
Research article**A numerical approach to approximate the solution of a quasilinear singularly perturbed parabolic convection diffusion problem having a non-smooth source term****Ruby¹, Vembu Shanthi¹ and Higinio Ramos^{2,3,*}**¹ Department of Mathematics, National Institute of Technology, Tiruchirappalli, Tamilnadu, India² Scientific Computing Group, University of Salamanca, 49029 Zamora, Spain³ Escuela Politécnica Superior de Zamora, Universidad de Salamanca, Avda. Requejo 33, 49029 Zamora, Spain*** Correspondence:** Email: higra@usal.es.

Abstract: The objective of the present paper is to solve a one-dimensional quasilinear parabolic singularly perturbed problem with a discontinuous source term. Due to the presence of such a discontinuity, an interior layer exists at the location of the discontinuity. The problem is solved by discretizing the spatial variable on a piecewise uniform Shishkin mesh using the standard upwind approach, while the backward Euler scheme is employed on a uniform mesh to discretize the time variable. The method is ε -uniformly convergent, providing first-order convergence in the time domain and almost first-order convergence in the spatial variable. To validate the theoretical findings, the scheme was tested by numerically solving two examples.

Keywords: convection diffusion; quasilinear; singular perturbation; Shishkin mesh; standard upwind scheme; backward Euler method; discontinuous source term; weak interior layer

Mathematics Subject Classification: 35B25, 65N06, 65N12, 65N15

1. Introduction

This work focuses on a quasilinear singularly perturbed convection diffusion parabolic initial-boundary value problem, where the source term exhibits a spatial discontinuity. The nonlinear convection coefficient, which characterizes the problem as quasilinear, is the novel aspect of the paper. Here, the source term has a jump discontinuity in space resulting in a weak interior layer appearing near the discontinuity, in addition to the presence of a boundary layer. The aim of this study is to propose a robust ε -uniform method, achieved by discretizing the spatial variable using the standard upwind

scheme on a piecewise Shishkin mesh and the temporal variable using the backward Euler scheme on a uniform mesh. The choice of the standard upwind scheme is motivated by its computational efficiency and simplicity, making it a suitable baseline solution for future comparisons with more advanced techniques. Problems of this type can be found in [1]. For studies on singular perturbation, one can refer to [2–5], in which various numerical and asymptotic methods are discussed. Regarding singularly perturbed problems with discontinuous terms, among all existing studies in the literature, we will refer only to a few related to our problem. The problem in [6] is time independent and has a strong interior layer because of the discontinuity in the source term and the convection coefficient. The problem was solved using the standard upwind scheme. In [7], a quasilinear convection-diffusion problem with discontinuous data resulting in strong interior layers was considered. Recently, a semilinear reaction-diffusion parabolic system with a discontinuous source term was studied in [8], while a system of quasilinear convection-diffusion problems with discontinuities in both the source term and the convection coefficient was investigated in [9]. Singularly perturbed problems occur in numerous scientific and engineering fields, including fluid dynamics [10, 11], chemical kinetics, population dynamics, control theory [12], and semiconductor modeling [13, 14].

Model problem

Let us first introduce some important notations:

$$\mathbb{D}_s = (0, 1), \quad \bar{\mathbb{D}}_s = [0, 1], \quad \zeta \in \mathbb{D}_s, \quad \mathbb{D}_s^- = (0, \zeta), \quad \mathbb{D}_s^+ = (\zeta, 1), \quad \mathbb{D}_t = (0, T], \quad \mathbb{D} = \mathbb{D}_s \times \mathbb{D}_t,$$

$$\bar{\mathbb{D}} = \bar{\mathbb{D}}_s \times \mathbb{D}_t, \quad \mathbb{D}^- = \mathbb{D}_s^- \times \mathbb{D}_t, \quad \mathbb{D}^+ = \mathbb{D}_s^+ \times \mathbb{D}_t, \quad \partial\mathbb{D} = \bar{\mathbb{D}} \setminus \mathbb{D}.$$

A function $u(s, t)$ is said to be Hölder continuous of order λ on Ω , with $\lambda \in (0, 1]$, if and only if $u \in C^0(\Omega)$ and

$$\sup_{(s_1, t_1), (s_2, t_2) \in \Omega} \frac{|u(s_1, t_1) - u(s_2, t_2)|}{((s_1 - s_2)^2 + |t_1 - t_2|)^{\lambda/2}}$$

is finite. This is denoted as $u \in C^\lambda(\Omega)$ (for more details, one can refer to [15]). For every integer $n \geq 0$,

$$u \in C^{n+\lambda}(\Omega) \text{ if } \frac{\partial^{k+m} u}{\partial s^k \partial t^m} \in C^\lambda(\Omega), \quad 0 \leq k + 2m \leq n.$$

Now, consider the following quasilinear singularly perturbed parabolic initial-boundary value problem where $z \in C^{1+\lambda}(\mathbb{D}) \cap C^{2+\lambda}(\mathbb{D}^- \cup \mathbb{D}^+)$:

$$Lz \equiv \varepsilon \frac{\partial^2 z}{\partial s^2}(s, t) + p(s, z(s, t)) \frac{\partial z}{\partial s}(s, t) - q(s)z(s, t) - r(s) \frac{\partial z}{\partial t}(s, t) = f(s, t), \quad (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \quad (1.1)$$

subject to the boundary and initial conditions

$$z(0, t) = z_l(t), \quad \forall t \in \mathbb{D}_t, \quad (1.2)$$

$$z(1, t) = z_r(t), \quad \forall t \in \mathbb{D}_t, \quad (1.3)$$

$$z(s, 0) = z_0(s), \quad \forall s \in \mathbb{D}_s. \quad (1.4)$$

It is assumed that the source term $f(s, t) \in C^{2+\lambda}(\mathbb{D}^- \cup \mathbb{D}^+)$, the convection coefficient $p(s, z) \in C^{2+\lambda}(\bar{\mathbb{D}})$, and the coefficients $q(s), r(s) \in C^{4+\lambda}(\bar{\mathbb{D}}_s)$ are smooth functions in their respective domains. The

perturbation parameter is assumed to be sufficiently small, such that $0 < \varepsilon \ll 1$. Furthermore, it is assumed that

$$\begin{cases} p(s, z) \geq \alpha > 0, \text{ on } \bar{\mathbb{D}}, \quad q(s) \geq \beta > 0, r(s) > 0 \text{ on } \bar{\mathbb{D}}_s, \\ |[f](\zeta, t)| \leq C, \end{cases} \quad (1.5)$$

where $s = \zeta$ is the line of discontinuity, and the source term has a jump discontinuity along $s = \zeta$. A weak interior layer appears on the right side of the discontinuity, and there is a boundary layer at $s = 0$. The initial and boundary conditions are assumed to be sufficiently smooth on $\bar{\mathbb{D}}$ and fulfill the compatibility conditions at the two corner points $(0, 0)$, $(1, 0)$ as follows:

$$\begin{cases} z_0(0) = z_l(0), \\ z_0(1) = z_r(0), \end{cases} \quad (1.6)$$

and

$$\begin{cases} \varepsilon \frac{\partial^2 z_0(0)}{\partial s^2} + p(0, z_0(0)) \frac{\partial z_0(0)}{\partial s} - q(0)z_0(0) - f(0, 0) = r(0) \frac{\partial z_l(0)}{\partial t}, \\ \varepsilon \frac{\partial^2 z_0(1)}{\partial s^2} + p(1, z_0(1)) \frac{\partial z_0(1)}{\partial s} - q(1)z_0(1) - f(1, 0) = r(1) \frac{\partial z_r(0)}{\partial t}. \end{cases} \quad (1.7)$$

Similarly, compatibility conditions are assumed to be fulfilled at the transition point $(\zeta, 0)$. Under these conditions, problem (1.1)–(1.5) has a unique solution $z \in C^{(1+\lambda)}(\mathbb{D}) \cap C^{(2+\lambda)}(\mathbb{D}^- \cap \mathbb{D}^+)$, as explained in [15–17].

This paper is organized as follows: In section 2, some analytical results are presented. Discretizations of the mesh and the solution are given in Section 3. Section 4 deals with error analysis, and in Section 5, two numerical examples are given to illustrate the performance of the proposed method. Finally, a conclusion of the paper is given in Section 6.

Throughout the paper, C denotes a generic constant that is independent of ε and the mesh parameters \mathcal{N} , \mathcal{M} . The maximum pointwise norm is used, denoted as $\|w\|_D = \max_{(s,t) \in D} |w(s, t)|$. When $D = \bar{\mathbb{D}}$, we will simply write $\|w\|$, avoiding the subscript.

2. Analytical results

In this section, some theoretical results are presented. Although the proofs of the following results are well known in the literature, we include them here for the sake of clarity and completeness.

Theorem 2.1. (*Minimum principle*): *Let $w \in C^0(\bar{\mathbb{D}}) \cap C^2(\mathbb{D}^- \cup \mathbb{D}^+)$ satisfy*

$$\begin{aligned} w(s, t) &\leq 0 \quad \text{on } \partial\mathbb{D}, \quad \left[\frac{\partial w}{\partial s} \right](\zeta, t) \geq 0 \quad t \in \mathbb{D}_t, \\ Lw(s, t) &\geq 0 \quad \text{for } (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+. \end{aligned}$$

Then,

$$w(s, t) \leq 0 \quad \text{for all } (s, t) \in \bar{\mathbb{D}}.$$

Proof. Consider a function ν , such that

$$w(s, t) = \begin{cases} \exp(-(\alpha_1(s - \zeta))/2\varepsilon)\nu(s, t), & s < \zeta, \\ \exp(-(\alpha_2(s - \zeta))/2\varepsilon)\nu(s, t), & s \geq \zeta, \end{cases}$$

where $\alpha \geq \alpha_1 \geq \alpha_2$. Let (s^*, t^*) be any point at which ν attains its maximum value in \mathbb{D} . If $\nu(s^*, t^*) \leq 0$, then there is nothing to prove. Now, we proceed by contradiction. Consider that $\nu(s^*, t^*) > 0$, then either $(s^*, t^*) \in (\mathbb{D}^- \cup \mathbb{D}^+)$ or $(s^*, t^*) = (\zeta, t^*)$. If $(s^*, t^*) \in (\mathbb{D}^- \cup \mathbb{D}^+)$, then

$$\nu_s = \nu_t = 0, \quad \nu_{ss} < 0.$$

Hence,

$$Lw \equiv \begin{cases} \exp(-(\alpha_1(s - \zeta))/2\varepsilon)\left(\varepsilon\nu_{ss} + (p - \alpha_1)\nu_s - \left(\frac{\alpha_1}{2\varepsilon}(p - \frac{\alpha_1}{2}) + q\right)\nu - r\nu_t\right) < 0, & s^* \leq \zeta, \\ \exp(-(\alpha_2(s - \zeta))/2\varepsilon)\left(\varepsilon\nu_{ss} + (p - \alpha_2)\nu_s - \left(\frac{\alpha_2}{2\varepsilon}(p - \frac{\alpha_2}{2}) + q\right)\nu - r\nu_t\right) < 0, & s^* \geq \zeta, \end{cases}$$

which is a contradiction.

If $(s^*, t^*) = (\zeta, t^*)$, then $[w_s](\zeta, t^*) = [\nu_s](\zeta, t^*) + \frac{\alpha_1 - \alpha_2}{2\varepsilon}\nu(\zeta, t) \leq 0$, which is also a contradiction. This completes the proof. \square

Theorem 2.2. *Let z be the solution of (1.1)–(1.5). Then, it holds:*

$$\|z\|_{\bar{\mathbb{D}}} \leq \frac{\|f\|_{\mathbb{D} \setminus (\zeta, t)}}{\alpha} + \max_{\partial\mathbb{D}} |z(s, t)|,$$

and

$$\|z^{(\kappa)}\|_{\mathbb{D} \setminus (\zeta, t)} \leq C\varepsilon^{-\kappa}, \quad 1 \leq \kappa \leq 3,$$

where the suffix (κ) denote the κ^{th} derivative of z with respect to s .

Proof. Consider the following barrier functions

$$\Psi^\pm(s, t) = -\frac{(1-s)\|f\|}{\alpha} - \max_{\partial\mathbb{D}} |z(s, t)| \pm z(s, t).$$

We note that $\Psi^\pm \in C^0(\bar{\mathbb{D}}) \cap C^2(\mathbb{D}^- \cup \mathbb{D}^+)$, and

$$\begin{aligned} \Psi^\pm(s, t) &\leq 0, \quad \text{on } \partial\mathbb{D}, \quad \left[\frac{\partial\Psi^\pm}{\partial s}\right](\zeta, t) \geq 0, \quad t > 0, \\ L\Psi^\pm(s, t) &\geq 0, \quad \text{for } (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+. \end{aligned}$$

Consequently, with the help of Theorem 2.1, the required result can be proved. By following a similar procedure as in [3], the bounds of the derivatives can be obtained. \square

Decomposition of the solution

We decompose the solution as the sum of a regular component x and a singular component y as:

$$z(s, t) = x(s, t) + y(s, t),$$

where the regular component $x \in C^0(\mathbb{D})$ satisfies the following problem

$$\begin{aligned} Lx &= f(s, t), \quad (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \\ x(s, 0) &= z(s, 0), \quad x(0, t) = z(0, t), \quad x(1, t) = z(1, t), \end{aligned} \quad (2.1)$$

and the singular component $y \in C^0(\mathbb{D})$ satisfies

$$\begin{aligned} Ly &= 0, \quad (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \\ y(s, 0) &= 0, \quad y(0, t) = z(0, t) - x(0, t), \quad y(1, t) = 0. \end{aligned} \quad (2.2)$$

Furthermore, the regular and singular components satisfy the conditions

$$[y](\zeta, t) = -[x](\zeta, t), \quad \left[\frac{\partial y}{\partial s} \right](\zeta, t) = -\left[\frac{\partial x}{\partial s} \right](\zeta, t). \quad (2.3)$$

The regular component x can be further decomposed as

$$x = x_0 + \varepsilon x_1 + \varepsilon^2 x_2, \quad (2.4)$$

where $x_0, x_1, x_2 \in C^0(\mathbb{D})$ and satisfy the following problems, respectively,

$$\begin{aligned} p(s, x_0) \frac{\partial x_0}{\partial s} + q(s)x_0 - r(s) \frac{\partial x_0}{\partial t} &= f(s, t), \quad (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \\ x_0(1, t) &= z(1, t), \quad x_0(s, 0) = z(s, 0), \end{aligned} \quad (2.5)$$

$$\begin{aligned} \varepsilon \frac{\partial^2 x_0}{\partial s^2} + p(s, x_0 + \varepsilon x_1) \frac{\partial(x_0 + \varepsilon x_1)}{\partial s} - q(x_0 + \varepsilon x_1) - r(s) \frac{\partial(x_0 + \varepsilon x_1)}{\partial t} &= f(s, t), \quad s \neq \zeta, \\ x_1(\zeta, t) &= x_1(1, t) = x_1(s, 0) = 0, \end{aligned} \quad (2.6)$$

$$\begin{aligned} \varepsilon \frac{\partial^2}{\partial s^2} (x_0 + \varepsilon x_1 + \varepsilon^2 x_2) + p(s, x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \frac{\partial}{\partial s} (x_0 + \varepsilon x_1 + \varepsilon^2 x_2) - q(s)(x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \\ - r(s) \frac{\partial}{\partial t} (x_0 + \varepsilon x_1 + \varepsilon^2 x_2) &= f(s, t), \end{aligned} \quad (2.7)$$

$$x_2(0, t) = x_2(\zeta, t) = x_2(1, t) = x_2(s, 0) = 0.$$

The singular component can be further decomposed as $y = y_1 + y_2$, where $y_1 \in C^2(\mathbb{D})$ is the boundary layer function satisfying

$$\begin{aligned} Ly_1 &= 0, \quad (s, t) \in \mathbb{D}, \\ y_1(0, t) &= z(0, t) - x(0, t), \quad y_1(1, t) = 0, \quad y_1(s, 0) = 0, \end{aligned} \quad (2.8)$$

and $y_2 \in C^0(\mathbb{D})$ is the weak interior layer function, which is the solution of

$$\begin{aligned} Ly_2 &= 0, (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \\ y_2(0, t) &= 0, y_2(1, t) = 0, y_2(s, 0) = 0, \\ \left[\frac{\partial y_2}{\partial s} \right](\zeta, t) &= - \left[\frac{\partial x}{\partial s} \right](\zeta, t). \end{aligned} \quad (2.9)$$

The method of upper and lower solutions will be useful to prove the existence and uniqueness of the regular component. By reversing all the inequalities, we can obtain the corresponding definition of lower solutions. In what follows, the upper and lower solutions will be denoted by Φ and ϕ , respectively.

Definition 2.3. *A function Φ is an upper solution of problem (2.1) if*

$$L\Phi \leq f(s, t), (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \quad (2.10)$$

$$\Phi(0, t) \geq z(0, t), \Phi(1, t) \geq z(1, t), \Phi(s, 0) \geq z(s, 0). \quad (2.11)$$

The solution of (2.1) satisfies Nagumo's condition on $\mathbb{D}^- \cup \mathbb{D}^+$ relative to the pair $\Psi(s, t) = \|f\| + \gamma s$, $\gamma := \sup_{\phi \leq z \leq \Phi} |p(s, z(s, t))|$. Thus, by using the result from [18, Thm 1.5.1], the existence of the regular component can be ensured by constructing lower and upper solutions. These components satisfy the bounds established in the following results.

Lemma 2.4. *For any integers κ, m satisfying $0 \leq \kappa+m \leq 2$, the solution x of (2.1) satisfies the following bounds:*

$$\begin{aligned} \|x\| &\leq C, (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \\ \left| \frac{\partial^{\kappa+m} x}{\partial s^\kappa \partial t^m}(s, t) \right| &\leq C(1 + \varepsilon^{2-\kappa}), (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+. \end{aligned}$$

Proof. Note that x_0 and its derivatives are independent of ε , and

$$\begin{aligned} \varepsilon \frac{\partial^2 x_0}{\partial s^2} + p(s, x_0 + \varepsilon x_1) \frac{\partial(x_0 + \varepsilon x_1)}{\partial s} - q(s)(x_0 + \varepsilon x_1) - r(s) \frac{\partial(x_0 + \varepsilon x_1)}{\partial t} &= f(s, t) \\ &= p(s, x_0) \frac{\partial x_0}{\partial s} + q(s)x_0 - r(s) \frac{\partial x_0}{\partial t}, \end{aligned}$$

$$x_1(0, t) = x_1(\zeta, t) = x_1(1, t) = 0,$$

which gives

$$\begin{aligned} \varepsilon \frac{\partial^2 x_0}{\partial s^2} + \varepsilon p(s, x_0 + \varepsilon x_1) \frac{\partial x_1}{\partial s} + (p(s, x_0 + \varepsilon x_1) - p(s, x_0)) \frac{\partial x_0}{\partial s} - \varepsilon q(s)x_1 - \varepsilon r(s) \frac{\partial x_1}{\partial t} &= 0, \\ \varepsilon \left(\frac{\partial^2 x_0}{\partial s^2} + p(s, x_0 + \varepsilon x_1) \frac{\partial x_1}{\partial s} + \frac{\partial x_0}{\partial s} \frac{\partial p(s, \hat{x})}{\partial x} x_1 - q(s)x_1 - r(s) \frac{\partial x_1}{\partial t} \right) &= 0, \end{aligned}$$

where $\hat{x}(s, t)$ is such that

$$p(s, x_0 + \varepsilon x_1) - p(s, x_0) = \varepsilon \frac{\partial p(s, \hat{x})}{\partial x} x_1.$$

Thus, the function x_1 satisfies the following problem

$$\begin{aligned} p(s, x_0 + \varepsilon x_1) \frac{\partial x_1}{\partial s} + \frac{\partial x_0}{\partial s} \frac{\partial p(s, \tilde{x})}{\partial x} x_1 - q(s)x_1 - r(s) \frac{\partial x_1}{\partial t} &= -\frac{\partial^2 x_0}{\partial s^2}, \\ x_1(1, t) &= x_0(s, 0) = 0. \end{aligned}$$

Hence, $|x_1| \leq C$ and $\left| \frac{\partial x_1}{\partial s} \right| \leq C$. Differentiating the above equation, we can also get the remaining bounds. Now, from Eqs (2.6) and (2.7), we obtain

$$\begin{aligned} &\varepsilon \frac{\partial^2}{\partial s^2} (x_0 + \varepsilon x_1 + \varepsilon^2 x_2) + p(s, x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \frac{\partial}{\partial s} (x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \\ &- q(s)(x_0 + \varepsilon x_1 + \varepsilon^2 x_2) - r(s) \frac{\partial}{\partial t} (x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \\ &= f(s, t) = \varepsilon \frac{\partial^2 x_0}{\partial s^2} + p(s, x_0 + \varepsilon x_1) \frac{\partial(x_0 + \varepsilon x_1)}{\partial s} - q(s)(x_0 + \varepsilon x_1) - r(s) \frac{\partial(x_0 + \varepsilon x_1)}{\partial t}. \end{aligned}$$

Hence, we get

$$\begin{aligned} &\varepsilon \frac{\partial^2 x_2}{\partial s^2} + p(s, x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \frac{\partial x_2}{\partial s} + \varepsilon^{-2} (p(s, x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \\ &- p(s, x_0 + \varepsilon x_1)) \frac{\partial(x_0 + \varepsilon x_1)}{\partial s} - q(s)x_2 - r(s) \frac{\partial x_2}{\partial t} = -\frac{\partial^2 x_1}{\partial s^2}, \end{aligned}$$

which gives

$$\varepsilon \frac{\partial^2 x_2}{\partial s^2} + p(s, x_0 + \varepsilon x_1 + \varepsilon^2 x_2) \frac{\partial x_2}{\partial s} + \left(\frac{\partial(x_0 + \varepsilon x_1)}{\partial s} \frac{\partial p(s, \tilde{x})}{\partial x} - q(s) \right) x_2 - r(s) \frac{\partial x_2}{\partial t} = -\frac{\partial^2 x_1}{\partial s^2},$$

$$x_2(0, t) = x_2(\zeta, t) = x_2(1, t) = x_2(s, 0) = 0.$$

where \tilde{x} is given by

$$p(s, x_0 + \varepsilon x_1 + \varepsilon^2 x_2) - p(s, x_0 + \varepsilon x_1) = \varepsilon^2 \frac{\partial p(s, \tilde{x})}{\partial x} x_2.$$

Hence, we have

$$\|x_2^{(\kappa)}\| \leq C\varepsilon^{-\kappa}.$$

Thus, by combining the bound of x_0 , x_1 , and x_2 , the proof is completed. \square

Lemma 2.5. *For any integers κ, m , the solutions y_1 and y_2 of problems (2.8) and (2.9), respectively, satisfy the following bounds:*

$$\begin{aligned} \left| \frac{\partial^{\kappa+m} y_1}{\partial s^\kappa \partial t^m}(s, t) \right| &\leq C\varepsilon^{-\kappa} \exp(-\alpha s/\varepsilon), \quad (s, t) \in \mathbb{D}, \quad 0 \leq \kappa + m \leq 3, \\ |y_2(s, t)| &\leq C\varepsilon, \quad (s, t) \in \mathbb{D}^- \cup \mathbb{D}^+, \\ \left| \frac{\partial^{\kappa+m} y_2}{\partial s^\kappa \partial t^m}(s, t) \right| &\leq \begin{cases} C\varepsilon^{1-\kappa} \exp(-\alpha s/\varepsilon), & (s, t) \in \mathbb{D}^-, \\ C\varepsilon^{1-\kappa} \exp(-\alpha(s - \zeta)/\varepsilon), & (s, t) \in \mathbb{D}^+. \end{cases} \quad 1 \leq \kappa + m \leq 3. \end{aligned}$$

Proof. To find the bounds of y_1 , choose C sufficiently large so that

$$|y_1(s, t)| \leq C \exp(-\alpha s/\varepsilon), \text{ for } (s, t) \in \partial\mathbb{D}.$$

Since $Ly_1 = 0$, and

$$\begin{aligned} L(C \exp(-\alpha s/\varepsilon)) &= C \exp(-\alpha s/\varepsilon)[\alpha^2 \varepsilon^{-1} - \alpha p(s, u) \varepsilon^{-1} - q(s)] \\ &\leq -C\beta \exp(-\alpha s/\varepsilon) \\ &\leq -Ly_1, \end{aligned}$$

then by the minimum principle [4], we get

$$|y_1(s, t)| \leq C \exp(-\alpha s/\varepsilon).$$

Now, taking the stretching variable $\tilde{s} = s/\varepsilon$ and following a procedure similar to that in [19], we can easily obtain bounds for the derivatives of y_1 .

To find the bounds for y_2 , we choose suitable barrier functions $\Psi(s, t) = -\psi(s, t) \pm y_2(s, t)$, where

$$\psi(s, t) = C \frac{\varepsilon}{\alpha} \begin{cases} 1, & (s, t) \in \mathbb{D}^-, \\ \exp(-(s - \zeta)\alpha/\varepsilon), & (s, t) \in \mathbb{D}^+. \end{cases}$$

Clearly, $\Psi(0, t) = -\psi(0, t) \leq 0$, $\Psi(1, t) = -\psi(1, t) \leq 0$, $\Psi(s, 0) = -\psi(s, 0) \leq 0$. Also,

$$L\Psi(s, t) = \begin{cases} Cq(s) \geq 0, & (s, t) \in \mathbb{D}^-, \\ C \exp(-(s - \zeta)\alpha/\varepsilon) \left(-\alpha + p(s, \psi) + \frac{\varepsilon}{\alpha} q \right) \geq 0, & (s, t) \in \mathbb{D}^+, \end{cases}$$

and $\left[\frac{\partial \Psi}{\partial s} \right](\zeta, t) = -\left[\frac{\partial \psi}{\partial s} \right](\zeta, t) \pm \left[\frac{\partial y_2}{\partial s} \right](\zeta, t) \geq 0$. Hence, Theorem 2.1 implies $\Psi(s, t) \leq 0$, for all $(s, t) \in \bar{\mathbb{D}}$, that is,

$$-\psi(s, t) \pm y_2(s, t) \leq 0 \implies |y_2(s, t)| \leq \psi(s, t) \leq C\varepsilon \text{ for } (s, t) \in \bar{\mathbb{D}}.$$

Hence, with the help of Theorem 2.1, we arrive at the required result.

Now, to get the derivative bounds, we follow a similar procedure as in [20]. The result is proven for \mathbb{D}^- and a similar procedure can be adopted for \mathbb{D}^+ . Take $\eta = (s - \zeta)/\varepsilon$ and set $\tilde{\mathbb{D}} = (0, \zeta\varepsilon^{-1}) \times (0, 1]$. Additionally, let $\tilde{y}_2(\eta, t) = y_2(s, t)$, and similarly define the transformed coefficients $\tilde{p}, \tilde{q}, \tilde{r}$. Then, Eq (2.9) becomes

$$(\tilde{y}_2)_{\eta\eta} + \tilde{p}(\eta, \tilde{y}_2)(\tilde{y}_2)_\eta - \varepsilon\tilde{q}(\eta)\tilde{y}_2 - \varepsilon\tilde{r}(\eta)(\tilde{y}_2)_t = 0, \text{ on } \tilde{\mathbb{D}}. \quad (2.12)$$

For each $\eta \in (0, \zeta\varepsilon^{-1})$ and each $\mu > 0$, let $R_{\eta, \mu}$ denote the rectangle

$$((\eta - \mu, \eta + \mu) \times (0, T]) \cap \tilde{\mathbb{D}},$$

and let $\bar{R}_{\eta, \mu}$ denote the closure of $R_{\eta, \mu}$ in the (η, t) -plane.

Now, since \tilde{y}_2 satisfies (2.9) and (2.12), by [15], for $0 \leq \kappa + 2m \leq 4$,

$$\max_{\tilde{R}_{\eta,1}} |D_{\eta}^{\kappa} D_t^m \tilde{y}_2| \leq C \max_{\tilde{R}_{\eta,2}} |\tilde{y}_2|. \quad (2.13)$$

Using bounds for y_2 , we have

$$|D_{\eta}^{\kappa} D_t^m \tilde{y}_2(\eta, \mu)| \leq C\varepsilon \exp\{-\alpha\eta\}.$$

Changing variables, this becomes

$$|D_s^{\kappa} D_t^m y_2(s, t)| \leq C\varepsilon^{1-\kappa} \exp(-\alpha s/\varepsilon).$$

□

3. Discretization of the problem

3.1. Mesh discretization

Let us subdivide the domain $\bar{\mathbb{D}}_s$ into

$$[0, \theta_1] \cup [\theta_1, \zeta] \cup [\zeta, \zeta + \theta_2] \cup [\zeta + \theta_2, 1], \quad (3.1)$$

where the transition points are defined as

$$\theta_1 = \min\left\{\frac{\zeta}{2}, \frac{2\varepsilon}{\alpha} \ln \mathcal{N}\right\}, \quad \theta_2 = \min\left\{\frac{1-\zeta}{2}, \frac{2\varepsilon}{\alpha} \ln \mathcal{N}\right\}. \quad (3.2)$$

We will divide each of the four intervals in $\mathcal{N}/4$ subintervals. The mesh points are denoted by

$$\mathbb{D}_s^{\mathcal{N}} = \mathbb{D}_s^{\mathcal{N}-} \cup \mathbb{D}_s^{\mathcal{N}+} = \{s_i : 1 \leq i \leq \mathcal{N}/2 - 1\} \cup \{s_i : \mathcal{N}/2 + 1 \leq i \leq \mathcal{N} - 1\}.$$

Clearly, $\bar{\mathbb{D}}_s^{\mathcal{N}} = \{s_i\}_{i=0}^{\mathcal{N}}$ where $s_0 = 0$, $s_{\mathcal{N}/2} = \zeta$ and $s_{\mathcal{N}} = 1$. The mesh width will be

$$h_i = \begin{cases} H_l = \frac{4\theta_1}{\mathcal{N}}, & i = 1, 2, \dots, \mathcal{N}/4, \\ h_l = \frac{4(\zeta - \theta_1)}{\mathcal{N}}, & i = \mathcal{N}/4 + 1, \mathcal{N}/4 + 2, \dots, \mathcal{N}/2, \\ H_r = \frac{4\theta_2}{\mathcal{N}}, & i = \mathcal{N}/2 + 1, \mathcal{N}/2 + 2, \dots, 3\mathcal{N}/4, \\ h_r = \frac{4(1 - \zeta - \theta_2)}{\mathcal{N}}, & i = 3\mathcal{N}/4 + 1, 3\mathcal{N}/4 + 2, \dots, \mathcal{N}. \end{cases}$$

Additionally, define $H_i = \frac{h_i + h_{i+1}}{2}$.

The time variable is discretized using a uniform mesh:

$$\mathbb{D}_t^{\mathcal{M}} = \left\{t_j = j \frac{T}{\mathcal{M}}\right\}, \quad j = 0, 1, 2, \dots, \mathcal{M}.$$

The piecewise uniform mesh is $\mathbb{D}^{\mathcal{N}, \mathcal{M}} = \mathbb{D}_s^{\mathcal{N}} \times \mathbb{D}_t^{\mathcal{M}}$. Let us also define some useful discrete domains $\bar{\mathbb{D}}^{\mathcal{N}, \mathcal{M}} = \bar{\mathbb{D}}_s^{\mathcal{N}} \times \mathbb{D}_t^{\mathcal{M}}$, $\partial\mathbb{D}^{\mathcal{N}, \mathcal{M}} = \bar{\mathbb{D}}^{\mathcal{N}, \mathcal{M}} \setminus \mathbb{D}^{\mathcal{N}, \mathcal{M}}$, $\mathbb{D}^{\mathcal{N}, \mathcal{M}-} = \mathbb{D}_s^{\mathcal{N}-} \times \mathbb{D}_t^{\mathcal{M}}$, $\mathbb{D}^{\mathcal{N}, \mathcal{M}+} = \mathbb{D}_s^{\mathcal{N}+} \times \mathbb{D}_t^{\mathcal{M}}$.

3.2. Discretization of the problem

The forward difference D_s^+ and backward difference D_s^- operators in space and the backward difference operator D_t^- in time are given by

$$D_s^+ Z_i^j = \frac{Z_{i+1}^j - Z_i^j}{h_{i+1}}, \quad D_s^- Z_i^j = \frac{Z_i^j - Z_{i-1}^j}{h_i},$$

$$D_t^- Z_i^j = \frac{Z_i^j - Z_i^{j-1}}{k}.$$

A difference operator of second order in space is described as:

$$\delta_s^2 Z_i^j = \frac{D_s^+ Z_i^j - D_s^- Z_i^j}{H_i}.$$

Using these operators, we obtain a discretization of problem (1.1)–(1.4) as:

$$\begin{aligned} Z_i^0 &= z_0(s_i), \quad \text{for } i = 1, \dots, \mathcal{N}. \\ \begin{cases} L_N^M Z_i^j \equiv \varepsilon \delta_s^2 Z_i^j + p_i^j D_s^+ Z_i^j - q_i Z_i^j - r_i D_t^- Z_i^j = f_i^j, & s_i \in \mathbb{D}_s^N, \\ D_s^+ Z_{N/2}^j = D_s^- Z_{N/2}^j, \\ Z_0^j = z_l(t_j), \quad Z_N^j = z_r(t_j), \end{cases} \\ &\quad \text{for } j = 1, 2, \dots, \mathcal{M}, \end{aligned} \quad (3.3)$$

where Z_i^j denotes an approximate value of $z(s_i, t_j)$, $p_i^j = p(s_i, Z_i^j)$, $q_i = q(s_i)$, $r_i = r(s_i)$, $f_i^j = f(s_i, t_j)$.

The finite difference operator L_N^M satisfies the following discrete minimum principle:

Theorem 3.1. (Discrete minimum principle) Let W be a mesh function defined on the discretized domain $\bar{\mathbb{D}}^{N,M}$. If $W \leq 0$ on $\partial\mathbb{D}^{N,M}$, $L_N^M W_i^j \geq 0$ on \mathbb{D}_s^N , and $D_s^+ W_{N/2}^j - D_s^- W_{N/2}^j \geq 0$, for $j = 1, 2, \dots, \mathcal{M}$, then $W_i^j \leq 0$ for all $(s_i, t_j) \in \bar{\mathbb{D}}^{N,M}$.

Proof. By following the methodology in [21], we proceed by contradiction. Let us assume that (s_n, t_m) is a point where W attains its maximum. If $W_n^m \leq 0$, then there is nothing to prove. Now, suppose that $W_n^m > 0$. Clearly, $n \neq 0, \mathcal{N}$, and hence, either $s_n \in \mathbb{D}_s^{N,-} \cup \mathbb{D}_s^{N,+}$ or $n = \mathcal{N}/2$. First consider the case $s_n \in \mathbb{D}_s^{N,-} \cup \mathbb{D}_s^{N,+}$. Then, $W_n^m - W_{n-1}^m \geq 0$, $W_{n+1}^m - W_n^m \leq 0$ and $W_n^m - W_{n-1}^m \geq 0$. Hence,

$$\begin{aligned} L_N^M W_n^m &= \varepsilon \delta_s^2 W_n^m + p_n^{m-1} D_s^+ W_n^m - q_n W_n^m - r_n D_t^- W_n^m \\ &= \frac{\varepsilon}{H_n} \left(\frac{W_{n+1}^m - W_n^m}{h_{n+1}} - \frac{W_n^m - W_{n-1}^m}{h_n} \right) + p_n^{m-1} \frac{W_{n+1}^m - W_n^m}{h_{n+1}} - q_n W_n^m - r_n \frac{W_n^m - W_{n-1}^m}{k} \\ &\leq 0, \end{aligned}$$

which is a contradiction. Now, the only possibility is $n = \mathcal{N}/2$, and we have

$$D_s^+ W_{N/2}^m - D_s^- W_{N/2}^m = \frac{W_{N/2+1}^m - W_{N/2}^m}{h_{N/2+1}} - \frac{W_{N/2}^m - W_{N/2-1}^m}{h_{N/2}} \leq 0,$$

which is again a contradiction. Hence, the proof is completed. \square

4. Error estimate

To get a bound of the nodal error $|(\mathbf{Z} - \mathbf{z})(s_i, t_j)|$, we will proceed as follows. First, we define the regular component X and the singular component Y of the discretized solution \mathbf{Z} . Later, the nodal errors are considered outside and within the layer using these mesh functions. We assume $\mathcal{M} = CN$ throughout the rest of the paper.

Define the discrete regular component X to be the solution of the following system:

$$\begin{aligned} L_N^M X &= f_i^j, \quad \text{for } i, j \in \mathbb{D}^{N, M}, \\ X(0, t_j) &= x(0, t_j), \quad X(\zeta, t_j) = x(\zeta, t_j), \quad X(1, t_j) = x(1, t_j), \quad X(s_i, 0) = x(s_i, 0). \end{aligned} \quad (4.1)$$

Define the discrete singular components Y to be the solution of the following system:

$$\begin{aligned} L_N^M Y &= 0, \quad \text{for } (s_i, t_j) \in \mathbb{D}^{N, M}, \quad Y(s_i, 0) = 0, \\ Y(0, t_j) &= y(0, t_j), \quad Y(1, t_j) = y(1, t_j), \\ [D_s Y(\zeta, t_j)] &= -[D_s X(\zeta, t_j)], \quad t_j \in \mathbb{D}^M, \end{aligned} \quad (4.2)$$

where for any mesh function W , the jump discontinuity along the line $s = \zeta$ is represented by:

$$[D_s W(\zeta, t_j)] = D_s^+ W(\zeta, t_j) - D_s^- W(\zeta, t_j).$$

The singular component can be further decomposed as $Y = Y_1 + Y_2$ where Y_1 is the boundary layer function satisfying

$$\begin{aligned} L_N^M Y_1 &= 0, \quad (s_i, t_j) \in \mathbb{D}^{N, M}, \\ Y_1(0, t_j) &= y(0, t_j) - x(0, t_j), \quad Y_1(1, t_j) = 0, \quad Y_1(s_i, 0) = 0, \end{aligned} \quad (4.3)$$

and Y_2 is the weak interior layer function, which solves the problem

$$\begin{aligned} L_N^M Y_2 &= 0, \quad (s_i, t_j) \in \mathbb{D}^{N, M}, \\ Y_2(0, t_j) &= 0, \quad Y_2(1, t_j) = 0, \quad Y_2(s_i, 0) = 0, \\ [D_s Y_2(\zeta, t_j)] &= -[D_s X(\zeta, t_j)] - [D_s Y_1(\zeta, t_j)]. \end{aligned} \quad (4.4)$$

So, the discretized solution can be expressed as

$$\mathbf{Z} = \mathbf{X} + \mathbf{Y} = \mathbf{X} + \mathbf{Y}_1 + \mathbf{Y}_2.$$

Lemma 4.1. *The regular component X of the discrete solution satisfies the following ε -uniform error estimate:*

$$|X(s_i, t_j) - x(s_i, t_j)| \leq \begin{cases} CN^{-1}(\zeta - s_i), & s_i \leq \zeta, \\ CN^{-1}(1 - s_i), & s_i \geq \zeta, \end{cases} \quad \text{where } t_j \in \mathbb{D}_t^M.$$

Proof. Let us first consider the case $s_i \leq \zeta$, we have

$$\begin{aligned} &\{\varepsilon \delta_s^2 + p(s_i, X) D_s^+ - q_i - r_i D_t^-\}(X - x) \\ &= \varepsilon \frac{\partial^2 x}{\partial s^2} + p(s_i, x) \frac{\partial x}{\partial s} - q(s_i)x - r(s_i) \frac{\partial x}{\partial t} - \{\varepsilon \delta_s^2 + p(s_i, X) D_s^+ - q_i - r_i D_t^-\}(x) \end{aligned}$$

$$\begin{aligned}
&= \varepsilon \left(\frac{\partial^2 x}{\partial s^2} - \delta_s^2 x \right) + p(s_i, X) \left(\frac{\partial x}{\partial s} - D_s^+ x \right) + (p(s_i, x) - p(s_i, X)) \frac{\partial x}{\partial s} - r(s_i) \left(\frac{\partial x}{\partial t} - D_t^- (x) \right) \\
&= \varepsilon \left(\frac{\partial^2 x}{\partial s^2} - \delta_s^2 x \right) + p(s_i, X) \left(\frac{\partial x}{\partial s} - D_s^+ x \right) + p_z(s_i, \xi_i^j) (x - X) \frac{\partial x}{\partial s} - r(s_i) \left(\frac{\partial x}{\partial t} - D_t^- (x) \right).
\end{aligned}$$

Introduce the linear difference operator

$$\mathcal{L}_x^{\mathcal{N}, \mathcal{M}} W := (\varepsilon \delta_s^2 + p(s_i, X) D_s^+ + p_z(s, \xi_i^j) \frac{\partial x}{\partial s} - r(s_i) D_t^-) W,$$

where ξ_i^j is defined implicitly by

$$p(s_i, x) - p(s_i, X) \equiv p_z(s_i, \xi_i^j) (x - X).$$

Note that $\|p_z(s, \xi_i^j) \frac{\partial x}{\partial s}\| \leq C$. Now,

$$\mathcal{L}_x^{\mathcal{N}, \mathcal{M}} (X - x) = \varepsilon \left(\delta_s^2 x - \frac{\partial x^2}{\partial^2 s} \right) + p(s_i, X) \left(D_s^+ x - \frac{\partial x}{\partial s} \right) + p_z(s_i, \xi_i^j) \frac{\partial x}{\partial s} (X - x) - r(s_i) D_t^- (X - x).$$

Then by using standard local truncation error estimates and by Lemma 2.4, we have

$$|\mathcal{L}_x^{\mathcal{N}, \mathcal{M}} (X - x)|(s_i, t_j) \leq \frac{\varepsilon}{3} (s_{i+1} - s_{i-1}) \left\| \frac{\partial x^3}{\partial^3 s} \right\| + \frac{p(s_i, X)}{2} (s_{i+1} - s_i) \left\| \frac{\partial x^2}{\partial^2 s} \right\| \leq C \mathcal{N}^{-1}.$$

Using the barrier functions $\psi^\pm(s, t) = -C \mathcal{N}^{-1} (\zeta - s_i) \pm (X - x)(s_i, t_j)$ and the discrete minimum principle, the proof can be completed in the usual way. A similar procedure can be used for $s_i \geq \zeta$, using a suitable barrier function and Theorem 3.1. \square

Lemma 4.2. *For $(s_i, t_j) \in \mathbb{D}^{\mathcal{N}, \mathcal{M}}$, the boundary layer function Y_1 satisfies*

$$|Y_1 - y_1| \leq C \mathcal{N}^{-1} \ln \mathcal{N}. \quad (4.5)$$

Additionally, for $s_i \geq \theta_1$, it holds that $|Y_1| \leq C \mathcal{N}^{-1}$.

Proof. Let us first consider the case $s \leq \zeta$. In the case when $\theta_1 \leq 1/4$ for the region away from the layer $[\theta_1, \zeta]$, using Lemma 2.5, we have

$$|y_1(s, t)| \leq C \exp(-\alpha s / \varepsilon) \leq C \exp(-\ln \mathcal{N}) = C \mathcal{N}^{-1}.$$

Now, consider the following transformation

$$Y_1 = \Lambda(s_i, t_j) \hat{Y}_1,$$

where $|\Lambda(s_i, t_j)| \leq C$ so that, for sufficiently large $\mathcal{N} \geq \mathcal{N}_0$ (\mathcal{N}_0 is independent of ε) and for sufficiently small $\varepsilon \leq \varepsilon_0$ (ε_0 is independent \mathcal{N})

$$\hat{L}_N^{\mathcal{M}} \hat{Y}_1 = \left\{ \varepsilon \delta_s^2 + \hat{p} D_s^- - \hat{q} - \hat{r} D_t^- \right\} \hat{Y}_1, \quad \hat{p} \geq \alpha, \quad \hat{q} \geq \beta, \quad \hat{r} > 0,$$

and $\hat{Y}_1(0, t) = |y_1(0, t)|$, $\hat{Y}_1(\zeta, t) = 0$. Now, we take

$$\Lambda(s_i, t_j) = \begin{cases} \prod_{k=1}^i (1 + \xi_1 \varepsilon^{-1} \mathcal{N}^{-1} h_k)^{-1}, & k < \mathcal{N}/4, \\ (1 + \xi_1 \varepsilon^{-1} \mathcal{N}^{-1} h_k)^{-\mathcal{N}/2} \prod_{k=\mathcal{N}/4}^{\mathcal{N}/2} (1 + \xi_2 h_k)^{-1}, & \mathcal{N}/4 \leq k \leq \mathcal{N}/2, \end{cases}$$

where $\xi_2 > \xi_1 > 0$ are suitably chosen constants. Let \tilde{Y}_1 be the solution of

$$\begin{aligned} \{\varepsilon \delta_s^2 + \alpha D_s^+ - q_i - r_i D_t^-\} \tilde{Y}_1 &= 0, \\ \tilde{Y}_1(0, t) &= |y_1(0, t)|, \quad \tilde{Y}_1(\zeta, t) = 0. \end{aligned}$$

Then, using the discrete comparison principle, we have

$$|\hat{Y}_1| \leq \tilde{Y}_1 \leq C \mathcal{N}^{-1}.$$

Hence,

$$|Y_1(s_i, t_j) - y_1(s_i, t_j)| \leq C \mathcal{N}^{-1}.$$

Now, for the layer region $[0, \theta_1]$, we have

$$\begin{aligned} L_{\mathcal{N}}^{\mathcal{M}}(Y_1 - y_1) &= L_{\mathcal{N}}^{\mathcal{M}} y_1 - L y_1 \\ &= \varepsilon (\delta_s^2 y_1 - (y_1)_{ss}) + p(s, Z) D_s^+ y_1 - p(s, z) (y_1)_s - q(s) (Y_1 - y_1) - r(s) (D_t^- y_1 - (y_1)_t) \\ &= \varepsilon (\delta_s^2 y_1 - (y_1)_{ss}) + (p(s, Z) - p(s, z)) (y_1)_s + p(s, Z) (D_s^+ y_1 - (y_1)_s) - r(s) (D_t^- y_1 - (y_1)_t) \\ &= \varepsilon (\delta_s^2 y_1 - (y_1)_{ss}) + (p(s, \xi_i^j) (Z - z)) (y_1)_s + p(s, Z) (D_s^+ y_1 - (y_1)_s) - r(s) (D_t^- y_1 - (y_1)_t). \end{aligned}$$

Let us define the linear discrete operator

$$L_{y_1}^{\mathcal{N}, \mathcal{M}} W \equiv (\varepsilon \delta_s^2 + p(s, \xi_i^j) (y_1)_s + p(s, Z) D_s^+ - r(s) D_t^-) W,$$

where ξ_i^j is given implicitly by $p_s(s, \xi_i^j) (Y_1 - y_1) \equiv p(s, Y_1) - p(s, y_1)$. Proving that

$$p(s, Z) - 4\varepsilon (p_s(s, \xi_i^j) (y_1)_s) > 0,$$

this linear discrete operator satisfies the maximum principle.

Now, adopting the methodology as in [3], using a suitable barrier function \bar{Y}_1 , which is the solution of the problem:

$$(\varepsilon \delta_s^2 + p(s, \xi_i^j) (y_1)_s + p(s, Z) D_s^+ - r(s) D_t^-) \bar{Y}_1 = 0, \quad (s_i, t_j) \in \mathbb{D}^-,$$

$$\bar{Y}_1(0, t_j) = 1, \quad \bar{Y}_1(\zeta, t_j) = 0, \quad \bar{Y}_1(s_i, 0) = 0,$$

and a discrete comparison principle, the proof is complete. Similarly, bounds for the case $s > \zeta$ can be proved easily. \square

Lemma 4.3. *The following ε -uniform bound holds:*

$$|[D_s Y_2(\zeta, t_j)]| \leq C(1 + \varepsilon^{-1} \mathcal{N}^{-1}).$$

Proof. Since

$$[D_s Y_2(\zeta, t_j)] = -[D_s X(\zeta, t_j)] - [D_s Y_1(\zeta, t_j)] = -(D_s^+ X(\zeta, t_j) - D_s^- X(\zeta, t_j)) - (D_s^+ Y_1(\zeta, t_j) - D_s^- Y_1(\zeta, t_j)),$$

it is enough to prove the bounds for each term in the right side of the equation. Now,

$$D_s^- X(\zeta, t_j) = D_s^-(X - x)(\zeta, t_j) + D_s^- x(\zeta, t_j)$$

and $\left\| \frac{\partial x}{\partial s} \right\|_{\mathbb{D}^-} \leq C$, which implies $|D_s^- x(\zeta, t_j)| \leq C$ and $|D_s^-(X - x)(\zeta, t_j)| = |(X - x)(\zeta - h_l, t_j)/h_l| \leq C\mathcal{N}^{-1}$ by Lemma 4.1. Hence,

$$[D_s^- X(\zeta, t_j)] \leq C. \quad (4.6)$$

Now, when $s_i \geq \zeta$,

$$D_s^+ X(s_i, t_j) = D_s^+(X - x)(s_i, t_j) + D_s^+ x(s_i, t_j),$$

and $\left\| \frac{\partial x}{\partial s} \right\|_{\mathbb{D}^+} \leq C$.

Also by applying Lemma 3.14 of [3], for $s_i \geq \zeta$, we get

$$\varepsilon |D_s^+(X - x)(s_i, t_j)| \leq C\mathcal{N}^{-1}, \quad (4.7)$$

which gives

$$D_s^+ X(s_i, t_j) \leq C\mathcal{N}^{-1}/\varepsilon + C = C(1 + \varepsilon^{-1}\mathcal{N}^{-1}). \quad (4.8)$$

Now, to get the bounds for Y_1 , note that $Y_1(s_i, t_j) \leq C\mathcal{N}^{-1}$, $s_i \geq \theta_1$, and hence

$$|D_s^- Y_1(\zeta, t_j)| \leq C. \quad (4.9)$$

Finally,

$$D_s^+ Y_1(\zeta, t_j) = D_s^+(Y_1 - y_1)(\zeta, t_j) + D_s^+ y_1(\zeta, t_j), \text{ and } \left\| \frac{\partial y_1}{\partial s} \right\|_{\mathbb{D}^+} \leq C\varepsilon^{-1} \exp(-\alpha\zeta/\varepsilon) \leq C.$$

Now, by following the arguments of Lemma 3.14 from [3] and using the bounds of the derivatives of y_1 , one can get

$$D_s^+(Y_1 - y_1)(s_i, t_j) \leq C, \quad s_i \geq \zeta. \quad (4.10)$$

Hence,

$$D_s^+ Y_1(\zeta, t_j) \leq C. \quad (4.11)$$

Now, by combining the Eqs (4.6), (4.8), (4.9), and (4.11), the proof is complete. \square

Concerning the interior layer function, the same bounds can be obtained, as stated in the following result.

Lemma 4.4. *For $(s_i, t_j) \in \mathbb{D}^{\mathcal{N}, \mathcal{M}}$, the interior layer function Y_2 satisfies*

$$|Y_2 - y_2| \leq C\mathcal{N}^{-1} \ln \mathcal{N}. \quad (4.12)$$

Proof. Let us first get the error estimates at the line of discontinuity $s = \zeta$. Recalling that $[x_s(\zeta, t)] + [(y_2)_s(\zeta, t)] = 0$, we obtain

$$\begin{aligned} [D_s(Y_2 - y_2)(\zeta, t)] &= [D_s Y_2(\zeta, t)] - [D_s y_2(\zeta, t)] = [x_s(\zeta, t)] + [(y_2)_s(\zeta, t)] + [D_s Y_2(\zeta, t)] - [D_s y_2(\zeta, t)] \\ &= [x_s(\zeta, t)] + [(y_2)_s(\zeta, t)] - [D_s X(\zeta, t_j)] - [D_s Y_1(\zeta, t_j)] - [D_s y_2(\zeta, t)] \\ &= ([x_s(\zeta, t)] - [D_s X(\zeta, t_j)]) + ([(y_2)_s(\zeta, t)] - [D_s y_2(\zeta, t)]) - [D_s Y_1(\zeta, t_j)]. \end{aligned} \quad (4.13)$$

Using [Lemma 3.14, [3]], we have

$$[x_s(\zeta, t)] - [D_s X(\zeta, t_j)] = x_s(\zeta^+, t) - D_s^+ X(\zeta, t_j) + x_s(\zeta^-, t) - D_s^- X(\zeta, t_j) = [D_s(x - X)(\zeta, t)] \leq C(\varepsilon N)^{-1}. \quad (4.14)$$

Also, it is

$$[(y_2)_s(\zeta, t)] - [D_s y_2(\zeta, t)] = (y_2)_s(\zeta^+, t) - D_s^+ Y_2(\zeta, t_j) + (y_2)_s(\zeta^-, t) - D_s^- Y_2(\zeta, t_j), \quad (4.15)$$

which implies

$$\begin{aligned} |[(y_2)_s(\zeta, t)] - [D_s y_2(\zeta, t)]| &\leq CH_r \left\| \frac{\partial y_2^2}{\partial^2 s} \right\|_{[\zeta, \zeta+H_r]} + Ch_l \left\| \frac{\partial y_2^2}{\partial^2 s} \right\|_{[\zeta-h_l, \zeta]} \\ &\leq C \frac{H_r}{\varepsilon} + C \frac{H_r}{\varepsilon} \exp(-\alpha/\varepsilon(\zeta - H_r)) \\ &\leq CH_r/\varepsilon + Ch_l \leq CN^{-1} \ln N. \end{aligned} \quad (4.16)$$

Finally,

$$|D_s y_1(\zeta, t)| \leq (h_l + H_r) |(y_1)_{ss}(\zeta - h_l)| \leq C \frac{h_l + H_r}{\varepsilon^2} \exp((-\alpha/\varepsilon)(\zeta - h_l)) \leq CN^{-1}.$$

Using inequalities (4.14)–(4.16) and Lemma 3.16 [3], we obtain from (4.13)

$$|[D_s(y_2 - Y_2)(\zeta, t_j)]| \leq C(\varepsilon N)^{-1} \ln N.$$

Consider the following transformation for the outer region $[\theta_1, \zeta]$,

$$Y_2(s_i, t_j) = \omega(s_i, t_j) \hat{Y}_2(s_i, t_j),$$

where $|\omega(s_i, t_j)| \leq C$ so that for $N \geq N_0$ and $\varepsilon \leq \varepsilon_0$ (where N_0 is independent of ε and ε_0 is independent of N),

$$\hat{\mathcal{L}}_N^M \hat{Y}_2 = \{\varepsilon \delta_s^2 + \hat{p} D_s^+ - \hat{q} - \hat{r} D_t^-\} \hat{Y}_2(s_i, t_j) = 0, \quad \hat{p} \geq \alpha, \quad \hat{q} \geq 0, \quad \hat{r} \geq 0,$$

$$\hat{Y}_2(0, t_j) = |Y_2(0, t_j)|, \quad \hat{Y}_2(1, t_j) = 0, \quad [D_s \hat{Y}_2(\zeta, t_j)] = -[D_s X(\zeta, t_j)] - [D_s Y_1(\zeta, t_j)], \quad \hat{Y}_2(s_i, 0) = 0.$$

By choosing a suitable barrier function $\omega(s_i, t_j)$ given by

$$\omega(s_i, t) = \begin{cases} \prod_{k=1}^i (1 + \hat{\xi}_1 \varepsilon^{-1} N^{-1} h_k)^{-1}, & k < N/4, \\ (1 + \hat{\xi}_1 \varepsilon^{-1} N^{-1} h_k)^{-N/2} \prod_{k=N/4}^{N/2} (1 + \hat{\xi}_2 h_k)^{-1}, & N/4 \leq k \leq N/2, \end{cases}$$

with $\hat{\xi}_2 > \hat{\xi}_1 > 0$ appropriate constants, the operator satisfies a comparison principle and we can get that

$$|\hat{Y}_2| \leq \tilde{Y}_2 \leq C\mathcal{N}^{-1},$$

where $\tilde{Y}_2(s_i, t_j)$ is the solution of the following problem

$$\{\varepsilon\delta_s^2 + \alpha D_s^+ - q - rD_t^-\}\tilde{Y}_2(s_i, t_j) = 0,$$

$$\tilde{Y}_2(0, t_j) = |y_2(0, t_j)|, \quad \tilde{Y}_2(1, t_j) = 0, \quad [D_s \tilde{Y}_2(\zeta, t_j)] = -[D_s X(\zeta, t_j)] - [D_s Y_1(\zeta, t_j)], \quad \tilde{Y}_2(s_i, 0) = 0.$$

Hence,

$$|Y_2(s_i, t_j) - y_2(s_i, t_j)| \leq C\mathcal{N}^{-1}.$$

Now, let us consider the truncation error in the layer regions:

$$\begin{aligned} L_N^M(Y_2 - y_2) &= Ly_2 - L_N^M y_2 \\ &= \varepsilon \left(\frac{\partial^2 y_2}{\partial s^2} - \delta_s^2 y_2 \right) + p(s_i, y_2) \frac{\partial y_2}{\partial s} - p(s_i, Y_2) D_s^+ y_2 - r \left(\frac{\partial y_2}{\partial t} - D_t^- (y_2) \right) \\ &= \varepsilon \left(\frac{\partial^2 y_2}{\partial s^2} - \delta_s^2 y_2 \right) + (p(s_i, y_2) - p(s_i, Y_2)) \frac{\partial y_2}{\partial s} + (p(s, Y_2) \left(\frac{\partial y_2}{\partial s} - D_s^+ y_2 \right) - r \left(\frac{\partial y_2}{\partial t} - D_t^- (y_2) \right)) \\ &= \varepsilon \left(\frac{\partial^2 y_2}{\partial s^2} - \delta_s^2 y_2 \right) + p_z(s_i, \bar{\xi}_i^j) (y_2 - Y_2) \frac{\partial y_2}{\partial s} + (p(s, Y_2) \left(\frac{\partial y_2}{\partial s} - D_s^+ y_2 \right) - r \left(\frac{\partial y_2}{\partial t} - D_t^- (y_2) \right)), \end{aligned}$$

where $\bar{\xi}_i^j$ is defined implicitly by

$$p(s_i, y_2) - p(s_i, Y_2) \equiv p_z(s_i, \bar{\xi}_i^j) (y_2 - Y_2).$$

Note that $\|p_z(s, \xi_i^j) \frac{\partial x}{\partial s}\| \leq C$. Now, we define the linear discrete operator

$$\mathcal{L}_{y_2}^{N,M} W \equiv (\varepsilon\delta_s^2 + p_z(s_i, \bar{\xi}_i^j) \frac{\partial y_2}{\partial s} + p(s_i, Y_2) D_s^+ - r(s_i) D_t^-) W.$$

The operator satisfies a discrete comparison principle, providing that the inequality

$$p^2(s_i, Y_2) - 4\varepsilon \left(p_z(s_i, \bar{\xi}_i^j) \frac{\partial y_2}{\partial s} \right) > 0$$

holds. Now, the operator $\mathcal{L}_{y_2}^{N,M}$ follows the linear case as in [22]. With the help of a comparison principle and a suitable barrier function $B_{i,j}^\pm = C\mathcal{N}^{-1} + C\mathcal{N}^{-1} \ln \mathcal{N} \pm (Y_2 - y_2)(s_i, t_j)$, we get the required result, as usually. \square

Now, we will state the main theoretical result of this paper.

Theorem 4.5. *Let z be the solution of problem (1.1) and Z be the numerical solution of (3.3). Then, the following error estimates hold:*

$$\|Z - z\|_{\mathbb{D}} \leq C\mathcal{N}^{-1}(\ln \mathcal{N}),$$

and

$$\varepsilon \left\| D_s^+ Z - \frac{\partial z}{\partial s} \right\|_{\mathbb{D}} \leq C\mathcal{N}^{-1} \ln \mathcal{N}.$$

Proof. Combining Lemmas 4.1, 4.2, and 4.4, the proof is completed.

To derive the error bound for approximations to the scaled derivative, by applying the arguments from [3, Section 3.5], separately on each subdomain $[0, \zeta] \times (0, T]$ and $[\zeta, 1] \times (0, T]$, we get

$$\begin{aligned}\varepsilon \left\| D_s^+ z - \frac{\partial z}{\partial s} \right\|_{\bar{\mathbb{D}}} &\leq C\mathcal{N}^{-1} \ln \mathcal{N}, \\ \varepsilon \left\| D_s^+ X - \frac{\partial X}{\partial s} \right\|_{\bar{\mathbb{D}}} &\leq C\mathcal{N}^{-1} \ln \mathcal{N}, \\ \varepsilon \left\| D_s^+ Y_1 - \frac{\partial y_1}{\partial s} \right\|_{\mathbb{D}^-} &\leq C\mathcal{N}^{-1} \ln \mathcal{N}\end{aligned}$$

and from the proof of Lemma 4.3, we obtain

$$\varepsilon \left\| D_s^+ Y_1 - \frac{\partial y_1}{\partial s} \right\|_{\mathbb{D}^+} \leq C\mathcal{N}^{-1} \ln \mathcal{N}.$$

Now, again by applying the arguments from [3, Section 3.5], for $s_i \geq \zeta$, we can get the following bounds

$$\varepsilon \left\| D_s^+ Y_2 - \frac{\partial y_2}{\partial s} \right\|_{\mathbb{D}^+} \leq C\mathcal{N}^{-1} \ln \mathcal{N}.$$

For, $s_i < \zeta$, define $\tilde{Y}_2(s_i, t_j) = Y_2(s_i, t_j) - Y_2(\zeta, t_j)$ with

$$\begin{aligned}L_N^M \tilde{Y}_2 &= 0, \text{ for } (s_i, t_j) \in \mathbb{D}^{N,M}, \\ \tilde{Y}_2(0, t_j) &= -Y_2(\zeta, t_j), \quad \tilde{Y}_2(\zeta, t_j) = 0, \quad \tilde{Y}_2(1, t_j) = -Y_2(\zeta, t_j).\end{aligned}$$

Then, for $s_i < \zeta$,

$$D_s^+(Y_2 - y_2)(s_i, t_j) = D_s^+(Y_2 - y_2)(\zeta, t_j) + D_s^+(\tilde{Y}_2 - \tilde{y}_2)(s_i, t_j) \quad (4.17)$$

where \tilde{y}_2 is defined analogously to \tilde{Y}_2 . Now, by adopting the methodology from [3], we can bound the second term $D_s^+(\tilde{Y}_2 - \tilde{y}_2)(s_i, t_j)$ on the right side of Eq (4.17), whereas the first term $D_s^+(Y_2 - y_2)(\zeta, t_j)$ of Eq (4.17) is already bounded. \square

5. Numerical results

To validate the theoretical results obtained in the previous sections, the proposed scheme has been employed for solving two test problems. In order to get numerical solutions, we have linearized system (3.3) and consider the following system

$$\begin{aligned}\mathcal{Z}_i^0 &= z_0(s_i), \quad \text{for } i = 1, \dots, \mathcal{N}, \\ \begin{cases} L_N^M \mathcal{Z}_i^j \equiv \varepsilon \delta_s^2 \mathcal{Z}_i^j + p_i^{j-1} D_s^+ \mathcal{Z}_i^j - q_i \mathcal{Z}_i^j - r_i D_t^- \mathcal{Z}_i^j = f_i^j, & s_i \in \mathbb{D}_s, \\ D_s^+ \mathcal{Z}_{N/2}^j = D_s^- \mathcal{Z}_{N/2}^j \\ \mathcal{Z}_0^j = z_l(t_j), \quad \mathcal{Z}_{\mathcal{N}}^j = z_r(t_j), \end{cases} \\ &\text{for } j = 1, 2, \dots, \mathcal{M},\end{aligned} \quad (5.1)$$

where we used a linearization technique as the one described in [23, 24]. The above system can be expressed in matrix form as

$$A \mathbf{Z}^j = \mathbf{Z}^{j-1}, \quad j = 1, 2, \dots, \mathcal{M},$$

where A is the coefficient matrix, which is an M-matrix of order $(\mathcal{N} - 1) \times (\mathcal{N} - 1)$, and $\mathbf{Z}^j = \{\mathcal{Z}_1^j, \mathcal{Z}_2^j, \dots, \mathcal{Z}_{N-1}^j\}$, $j = 1, 2, \dots, \mathcal{M}$.

Using $\mathbf{Z}^0 = \{\mathcal{Z}_1^0, \mathcal{Z}_2^0, \dots, \mathcal{Z}_{N-1}^0\} = \{z_0(s_1), z_0(s_2), \dots, z_0(s_{N-1})\}$, $\mathcal{Z}_0^j = z_l(t_j)$, $\mathcal{Z}_N^j = z_r(t_j)$, $j = 1, 2, \dots, \mathcal{M}$, the above system can be solved for \mathbf{Z}^j .

The double mesh principle [3] is applied to approximate the errors in the maximum norm. For simplicity, we take $\mathcal{M} = \mathcal{N}$. The errors are obtained through

$$E_{\varepsilon}^{\mathcal{N}, \mathcal{M}} = \max_{(s, t) \in \mathbb{D}} |\mathcal{Z}_{2i, 2j}^{2\mathcal{N}, 2\mathcal{M}} - \mathcal{Z}_{i, j}^{\mathcal{N}, \mathcal{M}}|,$$

where $\mathcal{Z}^{\mathcal{N}, \mathcal{M}}$ and $\mathcal{Z}^{2\mathcal{N}, 2\mathcal{M}}$ represent the numerical solutions to problem (1.1) on two meshes with \mathcal{N} and $2\mathcal{N}$ number of subintervals, respectively. The finer mesh has the mesh points of the coarse mesh along with their midpoints. The maximum pointwise error norm is determined by

$$E^{\mathcal{N}, \mathcal{M}} = \max_{\varepsilon} E_{\varepsilon}^{\mathcal{N}, \mathcal{M}}.$$

Furthermore, the approximate orders of convergence are calculated using the standard formula:

$$Q^{\mathcal{N}, \mathcal{M}} = \log_2 \left(\frac{E^{\mathcal{N}, \mathcal{M}}}{E^{2\mathcal{N}, 2\mathcal{M}}} \right).$$

The errors and orders of convergence associated with time and space are calculated separately by fixing \mathcal{N} and \mathcal{M} , respectively. The errors $E_{\mathcal{N}}$, $E_{\mathcal{M}}$ and the maximum pointwise error norms $E^{\mathcal{N}}$, $E^{\mathcal{M}}$ are obtained as

$$E_{\varepsilon}^{\mathcal{N}} = \max_{(s, t) \in \mathbb{D}} |\mathcal{Z}^{2\mathcal{N}, \mathcal{M}} - \mathcal{Z}^{\mathcal{N}, \mathcal{M}}|, \quad E^{\mathcal{N}} = \max_{\varepsilon} E_{\varepsilon}^{\mathcal{N}}.$$

$$E_{\varepsilon}^{\mathcal{M}} = \max_{(s, t) \in \mathbb{D}} |\mathcal{Z}^{\mathcal{N}, 2\mathcal{M}} - \mathcal{Z}^{\mathcal{N}, \mathcal{M}}|, \quad E^{\mathcal{M}} = \max_{\varepsilon} E_{\varepsilon}^{\mathcal{M}}.$$

The estimate of orders of convergence for space and time are calculated, respectively, by

$$Q^{\mathcal{N}} = \log_2 \left(\frac{E^{\mathcal{N}}}{E^{2\mathcal{N}}} \right), \quad Q^{\mathcal{M}} = \log_2 \left(\frac{E^{\mathcal{M}}}{E^{2\mathcal{M}}} \right).$$

Example 5.1. Consider problem (1.1) with the following data

$$\begin{aligned} p(s, z) &= \sin(z) + 1, \quad q(s) = \sin^2(s) + 1, \quad r(s) = 1, \\ z_0(s) &= s, \quad z_l(t) = z_0(0) = 0, \quad z_r(t) = z_0(1) = 1, \\ f(s, t) &= \begin{cases} -\sin(s)^2 - \sin(t), & s < \zeta, \\ \sin(s)^2 + \sin(t), & s > \zeta, \end{cases} \quad \text{where } \zeta = 0.5. \end{aligned}$$

Example 5.2. Consider problem (1.1) with the following data

$$\begin{aligned} p(s, z) &= z + 1, \quad q(s) = \log(s + 2), \quad r(s) = 1, \\ z_0(s) &= \log(s + 2), \quad z_l(t) = z_0(0) = \log(2), \quad z_r(t) = z_0(1) = \log(3), \\ f(s, t) &= \begin{cases} -(\log(s + 2) + t)^3, & s < \zeta, \\ (\log(s + 2) + t)^3, & s > \zeta, \end{cases} \quad \text{where } \zeta = 0.7. \end{aligned}$$

For $\varepsilon = 2^{-7}$ and $\mathcal{N} = 32$, solution graphs are given for both problems. From Figures 1 and 2, one can clearly observe that apart from a boundary layer at $s = 0$, a weak interior layer exists at the right side of the line of discontinuity. For $\varepsilon = 2^{-6}, \dots, 2^{-25}$, $\mathcal{N} = 2^5, \dots, 2^{10}$, Tables 1 and 2 provide the pointwise errors and convergence rates for Examples 5.1 and 5.2 with $\mathcal{N} = \mathcal{M}$. To analyse the rate of convergence in the space and time variable, Tables 3 and 4 are presented with $\mathcal{N} = 256$, $\mathcal{M} = 2^5, \dots, 2^{10}$ and Tables 5 and 6 are presented with $\mathcal{M} = 128$, $\mathcal{N} = 2^5, \dots, 2^{10}$. These tables demonstrate an almost first-order convergence in space and a first-order convergence in time. The numerical results presented in these tables corroborate the theoretical findings and confirm that the method achieves nearly first-order convergence.

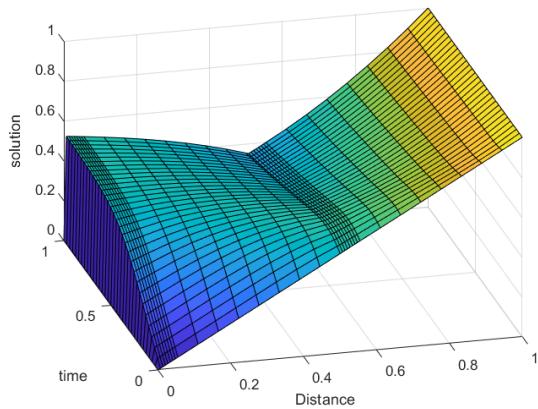


Figure 1. Solution graph for the value $\varepsilon = 2^{-7}$ and $\mathcal{N} = 32$ for Example 5.1.

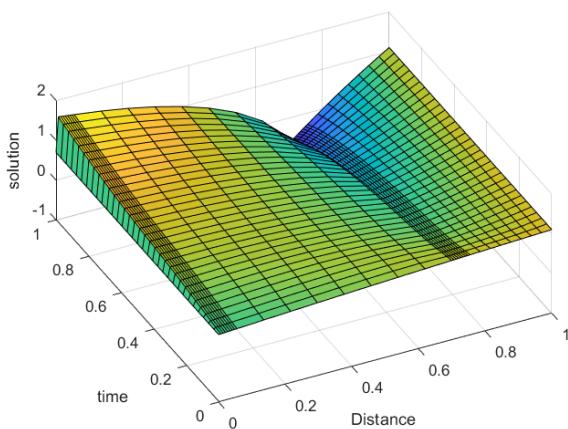


Figure 2. Solution graph for the value $\varepsilon = 2^{-7}$ and $\mathcal{N} = 32$ for Example 5.2.

Table 1. Maximum point-wise errors $E^{N,M}$ and orders of convergence $Q^{N,M}$ calculated for Example 5.1 for $N = M$.

ε	$M = 32$	$M = 64$	$M = 128$	$M = 256$	$M = 512$	$M = 1024$
	$N = 32$	$N = 64$	$N = 128$	$N = 256$	$N = 512$	$N = 1024$
2^{-6}	1.858e- 2	1.180e- 2	6.891e- 3	4.092e- 3	2.538e- 3	1.631e- 3
2^{-7}	2.013e- 2	1.308e- 2	7.730e- 3	4.550e- 3	2.788e- 3	1.767e- 3
2^{-8}	2.102e- 2	1.389e- 2	8.317e- 3	4.900e- 3	3.000e- 3	1.900e- 3
2^{-9}	2.149e- 2	1.433e- 2	8.649e- 3	5.108e- 3	3.129e- 3	1.982e- 3
2^{-10}	2.173e- 2	1.456e- 2	8.829e- 3	5.219e- 3	3.199e- 3	2.028e- 3
2^{-11}	2.185e- 2	1.468e- 2	8.922e- 3	5.276e- 3	3.235e- 3	2.051e- 3
2^{-12}	2.190e- 2	1.473e- 2	8.969e- 3	5.305e- 3	3.253e- 3	2.064e- 3
2^{-13}	2.193e- 2	1.476e- 2	8.992e- 3	5.319e- 3	3.262e- 3	2.070e- 3
2^{-14}	2.195e- 2	1.478e- 2	9.004e- 3	5.327e- 3	3.267e- 3	2.073e- 3
2^{-15}	2.196e- 2	1.479e- 2	9.010e- 3	5.330e- 3	3.269e- 3	2.074e- 3
2^{-16}	2.196e- 2	1.479e- 2	9.013e- 3	5.332e- 3	3.271e- 3	2.075e- 3
2^{-17}	2.196e- 2	1.479e- 2	9.014e- 3	5.333e- 3	3.271e- 3	2.076e- 3
.
.
.
2^{-25}	2.196e- 2	1.479e- 2	9.016e- 3	5.334e- 3	3.272e- 3	2.076e- 3
$E^{N,M}$	2.196e- 2	1.479e- 2	9.016e- 3	5.334e- 3	3.272e- 3	2.076e- 3
$Q^{N,M}$	0.5703	0.7143	0.7572	0.7052	0.6563	

Table 2. Maximum point-wise errors $E^{N,M}$ and orders of convergence $Q^{N,M}$ calculated for Example 5.2 for $N = M$.

ε	$M = 32$	$M = 64$	$M = 128$	$M = 256$	$M = 512$	$M = 1024$
	$N = 32$	$N = 64$	$N = 128$	$N = 256$	$N = 512$	$N = 1024$
2^{-6}	1.049e- 1	7.356e- 2	5.020e- 2	3.117e- 2	1.919e- 2	1.155e- 2
2^{-7}	7.026e- 2	4.641e- 2	3.060e- 2	2.022e- 2	1.330e- 2	8.602e- 3
2^{-8}	7.726e- 2	5.044e- 2	3.149e- 2	1.860e- 2	1.043e- 2	5.525e- 3
2^{-9}	8.332e- 2	5.558e- 2	3.502e- 2	2.130e- 2	1.221e- 2	6.668e- 3
2^{-10}	8.631e- 2	5.812e- 2	3.707e- 2	2.267e- 2	1.315e- 2	7.274e- 3
2^{-11}	8.779e- 2	5.938e- 2	3.89e- 2	2.335e- 2	1.362e- 2	7.577e- 3
2^{-12}	8.853e- 2	6.001e- 2	3.860e- 2	2.369e- 2	1.385e- 2	7.728e- 3
2^{-13}	8.890e- 2	6.033e- 2	3.885e- 2	2.386e- 2	1.397e- 2	7.809e- 3
2^{-14}	8.908e- 2	6.049e- 2	3.898e- 2	2.394e- 2	1.43e- 2	7.849e- 3
2^{-15}	8.917e- 2	6.056e- 2	3.904e- 2	2.398e- 2	1.406e- 2	7.870e- 3
2^{-16}	8.922e- 2	6.060e- 2	3.908e- 2	2.401e- 2	1.407e- 2	7.880e- 3
2^{-17}	8.924e- 2	6.062e- 2	3.909e- 2	2.402e- 2	1.408e- 2	7.885e- 3
.
.
.
2^{-25}	8.926e- 2	6.064e- 2	3.911e- 2	2.403e- 2	1.409e- 2	7.890e- 3
$E^{N,M}$	1.049e- 1	7.356e- 2	5.020e- 2	3.117e- 2	1.919e- 2	1.155e- 2
$Q^{N,M}$	0.5114	0.5512	0.6876	0.6998	0.7327	

Table 3. Maximum point-wise errors E^M and orders of convergence Q^M calculated for Example 5.1 with $N = 256$.

ε	$M = 32$	$M = 64$	$M = 128$	$M = 256$	$M = 512$	$M = 1024$
2^{-6}	1.536e-02	1.001e-02	6.071e-03	3.456e-03	1.874e-03	9.829e-04
2^{-7}	1.510e-02	9.732e-03	5.834e-03	3.283e-03	1.762e-03	9.167e-04
2^{-8}	1.492e-02	9.552e-03	5.685e-03	3.179e-03	1.698e-03	8.802e-04
2^{-9}	1.483e-02	9.453e-03	5.604e-03	3.124e-03	1.664e-03	8.617e-04
2^{-10}	1.477e-02	9.401e-03	5.563e-03	3.096e-03	1.648e-03	8.524e-04
2^{-11}	1.475e-02	9.374e-03	5.542e-03	3.082e-03	1.639e-03	8.477e-04
2^{-12}	1.473e-02	9.361e-03	5.531e-03	3.075e-03	1.635e-03	8.454e-04
2^{-13}	1.473e-02	9.354e-03	5.526e-03	3.072e-03	1.633e-03	8.443e-04
2^{-14}	1.472e-02	9.351e-03	5.523e-03	3.070e-03	1.632e-03	8.437e-04
2^{-15}	1.472e-02	9.349e-03	5.522e-03	3.069e-03	1.631e-03	8.434e-04
2^{-16}	1.472e-02	9.349e-03	5.521e-03	3.068e-03	1.631e-03	8.432e-04
2^{-17}	1.472e-02	9.348e-03	5.521e-03	3.068e-03	1.631e-03	8.432e-04
.
.
.
2^{-25}	1.472e-02	9.348e-03	5.521e-03	3.068e-03	1.631e-03	8.431e-04
E^M	1.536e-02	1.001e-02	6.071e-03	3.456e-03	1.874e-03	9.829e-04
Q^M	0.6180	0.7208	0.8127	0.8830	0.9311	

Table 4. Maximum point-wise errors E^M and orders of convergence Q^M calculated for Example 5.2 with $N = 256$.

ε	$M = 32$	$M = 64$	$M = 128$	$M = 256$	$M = 512$	$M = 1024$
2^{-6}	3.218e-02	1.799e-02	9.549e-03	4.925e-03	2.502e-03	1.359e-03
2^{-7}	3.252e-02	1.819e-02	9.660e-03	4.983e-03	2.531e-03	1.295e-03
2^{-8}	3.280e-02	1.835e-02	9.746e-03	5.027e-03	2.554e-03	1.287e-03
2^{-9}	3.296e-02	1.844e-02	9.792e-03	5.051e-03	2.566e-03	1.293e-03
2^{-10}	3.306e-02	1.849e-02	9.818e-03	5.064e-03	2.572e-03	1.296e-03
2^{-11}	3.311e-02	1.852e-02	9.831e-03	5.071e-03	2.576e-03	1.298e-03
2^{-12}	3.313e-02	1.853e-02	9.839e-03	5.074e-03	2.578e-03	1.299e-03
2^{-13}	3.315e-02	1.854e-02	9.843e-03	5.076e-03	2.579e-03	1.300e-03
2^{-14}	3.316e-02	1.854e-02	9.845e-03	5.078e-03	2.579e-03	1.300e-03
2^{-15}	3.316e-02	1.855e-02	9.846e-03	5.078e-03	2.579e-03	1.300e-03
2^{-16}	3.316e-02	1.855e-02	9.847e-03	5.078e-03	2.580e-03	1.300e-03
2^{-17}	3.316e-02	1.855e-02	9.847e-03	5.079e-03	2.580e-03	1.300e-03
.
.
.
2^{-25}	3.316e-02	1.855e-02	9.847e-03	5.079e-03	2.580e-03	1.300e-03
E^M	3.316e-02	1.855e-02	9.847e-03	5.079e-03	2.580e-03	1.300e-03
Q^M	0.8383	0.9135	0.9553	0.97773	0.9885	

Table 5. Maximum point-wise errors E^N and orders of convergence Q^N calculated for Example 5.1 with $\mathcal{M} = 128$.

ε	$\mathcal{N} = 32$	$\mathcal{N} = 64$	$\mathcal{N} = 128$	$\mathcal{N} = 256$	$\mathcal{N} = 512$	$\mathcal{N} = 1024$
2^{-6}	1.512e-02	9.353e-03	5.306e-03	2.818e-03	1.413e-03	6.875e-04
2^{-7}	1.716e-02	1.109e-02	6.632e-03	3.690e-03	1.941e-03	9.831e-04
2^{-8}	1.809e-02	1.191e-02	7.272e-03	4.142e-03	2.231e-03	1.155e-03
2^{-9}	1.855e-02	1.230e-02	7.590e-03	4.365e-03	2.376e-03	1.244e-03
2^{-10}	1.877e-02	1.250e-02	7.743e-03	4.474e-03	2.448e-03	1.288e-03
2^{-11}	1.888e-02	1.259e-02	7.821e-03	4.535e-03	2.485e-03	1.309e-03
2^{-12}	1.893e-02	1.263e-02	7.856e-03	4.561e-03	2.503e-03	1.321e-03
2^{-13}	1.896e-02	1.266e-02	7.872e-03	4.573e-03	2.512e-03	1.326e-03
2^{-14}	1.897e-02	1.267e-02	7.881e-03	4.579e-03	2.516e-03	1.328e-03
2^{-15}	1.898e-02	1.268e-02	7.887e-03	4.582e-03	2.518e-03	1.330e-03
2^{-16}	1.898e-02	1.268e-02	7.890e-03	4.583e-03	2.519e-03	1.330e-03
2^{-17}	1.898e-02	1.268e-02	7.891e-03	4.584e-03	2.519e-03	1.331e-03
.
.
.
2^{-25}	1.899e-02	1.269e-02	7.892e-03	4.585e-03	2.520e-03	1.331e-03
E^N	1.899e-02	1.269e-02	7.892e-03	4.585e-03	2.520e-03	1.331e-03
Q^N	0.5817	0.6848	0.7836	0.8636	0.9208	

Table 6. Maximum point-wise errors E^N and orders of convergence Q^N calculated for Example 5.2 with $\mathcal{M} = 128$.

ε	$\mathcal{N} = 32$	$\mathcal{N} = 64$	$\mathcal{N} = 128$	$\mathcal{N} = 256$	$\mathcal{N} = 512$	$\mathcal{N} = 1024$
2^{-6}	8.007e-02	5.905e-02	4.203e-02	2.649e-02	1.639e-02	9.775e-03
2^{-7}	6.695e-02	4.029e-02	2.331e-02	1.538e-02	1.060e-02	6.991e-03
2^{-8}	7.931e-02	5.062e-02	3.043e-02	1.752e-02	9.537e-03	4.940e-03
2^{-9}	8.530e-02	5.572e-02	3.438e-02	2.021e-02	1.127e-02	6.013e-03
2^{-10}	8.851e-02	5.824e-02	3.641e-02	2.157e-02	1.217e-02	6.577e-03
2^{-11}	9.051e-02	5.949e-02	3.743e-02	2.224e-02	1.264e-02	6.862e-03
2^{-12}	9.151e-02	6.012e-02	3.794e-02	2.258e-02	1.287e-02	7.007e-03
2^{-13}	9.201e-02	6.043e-02	3.819e-02	2.275e-02	1.299e-02	7.082e-03
2^{-14}	9.226e-02	6.059e-02	3.832e-02	2.284e-02	1.305e-02	7.120e-03
2^{-15}	9.239e-02	6.066e-02	3.838e-02	2.288e-02	1.308e-02	7.139e-03
2^{-16}	9.245e-02	6.070e-02	3.841e-02	2.290e-02	1.309e-02	7.148e-03
2^{-17}	9.248e-02	6.072e-02	3.843e-02	2.291e-02	1.310e-02	7.153e-03
.
.
.
2^{-25}	9.251e-02	6.074e-02	3.844e-02	2.292e-02	1.311e-02	7.158e-03
E^N	9.251e-02	6.074e-02	3.844e-02	2.292e-02	1.311e-02	7.158e-03
Q^N	0.6070	0.6600	0.7460	0.8065	0.8726	

6. Conclusions

A quasilinear one-dimensional parabolic convection-reaction-diffusion problem with a discontinuous source term has been considered. The solution exhibits a boundary layer at $s = 0$ and a weak interior layer to the right of the discontinuity. A numerical method is constructed to solve the problem, yielding an ε -uniform convergent numerical approximation to the solution. The method employs the standard upwind scheme on the spatial domain and the backward upwind scheme on the temporal domain. A Shishkin mesh is used to discretize the space while the time is discretized using a uniform mesh. The scheme achieves nearly first-order convergence in space and first-order convergence in time. Two numerical examples supporting the theoretical results are presented.

Author contributions

Ruby: Conceptualization, Formal analysis, Investigation, Methodology, Writing – original draft. Vembu Shanthi: Conceptualization, Formal analysis, Investigation, Methodology, Supervision, Writing – review & editing. Higinio Ramos: Conceptualization, Formal analysis, Investigation, Methodology, Supervision, Writing – review & editing. All authors have read and approved the final version of the manuscript for publication.

Use of Generative-AI tools declaration

The authors declare they have not used artificial intelligence (AI) tools in the creation of this article.

Acknowledgments

The first author is grateful for financial support in the form of research fellowship(File no. 09/0895(12534)/2021-EMR-I) the Council of Scientific and Industrial Research (CSIR), India.

Conflict of interest

Prof. Higinio Ramos is a Guest Editor of special issue “Numerical Analysis of Differential Equations with Real-world Applications” for AIMS Mathematics. Prof. Higinio Ramos was not involved in the editorial review and the decision to publish this article.

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

References

1. P. A. Markowich, C. A. Ringhofer, C. Schmeiser, *Semiconductor equations*, Springer Vienna, 2012.
2. E. P. Doolan, J. J. H. Miller, W. H. A. Schilders, *Uniform numerical methods for problems with initial and boundary layers*, Boole Press, 1980.

3. P. A. Farrell, A. Hegarty, J. J. H. Miller, E. O'Riordan, G. I. Shishkin, *Robust computational techniques for boundary layers*, Chapman and hall/CRC, 2000. <https://doi.org/10.1201/9781482285727>
4. J. J. H. Miller, E. O'Riordan, G. I. Shishkin, *Fitted numerical methods for singular perturbation problems: error estimates in the maximum norm for linear problems in one and two dimensions*, World scientific, 2012. <https://doi.org/10.1142/8410>
5. H. G. Roos, M. Stynes, L. Tobiska, *Robust numerical methods for singularly perturbed differential equations*, Springer Series in Computational Mathematics, Vol. 24, 2 Eds., Springer, Berlin, 2008. <https://doi.org/10.1007/978-3-540-34467-4>
6. P. A. Farrell, A. F. Hegarty, J. J. H. Miller, E. O'Riordan, G. I. Shishkin, Global maximum norm parameter-uniform numerical method for a singularly perturbed convection-diffusion problem with discontinuous convection coefficient, *Math. Comput. Model.*, **40** (2004), 1375–1392. <https://doi.org/10.1016/j.mcm.2005.01.025>
7. P. A. Farrell, E. O'Riordan, G. I. Shishkin, A class of singularly perturbed quasilinear differential equations with interior layers, *Math. Comp.*, **78** (2009), 103–127. <https://doi.org/10.1090/S0025-5718-08-02157-1>
8. S. C. S. Rao, A. K. Chaturvedi, Analysis and implementation of a computational technique for a coupled system of two singularly perturbed parabolic semilinear reaction–diffusion equations having discontinuous source terms, *Commun. Nonlinear Sci. Numer. Simul.*, **108** (2022), 106232. <https://doi.org/10.1016/j.cnsns.2021.106232>
9. Ruby, V. Shanthi, H. Ramos, Numerical analysis for a weakly coupled system of singularly perturbed quasilinear problem with non-smooth data, *J. Comput. Sci.*, **85** (2025), 102475. <https://doi.org/10.1016/j.jocs.2024.102475>
10. H. Schlichting, K. Gersten, *Boundary-layer theory*, 9 Eds., Springer Berlin, Heidelberg, 2017. <https://doi.org/10.1007/978-3-662-52919-5>
11. C. M. Bender, S. A. Orszag, *Advanced mathematical methods for scientists and engineers I: asymptotic methods and perturbation theory*, Springer Science & Business Media, 2013.
12. P. Kokotović, H. K. Khalil, J. O'Reilly, *Singular perturbation methods in control: analysis and design*, SIAM, 1999.
13. P. A. Markowich, *The stationary semiconductor device equations*, Springer Vienna, 1986. <https://doi.org/10.1007/978-3-7091-3678-2>
14. S. Selberherr, *Analysis and simulation of semiconductor devices*, Springer Science & Business Media, 2012.
15. O. A. Ladyženskaja, V. A. Solonnikov, N. N. Ural'ceva, Linear and quasi-linear equations of parabolic type, In: *Translations of mathematical monographs*, Vol. 23, American Mathematical Society, 1968.
16. E. O'Riordan, G. I. Shishkin, Singularly perturbed parabolic problems with non-smooth data, *J. Comput. Appl. Math.*, **166** (2004), 233–245. <https://doi.org/10.1016/j.cam.2003.09.025>

17. K. Mukherjee, S. Natesan, ε -Uniform error estimate of hybrid numerical scheme for singularly perturbed parabolic problems with interior layers, *Numer. Algor.*, **58** (2011), 103–141. <https://doi.org/10.1007/s11075-011-9449-6>

18. S. R. Bernfeld, V. Lakshmikantham, *An introduction to nonlinear boundary value problems*, New York: Academic Press, 1974.

19. J. J. H. Miller, E. O’Riordan, G. I. Shishkin, L. P. Shishkina Fitted mesh methods for problems with parabolic boundary layers, *Math. Proc. Royal Irish Acad.*, **98A** (1998), 173–190.

20. M. Stynes, E. O’Riordan, Uniformly convergent difference schemes for singularly perturbed parabolic diffusion-convection problems without turning points, *Numer. Math.*, **55** (1989), 521–544. <https://doi.org/10.1007/BF01398914>

21. P. A. Farrell, A. F. Hegarty, J. J. H. Miller, E. O’Riordan, G. I. Shishkin, Singularly perturbed convection-diffusion problems with boundary and weak interior layers, *J. Comput. Appl. Math.*, **166** (2004), 133–151. <https://doi.org/10.1016/j.cam.2003.09.033>

22. H. G Roos, M. Stynes, L. Tobiska, *Numerical methods for singularly perturbed differential equations: convection-diffusion and flow problems*, Springer Berlin, Heidelberg, 1996. <https://doi.org/10.1007/978-3-662-03206-0>

23. E. O’Riordan, J. Quinn, A linearised singularly perturbed convection-diffusion problem with an interior layer, *Appl. Numer. Math.*, **98** (2015), 1–17. <https://doi.org/10.1016/j.apnum.2015.08.002>

24. J. W. Thomas, *Numerical partial differential equations: finite difference methods*, Springer Science & Business Media, 2013.



AIMS Press

©2025 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (<https://creativecommons.org/licenses/by/4.0/>)