



Research article

Analyzing the continuity of the mild solution in finite element analysis of semilinear stochastic subdiffusion problems

Fang Cheng¹, Ye Hu^{2,*} and Mati ur Rahman^{3,4}

¹ School of Statistics and Applied Mathematics, Anhui University of Finance & Economics, Bengbu 233030, China

² Department of Mathematics and Artificial Intelligence, Lyuliang University, Lishi 033000, China

³ School of Mathematical Sciences, Jiangsu University, Zhenjiang 212013, China

⁴ Department of Computer Science and Mathematics, Lebanese American University, Beirut Lebanon

* **Correspondence:** Email: huye1015@163.com.

Abstract: This paper aimed to further introduce the finite element analysis of non-smooth data for semilinear stochastic subdiffusion problems driven by fractionally integrated additive noise. The mild solution of this stochastic model consisted of three different Mittag-Leffler functions. We analyzed the smoothness of the solution and utilized complex integration to approximate the error of the solution operator under non-smooth data. Consequently, optimal convergence estimates were obtained, and we also obtained the continuity conditions of the mild solution. Finally, the influence of the fractional parameters α and γ on the convergence rates were accurately demonstrated through numerical examples.

Keywords: stochastic time-fractional equation; nonsmooth data analysis; continuity; error estimate

Mathematics Subject Classification: 65C60, 65J15, 65M70, 65N35

1. Introduction

Fractional calculus has gained increasing attention due to its potential applications in various fields of science and engineering [19]. However, the study of fractional calculus has been predominantly focused on deterministic equations, using either deterministic or probabilistic methods. This approach limits the modeling of real-world phenomena where the propagation speed can be finite, as heat flow can be disrupted by material response. In contrast, the classical heat equation assumes infinite speed of heat flow. Recent studies have shown that materials with thermal memory can exhibit finite heat flow speed [2]. This is due to the convolution term in the definition of fractional derivatives and

integrals, which implies that the nearer past has a stronger influence on the present. Additionally, if the internal energy of the material is affected by past random effects, it can be modeled as fractionally integrated additive noise, represented as ${}_0I_t^\gamma \dot{W}(t)$ using the classical Wiener process. Here, ${}_0I_t^\gamma$ (or ${}_{RL}D_{0,t}^{-\gamma}u$) denotes the Riemann-Liouville fractional integral of order γ of the function u defined by

$${}_0I_t^\gamma u(t) \equiv {}_{RL}D_{0,t}^{-\gamma}u(t) = \frac{1}{\Gamma(\gamma)} \int_0^t (t-\sigma)^{\gamma-1} u(\sigma) d\sigma.$$

In this paper, we consider the time fractional semilinear stochastic partial differential equation driven by fractionally integrated additive noise,

$$\begin{cases} {}_cD_{0,t}^\alpha u(t) + Au(t) = f(u(t)) + {}_0I_t^\gamma \dot{W}(t), & 0 < t \leq T, \\ u(0) = u_0, \end{cases} \quad (1.1)$$

where $0 < \alpha < 1$, $0 \leq \gamma \leq 1$, $A = -\Delta$ is a self-adjoint, positive definite, not necessarily bounded operator on the Hilbert space with domain $D(A) = H^2(D) \cap H_0^1(D)$, where $D \subset \mathbb{R}^d$, $d = 1, 2, 3$ denotes a bounded convex polygonal domain. Here, $\dot{W}(t) = \frac{dW(t)}{dt}$ denotes the white noise, and the time fractional derivative ${}_cD_{0,t}^\alpha u$ with order $\alpha \in (0, 1)$ is defined as follows [12],

$${}_cD_{0,t}^\alpha u(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-\sigma)^{-\alpha} \frac{\partial u}{\partial \sigma} d\sigma,$$

where $\Gamma(\cdot)$ denotes the Gamma function.

In recent years, the numerical solution of fractional partial differential equations has become a major focus of research. This is because analytical expressions for such equations are generally difficult to obtain, prompting researchers to explore numerical methods instead. The presence of singular convolution kernels in fractional operators further complicates the solving process. To tackle this challenge, a plethora of remarkable mathematical techniques and innovative approaches have emerged in recent years. These techniques not only play a crucial role in theoretical investigations but also demonstrate great potential in practical applications. For example, a novel technique employing double reduction order and a newly constructed nonlinear compact difference operator has been developed to simulate nonlocal problems on graded meshes [13]. In a separate study, researchers [15] have devised a conservative, positivity-preserving, nonlinear finite volume scheme suitable for multi-term nonlocal Nagumo-type equations using distorted meshes. Additionally, [16] proposed a positivity-preserving finite volume scheme tailored for subdiffusion equations on nonconforming quadrilateral distorted meshes with hanging nodes. Furthermore, [20] addresses the numerical solution of the three-dimensional nonlocal evolution equation with a weakly singular kernel.

Many researchers are also dedicated to studying techniques for solving stochastic partial differential equations. Interested readers are encouraged to explore [5, 6, 14, 17] for further works in this area. In prior research, our focus revolved around weak convergence analysis of the L1 scheme for a stochastic subdiffusion problem, as well as leveraging the spectral method for strong approximation of stochastic semilinear subdiffusion and superdiffusion equations driven by fractionally integrated additive noise [8, 9]. In this article, we delve deeper into the analysis of non-smooth data in this model. Notably, our model problem (1.1), in contrast to the formulation studied in [6, 14], exhibits greater generality and necessitates the involvement of three distinct Mittag-Leffler solution operators (namely $E(t)$, $\bar{E}(t)$,

and $\tilde{E}(t)$, as detailed in Section 2) due to the presence of time fractional derivative and fractionally integrated additive noise.

Upon applying the Riemann-Liouville derivative operator ${}_{RL}D_{0,t}^{1-\alpha} := ({}_0I_t^\alpha)'$ on both sides of (1.1), it is then formally equivalent to a semilinear fractional Volterra type evolution equation

$$du(t) + {}_{RL}D_{0,t}^{1-\alpha} Au(t)dt = {}_{RL}D_{0,t}^{1-\alpha} f(u(t))dt + {}_{RL}D_{0,t}^{1-\alpha-\gamma} dW(t), \quad (1.2)$$

so the existence and uniqueness of a mild solution u can be proved according to the literature methods [1] analogously, even only under some assumptions, via a standard Banach fixed point argument.

The main contributions of the paper are the following:

(i) The paper introduces finite element analysis for semilinear stochastic subdiffusion problems driven by fractionally integrated additive noise. It explores the smoothness of the solution and employs complex integration techniques to approximate the error of the solution operator under non-smooth data.

(ii) The paper establishes the continuity conditions of the mild solution, providing insights into the behavior and regularity of the solution when dealing with non-smooth data. We accurately demonstrate the impact of the fractional parameters α and γ on the convergence rates through numerical examples, offering valuable insights into the sensitivity and dependence of the solution on these parameters.

The remaining sections of this paper are structured as follows. In the upcoming section, we lay out the framework for the stochastic partial differential equation (SPDE) and establish significant smoothing properties using Mittag-Leffler functions. Subsequently, we derive essential error estimates for deterministic subdiffusion under non-smooth data, followed by obtaining optimal convergence estimates for the finite element method in the presence of non-smooth data.

2. Preliminaries

It is known that if $A = -\Delta$ with homogenous Dirichlet boundary conditions, one has $A\varphi_k = \lambda_k\varphi_k$, $k \in \mathbb{N}$, where $0 < \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_k \leq \dots$, $\lim_{k \rightarrow \infty} \lambda_k = \infty$, and the eigenvectors $\{\varphi_k\}_{k=1}^\infty$ form an orthonormal basis for H . Let $H = L_2(D)$ be a separable Hilbert space with inner product (\cdot, \cdot) and norm $\|\cdot\|$. Let $(\Omega, \mathcal{F}, \mathbb{P}, \{\mathcal{F}_t\}_{t \geq 0})$ be a filtered probability space, with Bochner spaces $L_p(\Omega; H) = L_p((\Omega, \mathcal{F}, \mathbb{P}); H)$. Let \mathbb{E} denote the expectation (with respect to \mathbb{P}). We recall an abstract framework to describe the noise $W(t)$ in the model (1.1) more precisely. A Wiener process $W(t)$ with a covariance operator Q may be characterized by the Fourier type series as follows:

$$W(t) = \sum_{k=1}^{\infty} \mu_k^{\frac{1}{2}} \varphi_k \beta_k(t), \quad (2.1)$$

where Q is a bounded, linear, self-adjoint, positive definite operator on H , with the pairs of eigenvalue and eigenfunction $\{(\mu_k, \varphi_k)\}_{k=1}^\infty$. The $\{\beta_k(t)\}_{k=1}^\infty$ is a sequence of independently and identically distributed standard Brownian motions.

For any $\nu \in \mathbb{R}$, we introduce the space $\dot{H}^\nu(D) = D(A^{\frac{\nu}{2}})$ with norm $|v|_\nu^2 = \|A^{\frac{\nu}{2}} v\|^2 = \sum_{k=1}^{\infty} \lambda_k^\nu (v, \varphi_k)^2$, where $\{\varphi_k\}_{k=1}^\infty$ is the orthonormal basis in H .

Let $\mathcal{L} = \mathcal{L}(H)$ denote the space of all bounded linear operators on H and let $\mathcal{L}_2^0 = HS(Q^{\frac{1}{2}}(H), H)$ be the space of Hilbert-Schmidt operators from $Q^{\frac{1}{2}}(H)$ to H , i.e.,

$$\mathcal{L}_2^0 = \{T \in \mathcal{L}(H) : \sum_{k=1}^{\infty} \|TQ^{\frac{1}{2}}\varphi_k\|^2 < \infty\},$$

furnished with the norm $\|T\|_{\mathcal{L}_2^0}^2 = \sum_{k=1}^{\infty} \|TQ^{\frac{1}{2}}\varphi_k\|^2$, thus $\|T\|_{\mathcal{L}_2^0} = \|TQ^{\frac{1}{2}}\|_{HS} < \infty$, for $T \in \mathcal{L}_2^0$.

We also need to recall the Burkholder-Davis-Gundy inequality [11], for $p \geq 2$,

$$\left\| \int_0^t \phi(\sigma) dW(\sigma) \right\|_{L_p(\Omega; H)} \leq C_p \left\| \left(\int_0^t \|\phi(\sigma)\|_{\mathcal{L}_2^0}^2 d\sigma \right)^{\frac{1}{2}} \right\|_{L_p(\Omega; \mathbb{R})}, \quad (2.2)$$

for strongly measurable functions $\phi : [0, T] \rightarrow \mathcal{L}_2^0$. It's important to emphasize that when $p = 2$, it is the Ito isometry.

By using time fractional Duhamel's principle and Laplace transform, we can obtain the mild solution of (1.1) as follows:

$$u(t) = E(t)u_0 + \int_0^t \bar{E}(t-\sigma)f(u(\sigma))d\sigma + \int_0^t \tilde{E}(t-\sigma)dW(\sigma), \quad \mathbb{P} - a.s., \quad (2.3)$$

where

$$E(t) := E_{\alpha,1}(-t^\alpha A), \quad (2.4)$$

$$\bar{E}(t) := t^{\alpha-1} E_{\alpha,\alpha}(-t^\alpha A), \quad (2.5)$$

$$\tilde{E}(t) := t^{\alpha+\gamma-1} E_{\alpha,\alpha+\gamma}(-t^\alpha A). \quad (2.6)$$

The two parameter function of the Mittag-Leffler type plays a very important role in the fractional calculus [5]. We recall the following important properties of the Mittag-Leffler function essential in our analysis.

Lemma 2.1. *Let $0 < \alpha < 2$ and $\beta \in \mathbb{R}$ be arbitrary and $\frac{\pi\alpha}{2} < \mu < \min(\pi, \alpha\pi)$, then there exists a constant $C = C(\alpha, \beta, \mu)$ such that*

$$|E_{\alpha,\beta}(z)| \leq \begin{cases} C(1+|z|)^{-1}, & \beta - \alpha \notin \mathbb{Z}^-, \\ C(1+|z|)^{-2}, & \beta - \alpha \in \mathbb{Z}^-, \end{cases} \quad \mu \leq |\arg(z)| \leq \pi, \quad (2.7)$$

where the notation \mathbb{Z}^- denotes the set of nonpositive integers, i.e., $\mathbb{Z}^- = \{0, -1, -2, \dots\}$.

Throughout we always make the following standing assumption on the fractional orders α and γ , which is sufficient to ensure the well-posedness of Eq (1.1) (see [2, 5]).

Assumption 2.1. *Let $0 < \alpha < 1$, $0 \leq \gamma \leq 1$ with $\alpha + \gamma > \frac{1}{2}$.*

On the Assumption 2.1, for $\beta \in (0, \kappa]$, the regularity of noise can be characterized as follows (see Lemma A.1 of [5]):

$$\|A^{\frac{\beta-\kappa}{2}}\|_{\mathcal{L}_2^0} = \|A^{\frac{\beta-\kappa}{2}}Q^{\frac{1}{2}}\|_{HS} \leq C, \quad (2.8)$$

where (with $\varepsilon > 0$ small)

$$\kappa = \begin{cases} 2, & \text{if } \frac{1}{2} < \gamma < 1, \\ 2 - \varepsilon, & \text{if } \gamma = \frac{1}{2}, \\ 2 - \frac{1-2\gamma}{\alpha} - \varepsilon, & \text{if } 0 \leq \gamma < \frac{1}{2}. \end{cases} \quad (2.9)$$

Now we state the smoothing property of the operators $E(t)$, $\bar{E}(t)$, and $\tilde{E}(t)$.

Lemma 2.2. *There exists C such that for $t > 0$, we have*

$$\|A^s E(t)\| \leq Ct^{-\alpha s}, \quad s \in [0, 1], \quad (2.10)$$

$$\|A^{-s} \dot{E}(t)\| \leq Ct^{\alpha s - 1}, \quad s \in [0, 1], \quad (2.11)$$

$$\|A^s \bar{E}(t)\| \leq Ct^{(1-s)\alpha - 1}, \quad s \in [0, 1], \quad (2.12)$$

$$\|A^{-s} \dot{\bar{E}}(t)\| \leq Ct^{\alpha - 2}, \quad s \geq 0, \quad (2.13)$$

$$\|A^s \tilde{E}(t)\| \leq Ct^{(1-s)\alpha + \gamma - 1}, \quad s \in [0, 1], \quad (2.14)$$

$$\|A^{-s} \dot{\tilde{E}}(t)\| \leq Ct^{\alpha + \gamma - 2}, \quad s \geq 0. \quad (2.15)$$

Proof. We just prove (2.13). The other conclusions can be obtained by using the similar method. By Lemma 2.1, we have

$$\|A^{-s} \dot{\bar{E}}(t)\| = \|A^{-s} t^{\alpha - 2} E_{\alpha, \alpha - 1}(-At^\alpha)\| \leq \sup_{\lambda > 0, t \geq 0} \frac{\lambda^{-s} t^{\alpha - 2}}{(1 + t^\alpha \lambda)^2} \leq Ct^{\alpha - 2}.$$

This completes the proof of the lemma. \square

Remark 2.1. *Overall, the corresponding conclusions are the smoothing properties of the heat semigroup, when $\alpha \rightarrow 1$ and $\gamma \rightarrow 0$, see [17].*

In the qualitative theory of nonlinear PDEs, the subsequent Gronwall type inequalities play a very important role in error estimate.

Lemma 2.3. [3] *Let $T > 0$, $N \in \mathbb{N}$, $k = \frac{T}{N}$, and $t_n = nk$ for $0 \leq n \leq N$. If $\zeta_1, \dots, \zeta_N \geq 0$ satisfy for some $M_0, M_1 \geq 0$, and $\mu, \nu > 0$, the inequality*

$$\zeta_n \leq M_0(1 + t_n^{-1 + \mu}) + M_1 k \sum_{j=1}^{n-1} t_{n-j}^{-1 + \nu} \zeta_j, \quad 1 \leq n \leq N,$$

then there exists a constant $M_2 = M_2(\mu, \nu, M_1, T)$ such that $\zeta_n \leq M_0 M_2(1 + t_n^{-1 + \mu})$, $1 \leq n \leq N$.

To mimic Assumption 2.14 as stated in book [7], let \mathcal{P}_T be the σ -field of predictable stochastic processes, and $\mathcal{B}(S)$ be the Borel σ -field of S . Given the available options, we shall make some reasonable assumptions about the nonlinear parts.

Assumption 2.2. *The mapping $f : [0, T] \times \Omega \times H \rightarrow \dot{H}^{-1}$, $(t, \omega, h) \rightarrow f(t, \omega, h)$ is $\mathcal{P}_T \times \mathcal{B}(H)/\mathcal{B}(\dot{H}^{-1})$ measurable. When $\delta \in (0, \frac{1}{2})$, there exists a constant C that satisfies the following expression:*

$$\|f(t_1, \omega, h) - f(t_2, \omega, h)\|_{-1} \leq C(1 + \|h\|)(t_2 - t_1)^\delta, \quad (2.16)$$

for all $h \in H$, $0 \leq t_1 < t_2 \leq T$, $\omega \in \Omega$.

3. Non-smooth data analysis for stochastic problem

In this section, we formulate the Galerkin finite element methods for spatial discretization in combination with time discretization based on an exponential Euler type method for approximation of (1.1). Let \mathcal{T}_h be a regular shaped quasi-uniform triangulation of the domain D , and let $S_h \subset H_0^1(D)$ be the space of continuous piecewise linear functions on the triangulation \mathcal{T}_h . We define the L_2 -projection $P_h : H \rightarrow S_h$ by

$$(P_h u, \chi) = (u, \chi), \quad \chi \in S_h, \quad (3.1)$$

and the Ritz projection $R_h : H_0^1 \rightarrow S_h$ by

$$a(R_h u, \chi) = a(u, \chi), \quad \chi \in S_h,$$

where $a(u, \chi) = (\nabla u, \nabla \chi)$ is the associated bilinear form.

Note that by interpreting the righthand side of (3.1) as a duality pairing between $\dot{H}^1(D)$ and $\dot{H}^{-1}(D)$, one may extend P_h to be a bounded operator from $\dot{H}^{-1}(D)$ to S_h . It is well-known that the operators P_h and R_h have the following approximation properties [18].

Lemma 3.1. *The operators P_h and R_h satisfy*

$$\begin{aligned} \|P_h u - u\| + h \|\nabla(P_h u - u)\| &\leq Ch^q |u|_q, \quad \text{for } u \in \dot{H}^q, \quad q = 1, 2, \\ \|R_h u - u\| + h \|\nabla(R_h u - u)\| &\leq Ch^q |u|_q, \quad \text{for } u \in \dot{H}^q, \quad q = 1, 2. \end{aligned}$$

The semi-discrete Galerkin FEM scheme for (1.1) is to find $u_h(t) \in S_h$ such that

$$\begin{cases} {}^c D_{0,t}^\alpha u_h(t) + A_h u_h(t) = P_h f(u_h(t)) + {}_0 I_t^\gamma P_h \dot{W}(t), & 0 < t \leq T, \\ u_h(0) = P_h u_0, \end{cases} \quad (3.2)$$

where the discrete Laplacian A_h is defined by

$$A_h : S_h \rightarrow S_h, \quad (A_h \psi, \chi) = a(\psi, \chi), \quad \forall \psi, \chi \in S_h.$$

Naturally, we present the discrete analogues of operators $E_h(t)$, $\bar{E}_h(t)$, and $\widetilde{E}_h(t)$ as follows

$$E_h(t) := E_{\alpha,1}(-t^\alpha A_h), \quad (3.3)$$

$$\bar{E}_h(t) := t^{\alpha-1} E_{\alpha,\alpha}(-t^\alpha A_h), \quad (3.4)$$

$$\widetilde{E}_h(t) := t^{\alpha+\gamma-1} E_{\alpha,\alpha+\gamma}(-t^\alpha A_h). \quad (3.5)$$

Now, let $0 < t_0 < t_1 < \dots < t_M = T$ be a uniform partition of time interval $[0, T]$, with time step $\Delta t = t_{m+1} - t_m$, $m = 0, 1, \dots, M-1$. Hence, the fully discrete approximation of (1.1) is given by

$$U_h^m = E_h(t_m) P_h u_0 + \sum_{j=0}^{m-1} \int_{t_j}^{t_{j+1}} \bar{E}_h(t_m - \sigma) d\sigma (P_h F(U_h^j)) + \int_0^{t_m} \widetilde{E}_h(t_m - \sigma) P_h dW(\sigma), \quad (3.6)$$

with initial value $U_h^0 = P_h u_0$.

Let us introduce and prove some Lemmas that will play an important role later on.

Lemma 3.2. [10] Let $0 \leq \omega \leq \mu \leq 2$. For $\alpha \in (0, 1)$, there exists a constant C such that

$$\|(E(t) - E_h(t)P_h)v\| \leq Ch^\mu t^{-\alpha \frac{\mu-\omega}{2}} \|v\|_\omega, \quad \text{for } v \in \dot{H}^\omega.$$

We now turn to the non-smooth data error estimates of the approximations to $\bar{E}(t)g$, $g \in H$ in the semi-discrete case.

Lemma 3.3. Let $0 \leq s \leq 1$ and $0 \leq r \leq 2$ with $r + s \leq 2$. For $g \in H$, there holds

$$\|A^{\frac{s}{2}}(\bar{E}(t) - \bar{E}_h(t)P_h)g\| \leq Ch^{2-s-r} t^{\frac{r}{2}-1} \|g\|. \tag{3.7}$$

Proof. In the case $s = 0$, by the inverse Laplace transform, for any given $g \in H$, we have

$$\bar{E}(t)g = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\delta}} e^{zt} (z^\alpha + A)^{-1} g dz, \tag{3.8}$$

$$\bar{E}_h(t)P_h g = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\delta}} e^{zt} (z^\alpha + A_h)^{-1} P_h g dz \tag{3.9}$$

where $\Gamma_{\theta,\delta}^\tau = \{z \in \Gamma_{\theta,\delta} : |\Im z| \leq \frac{\pi}{\tau}\}$, and $\Gamma_{\theta,\delta} = \{z \in \mathbb{C} : z = re^{\pm i\theta}, r \geq \delta\} \cup \{z \in \mathbb{C} : z = \delta e^{i\phi}, |\phi| \leq \theta\}$, $\frac{\pi}{2} < \theta < \pi$, $\frac{\pi}{\tau} > \delta$.

Let us first show (3.7). For any fixed $g \in H$, we have, by (3.8) and (3.9),

$$\|(\bar{E}(t) - \bar{E}_h(t)P_h)g\| \leq C \int_{\Gamma_{\theta,\delta}} e^{\Re(z)t} \left\| ((z^\alpha + A)^{-1} - (z^\alpha + A_h)^{-1} P_h) g \right\| |dz|.$$

By [4], p. 820, we get $\|((z^\alpha + A)^{-1} - (z^\alpha + A_h)^{-1} P_h)g\| \leq Ch^2 \|g\|$, $\forall z \in \Gamma_{\theta,\delta}$. Hence, we have

$$\begin{aligned} \|(\bar{E}(t) - \bar{E}_h(t)P_h)g\| &\leq Ch^2 \|g\| \left(\int_{\{z \in \mathbb{C} : z = \delta e^{i\phi}, |\phi| \leq \theta\}} + \int_{\{z \in \mathbb{C} : z = re^{\pm i\theta}, r \geq \delta\}} \right) e^{\Re(z)t} |dz| \\ &= I + II. \end{aligned}$$

For I , with $z = \delta e^{i\phi}$, we have $\delta = t^{-1}$, where $t^{-1} < \frac{\pi}{\tau}$ for sufficiently small τ ,

$$I \leq Ch^2 \|g\| \int_{-\theta}^{\theta} e^{t\delta \cos \phi} \delta d\phi \leq Ch^2 \|g\| \delta \int_{-\theta}^{\theta} e^{\cos \phi} d\phi \leq Ch^2 t^{-1} \|g\|.$$

For II , with $z = re^{\pm i\theta}$, we have $r \geq \delta$, $\delta = t^{-1}$,

$$II \leq Ch^2 \|g\| \int_{\delta}^{\infty} e^{tr \cos \theta} dr \leq Ch^2 \|g\| \int_{t^{-1}}^{\infty} e^{-ctr} dr \leq Ch^2 \|g\| t^{-1} \int_c^{\infty} e^{-x} dx \leq Ch^2 t^{-1} \|g\|,$$

where c, C is some suitable positive constant.

Meanwhile, by (2.12) and the triangle inequality,

$$\|(\bar{E}(t) - \bar{E}_h(t)P_h)g\| \leq Ct^{\alpha-1} \|g\|.$$

Similarly, for $s = 1$, there holds

$$\|A^{\frac{1}{2}}(\bar{E}(t) - \bar{E}_h(t)P_h)g\| \leq Ch t^{-1} \|g\|.$$

Now the desired assertion follows by interpolation, which completes the proof. □

The next lemma gives an error estimate on $\tilde{E}(t)$. More details can be found in Lemma 4.4 of [5].

Lemma 3.4. *Let $0 \leq s \leq 1$ and $0 \leq r \leq 2$ with $r + s \leq 2$. For $g \in H$, there holds*

$$\|A^{\frac{s}{2}}(\tilde{E}(t) - \tilde{E}_h(t)P_h)g\| \leq Ch^{2-s-r}t^{\frac{\alpha}{2}+\gamma-1}\|g\|. \quad (3.10)$$

Based on the previous discussion, we are ready to prove the error estimates for the fully discrete approximation.

Theorem 3.1. *For $0 \leq r \leq 2$ and $r\alpha + 2\gamma > 1$, by the Assumptions 2.1 and 2.2, with $\nu \in [0, \beta]$, $\beta \in (0, \kappa]$, $\alpha + \gamma \in (\frac{1}{2}, 1)$ and $\alpha(2 - \kappa + \beta - \nu) + 2\gamma - 1 \in (0, 1)$, then there holds*

$$\sup_{t_m \in [0, T]} \|u(t_m) - U_h^m\|_{L_2(\Omega; H)} \leq C \left(h^{(2-r)\alpha} + h^{2-r} \max \left\{ t_m^{-\frac{(2-r)\alpha}{2}}, \ln \frac{t_m}{h^{2-r}} \right\} + \Delta t^{\alpha\nu} \right).$$

Proof. Due to (1.1), it is formally then equivalent to a nonlinear fractional Volterra type evolution equation. We can obtain $\sup_{t \in [0, T]} \|u(t)\|_{L_p(\Omega; \dot{H}^\nu)} \leq C$. Subtracting (3.6) from (2.3) and by taking norms, one obtains

$$\begin{aligned} \|u(t_m) - U_h^m\|_{L_2(\Omega; H)} &\leq \|E(t_m)u_0 - E_h(t_m)P_h u_0\|_{L_2(\Omega; H)} \\ &\quad + \left\| \int_0^{t_m} (\bar{E}(t_m - \sigma) - \bar{E}_h(t_m - \sigma)P_h)F(u(\sigma))d\sigma \right\|_{L_2(\Omega; H)} \\ &\quad + \left\| \int_0^{t_m} (\tilde{E}(t_m - \sigma) - \tilde{E}_h(t_m - \sigma)P_h)dW(\sigma) \right\|_{L_2(\Omega; H)} \\ &\quad + \left\| \sum_{j=0}^{m-1} \int_{t_j}^{t_{j+1}} \bar{E}(t_m - \sigma)P_h(F(u(\sigma)) - F(U_h^j))d\sigma \right\|_{L_2(\Omega; H)} \\ &= I_1 + I_2 + I_3 + I_4. \end{aligned}$$

Here, we note that I_1 , I_2 , and I_3 correspond to the spatial finite element discretization error, while I_4 corresponds to the temporal error.

The estimate of I_1 is a consequence of Lemma 3.2. For $\omega = 0$ and $\mu = 2 - r$, we get

$$I_1 = \|E(t_m)u_0 - E_h(t_m)P_h u_0\|_{L_2(\Omega; H)} \leq Ch^{2-r}t_m^{-\frac{(2-r)\alpha}{2}}.$$

For I_2 , by using (3.7) and $r > 0$, we have

$$\begin{aligned} I_2 &= \left\| \int_0^{t_m} (\bar{E}(t_m - \sigma) - \bar{E}_h(t_m - \sigma)P_h)F(u(\sigma))d\sigma \right\|_{L_2(\Omega; H)} \\ &\leq \int_0^{t_m} \|\bar{E}(t_m - \sigma) - \bar{E}_h(t_m - \sigma)P_h\| \cdot (1 + \|u(\sigma)\|_{L_2(\Omega; H)})d\sigma \\ &\leq C \int_0^{t_m} \|\bar{E}(t_m - \sigma) - \bar{E}_h(t_m - \sigma)P_h\|d\sigma = C \int_0^{t_m} \|\bar{E}(\sigma) - \bar{E}_h(\sigma)P_h\|d\sigma \\ &\leq \int_0^{t_m} Ch^{2-r}\sigma^{\frac{r\alpha}{2}-1}d\sigma = Ch^{2-r}t_m^{\frac{r\alpha}{2}}. \end{aligned}$$

In the case $r = 0$, similarly, we split I_2 into two terms,

$$I_2 \leq \int_0^{h^{2-r}} \|\bar{E}(\sigma) - \bar{E}_h(\sigma)\| d\sigma + \int_{h^{2-r}}^{t_m} \|\bar{E}(\sigma) - \bar{E}_h(\sigma)\| d\sigma = I_{21} + I_{22}.$$

For I_{21} , noting $\|\bar{E}(\sigma)\| \leq C\sigma^{\alpha-1}$ and $\|\bar{E}_h(\sigma)\| \leq C\sigma^{\alpha-1}$, then

$$I_{21} \leq \int_0^{h^{2-r}} \|\bar{E}(\sigma) - \bar{E}_h(\sigma)\| d\sigma \leq Ch^{(2-r)\alpha}.$$

For I_{22} , by (3.7) with $r = 0$, we derive

$$I_{22} \leq \int_{h^{2-r}}^{t_m} \|\bar{E}(\sigma) - \bar{E}_h(\sigma)\| d\sigma \leq \int_{h^{2-r}}^{t_m} Ch^{2-r}\sigma^{-1} d\sigma = Ch^{2-r}l_h,$$

where $l_h = \ln \frac{t_m}{h^{2-r}}$.

Now, we estimate I_3 . Using (3.10), by Itô's formula, for trace class and $r\alpha + 2\gamma > 1$, we obtain

$$\begin{aligned} I_3^2 &= \left\| \int_0^{t_m} (\bar{E}(t_m - \sigma) - \bar{E}_h(t_m - \sigma)P_h) dW(\sigma) \right\|_{L_2(\Omega; H)}^2 \\ &= \int_0^{t_m} \left\| (\bar{E}(t_m - \sigma) - \bar{E}_h(t_m - \sigma)P_h) \right\|_{\mathcal{L}_2^0}^2 d\sigma \leq C \int_0^{t_m} \left\| (\bar{E}(\sigma) - \bar{E}_h(\sigma)) \right\|_{\mathcal{L}_2^0}^2 d\sigma \\ &\leq C \int_0^{t_m} \left\| (\bar{E}(\sigma) - \bar{E}_h(\sigma)) Q^{\frac{1}{2}} \right\|_{HS}^2 d\sigma \leq C \int_0^{t_m} \left\| \bar{E}(\sigma) - \bar{E}_h(\sigma) \right\|^2 \cdot \left\| Q^{\frac{1}{2}} \right\|_{HS}^2 d\sigma \\ &\leq C \int_0^{t_m} (h^{2-r}\sigma^{\frac{r\alpha}{2} + \gamma - 1})^2 \cdot \left\| Q^{\frac{1}{2}} \right\|_{HS}^2 d\sigma \leq C \int_0^{t_m} (h^{2-r}\sigma^{\frac{r\alpha}{2} + \gamma - 1})^2 d\sigma \\ &\leq Ch^{4-2r} t_m^{r\alpha + 2\gamma - 1}. \end{aligned}$$

For I_4 , by Lemma 2.2, the idea is similar to [9] with different coefficient condition $\alpha(2 - \kappa + \beta - \nu) + 2\gamma - 1 > 0$, so we get $I_4 \leq \Delta t^{\alpha\nu}$. Finally, combining the Gronwall's inequality of Lemma 2.3, we complete the proof of Theorem 3.1. \square

Under the assumptions of smoothness of the solution operator and the nonlinear term, we can obtain the following conclusions.

Theorem 3.2. *Let $p \in [2, \infty)$ be given such that Assumptions 2.13–2.17 hold, then the unique mild solution u to (2.3) is continuous with respect to the norm $\|\cdot\|_{L_p(\Omega; \dot{H}^s)}$.*

Proof. According to the definition of continuity, our goal is to show that

$$\lim_{\substack{t_2 - t_1 \rightarrow 0 \\ t_1 < t_2}} \|u(t_2) - u(t_1)\|_{L_p(\Omega; \dot{H}^s)} = 0$$

with either t_1 or t_2 fixed.

Therefore, with the expression of a mild solution u , we divide $\|u(t_2) - u(t_1)\|_{L_p(\Omega; \dot{H}^s)}$ into five parts, and only need to show that each of the parts approaches zero when $t_2 \rightarrow t_1$. By utilizing the triangle inequality, we obtain

$$\|u(t_1) - u(t_2)\|_{L_p(\Omega; \dot{H}^s)} \leq \|E(t_1)u_0 - E(t_2)u_0\|_{L_p(\Omega; \dot{H}^s)}$$

$$\begin{aligned}
& + \left\| \int_{t_1}^{t_2} \bar{E}(t_2 - \sigma) f(u(\sigma)) d\sigma \right\|_{L_p(\Omega; \dot{H}^s)} \\
& + \left\| \int_0^{t_1} (\bar{E}(t_2 - \sigma) - \bar{E}(t_1 - \sigma)) f(u(\sigma)) d\sigma \right\|_{L_p(\Omega; \dot{H}^s)} \\
& + \left\| \int_{t_1}^{t_2} \tilde{E}(t_2 - \sigma) dW(\sigma) \right\|_{L_p(\Omega; \dot{H}^s)} \\
& + \left\| \int_0^{t_1} (\tilde{E}(t_2 - \sigma) - \tilde{E}(t_1 - \sigma)) dW(\sigma) \right\|_{L_p(\Omega; \dot{H}^s)} \\
& = M_1 + M_2 + M_3 + M_4 + M_5.
\end{aligned}$$

For the expression M_1 , we apply (2.11) from Lemma 2.2. One yields

$$\begin{aligned}
M_1 & = \left\| E(t_1)u_0 - E(t_2)u_0 \right\|_{L_p(\Omega; \dot{H}^s)} = \left\| A^{\frac{s}{2}} (E(t_1) - E(t_2)) u_0 \right\|_{L_p(\Omega; H)} \\
& = \left\| \int_{t_1}^{t_2} A^{\frac{s}{2}} \dot{E}(\tau) u_0 d\tau \right\|_{L_p(\Omega; H)} = \left\| \int_{t_1}^{t_2} A^{\frac{-\nu}{2}} \dot{E}(\tau) A^{\frac{\nu+s}{2}} u_0 d\tau \right\|_{L_p(\Omega; H)} \\
& \leq \|u_0\|_{L_p(\Omega; \dot{H}^{s+\nu})} \cdot \int_{t_1}^{t_2} \|A^{\frac{-\nu}{2}} \dot{E}(\tau)\| d\tau \leq C \|u_0\|_{L_p(\Omega; \dot{H}^{s+\nu})} \cdot \int_{t_1}^{t_2} \tau^{\frac{\alpha\nu}{2}-1} d\tau \\
& \leq C \|u_0\|_{L_p(\Omega; \dot{H}^{s+\nu})} \cdot (t_2^{\frac{\alpha\nu}{2}} - t_1^{\frac{\alpha\nu}{2}}),
\end{aligned}$$

and the validity of $\lim_{t_2-t_1} M_1 = 0$ here is obvious when $\nu \geq 0$.

For M_2 , we insert a node t_3 and split it into three parts.

$$\begin{aligned}
M_2 & = \left\| \int_{t_1}^{t_2} \bar{E}(t_2 - \sigma) f(u(\sigma)) d\sigma \right\|_{L_p(\Omega; \dot{H}^s)} \\
& \leq \left\| \int_{t_1}^{t_2} \bar{E}(t_2 - \sigma) f(u(\sigma)) d\sigma \right\|_{L_p(\Omega; \dot{H}^s)} \\
& \leq \left\| \int_{t_1}^{t_2} A^{\frac{s+1}{2}} \bar{E}(t_2 - \sigma) A^{\frac{-1}{2}} (f(u(\sigma)) - f(u(t_2))) d\sigma \right\|_{L_p(\Omega; H)} \\
& \quad + \left\| \int_{t_1}^{t_2} A^{\frac{s+1}{2}} \bar{E}(t_2 - \sigma) A^{\frac{-1}{2}} (f(u(t_2)) - f(u(t_3))) d\sigma \right\|_{L_p(\Omega; H)} \\
& \quad + \left\| \int_{t_1}^{t_2} A^{\frac{s+1}{2}} \bar{E}(t_2 - \sigma) A^{\frac{-1}{2}} f(u(t_3)) d\sigma \right\|_{L_p(\Omega; H)} \\
& = M_{21} + M_{22} + M_{23}.
\end{aligned}$$

For M_{21} , we apply (2.12) from Lemma 2.2, and Assumption 2.2. We obtain

$$\begin{aligned}
M_{21} & \leq \int_{t_1}^{t_2} \|A^{\frac{s+1}{2}} \bar{E}(t_2 - \sigma) A^{\frac{-1}{2}} (f(u(t_2)) - f(u(\sigma)))\|_{L_p(\Omega; H)} d\sigma \\
& \leq C \int_{t_1}^{t_2} (t_2 - \sigma)^{-\frac{s+1}{2}\alpha + \alpha - 1} (t_2 - \sigma)^\delta (1 + \|u(\sigma)\|_{L_p(\Omega; H)}) d\sigma \\
& \leq C (t_2 - t_1)^{\frac{1-s}{2}\alpha + \delta} (1 + \sup_{\sigma \in [0, T]} \|u(\sigma)\|_{L_p(\Omega; H)}).
\end{aligned}$$

Due to $\delta \in (0, \frac{1}{2})$, therefore $\frac{1-s}{2}\alpha + \delta > 0$, and thus when $t_2 \rightarrow t_1$, it is valid.

For M_{22} ,

$$\begin{aligned} M_{22} &= \left\| \int_{t_1}^{t_2} A^{\frac{s+1}{2}} \bar{E}(t_2 - \sigma) A^{-\frac{1}{2}} (f(u(t_2)) - f(u(t_3))) d\sigma \right\|_{L_p(\Omega; H)} \\ &\leq \int_{t_1}^{t_2} \|A^{\frac{s+1}{2}} \bar{E}(t_2 - \sigma) A^{-\frac{1}{2}} (f(u(t_2)) - f(u(t_3)))\|_{L_p(\Omega; H)} d\sigma, \end{aligned}$$

which is formally identical to M_{21} . When $t_2 \rightarrow t_3 \in [t_1, t_2]$ holds, the conclusion is valid.

For M_{23} , with the help of (2.12), one gets

$$\begin{aligned} M_{23} &= \left\| \int_{t_1}^{t_2} A^{\frac{s+1}{2}} \bar{E}(t_2 - \sigma) A^{-\frac{1}{2}} f(u(t_3)) d\sigma \right\|_{L_p(\Omega; H)} \\ &\leq \int_{t_1}^{t_2} \|A^{\frac{s+1-r}{2}} \bar{E}(t_2 - \sigma) A^{-\frac{1+r}{2}} f(u(t_3))\|_{L_p(\Omega; H)} d\sigma \\ &\leq C \int_{t_1}^{t_2} (t_2 - \sigma)^{\frac{1+r-s}{2}\alpha} d\sigma \cdot (1 + \sup_{\sigma \in [0, T]} \|Y(\sigma)\|_{L_p(\Omega; \dot{H}^r)}). \end{aligned}$$

To ensure the continuity property holds, it is sufficient to only have the integration in the last inequality be well-defined and satisfy the conditions of $\frac{1+r-s}{2}\alpha > 0$.

For M_3 , we prove that the following property holds first. One obtains

$$\begin{aligned} \|A^s \dot{\bar{E}}(t)\|^2 &= \|A^s t^{\alpha-2} E_{\alpha, \alpha-1}(-At^\alpha)\|^2 = |t^{\alpha-2} E_{\alpha, \alpha-1}(-At^\alpha)|_s^2 \\ &= \sum_{j=1}^{\infty} \lambda_j^{2s} t^{2(\alpha-2)} \cdot E_{\alpha, \alpha-1}(-t^\alpha \lambda_j)^2 \cdot (v, \varphi_j)^2 \\ &\leq t^{2[(1-s)\alpha-2]} \cdot \sum_{j=1}^{\infty} \frac{(t^\alpha \lambda_j)^{2s}}{(1+t^\alpha \lambda_j)^4} (v, \varphi_j)^2, \end{aligned}$$

and then we have

$$\begin{aligned} M_3 &= \left\| \int_0^{t_1} (\bar{E}(t_2 - \sigma) - \bar{E}(t_1 - \sigma)) f(u(\sigma)) d\sigma \right\|_{L_p(\Omega; \dot{H}^s)} \\ &\leq \int_0^{t_1} \|(\bar{E}(t_2 - \sigma) - \bar{E}(t_1 - \sigma)) f(u(\sigma))\|_{L_p(\Omega; \dot{H}^s)} d\sigma \\ &= \int_0^{t_1} \left\| \int_{t_1}^{t_2} A^{\frac{s+1-r}{2}} \dot{\bar{E}}(t_2 - \sigma) d\tau A^{-\frac{1+r}{2}} f(u(\sigma)) \right\|_{L_p(\Omega; H)} d\sigma \\ &\leq C \int_{t_1}^{t_2} |(t_2 - \sigma)^{\frac{1+r-s}{2}\alpha-1} - (t_1 - \sigma)^{\frac{1+r-s}{2}\alpha-1}| d\sigma \cdot (1 + \sup_{\sigma \in [0, T]} \|u(\sigma)\|_{L_p(\Omega; \dot{H}^r)}) \\ &\leq C(t_2 - t_1)^{\frac{1+r-s}{2}\alpha} \cdot (1 + \sup_{\sigma \in [0, T]} \|u(\sigma)\|_{L_p(\Omega; \dot{H}^r)}). \end{aligned}$$

Since $\frac{1+r-s}{2}\alpha > 0$ holds here, the conclusion is valid.

For M_4 , by applying (2.14) and the Ito formula, we can obtain

$$M_4 = \left\| \int_{t_1}^{t_2} \tilde{\bar{E}}(t_2 - \sigma) dW(\sigma) \right\|_{L_p(\Omega; \dot{H}^s)}$$

$$\begin{aligned}
&\leq C\|(\int_{t_1}^{t_2} \|A^{\frac{s}{2}}\tilde{E}(t_2-\sigma)\|_{L_2^0}^2 d\sigma)^{\frac{1}{2}}\|_{L_p(\Omega;R)} \\
&\leq C\|(\int_{t_1}^{t_2} (t_2-\sigma)^{2(-\frac{s}{2}\alpha+\alpha+\gamma-1)} d\sigma)^{\frac{1}{2}}\|_{L_p(\Omega;R)} \\
&\leq C(t_2-t_1)^{(1-\frac{s}{2})\alpha+\gamma-\frac{1}{2}}.
\end{aligned}$$

In order for the conclusion to hold, it is necessary for the conditions of $(1-\frac{s}{2})\alpha+\gamma-\frac{1}{2} > 0$ to be satisfied here.

For M_5 , according to the requirement, we first prove that the following property holds.

$$\begin{aligned}
\|A^s\tilde{E}(t)\|^2 &\leq \|A^s t^{\alpha+\gamma-2} E_{\alpha,\alpha+\gamma-1}(-At^\alpha)\|^2 = |t^{\alpha+\gamma-2} E_{\alpha,\alpha+\gamma-1}(-At^\alpha)|_{2s}^2 \\
&= \sum_{j=1}^{\infty} \lambda_j^{2s} t^{2(\alpha+\gamma-2)} E_{\alpha,\alpha+\gamma-1}(-t^\alpha \lambda_j)^2(v,\varphi)^2 \\
&\leq t^{2[(1-s)\alpha+\gamma-2]} \sum_{j=1}^{\infty} \frac{(\lambda_j t^\alpha)^{2s}}{(1+t^\alpha \lambda_j)^2} (v,\varphi)^2,
\end{aligned}$$

then we have

$$\begin{aligned}
M_5 &= \|\int_0^{t_1} (\tilde{E}(t_2-\sigma) - \tilde{E}(t_1-\sigma)) dW(\sigma)\|_{L_p(\Omega;H^s)} \\
&\leq C\|(\int_0^{t_1} \|A^{\frac{s}{2}}(\tilde{E}(t_2-\sigma) - \tilde{E}(t_1-\sigma))\|_{L_2^0}^2 d\sigma)^{\frac{1}{2}}\|_{L_p(\Omega;R)} \\
&= C\|(\int_0^{t_1} \int_{t_1}^{t_2} \|A^{\frac{s}{2}}\tilde{E}(\tau-\sigma)\|_{L_2^0}^2 d\sigma)^{\frac{1}{2}}\|_{L_p(\Omega;R)} \\
&\leq C\|(\int_0^{t_1} \int_{t_1}^{t_2} [(\tau-\sigma)^{(1-\frac{s}{2})\alpha+\gamma-2}]^2 d\tau d\sigma)^{\frac{1}{2}}\|_{L_p(\Omega;R)} \\
&= C\|(\int_0^{t_1} [(t_2-\sigma)^{(2-s)\alpha+2\gamma-3} - (t_1-\sigma)^{(2-s)\alpha+2\gamma-3}] d\sigma)^{\frac{1}{2}}\|_{L_p(\Omega;R)} \\
&\leq C(t_2-t_1)^{(1-\frac{s}{2})\alpha+\gamma-1}.
\end{aligned}$$

The last term of the inequality only needs to satisfy the condition of $(1-\frac{s}{2})\alpha+\gamma-\frac{1}{2} > 0$. Thus, the theorem is proved. \square

4. Numerical implementation

We conducted several numerical experiments in this section to validate our previous theoretical findings. To illustrate the theoretical results obtained in Theorem 3.1 for $\alpha \in (0, 1)$, we provide two numerical examples. Specifically, we set $T = 0.1$, $D = (0, 1)$, and $F(u) = \sin(u)$.

To explain the computer implementation of the full-discrete method (3.6), we make the assumption that the covariance operator Q possesses the same eigenfunctions as A , such that $Qv = \sum_{k=1}^{\infty} \mu_k(v, \varphi_k) \varphi_k$. Furthermore, we assume that $W(t)$ has the following Fourier series expansion: $W(t) = \sum_{k=1}^{\infty} \mu_k^{1/2} \varphi_k \beta_k(t)$.

Thus, the semi-discrete solution $U_{m,k,h}^N$ satisfies

$$U_{m,k,h}^N = E_{\alpha,1}(-t_m^\alpha A_h) P_h u_{0,k} + \sum_{j=0}^{m-1} \int_{t_j}^{t_{j+1}} (t_m - \sigma)^{\alpha-1} E_{\alpha,\alpha}(-(t_m - \sigma)^\alpha A_h) d\sigma \cdot P_h F_k(u(t_j)) \\ + \int_0^{t_m} (t_m - \sigma)^{\alpha+\gamma-1} E_{\alpha,\alpha+\gamma}(-(t_m - \sigma)^\alpha A_h) P_h \mu_k^{\frac{1}{2}} d\beta_k(\sigma),$$

where $u_{0,k} = (u_0, \varphi_k)$, $F_k(\cdot) = (F(\cdot), \varphi_k)$, and $\beta_k(\sigma)$ for $k = 1, \dots, N$ are mutually independent standard Brownian motions.

In our experiments, we investigated the dependence of the error estimates in Theorem 3.1 on the time step size Δt . To approximate the exact solution, we used the fully discrete solution with a small time step size of $\Delta t = 2^{-10}$ and a spatial step size of $\Delta x = 2^{-10}$. The theoretical rate of convergence is $\alpha\nu$ for $\nu \in (0, 1)$, which approaches α as $\nu \rightarrow 1$.

To conduct the experiments, we fixed a small spatial step size Δx and considered a sequence of moderate time step sizes $\Delta t_i = 2^{-i}$, $i = 2, \dots, 5$. We performed $M = 100$ simulations for each time step size Δt_i . In each simulation ω_j , $j = 1, 2, \dots, M$, we generated N_h independent Brownian motions $\beta_k(t)$, $k = 1, 2, \dots, N_h$. The initial data was set to $u_0 = 1$, and we defined the nonlinear operator $F(u) = \sin(u)$.

In Figure 1, we investigated the effects of different parameters α and γ . We observed that the numerical results are consistent with the theoretical results stated in Theorem 3.1. The numerical values vary slightly due to the limited range of α and γ . As the mesh size is refined, the numerical experiments support the theoretical findings in Theorem 3.1.

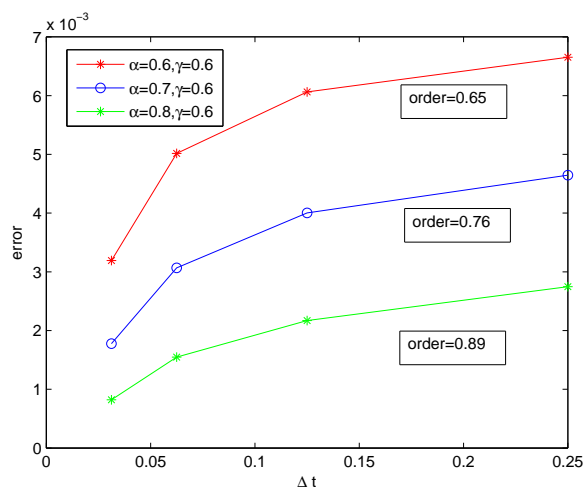


Figure 1. The $L_2(\Omega; H)$ errors and orders of convergence with $\alpha \in (0, 1)$.

In the following analysis, we focus on the spatial convergence. To this end, we consider a fixed number of space steps $M = 200$ and final time $T = 1$, and we obtain the reference solution at $N = 480$. Next, we compute the numerical results for different combinations of fractional orders α and γ , with trace class noise (with $m = 2$). Table 1 illustrates the results, which show that an $O(h^2)$ convergence order is observed for all combinations. We observed that the initial convergence rate changes slowly, but as the mesh is refined, the results are in excellent agreement with the theoretical predictions from Theorem 3.1.

Table 1. The $L^2(\Omega; H)$ -error with trace class noise ($m = 2$) at $T = 1$.

| γ | $\alpha \backslash M$ | 10 | 20 | 40 | 80 | 160 | order |
|----------|-----------------------|------------|------------|------------|------------|------------|-------------|
| 0.2 | 0.3 | 7.6512e-03 | 2.0123e-03 | 5.1534e-04 | 1.2745e-04 | 2.9646e-05 | 2.00 (2.00) |
| | 0.7 | 4.8203e-03 | 1.3214e-03 | 3.4625e-04 | 8.7435e-05 | 2.0347e-05 | 1.97 (2.00) |
| 0.6 | 0.3 | 2.3923e-03 | 6.2545e-04 | 1.5964e-04 | 3.9323e-05 | 9.0951e-06 | 2.01 (2.00) |
| | 0.7 | 2.2653e-03 | 5.9214e-04 | 1.5031e-04 | 3.7334e-05 | 8.6453e-06 | 2.00 (2.00) |
| 0.8 | 0.3 | 2.0211e-03 | 5.2732e-04 | 1.3312e-04 | 3.3132e-05 | 7.6553e-06 | 2.01 (2.00) |
| | 0.7 | 2.0132e-03 | 5.2513e-04 | 1.3321e-04 | 3.3051e-05 | 7.6249e-06 | 2.01 (2.00) |

5. Conclusions

In this study, we proposed a numerical method for solving stochastic semilinear subdiffusion equations driven by fractionally integrated additive noise. The temporal discretization relies on Mittag-Leffler integrators, while the spatial discretization is based on the finite element method. The effectiveness of the proposed method was demonstrated through illustrative examples that provided support for the theoretical analysis. In the next phase, we aim to enhance efficiency by incorporating the stochastic fractional system model. We also plan to preserve important physical properties and structures, such as positivity preservation, maximum principle, long time behavior, and investigate singular solutions. This includes studying the uniform L1 long time behavior of time discretization for time-fractional partial differential equations with non-smooth data. Additionally, we will develop a finite volume scheme for two-dimensional time-fractional stochastic Fokker-Planck equations on distorted meshes, which preserves the maximum principle.

Use of AI tools declaration

The authors declare they have not used Artificial Intelligence (AI) tools in the creation of this article.

Acknowledgments

This work is supported by the NSFC under Grant No. 12147115, the University Natural Science Research Project of Anhui Province under Grant No. 2023AH050254, the Fundamental Research Program of Shanxi Province under Grant No. 202103021224317, the Lvliang High Tech Research and Development Program under Grant No. 2023GXYP14, and the present job was also partially supported by the School Research Project of AUFE, Grant No. ACKYC21049.

Conflict of interest

The authors declare no conflict of interest.

References

1. B. Baeumer, M. Geissert, M. Kovács, Existence, uniqueness and regularity for a class of semilinear stochastic Volterra equations with multiplicative noise, *J. Differ. Equations*, **258** (2015), 535–554. <https://doi.org/10.1016/j.jde.2014.09.020>
2. Z. Q. Chen, K. H. Kim, P. Kim, Fractional time stochastic partial differential equations, *Stoch. Proc. Appl.*, **125** (2015), 1470–1499. <https://doi.org/10.1016/j.spa.2014.11.005>
3. C. M. Elliott, S. Larsson, Error estimates with smooth and nonsmooth data for a finite element method for the Cahn-Hilliard equation, *Math. Comp.*, **58** (1992), 603–630.
4. H. Fujita, T. Suzuki, *Evolution problems*, In Handbook of Numerical Analysis, North-Holland, Amsterdam, **2** (1991), 789–928.
5. B. T. Jin, Y. B. Yan, Z. Zhou, Numerical approximation of stochastic time-fractional diffusion, *ESAIM: M2AN*, **53** (2019), 1245–1268. <https://doi.org/10.1051/m2an/2019025>
6. M. Kovács, S. Larsson, F. Saedpanah, Mittag-Leffler Euler integrator for a stochastic fractional order equation with additive noise, *SIAM J. Numer. Anal.*, **58** (2020), 66–85. <https://doi.org/10.1137/18M1177895>
7. R. Kruse, *Strong and weak approximation of semilinear stochastic evolution equations*, Springer, Berlin, 2016.
8. Y. Hu, C. P. Li, Y. Yan, Weak convergence of the L1 scheme for a stochastic subdiffusion problem driven by fractionally integrated additive noise, *Appl. Numer. Math.*, **178** (2022), 192–215. <https://doi.org/10.1016/j.apnum.2022.04.004>
9. Y. Hu, Y. Yan, Shahzad Sarwar, Strong approximation of stochastic semilinear subdiffusion and superdiffusion driven by fractionally integrated additive noise, *Numer. Meth. Part. D. E.*, 2023. <https://doi.org/10.1002/num.23068>
10. X. C. Li, X. Y. Yang, Error estimates of finite element methods for stochastic fractional differential equations, *J. Comput. Math.*, **35** (2017), 346–362. <https://doi.org/10.4208/jcm.1607-m2015-0329>
11. G. D. Prato, J. Zabczyk, *Stochastic equations in infinite dimensions*, Cambridge University Press, Cambridge, 1992. <https://doi.org/10.1017/CBO9780511666223>
12. I. Podlubny, *Fractional differential equations*, Academic Press, San Diego, 1999.
13. J. W. Wang, X. X. Jiang, X. H. Yang, H. X. Zhang, A nonlinear compact method based on double reduction order scheme for the nonlocal fourth-order PDEs with Burgers' type nonlinearity, *J. Appl. Math. Comput.*, **70** (2024), 489–511. <https://doi.org/10.1007/s12190-023-01975-4>
14. X. J. Wang, R. S. Qi, A note on an accelerated exponential Euler method for parabolic SPDEs with additive noise, *Appl. Math. Lett.*, **46** (2015), 31–37. <https://doi.org/10.1016/j.aml.2015.02.001>
15. X. H. Yang, Z. M. Zhang, On conservative, positivity preserving, nonlinear FV scheme on distorted meshes for the multi-term nonlocal Nagumo-type equations, *Appl. Math. Lett.*, **150** (2024), 108972. <https://doi.org/10.1016/j.aml.2023.108972>
16. X. H. Yang, Z. M. Zhang, Q. Zhang, G. W. Yuan, Simple positivity-preserving nonlinear finite volume scheme for subdiffusion equations on general non-conforming distorted meshes, *Nonlinear Dyn.*, **108** (2022), 3859–3886. <https://doi.org/10.1007/s11071-022-07399-2>

17. Y. Yan, Galerkin finite element methods for stochastic parabolic partial differential equations, *SIAM J. Numer. Anal.*, **43** (2005), 1363–1384. <https://doi.org/10.1137/040605278>
18. V. Thomée, *Galerkin finite element methods for parabolic problems*, Springer-Verlag, Berlin, 2007.
19. S. G. Samko, A. A. Kilbas, O. I. Marichev, *Fractional integrals and derivatives: Theory and applications*, Gordon and Breach, New York, 1993.
20. H. X. Zhang, Y. Liu, X. H. Yang, An efficient ADI difference scheme for the nonlocal evolution problem in three-dimensional space, *J. Appl. Math. Comput.*, **69** (2023), 651–674. <https://doi.org/10.1007/s12190-022-01760-9>



AIMS Press

©2024 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (<http://creativecommons.org/licenses/by/4.0>)