



Research article

Numerical method for solving the subdiffusion differential equation with nonlocal boundary conditions

Murat A. Sultanov^{1,*}, Vladimir E. Misilov^{1,2,3} and Makhmud A. Sadybekov^{1,4}

¹ Department of Mathematics, Faculty of Natural Science, Khoja Akhmet Yassawi International Kazakh-Turkish University, B. Sattarhanov Street 29, Turkistan 160200, Kazakhstan

² Krasovskii Institute of Mathematics and Mechanics, Ural Branch of RAS, S. Kovalevskaya Street 16, Ekaterinburg 620108, Russia

³ Department of High Performance Computing Technologies, Institute of Natural Sciences and Mathematics, Ural Federal University, Mira Street 19, Ekaterinburg 620002, Russia

⁴ Institute of Mathematics and Mathematical Modeling, 125 Pushkin street, 050010 Almaty, Kazakhstan

* **Correspondence:** Email: murat.sultanov@ayu.edu.kz.

Abstract: This work was devoted to the construction of a numerical algorithm for solving the initial boundary value problem for the subdiffusion equation with nonlocal boundary conditions. For the case of not strongly regular boundary conditions, the well-known methods cannot be used. We applied an algorithm that consists of reducing the nonlocal problem to a sequential solution of two subproblems with local boundary conditions. The solution to the original problem was summed up from the solutions of the subproblems. To solve the subproblems, we constructed implicit difference schemes on the basis of the L1 formula for approximating the Caputo fractional derivative and central difference for approximating the space derivatives. Stability and convergence of the schemes were established. The Thomas algorithm was used to solve systems of linear algebraic equations. Numerical experiments were conducted to study the constructed algorithm. In terms of accuracy and stability, the algorithm performs well. The results of experiments confirmed that the convergence order of the method coincides with the theoretical one, $O(\tau^{2-\alpha} + h^2)$.

Keywords: differential equations; fractional derivative; subdiffusion equation; nonlocal problems; not strongly regular boundary conditions; boundary value problems; numerical algorithms

Mathematics Subject Classification: 35K20, 35R11, 65M06

1. Introduction

Fractional differential equations are widely used in various fields of science and engineering [1–3]. Among the popular fields, we can mention physics models, such as anomalous diffusion [4], viscoelastic media [5], propagation of acoustic waves in porous media [6], etc. Fractional calculus is used in control problems, to model dynamic systems in which the current state depends on the whole history [7], in signal processing [8, 9], in biology for the modeling of neurons [10], in epidemiology [11], in environmental sciences [12], and in economic models with memory [13].

Differential equations with nonlocal boundary conditions arise in problems of mathematical modeling of various processes such as heat transfer, chemical diffusion, hydrology, and biochemistry.

Boundary conditions of Samarskii-Ionkin type were originally considered in works [14] for the heat transfer problem in a tenuous plasma:

$$\begin{aligned}w_t(x, t) &= w_{xx}(x, t), \quad x < 0 < 1, \quad 0 < t < \infty, \\w(x, 0) &= \varphi(x), \quad w(0, t) = 0, \\&\int_0^1 w(x, t) dx = \text{const.}\end{aligned}$$

The last condition means that the total energy of the system is constant. This problem was further investigated in other works [15, 16]. The nonlocal condition was shown to be equivalent to

$$w_x(0, t) - w_x(1, t) = 0,$$

which means equal heat flows at the ends of the interval. The existence of the solution was proved.

Later, the heat equation with two-point boundary conditions of the general type was considered in [17]:

$$\begin{cases} a_1 w_x(0, t) + b_1 w_x(1, t) + a_0 w(0, t) + b_0 w(1, t) = 0, \\ c_1 w_x(0, t) + d_1 w_x(1, t) + c_0 w(0, t) + d_0 w(1, t) = 0. \end{cases} \quad (1.1)$$

Restricting the class of such conditions to strongly regular boundary conditions allows one to use the well-established theory for self-adjointed operators and well-known methods [18–20]. In the case of not strongly regular boundary conditions, the system of eigenfunctions of the spectral problem do not form a Riesz basis. Thus, these methods cannot be used. In this case, new methods or modifications of existing methods must be developed.

In [21], the heat problem with nonlocal boundary condition

$$\begin{cases} w_x(0, t) - w_x(1, t) - aw(1, t) = 0, \\ w(0, t) = 0, \end{cases} \quad 0 \leq t \leq T,$$

was considered. It was shown that these conditions are not strongly regular for $a \neq 0$. The existence, uniqueness, and stability of the solution was established.

Work [22] was devoted to solving an initial boundary value problem for a heat equation:

$$w_t(x, t) - w_{xx}(x, t) + q(x)w(x, t) = f(x, t)$$

with regular but not strongly regular boundary conditions of the general type (1.1). It was shown that for the case of the even potential function $q(x) = q(1 - x)$, the considered class of problems can always be reduced to the sequential solution of two similar problems with strongly regular boundary conditions. The proof does not depend on whether the system of eigenfunctions of a spectral problem forms a basis.

Note that the problem of establishing the basis properties of a system of eigenfunctions remains open so far. For special cases of not strongly regular boundary conditions, the basis property was confirmed in [23].

Problems for fractional differential equations with local and nonlocal boundary conditions have been the topic of many works [24–28]. There are many methods for numerical solutions of such problems, such as the finite difference methods [29], spectral methods [30, 31], etc. For various types of fractional derivatives, stability, convergence, and solvability of difference schemes were researched in [32].

One of the methods for solving the initial boundary value problem for a difference equation with nonlocal boundary conditions was developed in the works [22, 33] and applied to the nonlocal heat equation with not strongly regular boundary conditions. It consists in reducing the nonlocal difference problem to a sequential solution of two local difference problems.

In [34], a parallel numerical algorithm was constructed for solving the initial boundary value problem for the subdiffusion equation with the homogeneous Dirichlet condition. The stability of the difference scheme was established. Numerical experiments were performed to study the performance of parallel implementations. These results were used in [35] to develop a parallel numerical algorithm for solving the inverse problem of identifying the space-dependent source term in the two-dimensional fractional diffusion equation. To solve the inverse problem, a regularized iterative conjugate gradient method is used. At each iteration of the method, a finite difference scheme is used to solve the auxiliary direct initial boundary value problem.

Research on correctness and stability of direct and inverse problems of mathematical physics was performed in numerous works [36–41].

In this work, we construct a numerical method for solving the initial boundary value problem for the subdiffusion equation with nonlocal boundary conditions. We solve the problem by reducing it to two problems with local boundary conditions. The solution of the original problem is found as a sum of the solutions of the subproblems. These subproblems are solved by using implicit finite difference schemes. The schemes produce systems of linear algebraic equations (SLAE) that must be solved for each successive time step. To solve these SLAEs, we utilize the Thomas algorithm. We research the stability and convergence of the constructed difference schemes and obtain the stability estimates for the proposed method. To confirm the stability and convergence order of the developed method, we perform numerical experiments.

Our approach to the numerical solution of the nonlocal initial boundary value problem is used for the first time to solve the problem for a time-fractional equation. For simplicity, we consider the problem for the case of a zero initial value. The proposed numerical method and algorithm presented below may be applied for more generalized statements of the problem.

The rest of the article is structured as follows: Section 2 describes the statement of the initial boundary value problem for the subdiffusion differential equation with nonlocal boundary conditions. Section 3 is devoted to the numerical algorithm for solving the problem. It describes the construction

of the finite difference scheme for solving the subproblems. It also contains the results on researching the stability and convergence of the constructed finite difference scheme. Section 4 describes the performed numerical experiments and discusses their results. Section 5 concludes the article.

2. Problem statement

In this work, we consider numerical algorithms for solving the initial boundary value problem for the subdiffusion equation with nonlocal boundary conditions. Consider the equation

$$D_t^\alpha u(x, t) - u_{xx}(x, t) = f(x, t), \quad (2.1)$$

with initial conditions

$$u(x, 0) = 0, \quad 0 \leq x \leq 1, \quad (2.2)$$

and the following nonlocal boundary condition [21]

$$\begin{cases} u_x(0, t) - u_x(1, t) - au(1, t) = 0, \\ u(0, t) = 0, \end{cases} \quad 0 \leq t \leq T, \quad (2.3)$$

where $a > 0$.

Here, we consider the Caputo fractional derivative with order α in the form [42]

$$D_t^\alpha f(x, t) = \frac{1}{\Gamma(m - \alpha)} \int_0^t \frac{f^{(m)}(x, s)}{(t - s)^{\alpha - m + 1}} ds,$$

with $m = [\alpha] : \alpha \in (m - 1, m)$, $m \in N$, $x > 0$.

By attempting to solve the initial boundary value problem (2.1)–(2.3) by using the Fourier method, we obtain the spectral problem for operator l , which is defined by the differential expression and boundary conditions:

$$l(y) \equiv -y''(x) = \lambda y(x), \quad 0 < x < 1, \quad (2.4)$$

$$y'(0) = y'(1) + ay(1), \quad y(0) = 0. \quad (2.5)$$

Boundary conditions (2.5) are regular but not strongly regular. Therefore, the system of eigenfunctions of operator l is a complete system, but it does not form a basis in $L_2(0, 1)$. The spectral problem (2.4) and (2.5) has two series of eigenvalues:

$$\lambda_k^{(1)} = (2\pi k)^2, \quad k = 1, 2, \dots,$$

$$\lambda_k^{(2)} = (2\beta_k)^2, \quad k = 0, 1, 2, \dots,$$

where β_k are roots of the equation

$$\tan \beta = \frac{a}{2\beta}, \quad \beta > 0.$$

They satisfy the inequalities $\pi k < \beta_k < \pi k + \pi/2$, $k = 0, 1, 2, \dots$. For $\delta_k = \beta_k - \pi k$ with sufficiently large k , the following inequalities hold:

$$\frac{a}{2\pi k} \left(1 - \frac{1}{2\pi k}\right) < \delta_k < \frac{a}{2\pi k} \left(1 + \frac{1}{2\pi k}\right).$$

The eigenvalues of the problem are

$$y_k^{(1)}(x) = \sin(2\pi kx), \quad k = 1, 2, \dots, \quad y_k^{(2)}(x) = \sin(2\beta_k x), \quad k = 0, 1, 2, \dots$$

The system is almost normalized, but it does not even form a normal basis in $L_2(0, 1)$. Therefore, the direct application of the Fourier method is impossible.

3. Numerical algorithm for solving the problem

3.1. Reducing the nonlocal problem to two local problems

For the numerical solution of problem (2.1)–(2.3), we adapt the algorithm [22]. It was proposed for nonlocal problems with classical derivatives. The algorithm is based on reducing the nonlocal differential problem to the successive solving of two problems with Sturm-type boundary conditions. This method allows one to avoid studying correctness and stability of difference schemes for the original nonlocal problem, and replace it with studying local difference schemes for the auxiliary subproblems.

Let us represent the solution $u(x, t)$ as a sum of two functions $C(x, t)$ and $S(x, t)$, such as

$$u(x, t) = C(x, t) + S(x, t), \quad (3.1)$$

where

$$C(x, t) = \frac{u(x, t) + u(1-x, t)}{2}, \quad S(x, t) = \frac{u(x, t) - u(1-x, t)}{2}.$$

$C(x, t)$ is an even part of function $u(x, t)$ and $S(x, t)$ is an odd part of $u(x, t)$ on the interval $0 \leq x \leq 1$. Thus, these functions have the following properties:

$$C(x, t) = C(1-x, t), \quad S(x, t) = -S(1-x, t),$$

$$C_x(x, t) = -C_x(1-x, t), \quad S_x(x, t) = S_x(1-x, t).$$

Then, for the functions $C(x, t)$ and $S(x, t)$, we solve two initial boundary value problems.

Let us obtain the boundary conditions for the subproblems for $C(x, t)$ and $S(x, t)$. From (2.3), we have

$$\begin{cases} C_x(0, t) + S_x(0, t) - (C_x(1, t) + S_x(1, t)) - a(C(1, t) + S(1, t)) = 0, \\ C(0, t) + S(0, t) = 0, \end{cases} \quad 0 \leq t \leq T.$$

Using the properties of $C(x, t)$ and $S(x, t)$, we transform these conditions to contain only values for $x = 0$:

$$\begin{aligned} C_x(0, t) + S_x(0, t) - (-C_x(0, t) + S_x(0, t)) - a(C(0, t) - S(0, t)) &= 0, \\ S(0, t) &= -C(0, t), \end{aligned} \quad 0 \leq t \leq T;$$

$$\begin{aligned} C_x(0, t) - aC(0, t) &= 0, \\ S(0, t) &= -C(0, t), \end{aligned} \quad 0 \leq t \leq T.$$

Now, let us transform the conditions to contain only values for $x = 1$:

$$\begin{aligned} -C_x(1, t) + S_x(1, t) - \left(C_x(1, t) + S_x(1, t) \right) - a \left(C(1, t) + S(1, t) \right) &= 0, & 0 \leq t \leq T; \\ C(1, t) - S(1, t) &= 0, \end{aligned}$$

$$\begin{aligned} C_x(1, t) + aC(1, t) &= 0, \\ S(1, t) &= C(1, t), \end{aligned} \quad 0 \leq t \leq T.$$

Now, for $C(x, t)$, we have a problem with homogeneous boundary conditions of the third kind and a homogeneous initial condition:

$$\begin{cases} D_t^\alpha C(x, t) - C_{xx}(x, t) = f^C(x, t), \\ C(x, 0) = 0, & 0 \leq x \leq 1, \\ C_x(0, t) - aC(0, t) = 0, & 0 \leq t \leq T, \\ C_x(1, t) + aC(1, t) = 0, & 0 \leq t \leq T, \end{cases} \quad (3.2)$$

where $f^C(x, t) = (f(x, t) + f(1 - x, t))/2$.

For $S(x, t)$, we have the following problem with the Dirichlet boundary condition and the homogeneous initial condition:

$$\begin{cases} D_t^\alpha S(x, t) - S_{xx}(x, t) = f^S(x, t), \\ S(x, 0) = 0, & 0 \leq x \leq 1, \\ S(0, t) = -C(0, t), & 0 \leq t \leq T, \\ S(1, t) = C(0, t), & 0 \leq t \leq T, \end{cases} \quad (3.3)$$

where $f^S(x, t) = (f(x, t) - f(1 - x, t))/2$.

By construction, we can formulate the following statement:

Statement 1. *A solution of problem (2.1)–(2.3) can always be equivalently reduced to a sequential solution of two boundary value problems (3.2) and (3.3) with local boundary conditions.*

Thus, to solve problem (2.1)–(2.3), we need first to solve problem (3.2) for $C(x, t)$. Then, knowing the function $C(x, t)$, we solve problem (3.3), because its boundary conditions depend on $C(0, t)$. Finally, we obtain the sought function $u(x, t)$ by formula (3.1).

3.2. Discretization and difference schemes

Let us introduce regular grids for x and t with $M + 1$ and $N + 1$ points, respectively: $i = 0, \dots, M$, $h = 1/M$, $x_i = ih$, $n = 0, \dots, N$, $\tau = 1/N$, $t_n = n\tau$. Denote the values of grid functions at grid points as $u_{i,n} = u(x_i, t_n)$.

To approximate the Caputo fractional derivative of function $u(x, t)$ (and similarly for functions $C(x, t)$ and $S(x, t)$) at the time level n , we use the L1 formula [32, 43]:

$$\begin{aligned} D_t^\alpha(u_{i,n}) &\approx \sigma_{\alpha,\tau} \sum_{j=1}^n w_j^{(\alpha)} (u_{i,n-j+1} - u_{i,n-j}), \\ \sigma_{\alpha,\tau} &= \frac{1}{\Gamma(2 - \alpha)\tau^\alpha}, \quad w_j^{(\alpha)} = j^{1-\alpha} - (j-1)^{1-\alpha}, \\ n &= 1, \dots, N. \end{aligned} \quad (3.4)$$

Using formula (3.4) of order $O(\tau^{2-\alpha})$ and a central difference scheme of order $O(h^2)$ to construct an implicit difference scheme for Eq (3.2), we obtain

$$\sigma_{\alpha,\tau} \sum_{j=1}^n w_j^\alpha (C_{i,n-j+1} - C_{i,n-j}) = \frac{C_{i-1,n} - 2C_{i,n} + C_{i+1,n}}{h^2} + f_{i,n}^C + O(\tau^{2-\alpha} + h^2), \quad i = 0, \dots, M. \quad (3.5)$$

Omitting the small term $O(\tau^{2-\alpha} + h^2)$, we obtain the difference equations

$$\sigma_{\alpha,\tau} \sum_{j=1}^n w_j^\alpha (C_{i,n-j+1} - C_{i,n-j}) \approx \frac{C_{i-1,n} - 2C_{i,n} + C_{i+1,n}}{h^2} + f_{i,n}^C, \quad i = 0, \dots, M. \quad (3.6)$$

Then, let us rearrange

$$\begin{aligned} \sigma_{\alpha,\tau} (C_{i,n} - C_{i,n-1}) + \sigma_{\alpha,\tau} \sum_{j=2}^n w_j^{(\alpha)} (C_{i,n-j+1} - C_{i,n-j}) &= \frac{C_{i-1,n} - 2C_{i,n} + C_{i+1,n}}{h^2} + f_{i,n}^C, \\ -\frac{1}{h^2} C_{i-1,n} + \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) C_{i,n} - \frac{1}{h^2} C_{i+1,n} &= \sigma_{\alpha,\tau} \left(C_{i,n-1} - \sum_{j=2}^n w_j^{(\alpha)} (C_{i,n-j+1} - C_{i,n-j}) \right) + f_{i,n}^C. \end{aligned} \quad (3.7)$$

The boundary condition at the point $x_0 = 0$ is

$$C_x(0, t) - aC(0, t) = 0.$$

To approximate this condition with the order of approximation $O(h^2)$, let us obtain the derivative $C_x(0, t)$. Consider the Taylor series expansion

$$\frac{C(x_1, t) - C(x_0, t)}{h} = C_x(x_0, t) + \frac{h}{2} C_{xx}(x_0, t) + O(h^2).$$

Substituting $C_{xx}(x_0, t) = D_t^\alpha C(x_0, t) - f^C(x_0, t)$ from (3.2), we obtain

$$C_x(x_0, t) \approx \frac{C(x_1, t) - C(x_0, t)}{h} - \frac{h}{2} (D_t^\alpha C(x_0, t) - f^C(x_0, t)).$$

Substitute this expression for derivative $C_x(x_0, t)$ and formula (3.4) into the boundary condition

$$\frac{C_{1,n} - C_{0,t}}{h} - \frac{h}{2} \left(\sigma_{\alpha,\tau} \sum_{j=1}^n w_j^\alpha (C_{0,n-j+1} - C_{0,n-j}) - f_0^C \right) - aC_{0,n} = 0.$$

Let us rearrange the terms:

$$\begin{aligned} \frac{C_{1,n} - C_{0,t}}{h} - \frac{h}{2} \left(\sigma_{\alpha,\tau} (C_{0,n} - C_{0,n-1}) + \sigma_{\alpha,\tau} \sum_{j=2}^n w_j^{(\alpha)} (C_{0,n-j+1} - C_{0,n-j}) - f^C(0, t) \right) - aC_{0,n} &= 0, \\ \left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h} \right) C_{0,n} + \left(-\frac{2}{h^2} \right) C_{1,n} &= \sigma_{\alpha,\tau} \left(C_{0,n-1} - \sum_{j=2}^n w_j^{(\alpha)} (C_{0,n-j+1} - C_{0,n-j}) \right) + f_{0,n}^C. \end{aligned} \quad (3.8)$$

Similarly, for the boundary condition at the point $x_M = 1$, we obtain a difference equation:

$$\left(-\frac{2}{h^2}\right)C_{M-1,n} + \left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right)C_{M,n} = \sigma_{\alpha,\tau} \left(C_{M,n-1} - \sum_{j=2}^n w_j^{(\alpha)} (C_{M,n-j+1} - C_{M,n-j}) \right) + f_{M,n}^C. \quad (3.9)$$

Combining all the difference equations for points x_i , $i = 0, \dots, M$, we can form the system of $(M+1)$ linear algebraic equations in the matrix form

$$\begin{bmatrix} \left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right) & \left(-\frac{2}{h^2}\right) & \cdots & & \\ \left(-\frac{1}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) & \left(-\frac{1}{h^2}\right) & & \\ \vdots & \vdots & \ddots & \vdots & \\ & & \left(-\frac{1}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) & \left(-\frac{1}{h^2}\right) \\ & & \cdots & \left(-\frac{2}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right) \end{bmatrix} \cdot \begin{bmatrix} C_{0,n} \\ C_{1,n} \\ \vdots \\ C_{M-1,n} \\ C_{M,n} \end{bmatrix} = \begin{bmatrix} F_{0,n}^C \\ F_{1,n}^C \\ \vdots \\ F_{M-1,n}^C \\ F_M^C \end{bmatrix}, \quad (3.10)$$

where

$$F_{i,n}^C = \sigma_{\alpha,\tau} \left(C_{i,n-1} - \sum_{j=2}^n w_j^{(\alpha)} (C_{i,n-j+1} - C_{i,n-j}) \right) + f_{i,n}^C,$$

$$F_{i,1}^C = \sigma_{\alpha,\tau} C_{i,0} + f_{i,1}^C.$$

Similarly, we can obtain the difference equations for the function $S(x, t)$ and problem (3.3):

$$-\frac{1}{h^2}S_{i-1,n} + \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right)S_{i,n} - \frac{1}{h^2}S_{i+1,n} = \sigma_{\alpha,\tau} \left(S_{i,n-1} - \sum_{j=2}^n w_j^{(\alpha)} (S_{i,n-j+1} - S_{i,n-j}) \right) + f_{i,n}^S, \quad (3.11)$$

$$i = 1, \dots, M-1.$$

Due to the Dirichlet boundary conditions, we have a SLAE of smaller size, $M-1$:

$$\begin{bmatrix} \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) & \left(-\frac{1}{h^2}\right) & \cdots & & \\ \left(-\frac{1}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) & \left(-\frac{1}{h^2}\right) & & \\ \vdots & \vdots & \ddots & \vdots & \\ & & \left(-\frac{1}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) & \left(-\frac{1}{h^2}\right) \\ & & \cdots & \left(-\frac{1}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) \end{bmatrix} \cdot \begin{bmatrix} S_{1,n} \\ S_{2,n} \\ \vdots \\ S_{M-2,n} \\ S_{M-1,n} \end{bmatrix} = \begin{bmatrix} F_{1,n}^S - \frac{1}{h^2}C_{0,n} \\ F_{2,n}^S \\ \vdots \\ F_{M-2,n}^S \\ F_{M-1,n}^S + \frac{1}{h^2}C_{0,n} \end{bmatrix}, \quad (3.12)$$

where

$$F_{i,n}^S = \sigma_{\alpha,\tau} \left(S_{i,n-1} - \sum_{j=2}^n w_j^{(\alpha)} (S_{i,n-j+1} - S_{i,n-j}) \right) + f_{i,n}^S,$$

$$F_{i,1}^S = \sigma_{\alpha,\tau} S_{i,0} + f_{i,1}^S.$$

Thus, the numerical algorithm for solving problem (2.1)–(2.3) is the following:

At each sequential time level $n = 1, 2, \dots, N$:

- (1) Solve system (3.10) for $C_{i,n}$, $i = 0, \dots, M$.
- (2) Solve system (3.12) for $S_{i,n}$, $i = 1, \dots, M-1$.
- (3) Compute the values $u_{i,n} = C_{i,n} + S_{i,n}$, $i = 0, \dots, M$.

3.3. Solving the SLAEs

To solve SLAEs (3.10) and (3.12) with tridiagonal matrices, the Thomas algorithm (also known as the sweep method or elimination method) [44] is used. It is a direct method for systems with special matrices. For the tridiagonal systems (3.10) and (3.12), we denote the elements of the main diagonal of the matrices as q_i , and the elements of the lower and upper diagonals as p_i and r_i , respectively. The algorithm may be written as follows (in the case of system (3.10) for C):

- The forward elimination phase:

$$\begin{aligned}\alpha_0 &= r_0/q_0, \\ \beta_1 &= F_{0,n}^C/q_0, \\ \alpha_{i+1} &= r_i/(q_i - p_i\alpha_i), \quad i = 0, 1, \dots, M-1, M, \\ \beta_{i+1} &= (F_{i,n}^C + p_i\beta_i)/(q_i - p_i\alpha_i), \quad i = 1, 2, \dots, M, M+1.\end{aligned}\tag{3.13}$$

- The backward substitution phase:

$$\begin{aligned}C_{M,n} &= \beta_{M+1}, \\ C_{i,n} &= \alpha_{i+1}C_{i+1,n} + \beta_{i+1}, \quad i = M-1, M-2, \dots, 0.\end{aligned}\tag{3.14}$$

The conditions for correctness of this algorithm are the diagonal dominance in the matrix

$$|q_i| \geq |p_i| + |r_i|, \quad i = 0, 1, \dots, M.\tag{3.15}$$

This property obviously holds for matrices of systems (3.10) and (3.12).

3.4. Stability of the difference schemes

To justify the proposed algorithm, let us derive estimates of the stability of schemes (3.10) and (3.12) with respect to the initial data and the right-hand side.

Theorem 1. Suppose $\{C_{i,n} \mid 0 \leq i \leq M, 0 \leq n \leq N\}$ is the solution of the difference scheme (3.10). Then, the following estimate holds:

$$\|C_n\|_\infty \leq \frac{1}{\sigma_{\alpha,\tau} w_n^{(\alpha)}} \max_{1 \leq m \leq n} \|f_m\|_\infty, \quad 1 \leq n \leq N,\tag{3.16}$$

where $C_n = \{C_{0,n}, C_{1,n}, \dots, C_{M,n}\}$.

Proof. Rewrite (3.7) as follows:

$$\begin{aligned}-\frac{1}{h^2}C_{i-1,n} + \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right)C_{i,n} - \frac{1}{h^2}C_{i+1,n} &= \sigma_{\alpha,\tau} \left(C_{i,n-1} - \sum_{j=2}^n w_j^{(\alpha)} (C_{i,n-j+1} - C_{i,n-j}) \right) + f_{i,n}^C. \\ \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right)C_{i,n} &= \sigma_{\alpha,\tau} \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)}) C_{i,n-j} + \frac{1}{h^2}C_{i-1,n} + \frac{1}{h^2}C_{i+1,n} + f_{i,n}^C.\end{aligned}$$

Similarly, we can rewrite the difference equations for boundary conditions (3.8) and (3.9):

$$\left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right)C_{0,n} = \sigma_{\alpha,\tau} \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})C_{0,n-j} + \frac{2}{h^2}C_{1,n} + f_{0,n}^C,$$

$$\left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right)C_{M,n} = \sigma_{\alpha,\tau} \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})C_{M,n-j} + \frac{2}{h^2}C_{M-1,n} + f_{M,n}^C.$$

Suppose $\|C_n\|_\infty = |C_{l,n}|$ for some inner point $0 < l < M$. Note that $(w_j^{(\alpha)} - w_{j+1}^{(\alpha)}) > 0$ and $\sigma_{\alpha,\tau} > 0$. Taking the absolute value of both sides and using the triangle inequality, we obtain

$$\left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right)\|C_n\|_\infty \leq \sigma_{\alpha,\tau} \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})\|C_{n-j}\|_\infty + \frac{1}{h^2}\|C_n\|_\infty + \frac{1}{h^2}\|C_n\|_\infty + \|f_n^C\|_\infty.$$

We can simplify it to

$$\|C_n\|_\infty \leq \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})\|C_{n-j}\|_\infty + \frac{1}{\sigma_{\alpha,\tau}}\|f_n\|_\infty.$$

For the cases of $\|C_n\|_\infty = |C_{0,n}|$ or $\|C_n\|_\infty = |C_{M,n}|$, we can obtain the same inequality:

$$\left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right)\|C_n\|_\infty \leq \sigma_{\alpha,\tau} \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})\|C_{n-j}\|_\infty + \frac{2}{h^2}\|C_n\|_\infty + \|f_n^C\|_\infty,$$

$$\left(\sigma_{\alpha,\tau} + \frac{2a}{h}\right)\|C_n\|_\infty \leq \sigma_{\alpha,\tau} \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})\|C_{n-j}\|_\infty + \|f_n^C\|_\infty,$$

$$\|C_n\|_\infty \leq \frac{\sigma_{\alpha,\tau}}{\sigma_{\alpha,\tau} + \frac{2a}{h}} \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})\|C_{n-j}\|_\infty + \frac{1}{\sigma_{\alpha,\tau} + \frac{2a}{h}}\|f_n^C\|_\infty \leq \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})\|C_{n-j}\|_\infty + \frac{1}{\sigma_{\alpha,\tau}}\|f_n\|_\infty.$$

Next, let us use the mathematical induction method. Let

$$A_n = \frac{1}{\sigma_{\alpha,\tau}w_n^{(\alpha)}} \max_{1 \leq m \leq n} \|f_m\|_\infty.$$

Note that since $w_n^{(\alpha)} < w_{n-1}^{(\alpha)}$, $A_n > A_{n-1}$.

The base case. For $n = 1$, we have

$$\|C_1\|_\infty \leq \frac{1}{\sigma_{\alpha,\tau}}\|f_1\|_\infty \leq A_1.$$

The induction step. Suppose $\|C_k\|_\infty \leq A_k$ for $1 \leq k \leq n-1$. Consider the case of $k = n$.

$$\begin{aligned} \|C_n\|_\infty &\leq \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})\|C_{n-j}\|_\infty + \frac{1}{\sigma_{\alpha,\tau}}\|f_n\|_\infty \leq \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})A_{n-j} + \frac{1}{\sigma_{\alpha,\tau}}\|f_n\|_\infty \\ &\leq \sum_{j=1}^{n-1} (w_j^{(\alpha)} - w_{j+1}^{(\alpha)})A_n + \frac{1}{\sigma_{\alpha,\tau}}\|f_n\|_\infty \leq (w_1^{(\alpha)} - w_n^{(\alpha)})A_n + \frac{1}{\sigma_{\alpha,\tau}}\|f_n\|_\infty \\ &= A_n - w_n^{(\alpha)} \left(A_n - \frac{1}{\sigma_{\alpha,\tau}w_n^{(\alpha)}}\|f_n\|_\infty \right) \leq A_n. \end{aligned}$$

Thus, by induction, inequality (3.16) is true for $k = 1, 2, \dots, n$. \square

Similarly, we can obtain an estimate for S .

Theorem 2. Suppose $\{S_{i,n} \mid 0 \leq i \leq M, 0 \leq n \leq N\}$ is the solution of the difference scheme (3.12). Then, the following estimate holds:

$$\|S_n\|_\infty \leq \frac{1}{\sigma_{\alpha,\tau} W_n^{(\alpha)}} \max_{1 \leq m \leq n} \|f_m\|_\infty, \quad 1 \leq n \leq N. \quad (3.17)$$

Proof. The proof is similar to Theorem 1. For the boundary cases, we utilize the inequality (3.16). \square

Now, we can obtain an estimate for approximate solution $u_n = C_n + S_n$, $0 < n \leq N$, to the base problem (2.1)–(2.3).

Theorem 3. Suppose $\{C_{i,n}, S_{i,n} \mid 0 \leq i \leq M, 0 \leq n \leq N\}$ are the solutions of the difference schemes (3.10) and (3.12).

Then, the following estimate holds:

$$\|u_n\|_\infty \leq \frac{2}{\sigma_{\alpha,\tau} W_n^{(\alpha)}} \max_{1 \leq m \leq n} \|f_m\|_\infty, \quad 1 \leq n \leq N. \quad (3.18)$$

Proof. Let us utilize the triangle inequality. Since $u_n = C_n + S_n$, then $\|u_n\|_\infty \leq \|C_n\|_\infty + \|S_n\|_\infty$.

Now, by inequalities (3.16) and (3.17), we have

$$\|u_n\|_\infty \leq \frac{1}{\sigma_{\alpha,\tau} W_n^{(\alpha)}} \max_{1 \leq m \leq n} \|f_m\|_\infty + \frac{1}{\sigma_{\alpha,\tau} W_n^{(\alpha)}} \max_{1 \leq m \leq n} \|f_m\|_\infty = \frac{2}{\sigma_{\alpha,\tau} W_n^{(\alpha)}} \max_{1 \leq m \leq n} \|f_m\|_\infty.$$

\square

Thus, the difference schemes (3.10) and (3.12) are unconditionally stable with respect to the right-hand side.

3.5. Convergence of the difference schemes

In [32], it is shown that the approximation schemes used in (3.10) and (3.12) have a truncation error of order $O(\tau^{2-\alpha} + h^2)$, i.e., for the error $r_{i,n}$, there exists a constant $c > 0$, such that

$$|r_{i,n}| \leq c(\tau^{2-\alpha} + h^2), \quad 0 \leq i \leq M, 0 \leq n \leq N. \quad (3.19)$$

Theorem 4. Suppose $\{C_{i,n} \mid 0 \leq i \leq M, 0 \leq n \leq N\}$ is the exact solution of problem (3.2) and $\{\overline{C_{i,n}} \mid 0 \leq i \leq M, 0 \leq n \leq N\}$ is the solution of difference scheme (3.10). Consider the error

$$e^C_{i,n} = C_{i,n} - \overline{C_{i,n}}, \quad 0 \leq i \leq M, 0 \leq n \leq N.$$

Then, for some $c_1 > 0$, it holds that

$$\|e^C_n\|_\infty \leq c_1 T^\alpha \Gamma(1-\alpha) (\tau^{2-\alpha} + h^2), \quad 1 \leq n \leq N. \quad (3.20)$$

Proof. Let us subtract Eq (3.6) from (3.5), and, similarly, subtract boundary difference Eqs (3.8) and (3.9) from the boundary conditions. We obtain the difference scheme for error e_n^C :

$$\begin{bmatrix} \left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right) & \left(-\frac{2}{h^2}\right) & \cdots & & & \\ \left(-\frac{1}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) & \left(-\frac{1}{h^2}\right) & & & \\ \vdots & \vdots & \ddots & \vdots & & \\ & \left(-\frac{1}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2}\right) & \left(-\frac{1}{h^2}\right) & & \\ \cdots & \cdots & \left(-\frac{2}{h^2}\right) & \left(\sigma_{\alpha,\tau} + \frac{2}{h^2} + \frac{2a}{h}\right) & & \end{bmatrix} \cdot \begin{bmatrix} e_{0,n}^C \\ e_{1,n}^C \\ \vdots \\ e_{M-1,n}^C \\ e_{M,n}^C \end{bmatrix} = \begin{bmatrix} R_{0,n}^C \\ R_{1,n}^C \\ \vdots \\ R_{M-1,n}^C \\ R_M^C \end{bmatrix},$$

where

$$R_{i,n}^C = \sigma_{\alpha,\tau} \left(e_{i,n-1}^C - \sum_{j=2}^n w_j^{(\alpha)} (e_{i,n-j+1}^C - e_{i,n-j}^C) \right) + r_{i,n}^C,$$

$$R_{i,1}^C = \sigma_{\alpha,\tau} e_{i,0}^C + r_{i,1}^C.$$

We can apply Theorem 1 to this scheme and obtain

$$\|e_n^C\|_\infty \leq \frac{1}{\sigma_{\alpha,\tau} w_n^{(\alpha)}} \max_{1 \leq m \leq n} \|r_m\|_\infty, \quad 1 \leq n \leq N.$$

By (3.19), for some constant $c_1 > 0$, we have

$$\|e_n^C\|_\infty \leq \frac{c_1}{\sigma_{\alpha,\tau} w_n^{(\alpha)}} (\tau^{2-\alpha} + h^2).$$

For the coefficients $w_j^{(\alpha)}$, the following property holds (see [32, p. 32]):

$$(1 - \alpha)j^{-\alpha} < w_j^{(\alpha)} < (1 - \alpha)(j - 1)^{-\alpha}.$$

Thus,

$$\begin{aligned} \|e_n^C\|_\infty &\leq \frac{c_1}{\sigma_{\alpha,\tau}(1 - \alpha)n^{-\alpha}} (\tau^{2-\alpha} + h^2) \leq \frac{c_1 \Gamma(2 - \alpha) \tau^\alpha}{(1 - \alpha)n^{-\alpha}} (\tau^{2-\alpha} + h^2) \\ &\leq \frac{c_1 (1 - \alpha) \Gamma(1 - \alpha) \tau^\alpha}{(1 - \alpha)n^{-\alpha}} (\tau^{2-\alpha} + h^2) \leq c_1 t_n^\alpha \Gamma(1 - \alpha) (\tau^{2-\alpha} + h^2) \\ &\leq c_1 T^\alpha \Gamma(1 - \alpha) (\tau^{2-\alpha} + h^2), \quad 1 \leq n \leq N. \end{aligned}$$

□

Similarly, we can obtain an estimate for the error $e^{S_{i,n}} = S_{i,n} - \overline{S_{i,n}}$, $0 \leq i \leq M$, $0 \leq n \leq N$.

Theorem 5. Suppose $\{S_{i,n} | 0 \leq i \leq M, 0 \leq n \leq N\}$ is the exact solution of problem (3.3) and $\{\overline{S_{i,n}} | 0 \leq i \leq M, 0 \leq n \leq N\}$ is the solution of the difference scheme (3.12).

Then, for some $c_2 > 0$, it holds that

$$\|e_n^S\|_\infty \leq c_2 T^\alpha \Gamma(1 - \alpha) (\tau^{2-\alpha} + h^2), \quad 1 \leq n \leq N. \quad (3.21)$$

Now, consider the error between the solution of the original problem and the approximate solution obtained by the proposed numerical algorithm.

Theorem 6. Suppose $\{u_{i,n} | 0 \leq i \leq M, 0 \leq n \leq N\}$ is the exact solution of problem (2.1)–(2.3) and $\{\overline{C}_{i,n}, \overline{S}_{i,n} | 0 \leq i \leq M, 0 \leq n \leq N\}$ are the solutions of the difference schemes (3.10) and (3.12). Let

$$e_{i,n} = u_{i,n} - \overline{u}_{i,n}, \quad 0 \leq i \leq M, \quad 0 \leq n \leq N,$$

where $u_{i,n} = C_{i,n} + S_{i,n}$ and $\overline{u}_{i,n} = \overline{C}_{i,n} + \overline{S}_{i,n}$.

Then, for some $c_3 > 0$, it holds that

$$\|e_n\|_\infty \leq c_3 T^\alpha \Gamma(1 - \alpha) (\tau^{2-\alpha} + h^2), \quad 1 \leq n \leq N. \quad (3.22)$$

Proof. By Theorems 4 and 5, we have

$$\begin{aligned} \|e_n\|_\infty &= \|u_n - \overline{u}_n\|_\infty = \|C_n - \overline{C}_n + S_n - \overline{S}_n\|_\infty \leq \|C_n - \overline{C}_n\|_\infty + \|S_n - \overline{S}_n\|_\infty \\ &= \|e_n^C\|_\infty + \|e_n^S\|_\infty \leq c_1 T^\alpha \Gamma(1 - \alpha) (\tau^{2-\alpha} + h^2) + c_2 T^\alpha \Gamma(1 - \alpha) (\tau^{2-\alpha} + h^2) \\ &\leq c_3 T^\alpha \Gamma(1 - \alpha) (\tau^{2-\alpha} + h^2). \end{aligned}$$

The proof is complete. □

In the next section, we perform numerical experiments to confirm the convergence order.

4. Numerical experiments

The test problem is as follows:

$$D_t^\alpha u(x, t) - u_{xx}(x, t) = f(x, t),$$

$$u(x, 0) = 0, \quad 0 \leq x \leq 1,$$

$$\begin{cases} u_x(0, t) - u_x(1, t) + au(1, t) = 0, \\ u(0, t) = 0, \end{cases} \quad 0 \leq t \leq 1,$$

where

$$f(x, t) = \frac{\Gamma(m+1)}{\Gamma(m+1-\alpha)} t^{m-\alpha} [x^2(1-x)^2 - ax^2 + (2+a)x] + 2t^m [6x(1-x) - 1 + a].$$

The analytical solution is $u(x, t) = t^m [x^2(1-x)^2 - ax^2 + (2+a)x]$.

For subproblems $C(x, t)$ and $S(x, t)$, the right-hand terms are

$$f^C(x) = \frac{\Gamma(m+1)}{\Gamma(m+1-\alpha)} t^{m-\alpha} [x^2(1-x)^2 + ax(1-x) + 1] + 2t^m [6x(1-x) - 1 + a],$$

$$f^S(x) = \frac{\Gamma(m+1)}{\Gamma(m+1-\alpha)} t^{m-\alpha} [2x - 1].$$

The analytical solutions for $S(x, t)$ and $C(x, t)$ are

$$C(x, t) = t^m (x^2(1-x)^2 + ax(1-x) + 1),$$

$$S(x, t) = t^m (2x - 1).$$

4.1. Experiment 1: convergence in time

Experiment 1 was performed for the parameter values $\alpha = 0.5$, $m = 2$, and $a = 1$. To estimate the order of convergence in time, we solved the test problem for grid sizes $N = 64, 128, 256, 512$, and 1024 . To reduce the impact of the space step $h = 1/M$, we used the fixed value $M = 4096$. Thus, $\tau^{2-\alpha} > h^2$ for all values of $\tau = 1/N$.

Table 1 shows the absolute errors $\delta_1(\tau) = \|u(1, t) - \overline{u_\tau(1, t)}\|_\infty$ of the approximate solution $\overline{u_\tau(1, t)}$ obtained for step size τ , as well as the mixed relative errors $\delta_2(\tau) = \frac{\|u(1, t) - \overline{u_\tau(1, t)}\|_\infty}{1 + \|u(1, t)\|_\infty}$ for various dimensions N of the computational grid. It also contains the convergence orders $\frac{\log(\delta(\tau_1)/\delta(\tau_2))}{\log(\tau_1/\tau_2)}$ for both errors.

Table 1. Errors of the solutions and order of convergence in time for various grid sizes.

N	Absolute error δ_1	Order	Mixed relative error δ_2	Order
64	7.55×10^{-4}	—	2.52×10^{-4}	—
128	2.69×10^{-4}	1.49	8.98×10^{-5}	1.49
256	9.59×10^{-5}	1.49	3.20×10^{-5}	1.49
512	3.40×10^{-5}	1.49	1.13×10^{-5}	1.49
1024	1.20×10^{-5}	1.50	4.00×10^{-6}	1.50

4.2. Experiment 2: convergence in space

Experiment 2 was performed for the parameter values $\alpha = 0.5$, $m = 2$, and $a = 1$. To estimate the order of convergence in space, we solved the test problem for grid sizes $M = 64, 128, 256, 512$, and 1024 . To reduce the impact of the time step $\tau = 1/N$, we used the fixed value $N = 16,384$. Thus, $h^2 > \tau^{2-\alpha}$ for all values of $h = 1/M$.

Table 2 shows the absolute errors $\delta_1(h) = \|u(x, 1) - \overline{u_h(x, 1)}\|_\infty$ of the approximate solution $\overline{u_h(x, 1)}$ obtained for step size h , as well as the mixed relative errors $\delta_2(h) = \frac{\|u(x, 1) - \overline{u_h(x, 1)}\|_\infty}{1 + \|u(x, 1)\|_\infty}$ for various dimensions M of the computational grid. It also contains the convergence orders $\frac{\log(\delta(h_1)/\delta(h_2))}{\log(h_1/h_2)}$ for both errors.

Table 2. Errors of the solutions and order of convergence in space for various grid sizes.

M	Absolute error δ_1	Order	Mixed relative error δ_2	Order
64	3.19×10^{-4}	—	1.06×10^{-4}	—
128	7.97×10^{-5}	2.00	2.66×10^{-5}	2.00
256	1.98×10^{-5}	2.01	6.60×10^{-6}	2.01
512	4.80×10^{-6}	2.04	1.60×10^{-6}	2.04
1024	1.06×10^{-6}	2.18	3.53×10^{-7}	2.18

Figure 1 shows the exact and approximate solutions $u(x, 1)$ at the final instant $t = T = 1$, as well as the auxiliary solutions $S(x, 1)$ and $C(x, 1)$ obtained for the grid sizes $M = 4096$ and $N = 1024$. Solid graphs are exact solutions and dots are approximate ones.

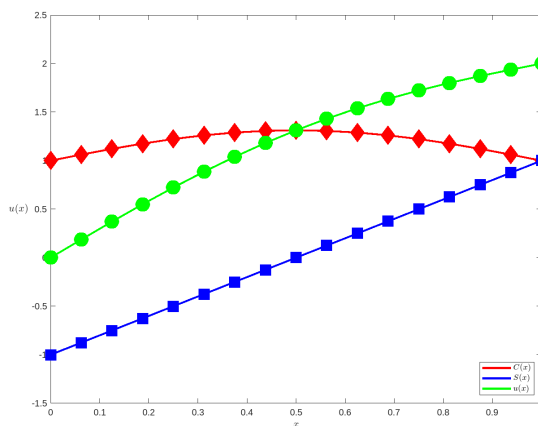


Figure 1. Solutions at the final instant $t = T = 1$. The blue graph is $S(x, 1)$, the red one is $C(x, 1)$, and the green one is $u(x, 1)$. Solid lines represent the exact solution and dots are the approximate solution.

Figures 2(a)–(c) show the surface plots for $C(x, t)$, $S(x, t)$, and $u(x, t)$ obtained for the grid sizes $M = 4096$ and $N = 1024$. Solid graphs are the exact solutions and dots are the approximate ones. Figure 2(d) shows the plot of error $\delta(x, t) = u(x, 1) - \overline{u(x, 1)}$.

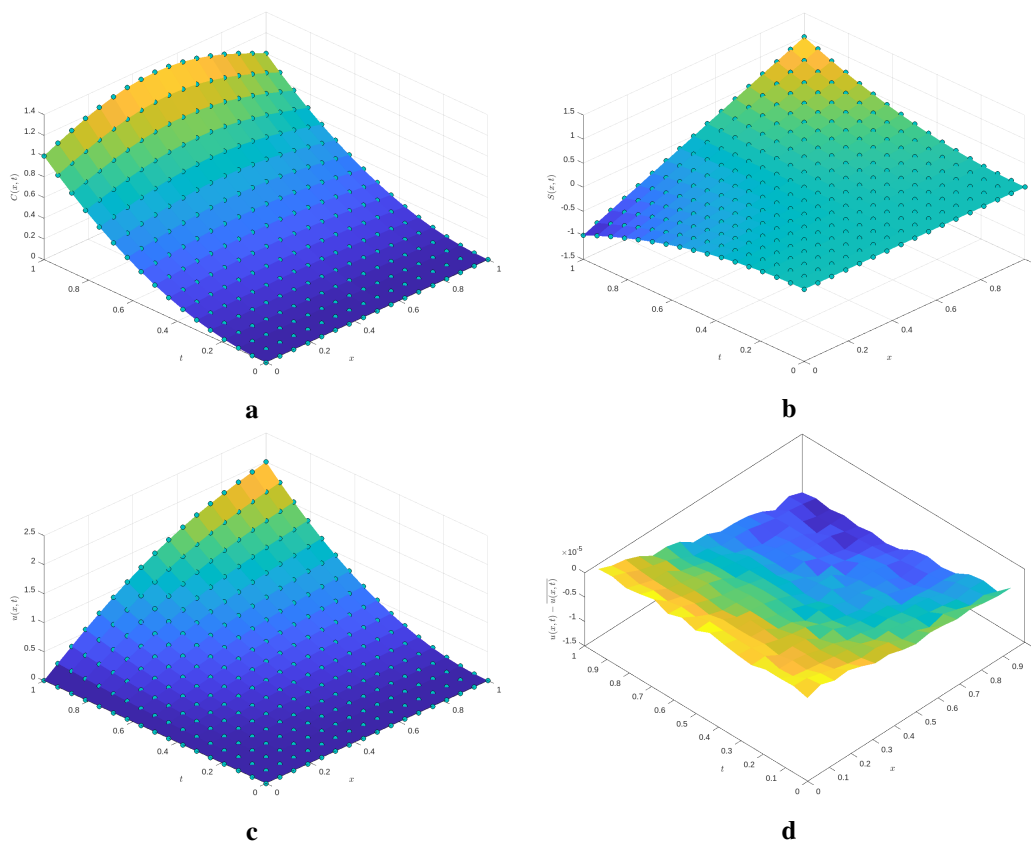


Figure 2. Results of numerical experiments for grid size $M = 4096$ and $N = 1024$. A solid surface is the exact solution and dots are the approximate one. (a) Surface plot of $C(x, t)$; (b) surface plot of $S(x, t)$; (c) surface plot of $u(x, t)$; (d) surface plot of error $\delta = (u(x, t) - u(x, 1))$.

Let us discuss the results of Experiments 1 and 2.

- (1) In Experiments 1 and 2, the finer the grid is for x or t , the lower the errors of the solutions. This confirms the stability and convergence of the difference schemes (3.10) and (3.12).
- (2) In Experiment 1, the computed convergence order is 1.5 in time for $\alpha = 0.5$.
- (3) In Experiment 2, the convergence order is 2 in space.
- (4) This coincides with the theoretical convergence order for the difference schemes based on the L1 formula for the Caputo time derivative and central difference scheme for space.
- (5) In the future, we plan to utilize higher-order approximations for the Caputo fractional derivative [45–47].

5. Conclusions

In this work, we have constructed a numerical algorithm for solving the initial boundary value problem for the subdiffusion equation with nonlocal boundary conditions. The algorithm consists of reducing the nonlocal problem to the successive solving of two subproblems with local boundary conditions. The solution of the original problem has been found as a sum of the solutions of the subproblems. To solve the subproblems, the implicit difference schemes have been constructed. For solving systems of linear algebraic equations, we have applied the Thomas algorithm. We have established the unconditional stability and obtained the convergence estimates of the difference schemes. We have conducted numerical experiments to study the constructed algorithm. The numerical results show the stability and convergence of the algorithm. The experiments confirm that the order of convergence of the proposed method coincides with the theoretical one.

In the future, we plan to adapt our approach to subdiffusion equations with more generalized formulations (1.1) of not strongly regular boundary conditions and nonzero initial conditions. We plan to adapt our approach to two-dimensional problems. Our approach for solving the initial boundary value problems may also be utilized in numerical algorithms for solving the inverse problems for nonlocal subdiffusion equations.

Author contributions

Murat A. Sultanov: Methodology, investigation, writing–review and editing, supervision, project administration, funding acquisition; Vladimir E. Misilov: Methodology, investigation, writing–original draft preparation, writing–review and editing; Makhmud A. Sadybekov: Methodology, validation, investigation, writing–review and editing, supervision. All authors have read and agreed to the published version of the manuscript.

Use of Generative-AI tools declaration

The authors declare they have not used Artificial Intelligence (AI) tools in the creation of this article.

Acknowledgments

This work was financially supported by the Ministry of Science and Higher Education of the Republic of Kazakhstan (project No. AP19676663).

Conflict of interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

References

1. H. Sun, Y. Zhang, D. Baleanu, W. Chen, Y. Chen, A new collection of real world applications of fractional calculus in science and engineering, *Commun. Nonlinear Sci.*, **64** (2018), 213–231. <https://doi.org/10.1016/j.cnsns.2018.04.019>
2. A. A. Kilbas, H. M. Srivastava, J. J. Trujillo, *Theory and applications of fractional differential equations*, Amsterdam: Elsevier, 2006.
3. J. T. Machado, V. Kiryakova, F. Mainardi, Recent history of fractional calculus, *Commun. Nonlinear Sci.*, **16** (2011), 1140–1153. <https://doi.org/10.1016/j.cnsns.2010.05.027>
4. R. Metzler, J. H. Jeon, A. G. Cherstvy, E. Barkai, Anomalous diffusion models and their properties: non-stationarity, non-ergodicity, and ageing at the centenary of single particle tracking, *Phys. Chem. Chem. Phys.*, **16** (2014), 24128–24164. <https://doi.org/10.1039/C4CP03465A>
5. M. J. Buckingham, On pore-fluid viscosity and the wave properties of saturated granular materials including marine sediments, *J. Acoust. Soc. Am.*, **122** (2007), 1486–1501. <https://doi.org/10.1121/1.2759167>
6. W. Chen, S. Hu, W. Cai, A causal fractional derivative model for acoustic wave propagation in lossy media, *Arch. Appl. Mech.*, **86** (2016), 529–539. <https://doi.org/10.1007/s00419-015-1043-2>
7. M. Du, Z. Wang, H. Hu, Measuring memory with the order of fractional derivative, *Sci. Rep.*, **3** (2013), 3431. <https://doi.org/10.1038/srep03431>
8. J. Zhang, Z. Wei, L. Xiao, Adaptive fractional-order multi-scale method for image denoising, *J. Math. Imaging Vis.*, **43** (2012), 39–49. <https://doi.org/10.1007/s10851-011-0285-z>
9. Q. Yang, D. Chen, T. Zhao, Y. Chen, Fractional calculus in image processing: a review, *Fract. Calc. Appl. Anal.*, **19** (2016), 1222–1249. <https://doi.org/10.1515/fca-2016-0063>
10. R. Magin, Fractional calculus in bioengineering, part 1, *Crit. Rev. Biomed. Eng.*, **32** (2004), 1–104. <https://doi.org/10.1615/CritRevBiomedEng.v32.i1.10>
11. C. Pinto, A. Carvalho, Fractional complex-order model for HIV infection with drug resistance during therapy, *J. Vib. Control*, **22** (2016), 2222–2239. <https://doi.org/10.1177/1077546315574964>
12. Y. Zhang, B. Baeumer, L. Chen, D. Reeves, H. Sun, A fully subordinated linear flow model for hillslope subsurface stormflow, *Water Resour. Res.*, **53** (2017), 3491–3504. <https://doi.org/10.1002/2016WR020192>

13. V. E. Tarasov, V. V. Tarasova, Long and short memory in economics: fractional-order difference and differentiation, *IRA-International Journal of Management and Social Sciences*, **5** (2016), 327–334. <https://dx.doi.org/10.21013/jmss.v5.n2.p10>
14. A. V. Bitsadze, A. A. Samarskii, Some elementary generalizations of linear elliptic boundary value problems, *Dokl. Akad. Nauk SSSR*, **185** (1969), 739–740.
15. V. A. Il'in, Existence of a reduced system of eigen- and associated functions for a nonselfadjoint ordinary differential operator, *Trudy Mat. Inst. Steklov.*, **142** (1976), 148–155.
16. N. I. Ionkin, The solution of a certain boundary value problem of the theory of heat conduction with a nonclassical boundary condition, *Differ. Uravn.*, **13** (1977), 294–304.
17. N. I. Ionkin, E. I. Moiseev, A problem for a heat equation with two-point boundary conditions, *Differ. Uravn.*, **15** (1979), 1284–1295.
18. V. P. Mikhailov, On Riesz bases in $L_2(0, 1)$, *Dokl. Akad. Nauk SSSR*, **144** (1962), 981–984.
19. M. A. Naimark, *Linear differential operators*, London: Harrap, 1967.
20. N. Dunford, J. Schwartz, *Linear operators, part III: spectral operators*, New York: John Wiley & Sons Inc, 1957.
21. A. Mokin, On a family of initial-boundary value problems for the heat equation, *Diff. Equat.*, **45** (2009), 126–141. <https://doi.org/10.1134/S0012266109010133>
22. M. A. Sadybekov, Initial-boundary value problem for a heat equation with not strongly regular boundary conditions, In: *Functional analysis in interdisciplinary applications*, Cham: Springer, 2017, 330–348. https://doi.org/10.1007/978-3-319-67053-9_32
23. A. S. Makin, On spectral decompositions corresponding to non-self-adjoint Sturm-Liouville operators, *Dokl. Math.*, **73** (2006), 15–18. <https://doi.org/10.1134/S1064562406010042>
24. A. V. Gulin, On the spectral stability in subspaces for difference schemes with nonlocal boundary conditions, *Diff. Equat.*, **49** (2013), 815–823. <https://doi.org/10.1134/S0012266113070045>
25. R. Ashurov, Y. Fayziev, On the nonlocal problems in time for time-fractional subdiffusion equations, *Fractal Fract.*, **6** (2022), 41. <https://doi.org/10.3390/fractalfract6010041>
26. E. Ozbilge, F. Kanca, E. Özbilge, Inverse problem for a time fractional parabolic equation with nonlocal boundary conditions, *Mathematics*, **10** (2022), 1479. <https://doi.org/10.3390/math10091479>
27. M. A. Almalahi, M. S. Abdo, S. K. Panchal, Periodic boundary value problems for fractional implicit differential equations involving Hilfer fractional derivative, *Probl. Anal. Issues Anal.*, **9** (2020), 16–44. <https://doi.org/10.15393/j3.art.2020.7410>
28. M. A. Almalahi, A. B. Ibrahim, A. Almutairi, O. Bazighifan, T. A. Aljaaidi, J. A. Awrejcewicz, A qualitative study on second-order nonlinear fractional differential evolution equations with generalized ABC operator, *Symmetry*, **14** (2022), 207. <https://doi.org/10.3390/sym14020207>
29. A. A. Alikhanov, Stability and convergence of difference schemes approximating a two-parameter nonlocal boundary value problem for time-fractional diffusion equation, *Comput. Math. Model.*, **26** (2015), 252–272. <https://doi.org/10.1007/s10598-015-9271-4>

30. N. B. Kerimov, M. I. Ismailov, Direct and inverse problems for the heat equation with a dynamic-type boundary condition, *IMA J. Appl. Math.*, **80** (2015), 1519–1533. <https://doi.org/10.1093/imamat/hxv005>
31. N. H. Tuan, N. A. Triet, N. H. Luc, N. D. Phuong, On a time fractional diffusion with nonlocal in time conditions, *Adv. Differ. Equ.*, **2021** (2021), 204. <https://doi.org/10.1186/s13662-021-03365-1>
32. Z. Sun, G. Gao, *Fractional differential equations: finite difference methods*, Berlin: De Gruyter, 2020. <https://doi.org/10.1515/9783110616064>
33. M. A. Sadybekov, I. N. Pankratova, Correct and stable algorithm for numerical solving nonlocal heat conduction problems with not strongly regular boundary conditions, *Mathematics*, **10** (2022), 3780. <https://doi.org/10.3390/math10203780>
34. M. A. Sultanov, E. N. Akimova, V. E. Misilov, Y. Nurlanuly, Parallel direct and iterative methods for solving the time-fractional diffusion equation on multicore processors, *Mathematics*, **10** (2022), 323. <https://doi.org/10.3390/math10030323>
35. E. N. Akimova, M. A. Sultanov, V. E. Misilov, Y. Nurlanuly, Parallel algorithm for solving the inverse two-dimensional fractional diffusion problem of identifying the source term, *Fractal Fract.*, **7** (2023), 801. <https://doi.org/10.3390/fractalfract7110801>
36. A. A. Alikhanov, A time-fractional diffusion equation with generalized memory kernel in differential and difference settings with smooth solutions, *Comput. Meth. Appl. Math.*, **17** (2017), 647–660. <https://doi.org/10.1515/cmam-2017-0035>
37. M. A. Sultanov, M. I. Akylbaev, R. Ibragimov, Conditional stability of a solution of a difference scheme for an ill-posed Cauchy problem, *Electron. J. Differ. Eq.*, **2018** (2018), 33.
38. K. Cao, D. Lesnic, M. I. Ismailov, Determination of the time-dependent thermal grooving coefficient, *J. Appl. Math. Comput.*, **65** (2021), 199–221. <https://doi.org/10.1007/s12190-020-01388-7>
39. N. H. Luc, D. Baleanu, R. P. Agarwal, L. D. Long, Identifying the source function for time fractional diffusion with non-local in time conditions, *Comp. Appl. Math.*, **40** (2021), 159. <https://doi.org/10.1007/10.1007/s40314-021-01538-y>
40. M. Slodička, Uniqueness for an inverse source problem of determining a space-dependent source in a non-autonomous time-fractional diffusion equation, *Fract. Calc. Appl. Anal.*, **23** (2020), 1702–1711. <https://doi.org/10.1515/fca-2020-0084>
41. E. N. Akimova, V. E. Misilov, M. A. Sultanov, Regularized gradient algorithms for solving the nonlinear gravimetry problem for the multilayered medium, *Math. Method. Appl. Sci.*, **45** (2022), 8760–8768. <https://doi.org/10.1002/mma.7012>
42. Y. Zhang, A finite difference method for fractional partial differential equation, *Appl. Math. Comput.*, **215** (2009), 524–529. <https://doi.org/10.1016/j.amc.2009.05.018>
43. I. Podlubny, *Fractional differential equations*, London: Academic Press, 1999.
44. A. Samarskii, E. Nikolaev, *Numerical methods for grid equations, volume I: direct methods*, Basel: Birkhäuser, 1989. <https://doi.org/10.1007/978-3-0348-9272-8>
45. Y. Dimitrov, R. Miryanov, V. Todorov, Asymptotic expansions and approximations for the Caputo derivative, *Comp. Appl. Math.*, **37** (2018), 5476–5499. <https://doi.org/10.1007/s40314-018-0641-3>

-
46. G. Gao, Z. Sun, H. Zhang, A new fractional numerical differentiation formula to approximate the Caputo fractional derivative and its applications, *J. Comput. Phys.*, **259** (2014), 33–50. <https://doi.org/10.1016/j.jcp.2013.11.017>
47. A. A. Alikhanov, A new difference scheme for the time fractional diffusion equation, *J. Comput. Phys.*, **280** (2015), 424–438. <https://doi.org/10.1016/j.jcp.2014.09.031>



AIMS Press

©2024 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (<https://creativecommons.org/licenses/by/4.0>)