



Research article

Qualitative results and numerical approximations of the (k, ψ) -Caputo proportional fractional differential equations and applications to blood alcohol levels model

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Abstract: The initial value problem in Cauchy-type under the (k, ψ) -Caputo proportional fractional operators was our focus in this paper. An extended Gronwall inequality and its properties were analyzed. The existence and uniqueness results were proven utilizing the fixed point theory of Banach's and Leray-Schauder's types. The qualitative analysis included results for Ulam-Mittag-Leffler stability, which was also investigated. Using a decomposition principle, a novel numerical technique was presented for the (k, ψ) -Caputo proportional fractional derivative operator. Finally, theoretical results were supported with numerical examples to demonstrate their practical application, especially to blood alcohol level problems.

Keywords: Gronwall inequality; existence and uniqueness; blood alcohol dynamic model; (k, ψ) -Caputo proportional fractional derivative; Ulam-Mittag-Leffler stability

Mathematics Subject Classification: 26A33, 26D10, 34A08, 34B10, 33E12

1. Introduction

Fractional calculus has been a popular topic among scholars for almost three centuries. It is a branch of mathematical analysis that deviates from classical calculus and extends the concepts of derivatives and integrals to fractional-order. Differential equations under fractional-order are known as fractional differential equations (FDEs), and they are used in various domains of real-world

phenomena problems. It can be distinguished based on various criteria, such as the linearity of the equation (linear or nonlinear). Its distinguishing traits, including flexibility, memory, and hereditary properties, make it a powerful tool for modeling, evaluating, and regulating complex systems in many disciplines, including applied science and engineering (see the books of Podlubny [1], Hilfer [2], Kilbas, *et al.* [3], and Diethelm [4]), and a comprehensive work on boundary value problems with fractional-order [5].

It is well known that fractional derivative operators (FDOs) are often defined in the sense of fractional integral operators (FIOs) under the gamma function. In the field of fractional calculus, various definitions of fractional derivatives have emerged, such as Riemann-Liouville (RL), Caputo, Hadamard, Erdélyi-Kober, Katugampola, Hilfer, proportional, and so on, and each has different uses. The RL-FIOs and RL-FDOs concerning a function ψ are studied as in [3]. These were used to develop the ψ -Hilfer-FDO [6]. One FDO gaining attention is the proportional FDO [7]. This operator has exponential functions in the kernels, which are more advantageous than other fractional operators. For work regarding qualitative theory on proportional FDEs, see [8–14] and references cited therein. In parallel, the k -gamma function defines and develops fractional calculus concepts, especially the k -RL-FIOs and the k -RL-FDOs [15, 16]. After that, in 2018, the (k, ψ) -RL fractional operators were proposed by [17]. Later, in 2021, Kucche and Mali proposed FDO under the (k, ψ) -Hilfer type [18], which attracted many scholars' curiosity. There are some other intriguing pieces, including those of Aljaaidi *et al.* [19], which proved some properties of the (k, ψ) -proportional fractional operator ((k, ψ) -PFO) and presented a new technique to investigate the k -Pólya-Szegő integral inequalities in 2021. In the same year, Boucenna *et al.* [20] studied the existence and uniqueness of results and a numerical technique for the Caputo proportional fractional Cauchy-type problem. While a numerical technique based on the decomposition formula was used to solve the fractional Caputo-Katugampola derivative operator [21], the analytical solutions of fractional differential equations are difficult and complex, so numerical techniques are popular for solving these equations [22–24]. Sudsutad *et al.* [25] developed the (k, ψ) -Hilfer-PFO, which cooperates the (k, ψ) -proportional fractional derivative of RL's and Caputo's senses ((k, ψ) -RL-PFDO / (k, ψ) -Caputo-PFDO). They constructed the Laplace transform concerning a function ψ of the proposed operator and applied it to solving the initial value Cauchy-type problems. In the extension, a Cauchy-type problem for FDEs is commonly defined as an initial value problem in which the goal is to find a function corresponding to the equation and the provided beginning conditions. Furthermore, numerous researchers have popularly expressed their robust findings by using a realistic application of fractional calculus in the context of diverse operators for various real-world issues such as medicine, engineering, electrical, science, and finance. For example, in 2019, Qureshi *et al.* [26] proposed the fractional modeling of the blood ethanol concentration using three fractional operators such as Caputo, Atangana-Baleanu, and Caputo-Fabrizio with the real data. In 2021, Norouzi and N'Guérékata [27] studied FDEs in the sense of ψ -Hilfer-FDO and used the financial crisis as an application. In 2022, Awadalla *et al.* [28] used a FDE in the sense of ψ -Caputo-FDO to study the model of drug concentration. In 2024, Wanassi and Torres [29] utilized the blood alcohol model as an application of the fractional analysis based on the ψ -Caputo-FDO. We refer the reader for more works in [30–36]. On the other hand, Ulam stability is a popularly efficient tool for ensuring that the approximate solutions generated by numerical methods remain close to the exact solutions. Numerous researchers have provided this to analyze the mathematical stability of solutions in various fields, including fractional calculus and control theory.

It was initially created by Ulam in 1940 [37]. The following year, Hyer designed the Ulam-Hyers stability [38]. Presently, stability in the context of the Ulam's type has developed into various stabilities, such as Ulam-Hyers-Rassias stability [39] and Ulam-Hyer-Mittag-Leffler (UH-ML) stability, which are extensions of Ulam stability that incorporates the Mittag-Leffler function [40], and so on.

Motivated by the works [9, 20, 25], our major goals are to investigate the qualitative results and numerical approximations for the following Cauchy-type problem under (k, ψ) -Caputo-PFDO

$$\begin{cases} {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} u(\tau) = f(\tau, u(\tau)), & 0 < \alpha \leq 1, \quad 0 < \rho \leq 1, \\ u(a) = u_a, \quad u_a \in \mathbb{R}, \quad \tau \in [a, b], \quad 0 \leq a < b < +\infty, \end{cases} \quad (1.1)$$

where ${}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi}$ is the (k, ψ) -Caputo-PFDO of α and $f \in C([a, b] \times \mathbb{R}, \mathbb{R})$. To the best of our knowledge, this problem has yet to be considered. Specifically, we provide an extended Gronwall inequality under the (k, ψ) -PFOs to establish bounds on solutions, which is crucial in stability analysis. We investigate the existence and uniqueness of results utilizing the standard fixed point theory of Banach's and Leray-Schauder's types. Moreover, various UH-MT stability results are studied utilizing nonlinear functional analysis techniques. In addition, we develop a novel numerical approach based on a decomposition formula for solving the Cauchy-type problems. In the end, a blood alcohol level problem is presented as an application to the proposed system. The following structure governs the remaining sections of this work: In Sect. 2, we introduce the fundamental principles and some properties of the (k, ψ) -PFDO and the (k, ψ) -PFIO. The required lemmas utilized throughout this paper are also presented. In addition, we look at the extended Gronwall inequality under the aforementioned operators. The qualitative results for a Cauchy-type problem are accomplished in the remaining sections. In Section 3, we investigate the existence of the solution using a fixed point theory of Leray-Schauder's type, while the uniqueness of the solution is proved using Banach's contraction mapping principle. Then, a variety of UH-ML stability results are established to ensure the results. In Section 4, we demonstrate the numerical approach based on a decomposition formula for solving the Cauchy-type problems under the (k, ψ) -Caputo-PFDO. Some illustrative examples, particularly the blood alcohol level problem, are provided to help the accuracy of the theoretical results found in Section 5. In the final section, we outline the course of our work.

2. Preliminaries

Suppose that $C([a, b], \mathbb{R})$ is the Banach space of the continuous function u on $[a, b]$ supplemented with the supremum norm $\|u\| = \sup_{\tau \in [a, b]} \{|u(\tau)|\}$. The space of the n -times absolutely continuous function u on $[a, b]$ is defined by $\mathcal{AC}^n([a, b], \mathbb{R}) = \{u : [a, b] \rightarrow \mathbb{R}; u^{(n-1)} \in \mathcal{AC}([a, b], \mathbb{R})\}$ and $L^p([a, b], \mathbb{R})$ is the Banach space of all Lebesgue measurable $g : [a, b] \rightarrow \mathbb{R}$ supplemented with $\|g\|_{L^p} < +\infty$. Assume that $\psi : [a, b] \rightarrow \mathbb{R}$ is a strictly increasing continuous function under $\psi'(\tau)$ not equal zero. For ease of calculation through this work, we provide a notation as follows:

$${}^{\rho}_{k} \Psi^{\frac{\alpha}{k}-1}_{\psi}(\tau, s) = e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1}. \quad (2.1)$$

Now, we recall some definitions and lemmas of the (k, ψ) -PFDO and the (k, ψ) -PFIO that will be applied in this work.

Definition 2.1. ([25]). Let $\alpha, k \in \mathbb{R}^+, 0 < \rho \leq 1$, and $f \in L^1([a, b], \mathbb{R})$. Then, the (k, ψ) -RL-PFIO of α of f is defined by

$${}_{a,k} \mathcal{I}^{\alpha,\rho;\psi} f(\tau) = \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_k^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s) ds,$$

where $\Gamma_k(z) = \int_0^\infty s^{z-1} e^{-\frac{s}{k}} ds, z \in \mathbb{C}, \operatorname{Re}(z) > 0$ and

$$\Gamma_k(z+k) = z \Gamma_k(z), \quad \Gamma_k(k) = 1, \quad \Gamma_k(z) = k^{\frac{z}{k}-1} \Gamma\left(\frac{z}{k}\right), \quad \Gamma(z) = \lim_{k \rightarrow 1} \Gamma_k(z). \quad (2.2)$$

Definition 2.2. ([25]). Let $\alpha, k \in \mathbb{R}^+, 0 < \rho \leq 1, f \in C([a, b], \mathbb{R}), \psi(\tau) \in C^n([a, b], \mathbb{R})$ with $\psi'(\tau) \neq 0$, and $n = 1, 2, \dots$, such that $n = \lfloor \alpha/k \rfloor + 1$. Then, the (k, ψ) -RL-PFDO for α of f is defined by

$${}_{a,k}^{\text{RL}} \mathcal{D}^{\alpha,\rho;\psi} f(\tau) = {}_k \mathcal{D}^{n,\rho;\psi} \left({}_{a,k} \mathcal{I}^{nk-\alpha,\rho;\psi} f(\tau) \right) = \frac{{}_k \mathcal{D}^{n,\rho;\psi}}{\rho^{\frac{nk-\alpha}{k}} k \Gamma_k(nk-\alpha)} \int_a^\tau \rho \Psi_k^{\frac{nk-\alpha}{k}-1}(\tau, s) \psi'(\tau) f(s) ds,$$

where ${}_k \mathcal{D}^{n,\rho;\psi} = \underbrace{{}_k \mathcal{D}^{\rho;\psi} \mathcal{D}^{\rho;\psi} \dots \mathcal{D}^{\rho;\psi}}_{n \text{ times}}$ and ${}_k \mathcal{D}^{1,\rho;\psi} f(\tau) = {}_k \mathcal{D}^{\rho;\psi} f(\tau) = (1-\rho)f(\tau) + k\rho \frac{f'(\tau)}{\psi'(\tau)}$.

Definition 2.3. ([25]). Let $\alpha, k \in \mathbb{R}^+, 0 < \rho \leq 1, f \in C^n(\mathcal{J}, \mathbb{R}), \psi(\tau) \in C^n([a, b], \mathbb{R})$ with $\psi'(\tau) \neq 0$, and $n \in 1, 2, \dots$, such that $n = \lfloor \alpha/k \rfloor + 1$. Then, the (k, ψ) -Caputo-PFDO for α of f is defined by

$${}_{a,k}^C \mathcal{D}^{\alpha,\rho;\psi} f(\tau) = {}_{a,k} \mathcal{I}^{nk-\alpha,\rho;\psi} \left({}_k \mathcal{D}^{n,\rho;\psi} f(\tau) \right) = \frac{1}{\rho^{\frac{nk-\alpha}{k}} k \Gamma_k(nk-\alpha)} \int_a^\tau \rho \Psi_k^{\frac{nk-\alpha}{k}-1}(\tau, s) \psi'(s) \left({}_k \mathcal{D}^{n,\rho;\psi} f(s) \right) ds.$$

Next, we give some important properties that are applied in this work.

Lemma 2.4. ([25]). Let $\alpha, \delta \in \mathbb{R}^+ \cup \{0\}, k, \eta \in \mathbb{R}^+, 0 < \rho \leq 1, \omega \in \mathbb{R}, \omega/k > -1$, and $n = \lfloor \omega/k \rfloor + 1$. Then,

- (i) ${}_{a,k} \mathcal{I}^{\alpha,\rho;\psi} \left[\rho \Psi_k^{\frac{\omega}{k}-1}(\tau, a) \right] = \frac{\Gamma_k(\omega)}{\rho^{\frac{\omega}{k}} \Gamma_k(\omega+\alpha)} \rho \Psi_k^{\frac{\omega+\alpha}{k}-1}(\tau, a)$.
- (ii) ${}_{a,k}^C \mathcal{D}^{\alpha,\rho;\psi} \left[\rho \Psi_k^{\frac{\omega}{k}-1}(\tau, a) \right] = \frac{\rho^{\frac{\omega}{k}} \Gamma_k(\omega) \rho \Psi_k^{\frac{\omega-\alpha}{k}-1}(\tau, a)}{\Gamma_k(\omega-\alpha) k}$. Particularly, for $m = 0, 1, \dots, n-1$, we obtain ${}_{a,k}^C \mathcal{D}^{\alpha,\rho;\psi} \left[\rho \Psi_k^m(\tau, a) \right] = 0$.
- (iii) ${}_{a,k} \mathcal{I}^{\alpha,\rho;\psi} \left({}_{a,k} \mathcal{I}^{\delta,\rho;\psi} f(\tau) \right) = {}_{a,k} \mathcal{I}^{\delta+\alpha,\rho;\psi} f(\tau) = {}_{a,k} \mathcal{I}^{\delta,\rho;\psi} \left({}_{a,k} \mathcal{I}^{\alpha,\rho;\psi} f(\tau) \right)$.
- (iv) ${}_{a,k}^C \mathcal{D}^{\omega,\rho;\psi} \left({}_{a,k} \mathcal{I}^{\eta,\rho;\psi} f(\tau) \right) = {}_{a,k} \mathcal{I}^{\eta-\omega,\rho;\psi} f(\tau)$, where $\omega \in \mathbb{R}^+, \omega < k$, and $\eta > nk$.
- (v) ${}_{a,k} \mathcal{I}^{\alpha,\rho;\psi} \left({}_{a,k}^C \mathcal{D}^{\alpha,\rho;\psi} f(\tau) \right) = f(\tau) - \sum_{i=1}^n \frac{\rho \Psi_k^{\frac{n-i}{k}-1}(\tau, a)}{(\rho k)^{n-i} (n-i)!} {}_k \mathcal{D}^{n-i,\rho;\psi} f(a), \alpha \in (n-1, n]$.

Now, we prove an extended Gronwall inequality under the (k, ψ) -PFOs, and its properties are analyzed in the below Theorem.

Theorem 2.5. (An Extended (k, ψ) -Proportional Fractional Gronwall Inequality). Suppose that $\alpha > 0, k > 0, \rho \in (0, 1]$, and $\psi \in C^1([a, b], \mathbb{R})$ is an increasing function so that $\psi'(\tau)$ not equal zero for every $\tau \in [a, b]$. Assume that the following hypotheses hold:

- (\mathcal{H}_1) The two non-negative functions $u(\tau)$ and $v(\tau)$ are locally integrable on $[a, b]$;
- (\mathcal{H}_2) The function $w(\tau)$ is a non-decreasing, non-negative, and continuous function defined on $[a, b]$ such that $w(\tau) \leq w^* \in \mathbb{R}$.

If

$$u(\tau) \leq v(\tau) + \frac{\Gamma_k(\alpha)}{k} w(\tau) {}_{a,k} \mathcal{I}^{\alpha, \rho; \psi} u(\tau), \quad (2.3)$$

then, for any $\tau \in [a, b]$, we obtain

$$u(\tau) \leq v(\tau) + \int_a^\tau \left[\sum_{n=1}^{\infty} \frac{[\Gamma_k(\alpha)w(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^{n+1} \Gamma_k(n\alpha)} \rho \Psi_{\psi}^{\frac{n\alpha}{k}-1}(\tau, s) \psi'(s) v(s) \right] ds. \quad (2.4)$$

Proof. First, we define an operator

$$\mathcal{A}u(\tau) = \frac{w(\tau)}{\rho^{\frac{\alpha}{k}} k^2} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) u(s) ds. \quad (2.5)$$

From (2.5), we have $u(\tau) \leq v(\tau) + \mathcal{A}u(\tau)$. By applying the monotonicity of the operator \mathcal{A} , we obtain inequalities as follows:

$$u(\tau) \leq v(\tau) + \mathcal{A}u(\tau) \leq v(\tau) + \mathcal{A}v(\tau) + \mathcal{A}^2 u(\tau) \leq \sum_{i=0}^2 \mathcal{A}^i v(\tau) + \mathcal{A}^3 u(\tau). \quad (2.6)$$

By applying iterative technique, $n = 1, 2, \dots$, which yields that $u(\tau) \leq \sum_{i=0}^{n-1} \mathcal{A}^i v(\tau) + \mathcal{A}^n u(\tau)$, $\tau \in [a, b]$, where $\mathcal{A}^0 v(\tau) = v(\tau)$. Next, we claim that

$$\mathcal{A}^n u(\tau) \leq \int_a^\tau \frac{[\Gamma_k(\alpha)w(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^{n+1} \Gamma_k(n\alpha)^k} \rho \Psi_{\psi}^{\frac{n\alpha}{k}-1}(\tau, s) \psi'(s) u(s) ds, \quad (2.7)$$

and $\mathcal{A}^n u(\tau) \rightarrow 0$ as $n \rightarrow \infty$ for any $\tau \in [a, b]$. If $n = 1$, we get that the inequality (2.7) holds. Next, assume that the inequality (2.7) is held under $n = m$, that is

$$\mathcal{A}^m u(\tau) \leq \int_a^\tau \frac{[\Gamma_k(\alpha)w(\tau)]^m}{\rho^{\frac{m\alpha}{k}} k^{m+1} \Gamma_k(m\alpha)^k} \rho \Psi_{\psi}^{\frac{m\alpha}{k}-1}(\tau, s) \psi'(s) u(s) ds.$$

If $n = m + 1$, using the induction procedure, we have

$$\begin{aligned} \mathcal{A}^{m+1} u(\tau) &= \mathcal{A}(\mathcal{A}^m u(\tau)) \\ &\leq \mathcal{A} \left(\int_a^\tau \frac{[\Gamma_k(\alpha)w(\tau)]^m}{\rho^{\frac{m\alpha}{k}} k^{m+1} \Gamma_k(m\alpha)^k} \rho \Psi_{\psi}^{\frac{m\alpha}{k}-1}(\tau, s) \psi'(s) u(s) ds \right) \\ &\leq \frac{w(\tau)}{\rho^{\frac{\alpha}{k}} k^2} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) \left(\int_a^s \frac{[\Gamma_k(\alpha)w(s)]^m}{\rho^{\frac{m\alpha}{k}} k^{m+1} \Gamma_k(m\alpha)^k} \rho \Psi_{\psi}^{\frac{m\alpha}{k}-1}(s, r) \psi'(r) u(r) dr \right) ds. \end{aligned}$$

Since the function $w(\tau)$ is a non-decreasing, $w(s) \leq w(\tau)$, for every $s \leq \tau$ and taking $z = (\psi(s) - \psi(r))/(\psi(\tau) - \psi(r))$, then

$$\begin{aligned} \mathcal{A}^{m+1} u(\tau) &\leq \left(\frac{w^{m+1}(\tau) (\Gamma_k(\alpha))^m}{\rho^{\frac{(m+1)\alpha}{k}} k^{m+3} \Gamma_k(m\alpha)} \right) \int_a^\tau \int_a^s e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} e^{\frac{\rho-1}{k\rho}(\psi(s)-\psi(r))} \\ &\quad \times (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} (\psi(s) - \psi(r))^{\frac{m\alpha}{k}-1} \psi'(s) \psi'(r) u(r) dr ds \\ &= \left(\frac{w^{m+1}(\tau) (\Gamma_k(\alpha))^m}{\rho^{\frac{(m+1)\alpha}{k}} k^{m+3} \Gamma_k(m\alpha)} \right) \int_a^\tau e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(r))} \psi'(r) u(r) \end{aligned}$$

$$\begin{aligned}
& \times \left(\int_r^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} (\psi(s) - \psi(r))^{\frac{m\alpha}{k}-1} \psi'(s) ds \right) dr \\
& = \left(\frac{w^{m+1}(\tau) (\Gamma_k(\alpha))^m}{\rho^{\frac{(m+1)\alpha}{k}} k^{m+2} \Gamma_k(m\alpha)} \right) \int_a^\tau \rho \Psi_{\psi}^{\frac{(m+1)\alpha}{k}-1}(\tau, s) \psi'(r) u(r) \left(\frac{1}{k} \int_0^1 (1-z)^{\frac{\alpha}{k}-1} z^{\frac{m\alpha}{k}-1} dz \right) dr \\
& = \left(\frac{w^{m+1}(\tau) (\Gamma_k(\alpha))^{m+1}}{\rho^{\frac{(m+1)\alpha}{k}} k^{m+2} \Gamma_k((m+1)\alpha)} \right) \int_a^\tau \rho \Psi_{\psi}^{\frac{(m+1)\alpha}{k}-1}(\tau, s) \psi'(r) u(r) dr.
\end{aligned}$$

Since $w \in C([a, b], \mathbb{R})$, then there is $w^* \in \mathbb{R}$ such that $w(\tau) \leq w^*$, for all $\tau \in [a, b]$, one has

$$\mathcal{A}^n u(\tau) \leq \left(\frac{w^* \Gamma_k(\alpha)}{\rho^{\frac{\alpha}{k}} k} \right)^n \frac{1}{k \Gamma_k(n\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{n\alpha}{k}-1}(\tau, s) \psi'(r) u(r) dr. \quad (2.8)$$

Since $0 \leq e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(r))} \leq 1$, and u is non-negative and locally integrable on $[a, b]$. Hence, u is bounded on $[a, b]$, and there is $\mathcal{M}_u \in \mathbb{R}$ such that $|u(\tau)| \leq \mathcal{M}_u$. The inequality (2.8) can be obtained

$$\mathcal{A}^n u(\tau) \leq \left(\frac{w^* \Gamma_k(\alpha) (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} k} \right)^n \frac{\mathcal{M}_u}{n\alpha \Gamma_k(n\alpha)}. \quad (2.9)$$

Applying the Stirling's formula, that is $n! \sim \sqrt{2\pi n} (n/e)^n$, into (2.9) with (2.2), we obtain that

$$\mathcal{A}^n u(\tau) \leq \mathcal{M}_u \sqrt{\frac{k}{2\pi\alpha}} \left(\frac{\theta^n}{n^{\frac{n\alpha}{k} + \frac{1}{2}}} \right), \quad \theta := \frac{w^* \Gamma_k(\frac{\alpha}{k}) (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} k^2} \left(\frac{ke}{\alpha} \right)^{\frac{\alpha}{k}}. \quad (2.10)$$

where $\theta, \rho, k, \alpha \in \mathbb{R}^+$. This yields that, if $n \rightarrow \infty$, we obtain that $\mathcal{A}^n u(\tau) \rightarrow 0$. Then,

$$u(\tau) \leq v(\tau) + \int_a^\tau \left[\sum_{n=1}^{\infty} \frac{[\Gamma_k(\alpha) w(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^{n+1} \Gamma_k(n\alpha)} \rho \Psi_{\psi}^{\frac{n\alpha}{k}-1}(\tau, s) \psi'(s) v(s) \right] ds.$$

The inequality (2.4) is achieved.

Corollary 2.6. Let $\alpha, k \in \mathbb{R}^+$, $0 < \rho \leq 1$, and $\psi \in C^1([a, b], \mathbb{R})$ be an increasing function such that $\psi'(\tau) \neq 0$, $\tau \in [a, b]$. Let $u(\tau)$ and $v(\tau)$ be two non-negative locally integrable functions on $[a, b]$, and $w(\tau) \equiv M \geq 0$. If

$$u(\tau) \leq v(\tau) + \frac{M \Gamma_k(\alpha)}{k} \mathcal{I}_{a,k}^{\alpha, \rho; \psi} u(\tau), \quad (2.11)$$

then,

$$u(\tau) \leq v(\tau) + \int_a^\tau \left[\sum_{n=1}^{\infty} \frac{[M \Gamma_k(\alpha)]^n}{\rho^{\frac{n\alpha}{k}} k^{n+1} \Gamma_k(n\alpha)} \rho \Psi_{\psi}^{\frac{n\alpha}{k}-1}(\tau, s) \psi'(s) v(s) \right] ds. \quad (2.12)$$

Corollary 2.7. Assume all conditions in Theorem 2.5 are held, and the function $v(\tau)$ is non-decreasing on $\tau \in [a, b]$. Hence, we obtain the following inequality:

$$u(\tau) \leq v(\tau) \mathbb{E}_{k,\alpha,k} \left((\rho^{\frac{\alpha}{k}} k)^{-1} \Gamma_k(\alpha) w(\tau) (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right), \quad (2.13)$$

where

$$\mathbb{E}_{k,\alpha,\beta}(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma_k(n\alpha + \beta)}, \quad z \in \mathbb{R}, \quad \alpha, \beta \in \mathbb{C}, \quad \operatorname{Re}(\alpha) > 0, \quad k > 0.$$

Proof. From the inequality (2.12) with the help of the non-decreasing property of $v(\tau)$, that is $v(s) \leq v(\tau)$ for any $\tau \in [a, b]$, we obtain that

$$\begin{aligned} u(\tau) &\leq v(\tau) + \int_a^\tau \left[\sum_{n=1}^{\infty} \frac{[\Gamma_k(\alpha)w(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^{n+1} \Gamma_k(n\alpha)^k} \rho \Psi_{\psi}^{\frac{n\alpha}{k}-1}(\tau, s) \psi'(s) v(s) \right] ds \\ &\leq v(\tau) \left(1 + \int_a^\tau \left[\sum_{n=1}^{\infty} \frac{[\Gamma_k(\alpha)z(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^{n+1} \Gamma_k(n\alpha)^k} \rho \Psi_{\psi}^{\frac{n\alpha}{k}-1}(\tau, s) \psi'(s) \right] ds \right). \end{aligned}$$

Since $0 \leq e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} \leq 1$, which yields that

$$\begin{aligned} u(\tau) &\leq v(\tau) \left(1 + \sum_{n=1}^{\infty} \frac{[\Gamma_k(\alpha)w(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^{n+1} \Gamma_k(n\alpha)^k} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{n\alpha}{k}-1} \psi'(s) ds \right) \\ &= v(\tau) \left(1 + \sum_{n=1}^{\infty} \frac{[\Gamma_k(\alpha)z(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^n \Gamma_k(n\alpha + k)} (\psi(\tau) - \psi(a))^{\frac{n\alpha}{k}} \right) \\ &= v(\tau) \sum_{n=0}^{\infty} \frac{[\Gamma_k(\alpha)w(\tau)]^n}{\rho^{\frac{n\alpha}{k}} k^n \Gamma_k(n\alpha + k)} (\psi(\tau) - \psi(a))^{\frac{n\alpha}{k}}. \end{aligned}$$

Applying Definition 2.7, the inequality (2.13) is obtained.

3. Existence and stability results

Here, we analyze the qualitative results for the proposed problem, including existence and uniqueness results and various results for Ulam's stability. First, we prove an integral equation is equivalent to the proposed problem (1.1).

Lemma 3.1. *Suppose that $f \in C([a, b] \times \mathbb{R}, \mathbb{R})$, $\alpha, \rho \in (0, 1]$ and $k \in \mathbb{R}^+$. Then, the proposed problem (1.1) can be stated equivalently as*

$$u(\tau) = u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, u(s)) ds. \quad (3.1)$$

Proof. Let u be a solution to the problem (1.1). Taking the operator ${}_{a,k}\mathcal{I}^{\alpha,\rho;\psi}$ into the proposed problem (1.1) and applying (v) in Lemma 2.4, which implies that

$$u(\tau) = u(a) e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, u(s)) ds. \quad (3.2)$$

By using $u(a) = u_a$ into (3.2) and inserting the obtained value into itself, we get the achieved (3.1).

On the other hand, by applying (ii) and (iv) in Lemma 2.4 into (3.1), it follows that

$$\begin{aligned} {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} u(\tau) &= {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} \left[u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} \right] + {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} \left({}_{a,k}\mathcal{I}^{\alpha,\rho;\psi} f(\tau, u(\tau)) \right) \\ &= u_a {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} \left[\rho \Psi_{\psi}^0(\tau, s) \right] + {}_{a,k}\mathcal{I}^{\alpha-\alpha,\rho;\psi} f(\tau, u(\tau)) \\ &= f(\tau, u(\tau)). \end{aligned}$$

The proof is done.

3.1. Existence and uniqueness results

From Lemma 3.1, we provide the operator $Q : C([a, b], \mathbb{R}) \rightarrow C([a, b], \mathbb{R})$ by

$$(Qu)(\tau) = u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho^{\frac{\alpha}{k}} \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, u(s)) ds. \quad (3.3)$$

The first result is based on the Banach's contraction mapping [41].

Theorem 3.2. *Suppose that $\alpha, \rho \in (0, 1]$, $k \in \mathbb{R}^+$, and $f \in C([a, b] \times \mathbb{R}, \mathbb{R})$. Suppose that*

(\mathcal{A}_1) *There is a positive real constant \mathcal{L} such that*

$$|f(\tau, u) - f(\tau, v)| \leq \mathcal{L} |u(\tau) - v(\tau)|, \quad \tau \in [a, b], \quad u, v \in \mathbb{R}.$$

Then, the proposed problem (1.1) has a unique solution on $[a, b]$, provided that

$$\frac{\mathcal{L}(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} < 1. \quad (3.4)$$

Proof. Assume that $\mathcal{B}_{r_1} := \{u \in C([a, b], \mathbb{R}) : \|u\| \leq r_1\}$ is a closed bounded and convex subset of $C([a, b], \mathbb{R})$, where the radius r_1 corresponds to the following condition

$$r_1 \geq \left(u_a + \frac{\mathcal{F}(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \right) \left(1 - \frac{\mathcal{L}(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \right)^{-1}, \quad \mathcal{F} = \sup_{\tau \in [a, b]} |f(\tau, 0)| < +\infty.$$

We prove that (i) $Q\mathcal{B}_{r_1}$ is bounded, i.e. $Q\mathcal{B}_{r_1} \subseteq \mathcal{B}_{r_1}$ and (ii) Q is a contraction.

Step (i). We show that $Q\mathcal{B}_{r_1} \subseteq \mathcal{B}_{r_1}$.

For any $u \in \mathcal{B}_{r_1}$ with the property of $0 \leq e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} \leq 1$ and the assumption (\mathcal{A}_1) , we have

$$\begin{aligned} |(Qu)(\tau)| &= \sup_{\tau \in \mathcal{J}} \left\{ u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho^{\frac{\alpha}{k}} \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, u(s)) ds \right\} \\ &\leq u_a + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} \psi'(s) |f(s, u(s))| ds \\ &\leq u_a + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} \psi'(s) [|f(s, u(s)) - f(s, 0)| + |f(s, 0)|] ds \\ &\leq u_a + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} \psi'(s) [\mathcal{L} \|u\| + \mathcal{F}] ds \\ &\leq u_a + \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} [\mathcal{L} r_1 + \mathcal{F}] \leq r_1. \end{aligned}$$

This implies that $Q\mathcal{B}_{r_1} \subseteq \mathcal{B}_{r_1}$.

Step (ii). We show that Q is a contraction.

Let $u, v \in C([a, b], \mathbb{R})$. Then, for every $\tau \in [a, b]$, we get

$$|(Qu)(\tau) - (Qv)(\tau)| \leq \sup_{\tau \in \mathcal{J}} \left\{ \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho^{\frac{\alpha}{k}} \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) [f(s, u(s)) - f(s, v(s))] ds \right\}$$

$$\begin{aligned}
&\leq \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} \psi'(s) |f(s, u(s)) - f(s, v(s))| ds \\
&\leq \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} \psi'(s) \mathcal{L} |u(s) - v(s)| ds \\
&\leq \frac{\mathcal{L} \|u - v\|}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} \psi'(s) ds \\
&= \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \mathcal{L} \|u - v\|.
\end{aligned}$$

Then, $\|Qu - Qv\| \leq \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \mathcal{L} \|u - v\|$. Since the condition (3.4) holds, which implies that Q has a contraction property. Hence, by [41], Q has a unique fixed point. Then, there is a unique solution of the proposed problem (1.1) on $[a, b]$.

Next, the existence result is achieved by utilizing the fixed point theory of Leray-Schauder's [42].

Theorem 3.3. *Suppose that*

(\mathcal{A}_2) *There are a continuous non-decreasing function $h : [a, \infty) \rightarrow (0, \infty)$ and a function $\phi \in C([a, b], \mathbb{R}^+)$ such that the following condition:*

$$|f(\tau, u(\tau))| \leq \phi(\tau)h(\|u\|), \quad \forall (\tau, u) \in [a, b] \times \mathbb{R}.$$

(\mathcal{A}_3) *There is a number $\mathcal{N} \in \mathbb{R}^+$ such that the following inequality:*

$$\mathcal{N} \left(\frac{\|\phi\|h(\mathcal{N}) (\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} + u_a \right)^{-1} > 1.$$

Then the proposed problem (1.1) has at least one solution on \mathcal{J} .

Proof. Assume that Q is given by (3.3). The technique is done in three phases.

Step (i). We show that Q maps bounded sets (balls) into bounded sets in $C([a, b], \mathbb{R})$.

Assume that $\mathcal{B}_{r_2} = \{u \in C([a, b], \mathbb{R}) : \|u\| \leq r_2\}$ is a bounded ball in $C([a, b], \mathbb{R})$ and $\|\phi\| = \sup_{\tau \in [a, b]} |\phi(\tau)|$. By applying the fact of $0 \leq e^{\frac{\rho-1}{k\rho}(\psi(\tau) - \psi(a))} \leq 1$ and (\mathcal{A}_2), for every $\tau \in [a, b]$, which yields that

$$\begin{aligned}
|(Qu)(\tau)| &\leq u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau) - \psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) |f(s, u(s))| ds \\
&\leq u_a + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}-1} \psi'(s) \phi(s) h(\|u\|) ds \\
&\leq u_a + \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \|\phi\| h(r_2) := \mathcal{K}.
\end{aligned}$$

Step (ii). We show that Q maps bounded sets into equicontinuous sets of $C([a, b], \mathbb{R})$.

Assume that $\tau_1, \tau_2 \in [a, b]$ with $\tau_1 < \tau_2$, and $u \in \mathcal{B}_{r_2}$. Hence

$$|(Qu)(\tau_2) - (Qu)(\tau_1)|$$

$$\begin{aligned}
&\leq u_a \left| e^{\frac{\rho-1}{k\rho}(\psi(\tau_2)-\psi(a))} - e^{\frac{\rho-1}{k\rho}(\psi(\tau_1)-\psi(a))} \right| + \left| \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^{\tau_2} \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau_2, s) \psi'(s) f(s, u(s)) ds \right. \\
&\quad \left. - \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^{\tau_1} \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau_1, s) \psi'(s) f(s, u(s)) ds \right| \\
&\leq u_a \left| e^{\frac{\rho-1}{k\rho}(\psi(\tau_2)-\psi(a))} - e^{\frac{\rho-1}{k\rho}(\psi(\tau_1)-\psi(a))} \right| + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_{\tau_1}^{\tau_2} \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau_2, s) \psi'(s) |f(s, u(s))| ds \\
&\quad + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^{\tau_1} \left| \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau_2, s) - \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau_1, s) \right| \psi'(s) |f(s, u(s))| ds \\
&\leq u_a \left| e^{\frac{\rho-1}{k\rho}(\psi(\tau_2)-\psi(a))} - e^{\frac{\rho-1}{k\rho}(\psi(\tau_1)-\psi(a))} \right| + \frac{\|\phi\| h(r_2)}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} (\psi(\tau_2) - \psi(\tau_1))^{\frac{\alpha}{k}} \\
&\quad + \frac{\|\phi\| h(r_2)}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \left| (\psi(\tau_2) - \psi(a))^{\frac{\alpha}{k}} - (\psi(\tau_2) - \psi(\tau_1))^{\frac{\alpha}{k}} - (\psi(\tau_1) - \psi(a))^{\frac{\alpha}{k}} \right|. \tag{3.5}
\end{aligned}$$

Clearly, the right-hand side of (3.5) tends to zero independently of $u \in \mathcal{B}_{r_2}$ as $\tau_2 \rightarrow \tau_1$. Then, by utilizing the Arzelá-Ascoli theorem, we get that $Q : C([a, b], \mathbb{R}) \rightarrow C([a, b], \mathbb{R})$ is completely continuous.

Step (iii). We show that there exists an open set $D \subseteq C([a, b], \mathbb{R})$ with $u \neq \lambda Q(u)$, $0 < \lambda < 1$ and $u \in \partial D$.

Assume that $u \in C([a, b], \mathbb{R})$ is the solution of $u = \lambda Qu$, $0 < \lambda < 1$. For any $\tau \in [a, b]$, we obtain that

$$|u(\tau)| = |\lambda(Qu)(\tau)| \leq u_a + \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \|\phi\| h(\|u\|).$$

Then

$$\|u\| \leq u_a + \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \|\phi\| h(\|u\|).$$

Consequently,

$$\|u\| \left(\frac{\|\phi\| h(\|u\|) (\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} + u_a \right)^{-1} \leq 1.$$

In view of (\mathcal{A}_3) , there exists \mathcal{N} such that $\|u\| \neq \mathcal{N}$. Let us take $D = \{u \in C([a, b], \mathbb{R}) : \|u\| < \mathcal{N}\}$ and $U = D \cap \mathcal{B}_{r_2}$.

Finally, $Q : \bar{U} \rightarrow C([a, b], \mathbb{R})$ is continuous and completely continuous. Using the choice of U , there is no $0 \in \partial U$ so that $u = \lambda(Qu)$ for some $0 < \lambda < 1$. Hence, by [42], we obtain that Q has a fixed point $u \in \bar{U}$, which is a solution of the problem (1.1).

3.2. Ulam-Mittag-Leffler stability results

Next, we establish various of UH-ML stability results for the proposed problem (1.1).

Definition 3.4. The proposed problem (1.1) is called UH-ML stable, if there is a number $\mathfrak{C}_f \in \mathbb{R}^+$ so that for any $\epsilon > 0$ and $w \in C([a, b], \mathbb{R}^+)$ of

$$\left| {}_{a,k}^C \mathfrak{D}^{\alpha, \rho; \psi} w(\tau) - f(\tau, w(\tau)) \right| \leq \epsilon, \quad \tau \in [a, b], \tag{3.6}$$

there exists $u \in C([a, b], \mathbb{R}^+)$ of the proposed problem (1.1) under the assumption

$$|w(\tau) - u(\tau)| \leq \mathfrak{C}_f \in \mathbb{E}_{k, \alpha, k} \left(\kappa_f (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right), \quad \kappa_f \geq 0, \quad \tau \in [a, b]. \tag{3.7}$$

Definition 3.5. The proposed problem (1.1) is called generalized UH-ML stable, if there is a function $\mathcal{G}_f \in C(\mathbb{R}^+, \mathbb{R}^+)$ via $\mathcal{G}_f(0) = 0$, such that for any $\epsilon > 0$ and $w \in C([a, b], \mathbb{R}^+)$ of

$$\left| {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} w(\tau) - f(\tau, w(\tau)) \right| \leq \mathcal{G}_f(\tau), \quad \tau \in [a, b], \quad (3.8)$$

there exists $u \in C([a, b], \mathbb{R}^+)$ of the proposed problem (1.1) under the assumption

$$|w(\tau) - u(\tau)| \leq \mathcal{G}_f(\epsilon) \mathbb{E}_{k,\alpha,k} \left(\kappa_f (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right), \quad \kappa_f \geq 0, \quad \tau \in [a, b]. \quad (3.9)$$

Definition 3.6. The proposed problem (1.1) is called Ulam-Hyers-Rassias-Mittag-Leffler (UHR-ML) stable with respect to another function $\Phi(\tau)$, if there is a number $\mathfrak{C}_{f_\Phi} \in \mathbb{R}^+$ so that for any $\epsilon > 0$ and $w \in C([a, b], \mathbb{R}^+)$ of

$$\left| {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} w(\tau) - f(\tau, w(\tau)) \right| \leq \epsilon \Phi(\tau), \quad \tau \in [a, b], \quad (3.10)$$

there exists $u \in C([a, b], \mathbb{R}^+)$ of the proposed problem (1.1) under the assumption

$$|w(\tau) - u(\tau)| \leq \mathfrak{C}_{f_\Phi} \epsilon \Phi(\tau) \mathbb{E}_{k,\alpha,k} \left(\kappa_{f_\Phi} (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right), \quad \kappa_{f_\Phi} \geq 0, \quad \tau \in [a, b]. \quad (3.11)$$

Definition 3.7. The proposed problem (1.1) is called generalized UHR-ML stable concerning function $\Phi(\tau)$ so that for any $\epsilon > 0$ and $w \in C([a, b], \mathbb{R}^+)$ of

$$\left| {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} w(\tau) - f(\tau, w(\tau)) \right| \leq \Phi(\tau), \quad \tau \in [a, b], \quad (3.12)$$

there exists $u \in C([a, b], \mathbb{R}^+)$ of the proposed problem (1.1) under the assumption

$$|w(\tau) - u(\tau)| \leq \mathfrak{C}_{f_\Phi} \Phi(\tau) \mathbb{E}_{k,\alpha,k} \left(\kappa_{f_\Phi} (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right), \quad \kappa_{f_\Phi} \geq 0, \quad \tau \in [a, b]. \quad (3.13)$$

Remark 3.8. Assume that $w \in C([a, b], \mathbb{R})$ is the solution of (3.6) if and only if there is $u_w \in C([a, b], \mathbb{R})$, which depends on w , so that (i). $|u_w(\tau)| \leq \epsilon$, $\tau \in [a, b]$; (ii). ${}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} w(\tau) = f(\tau, w(\tau)) + u_w(\tau)$, $\tau \in [a, b]$.

Remark 3.9. Assume that $w \in C([a, b], \mathbb{R})$ is the solution of (3.10) if and only if there is $v_w \in C([a, b], \mathbb{R})$, that is depends on w , such that (i). $|v_w(\tau)| \leq \epsilon \Phi(\tau)$, $\tau \in [a, b]$; (ii). ${}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} z(\tau) = f(\tau, w(\tau)) + v_w(\tau)$, $\tau \in [a, b]$.

Theorem 3.10. Suppose that $f \in C([a, b] \times \mathbb{R}, \mathbb{R})$, (\mathcal{A}_1) , and (3.2) hold. Then, the proposed problem (1.1) is UH-ML stable and consequently generalized UH-ML stable on $[a, b]$.

Proof. Let $\epsilon > 0$ and $w \in C([a, b], \mathbb{R})$ be a solution of (3.6). From (ii) in Remark 3.8, we have

$$\begin{cases} {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} w(\tau) = f(\tau, w(\tau)) + u_w(\tau), & \tau \in (a, b], \\ w(a) = w_a, & w_a \in \mathbb{R}. \end{cases} \quad (3.14)$$

Applying Lemma 3.1, the solution of the problem (3.14) is

$$\begin{aligned} w(\tau) &= w_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho {}^{\rho} \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, w(s)) ds \\ &+ \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho {}^{\rho} \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) u_w(s) ds. \end{aligned} \quad (3.15)$$

Let $u \in C([a, b], \mathbb{R})$ be a solution of (1.1). Hence,

$$u(\tau) = u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho^{\frac{\alpha}{k}} \Psi_\psi^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, u(s)) ds. \quad (3.16)$$

By applying the property $|u - v| \leq |u| + |v|$, $\tau \in \mathcal{J}$ with (3.15)-(3.16), one has

$$\begin{aligned} |w(\tau) - u(\tau)| &\leq \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho^{\frac{\alpha}{k}} \Psi_\psi^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) |f(s, w(s)) - f(s, u(s))| ds \\ &\quad + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho^{\frac{\alpha}{k}} \Psi_\psi^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) |u_w(s)| ds. \end{aligned} \quad (3.17)$$

By using the property $0 \leq e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} \leq 1$, (\mathcal{H}_1) , and (i) in Remark 3.8, the inequality (3.17) can be written

$$|w(\tau) - u(\tau)| \leq \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \epsilon + \frac{\Gamma_k(\alpha)}{k} \left(\frac{k\mathcal{L}}{\Gamma_k(\alpha)} \right)_{a,k} \mathcal{I}^{\alpha, \rho; \psi} |w(\tau) - u(\tau)|. \quad (3.18)$$

From Theorem 2.5 and Corollary 2.7, which implies that

$$|w(\tau) - u(\tau)| \leq \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \epsilon \mathbb{B}_{k, \alpha, k} \left(\mathcal{L} \rho^{-\frac{\alpha}{k}} (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right).$$

By setting a constant $\mathfrak{C}_f := \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)}$ and $\kappa_f := \mathcal{L} \rho^{-\frac{\alpha}{k}}$, which yields that

$$|w(\tau) - u(\tau)| \leq \mathfrak{C}_f \epsilon \mathbb{B}_{k, \alpha, k} \left(\kappa_f (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right).$$

Hence, the proposed problem (1.1) is UH-ML stable. In addition, by setting $\mathcal{G}_f(\epsilon) = \mathfrak{C}_f \epsilon$ under the condition $\mathcal{G}_f(0) = 0$, then

$$|w(\tau) - u(\tau)| \leq \mathcal{G}_f(\epsilon) \mathbb{B}_{k, \alpha, k} \left(\kappa_f (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right).$$

Therefore, the solution of the proposed problem (1.1) is generalized UH-ML stable.

Now, we give the required assumption that is used in Theorem 3.11.

(\mathcal{P}_1) Suppose that $\Phi \in C([a, b], \mathbb{R})$ is a non-decreasing function. There exists $\chi_\Phi > 0$ so that

$$\frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho^{\frac{\alpha}{k}} \Psi_\psi^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) \Phi(s) ds \leq \chi_\Phi \Phi(\tau), \quad \tau \in [a, b].$$

Theorem 3.11. Suppose that $f \in C([a, b] \times \mathbb{R}, \mathbb{R})$, (\mathcal{A}_1) , and (3.2) hold. Then, the proposed problem (1.1) is UHR-ML stable and consequently generalized UHR-ML stable on $[a, b]$.

Proof. Let $\epsilon > 0$ and $w \in C([a, b], \mathbb{R})$ be the solution of (3.6). From (ii) in Remark 3.9, it follows from

$$\begin{cases} {}^C_{a,k} \mathfrak{D}^{\alpha, \rho; \psi} w(\tau) = f(\tau, w(\tau)) + v_w(\tau), & \tau \in (a, b], \\ w(a) = w_a, & w_a \in \mathbb{R}. \end{cases} \quad (3.19)$$

From Lemma 3.1, the solution of the problem (3.19) is defined by

$$\begin{aligned} w(\tau) &= w_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, w(s)) ds \\ &\quad + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) v_w(s) ds. \end{aligned} \quad (3.20)$$

Let $u \in C([a, b], \mathbb{R})$ be the solution of (1.1). It follows that

$$u(\tau) = u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) f(s, u(s)) ds. \quad (3.21)$$

By applying the property $|u - v| \leq |u| + |v|$, for $\tau \in [a, b]$ with (3.20)-(3.21), which yields that

$$\begin{aligned} |w(\tau) - u(\tau)| &\leq \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) |f(s, w(s)) - f(s, u(s))| ds \\ &\quad + \frac{1}{\rho^{\frac{\alpha}{k}} k \Gamma_k(\alpha)} \int_a^\tau \rho \Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, s) \psi'(s) |v_w(s)| ds. \end{aligned} \quad (3.22)$$

By using $0 \leq e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(a))} \leq 1$, (\mathcal{A}_1) , (\mathcal{P}_1) , and (i) in Remark 3.9, the inequality (3.22) can be written

$$|w(\tau) - u(\tau)| \leq \epsilon \chi_\Phi \Phi(\tau) + \frac{\Gamma_k(\alpha)}{k} \left(\frac{k\mathcal{L}}{\Gamma_k(\alpha)} \right)_{a,k} \mathcal{I}^{\alpha, \rho; \psi} |w(\tau) - u(\tau)|. \quad (3.23)$$

Applying Theorem 2.5 and Corollary 2.7, which implies that

$$|w(\tau) - u(\tau)| \leq \epsilon \chi_\Phi \Phi(\tau) \mathbb{E}_{k, \alpha, k} \left(\mathcal{L} \rho^{-\frac{\alpha}{k}} (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right).$$

By choosing $\mathfrak{C}_{f_\Phi} := \chi_\Phi$ and $\kappa_f := \mathcal{L} \rho^{-\frac{\alpha}{k}}$, then

$$|w(\tau) - u(\tau)| \leq \mathfrak{C}_{f_\Phi} \epsilon \Phi(\tau) \mathbb{E}_{k, \alpha, k} \left(\kappa_f (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right).$$

Thus, the proposed problem (1.1) is UHR-ML stable. Additionally, by setting $\epsilon = 1$, one has

$$|w(\tau) - u(\tau)| \leq \mathfrak{C}_{f_\Phi} \Phi(\tau) \mathbb{E}_{k, \alpha, k} \left(\kappa_f (\psi(\tau) - \psi(s))^{\frac{\alpha}{k}} \right).$$

Then, the solution of the proposed problem (1.1) is generalized UHR-ML stable.

4. Numerical approach

Now, we construct an approximation form for the (k, ψ) -PFDO in Caputo's sense under $\alpha \in (0, 1]$ of a function $u(\tau)$. We may generate a sequence of $N + 1$ equations with $N + 1$ conditions for a given (k, ψ) proportional fractional Cauchy-type problem under this tool. A sequence (u_N) of solutions to such systems eventually leads to the solution of the proposed problem.

Theorem 4.1. Assume that $N = 1, 2, \dots$, and $u \in \mathcal{AC}^2([a, b], \mathbb{R})$. Let

$$A_N = \frac{1}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \sum_{i=0}^N \frac{\Gamma\left(i + \frac{\alpha}{k} - 1\right)}{i! \Gamma\left(\frac{\alpha}{k} - 1\right)}, \quad (4.1)$$

$$B_{N,i} = \frac{\Gamma\left(i + \frac{\alpha}{k} - 1\right)}{\rho^{1-\frac{\alpha}{k}}(i-1)!\Gamma_k(2k-\alpha)\Gamma\left(\frac{\alpha}{k}-1\right)}, \quad i = 1, 2, \dots, N, \quad (4.2)$$

and let $\mathcal{V}_i : [a, b] \rightarrow \mathbb{R}$ be a function, which is given by

$$\mathcal{V}_i(\tau) = \int_a^\tau (\psi(s) - \psi(a))^{i-1} \psi'(s) e^{\frac{1-\rho}{k\rho}\psi(s)} {}_k\mathfrak{D}^{\rho;\psi} u(s) ds, \quad i = 1, 2, \dots, N. \quad (4.3)$$

Hence,

$$\begin{aligned} {}_C^C_{a,k}\mathfrak{D}^{\alpha,\rho;\psi} u(\tau) &\approx A_N(\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}} {}_k\mathfrak{D}^{\rho;\psi} u(\tau) \\ &\quad - e^{\frac{\rho-1}{k\rho}\psi(\tau)} \sum_{i=1}^N B_{N,i}(\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}-i} \mathcal{V}_i(\tau) + E_{tr}(\tau), \end{aligned} \quad (4.4)$$

where ${}_k\mathfrak{D}^{\rho;\psi} u(\tau) = (1-\rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)}$ and $\lim_{N \rightarrow \infty} E_{tr}(\tau) = 0$, $\tau \in [a, b]$.

Proof. By using Definition 2.3 and ${}_k\mathfrak{D}^{\rho;\psi} u(\tau) = (1-\rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)}$, for any $\alpha \in (0, 1]$, we have

$$\begin{aligned} &{}_C^C_{a,k}\mathfrak{D}^{\alpha,\rho;\psi} u(\tau) \\ &= \frac{1}{k\rho^{1-\frac{1}{k}}\Gamma_k(k-\alpha)} \int_a^\tau {}_k\Psi_{\psi}^{-\frac{\alpha}{k}}(\tau, s) \psi'(s) ({}_k\mathfrak{D}^{\rho;\psi} u(s)) ds \\ &= \frac{1}{k\rho^{1-\frac{\alpha}{k}}\Gamma_k(k-\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{-\frac{\alpha}{k}} \psi'(s) e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} \left[(1-\rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] ds. \end{aligned} \quad (4.5)$$

Changing new variables

$$x(s) = e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} \left[(1-\rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] \quad \text{and} \quad y'(s) = (\psi(\tau) - \psi(s))^{-\frac{\alpha}{k}} \psi'(s),$$

and by helping the integrating by parts technique, the equation (4.5) can be solved as

$$\begin{aligned} &{}_C^C_{a,k}\mathfrak{D}^{\alpha,\rho;\psi} u(\tau) \\ &= \frac{{}_k\Psi_{\psi}^{\frac{\alpha}{k}-1}(\tau, a)}{\rho^{1-\frac{\alpha}{k}}\Gamma_k(2k-\alpha)} \left[(1-\rho)u(a) + k\rho \frac{u'(a)}{\psi'(a)} \right] \\ &\quad + \frac{1}{\rho^{1-\frac{\alpha}{k}}\Gamma_k(2k-\alpha)} \int_a^\tau (\psi(\tau) - \psi(s))^{1-\frac{\alpha}{k}} \frac{d}{ds} \left(e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} \left[(1-\rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] \right) ds. \end{aligned} \quad (4.6)$$

Applying the Newton's generalized binomial theorem, it follows that

$$\begin{aligned} (\psi(\tau) - \psi(s))^{1-\frac{\alpha}{k}} &= (\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}} \left(1 - \frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \right)^{1-\frac{\alpha}{k}} \\ &= (\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}} \sum_{i=0}^N \frac{\Gamma\left(i + \frac{\alpha}{k} - 1\right)}{i!\Gamma\left(\frac{\alpha}{k} - 1\right)} \left(\frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \right)^i \end{aligned}$$

$$+(\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}} \sum_{i=N+1}^{\infty} \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} \left(\frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \right)^i. \quad (4.7)$$

Inserting (4.7) into (4.6), we have

$$\begin{aligned} {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} u(\tau) &= \frac{\rho \Psi_k^{\frac{\alpha}{k}-1}(\tau, a)}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \left[(1 - \rho)u(a) + k\rho \frac{u'(a)}{\psi'(a)} \right] \\ &+ \frac{1}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \int_a^\tau (\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}} \sum_{i=0}^N \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} \left(\frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \right)^i \\ &\times \frac{d}{ds} \left(e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} \left[(1 - \rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] \right) ds \\ &+ \frac{1}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \int_a^\tau (\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}} \sum_{i=N+1}^{\infty} \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} \left(\frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \right)^i \\ &\times \frac{d}{ds} \left(e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} \left[(1 - \rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] \right) ds. \end{aligned} \quad (4.8)$$

where

$$E_{tr}(\tau) = \frac{(\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}}}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \int_a^\tau R_N(s) \frac{d}{ds} \left(e^{\frac{\rho-1}{k\rho}(\psi(\tau)-\psi(s))} \left[(1 - \rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] \right) ds, \quad (4.9)$$

$$R_N(s) = \sum_{i=N+1}^{\infty} \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} \left(\frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \right)^i. \quad (4.10)$$

Then, we have

$$\begin{aligned} {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} u(\tau) &= \frac{(\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}}}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \left[(1 - \rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)} \right] \\ &+ \frac{e^{\frac{\rho-1}{k\rho}\psi(\tau)} (\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}}}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \sum_{i=1}^N \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} (\psi(\tau) - \psi(a))^i \\ &\times \int_a^\tau (\psi(s) - \psi(a))^i \frac{d}{ds} \left(e^{\frac{1-\rho}{k\rho}\psi(s)} \left[(1 - \rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] \right) ds + E_{tr}(\tau). \end{aligned}$$

By using integrating by parts with changing variable $x(s) = (\psi(s) - \psi(a))^i$ and $y'(s) = \frac{d}{ds} (e^{\frac{1-\rho}{k\rho}\psi(s)} [(1 - \rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)}])$, which yields that

$$\begin{aligned} {}^C_{a,k} \mathfrak{D}^{\alpha,\rho;\psi} u(\tau) &= \frac{(\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}}}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \left[(1 - \rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)} \right] \left(1 + \sum_{i=1}^N \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} \right) \\ &- \frac{e^{\frac{\rho-1}{k\rho}\psi(\tau)}}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha)} \sum_{i=1}^N \frac{\Gamma(i + \frac{\alpha}{k} - 1) (\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}-i}}{(i-1)! \Gamma(\frac{\alpha}{k} - 1)} \end{aligned}$$

$$\times \int_a^\tau (\psi(s) - \psi(a))^{i-1} \psi'(s) e^{\frac{1-\rho}{k\rho} \psi(s)} \left[(1-\rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] ds + E_{tr}(\tau),$$

Next, we study the term $E_{tr}(\tau)$, which is the error caused by truncation error. Finally, we show that $E_{tr}(\tau) \rightarrow 0$ as $N \rightarrow \infty$, $\tau \in [a, b]$, and to prove this, we provide an upper bound for the error. Since $0 \leq \frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \leq 1$ for $s, \tau \in \mathcal{J}$, The equation (4.10) can be computed that

$$|R_N(s)| \leq \sum_{i=N+1}^{\infty} \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} \left(\frac{\psi(s) - \psi(a)}{\psi(\tau) - \psi(a)} \right)^i \leq \sum_{i=N+1}^{\infty} \frac{\Gamma(i + \frac{\alpha}{k} - 1)}{i! \Gamma(\frac{\alpha}{k} - 1)} \leq \sum_{i=N+1}^{\infty} \frac{i^{\frac{\alpha}{k}-2}}{\Gamma(\frac{\alpha}{k} - 1)} \leq \int_N^{\infty} \frac{s^{\frac{\alpha}{k}-2}}{\Gamma(\frac{\alpha}{k} - 1)} ds,$$

which implies that

$$|R_N(s)| \leq \frac{1}{N^{1-\frac{\alpha}{k}} \left(1 - \frac{\alpha}{k}\right) \Gamma\left(\frac{\alpha}{k} - 1\right)}. \quad (4.11)$$

Taking $\mathcal{M}(\tau) = \max_{s \in [a, \tau]} \left| \frac{d}{ds} \left(e^{\frac{\rho-1}{k\rho}(\psi(\tau) - \psi(s))} \left[(1-\rho)u(s) + k\rho \frac{u'(s)}{\psi'(s)} \right] \right) \right|$ and (4.11) into (4.9), we obtain the following upper bound

$$|E_{tr}(\tau)| \leq \frac{\mathcal{M}(\tau)(\psi(\tau) - \psi(a))^{1-\frac{\alpha}{k}}(\tau - a)}{\rho^{1-\frac{\alpha}{k}} \Gamma_k(2k - \alpha) N^{1-\frac{\alpha}{k}} \left(1 - \frac{\alpha}{k}\right) \Gamma\left(\frac{\alpha}{k} - 1\right)}. \quad (4.12)$$

The right-sided of (4.12) tends to zero as $N \rightarrow \infty$, $\tau \in [a, b]$. The proof is done.

From the property (iv) in Lemma 2.4, the integral equation (3.1) in Lemma 3.1 can be reformed as

$$u(\tau) = u_a e^{\frac{\rho-1}{k\rho}(\psi(\tau) - \psi(a))} + {}_C \mathfrak{D}_{a,k}^{k-\alpha, \rho; \psi} \left(\frac{1}{\rho k} \int_a^\tau e^{\frac{\rho-1}{k\rho}(\psi(\tau) - \psi(s))} \psi'(s) f(s, u(s)) ds \right). \quad (4.13)$$

By applying Theorem 4.1, the equation (4.13) can be re-written as

$$\begin{aligned} u(s) &= u_a e^{\frac{\rho-1}{k\rho}(\psi(s) - \psi(a))} + A_N^* (\psi(s) - \psi(a))^{\frac{\alpha}{k}} f(s, u(s)) \\ &\quad - e^{\frac{\rho-1}{k\rho} \psi(s)} \sum_{i=1}^N B_{N,i}^* (\psi(s) - \psi(a))^{\frac{\alpha}{k} - i} \mathcal{V}_i^*(s) + E_{tr}^*(s), \end{aligned} \quad (4.14)$$

where

$$\begin{aligned} A_N^* &= \frac{1}{\rho^{\frac{\alpha}{k}} \Gamma_k(k + \alpha)} \sum_{i=0}^N \frac{\Gamma\left(i - \frac{\alpha}{k}\right)}{i! \Gamma\left(-\frac{\alpha}{k}\right)}, \\ B_{N,i}^* &= \frac{\Gamma\left(i - \frac{\alpha}{k}\right)}{\rho^{\frac{\alpha}{k}} (i-1)! \Gamma_k(k + \alpha) \Gamma\left(-\frac{\alpha}{k}\right)}, \quad i = 1, 2, \dots, N, \\ \mathcal{V}_i^*(s) &= \int_a^s (\psi(r) - \psi(a))^{i-1} \psi'(r) e^{\frac{1-\rho}{k\rho} \psi(r)} f(r, u(r)) dr, \\ |E_{tr}^*(s)| &\leq \frac{\mathcal{M}(s)(s-a)(\psi(s) - \psi(a))^{\frac{\alpha}{k}}}{N^{\frac{\alpha}{k}} \left(\frac{\alpha}{k}\right) \Gamma\left(-\frac{\alpha}{k}\right) \rho^{\frac{\alpha}{k}} \Gamma_k(k + \alpha)}. \end{aligned}$$

In order to obtain the formula of an approximated solution, $u_N(\tau)$, we truncate the formula up to order N , getting

$$\begin{aligned} u_N(s) &= u_a e^{\frac{\rho-1}{k\rho}(\psi(s)-\psi(a))} + A_N^*(\psi(s) - \psi(a))^{\frac{\alpha}{k}} f(s, u_N(s)) \\ &\quad - e^{\frac{\rho-1}{k\rho}\psi(s)} \sum_{i=1}^N B_{N,i}^*(\psi(s) - \psi(a))^{\frac{\alpha}{k}-i} \mathcal{V}_{i,N}^*(s), \end{aligned} \quad (4.15)$$

where

$$\mathcal{V}_{i,N}^*(s) = \int_a^s (\psi(r) - \psi(a))^{i-1} \psi'(r) e^{\frac{1-\rho}{k\rho}\psi(r)} f(r, u_N(r)) dr, \quad i = 1, 2, \dots, N.$$

Note that $u_N(a) = u(a)$.

Theorem 4.2. Assume that $f \in C([a, b] \times \mathbb{R}, \mathbb{R})$ which verifies (\mathcal{A}_1) , and assume that u and u_N as in (4.14) and (4.15), respectively, for $N \in \mathbb{N}$. Also, suppose that $b \in \mathbb{R}$ is a real in the open interval

$$\psi(a) < \psi(b) < \psi(a) + \rho k \left(\frac{\Gamma(\frac{\alpha}{k} + 1)}{\mathcal{L}} \right)^{\frac{k}{\alpha}}. \quad (4.16)$$

Then, $u_N(\tau) \rightarrow u(\tau)$ as $N \rightarrow \infty$, for all $\tau \in [a, b]$.

Proof. From the equations (4.14) and (4.15), for any $\tau \in [a, b]$, we have

$$\begin{aligned} |x_N(\tau) - x(\tau)| &\leq |A_N^*| (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}} |f(\tau, x_N(\tau)) - f(\tau, x(\tau))| \\ &\quad + e^{\frac{\rho-1}{k\rho}\psi(\tau)} \sum_{i=1}^N |B_{N,i}^*| (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}-i} |\mathcal{V}_{i,N}^*(\tau) - \mathcal{V}_i^*(\tau)| + |E_{tr}^*(\tau)|. \end{aligned} \quad (4.17)$$

Now, we define $\delta_{u_N} := \max_{\tau \in [a, b]} |u_N(\tau) - u(\tau)|$. Then, we have

$$\begin{aligned} |f(\tau, u_N(\tau)) - f(\tau, u(\tau))| &\leq \mathcal{L} |u_N(\tau) - u(\tau)| \leq \mathcal{L} \delta_{u_N}, \\ |\mathcal{V}_{i,N}^*(\tau) - \mathcal{V}_i^*(\tau)| &\leq \int_a^\tau (\psi(s) - \psi(a))^{i-1} \psi'(s) e^{\frac{1-\rho}{k\rho}\psi(s)} |f(s, u_N(s)) - f(s, u(s))| ds \\ &\leq \mathcal{L} \delta_{u_N} \int_a^\tau (\psi(s) - \psi(a))^{i-1} \psi'(s) e^{\frac{1-\rho}{k\rho}\psi(s)} ds \\ &\leq \frac{\mathcal{L} \delta_{u_N}}{i} (\psi(\tau) - \psi(a))^i. \end{aligned} \quad (4.19)$$

By using the formula (3) in [43], we have

$$|A_N^*| = \frac{1}{\rho^{\frac{\alpha}{k}} \Gamma_k(k + \alpha)} \left| \sum_{i=0}^N \frac{\Gamma(i - \frac{\alpha}{k})}{\Gamma(-\frac{\alpha}{k}) i!} \right| = \frac{\rho^{-\frac{\alpha}{k}} \Gamma(N + 1 - \frac{\alpha}{k})}{\Gamma_k(k + \alpha) \Gamma(-\frac{\alpha}{k}) |\frac{\alpha}{k} \Gamma(N + 1)|} \leq \frac{k^{1-\frac{\alpha}{k}} \Gamma(N + 1 - \frac{\alpha}{k})}{\alpha \rho^{\frac{\alpha}{k}} \pi \Gamma(N + 1)}, \quad (4.20)$$

and using (4.19), it follows that

$$\sum_{i=1}^N |B_{N,i}^*| (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}-i} |\mathcal{V}_{i,N}^*(\tau) - \mathcal{V}_i^*(\tau)|$$

$$\begin{aligned}
&\leq \frac{\mathcal{L}\delta_{x_N}(\psi(\tau) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\Gamma_k(k + \alpha)|\Gamma(-\frac{\alpha}{k})|} \left| \sum_{i=1}^N \frac{\Gamma(i - \frac{\alpha}{k})}{i!} \right| \\
&\leq \frac{\mathcal{L}\delta_{x_N}(\psi(\tau) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\Gamma_k(k + \alpha)|\Gamma(-\frac{\alpha}{k})|} \left(\frac{\Gamma(N + 1 - \frac{\alpha}{k})}{\frac{\alpha}{k}\Gamma(N + 1)} + \left| \Gamma(-\frac{\alpha}{k}) \right| \right) \\
&\leq \frac{\mathcal{L}\delta_{x_N}(\psi(\tau) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}} \left(\frac{k^{1-\frac{\alpha}{k}}\Gamma(N + 1 - \frac{\alpha}{k})}{\alpha\pi\Gamma(N + 1)} + \frac{k^{1-\frac{\alpha}{k}}}{\alpha\Gamma(\frac{\alpha}{k})} \right). \tag{4.21}
\end{aligned}$$

Substitution (4.18), (4.20), and (4.21) into the inequality (4.17), for any $\tau \in [a, b]$, one has

$$\begin{aligned}
|u_N(\tau) - u(\tau)| &\leq \frac{\mathcal{L}\delta_{u_N}k^{1-\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\alpha\pi} (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}} \frac{\Gamma(N + 1 - \frac{\alpha}{k})}{\Gamma(N + 1)} + |E_{tr}^*(\tau)| \\
&\quad + e^{\frac{\rho-1}{k\rho}\psi(\tau)} \frac{\mathcal{L}\delta_{u_N}(\psi(\tau) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}} \left(\frac{k^{1-\frac{\alpha}{k}}\Gamma(N + 1 - \frac{\alpha}{k})}{\alpha\pi\Gamma(N + 1)} + \frac{k^{1-\frac{\alpha}{k}}}{\alpha\Gamma(\frac{\alpha}{k})} \right) \\
&\leq \frac{\mathcal{L}\delta_{u_N}k^{1-\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\alpha} (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}} \left(\frac{2\Gamma(N + 1 - \frac{\alpha}{k})}{\pi\Gamma(N + 1)} + \frac{1}{\Gamma(\frac{\alpha}{k})} \right) + |E_{tr}^*(\tau)|. \tag{4.22}
\end{aligned}$$

Applying the maximum into the inequality (4.22) over $\tau \in [a, b]$, we obtain that

$$\delta_{u_N} \leq \frac{\mathcal{L}\delta_{u_N}k^{1-\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\alpha} (\psi(b) - \psi(a))^{\frac{\alpha}{k}} \left(\frac{2\Gamma(N + 1 - \frac{\alpha}{k})}{\pi\Gamma(N + 1)} + \frac{1}{\Gamma(\frac{\alpha}{k})} \right) + \max_{\tau \in [a, b]} |E_{tr}^*(\tau)|. \tag{4.23}$$

Clearly, $\lim_{N \rightarrow \infty} |E_{tr}^*(\tau)| = 0$. In addition, by applying the Stirling's formula [44], we have

$$\lim_{N \rightarrow \infty} \frac{\Gamma(N + 1 - \frac{\alpha}{k})}{\Gamma(N + 1)} = 0. \tag{4.24}$$

Then, taking $N \rightarrow \infty$ in (4.23), one has

$$\lim_{N \rightarrow \infty} \delta_{u_N} \left(1 - \frac{\mathcal{L}(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}k^{\frac{\alpha}{k}}\Gamma(\frac{\alpha}{k} + 1)} \right) \leq 0.$$

By the condition (4.16), we get that $\delta_{u_N} \rightarrow 0$ as $N \rightarrow \infty$. From (4.23), which implies that

$$\delta_{u_N} \left[1 - \frac{\mathcal{L}k^{1-\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\alpha} (\psi(b) - \psi(a))^{\frac{\alpha}{k}} \left(\frac{2\Gamma(N + 1 - \frac{\alpha}{k})}{\pi\Gamma(N + 1)} + \frac{1}{\Gamma(\frac{\alpha}{k})} \right) \right] \leq \max_{\tau \in \mathcal{J}} |E_{tr}^*(\tau)|.$$

Applying the condition (4.16), we obtain the following result

$$0 < \frac{\mathcal{L}k^{1-\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\alpha} (\psi(b) - \psi(a))^{\frac{\alpha}{k}} < \Gamma\left(\frac{\alpha}{k}\right).$$

Then,

$$-\frac{2\Gamma(\frac{\alpha}{k})\Gamma(N + 1 - \frac{\alpha}{k})}{\pi\Gamma(N + 1)} < 1 - \frac{\mathcal{L}k^{1-\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}}\alpha} (\psi(b) - \psi(a))^{\frac{\alpha}{k}} \left(\frac{2\Gamma(N + 1 - \frac{\alpha}{k})}{\pi\Gamma(N + 1)} + \frac{1}{\Gamma(\frac{\alpha}{k})} \right) < 1.$$

For N sufficiently large number, we obtain the property (4.24). Therefore, there exists a function \mathcal{B} (depending on α, ρ, k, τ but independent of N), such that

$$\delta_{u_N} \leq \mathcal{B}_\psi(\alpha, k, \rho, \tau) N^{-\frac{\alpha}{k}},$$

where

$$\mathcal{B}_\psi(\alpha, k, \rho, \tau) = \frac{\mathcal{M}(\tau)(\psi(b) - \psi(a))^{\frac{\alpha}{k}}(b - a)}{\left(\frac{\alpha}{k}\right)\Gamma\left(-\frac{\alpha}{k}\right)\Gamma_k(k + \alpha)} \left(\frac{\alpha\Gamma\left(\frac{\alpha}{k}\right)}{\rho^{\frac{\alpha}{k}}\alpha\Gamma\left(\frac{\alpha}{k}\right) - \mathcal{L}k^{1-\frac{\alpha}{k}}(\psi(b) - \psi(a))^{\frac{\alpha}{k}}} \right). \quad (4.25)$$

Remark 4.3. Under Theorem 4.1 and Theorem 4.2, we have the following results:

- (i) If $k = 1$ and $\psi(\tau) = \tau$ then, Theorem 4.1 and Theorem 4.2 reduce to the results as in [20].
- (ii) If $k > 1, \rho = 1$, and $\psi(\tau) = \tau^\mu$, where $\mu > 0$ then, Theorem 4.1 and Theorem 4.2 reduce to the results as in [21].
- (iii) If $k > 1, \rho = 1$ and $\psi(\tau) = \tau$ then, Theorem 4.1 and Theorem 4.2 reduce to the results as in [22, 23].
- (iv) If $k > 1, \rho = 1$, and $\psi(\tau) = \tau^\mu$, where $\mu \rightarrow 0^+$ then, Theorem 4.1 and Theorem 4.2 reduce to the results as in [24].

5. Applications

This section provides two different numerical examples of applications to verify the theoretical results of our main results.

5.1. A numerical example

Example 5.1. Consider the following Cauchy-type problem under the (k, ψ) -Caputo-PFDO:

$$\begin{cases} {}^C_{0,k}\mathfrak{D}^{\alpha,\rho;\psi}u(\tau) = \frac{\rho^{\frac{\alpha}{k}}\Gamma_k(\omega + k)}{\Gamma_k(\omega + k - \alpha)^k}\rho\Psi_{\psi}^{\frac{\omega-\alpha}{k}}(\tau, 0) + \frac{1}{2^{\tau+3}k}\rho\Psi_{\psi}^{\frac{\omega}{k}}(\tau, 0) - \frac{1}{2^{\tau+3}}u(\tau), & \tau \in [0, 1], \\ u(0) = 0, & \alpha \in (0, 1], \quad \rho \in (0, 1], \quad k > 0. \end{cases} \quad (5.1)$$

The exact solution of the problem (5.1) is provided by $u(\tau) = \rho\Psi_{\psi}^{\frac{\omega}{k}}(\tau, 0)$, where $\omega \in \mathbb{R}$ and $\omega > k$.

From the problem (5.1), we get

$$f(\tau, u(\tau)) = \frac{\rho^{\frac{\alpha}{k}}\Gamma_k(\omega + k)}{\Gamma_k(\omega + k - \alpha)^k}\rho\Psi_{\psi}^{\frac{\omega-\alpha}{k}}(\tau, 0) + \frac{1}{2^{\tau+3}k}\rho\Psi_{\psi}^{\frac{\omega}{k}}(\tau, 0) - \frac{1}{2^{\tau+3}}u(\tau).$$

For every $u, v \in \mathbb{R}, \tau \in [0, 1]$, we get that

$$|f(\tau, u(\tau)) - f(\tau, v(\tau))| \leq \frac{1}{2^{\tau+3}}|u(\tau) - v(\tau)| \leq \frac{1}{8}|u(\tau) - v(\tau)|. \quad (5.2)$$

It is noticed that the assumption (\mathcal{A}_1) is held with $\mathcal{L} = 1/8$. By setting $\alpha = 22/25, \rho = 1/4, k = 5/4, \omega = 3/2$, and $\psi(\tau) = \tau$, then $(\mathcal{L}(\psi(b) - \psi(a))^{\frac{\alpha}{k}})/(\rho^{\frac{\alpha}{k}}\Gamma_k(\alpha + k)) \approx 0.311730513 < 1$. Since all

assumptions in Theorem 3.2 are satisfied, the proposed problem (5.1) has a unique solution on $[0, 1]$. From Theorem 3.10, we can compute that

$$\mathfrak{C}_f := \frac{(\psi(b) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(\alpha + k)} \approx 2.493844103 > 0, \quad \text{and} \quad \kappa_f := \mathcal{L}\rho^{-\frac{\alpha}{k}} \approx 0.331711285 \geq 0.$$

Then, the proposed problem (5.1) is UH-ML stable on $[0, 1]$. If we set $\mathcal{G}_f(\epsilon) = \mathfrak{C}_f \epsilon$ under the condition $\mathcal{G}_f(0) = 0$, then we obtain the proposed problem (5.1) is generalized UH-ML stable on $[0, 1]$. In addition, by setting a non-decreasing function $\Phi(\tau) = \rho \Psi_{\psi}^{\frac{3}{k}}(\tau, a)$, we have

$${}_{a,k} \mathcal{I}^{\alpha, \rho; \psi} \Phi(\tau) = \frac{\Gamma_k(3) \rho \Psi_{\psi}^{\frac{3+\alpha}{k}-1}(\tau, a)}{\rho^{\frac{\alpha}{k}} \Gamma_k(3 + \alpha)} = \frac{\Gamma_k(3) (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}}}{\rho^{\frac{\alpha}{k}} \Gamma_k(3 + \alpha)} \Phi(\tau).$$

This yields that $\mathfrak{C}_{f_{\Phi}} := \chi_{\Phi} = (\Gamma_k(3)) / (\rho^{\frac{\alpha}{k}} \Gamma_k(3 + \alpha)) (\psi(\tau) - \psi(a))^{\frac{\alpha}{k}} \approx 1.276974061 > 0$ and $\kappa_{f_{\Phi}} := \mathcal{L}\rho^{-\frac{\alpha}{k}} \approx 0.331711285 \geq 0$. Hence, the proposed problem (5.1) is UHR-ML stable on $[0, 1]$. If we take $\epsilon = 1$, then the proposed problem (5.1) is generalized UHR-ML stable on $[0, 1]$.

To achieve the numerical approximation of the proposed problem (5.1), we apply (4.4) in Theorem 4.1. Then, the operator ${}_{0,k}^C \mathfrak{D}^{\alpha, \rho; \psi} u(\tau)$ can be replaced as follows:

$$\begin{aligned} {}_{0,k}^C \mathfrak{D}^{\alpha, \rho; \psi} u(\tau) &\approx A_N (\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}} \left[(1 - \rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)} \right] \\ &\quad - e^{\frac{\rho-1}{k\rho}\psi(\tau)} \sum_{i=1}^N B_{N,i} (\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}-i} \mathcal{V}_i(\tau), \end{aligned} \quad (5.3)$$

where A_N , $B_{N,i}$, and $\mathcal{V}_i(\tau)$ are given by (4.1), (4.2), and (4.3), respectively, in Theorem 4.1. Then

$$\begin{cases} \mathcal{V}'_i(\tau) = (\psi(\tau) - \psi(0))^{i-1} \psi'(\tau) e^{\frac{1-\rho}{k\rho}\psi(\tau)} \left[(1 - \rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)} \right], & i = 1, 2, \dots, N, \\ \mathcal{V}_i(0) = 0. \end{cases} \quad (5.4)$$

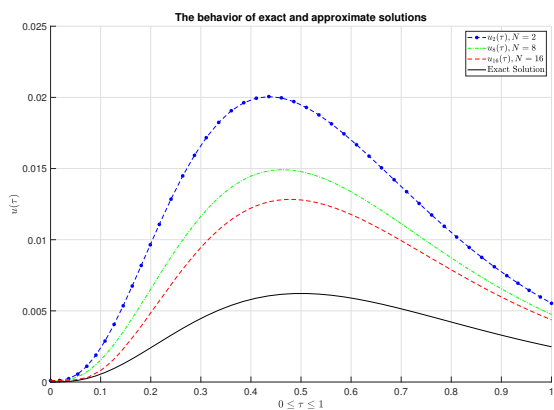
By applying (5.3) and (5.4), then

$$\begin{cases} A_N (\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}} \left[(1 - \rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)} \right] \\ - e^{\frac{\rho-1}{k\rho}\psi(\tau)} \sum_{i=1}^N B_{N,i} (\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}-i} \mathcal{V}_i(\tau) \\ = \frac{\rho^{\frac{\alpha}{k}} \Gamma_k(\omega + k)}{\Gamma_k(\omega + k - \alpha)^k} \rho \Psi_{\psi}^{\frac{\omega-\alpha}{k}}(\tau, 0) + \frac{1}{2^{\tau+3k}} \rho \Psi_{\psi}^{\frac{\alpha}{k}}(\tau, 0) - \frac{1}{2^{\tau+3}} u(\tau), & \tau \in [0, 1], \\ \mathcal{V}'_i(\tau) = (\psi(\tau) - \psi(0))^{i-1} \psi'(\tau) e^{\frac{1-\rho}{k\rho}\psi(\tau)} \left[(1 - \rho)u(\tau) + k\rho \frac{u'(\tau)}{\psi'(\tau)} \right], & i = 1, 2, \dots, N, \\ u(0) = 0, \quad \mathcal{V}_i(0) = 0. \end{cases} \quad (5.5)$$

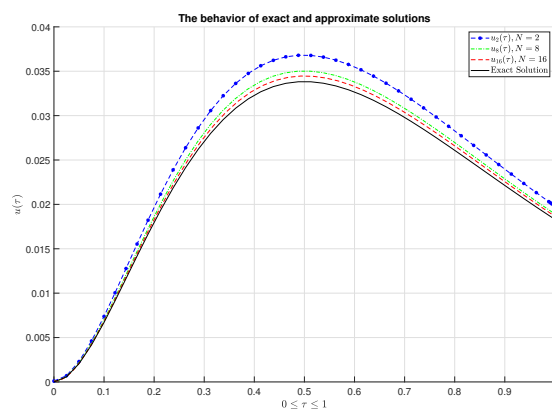
The proposed system (5.5) can solve the numerical solution by applying *ode45* in MATLAB software, which applies the explicit Runge-Kutta technique to achieve the approximated solution. The

exact solution (continuous line) and numerical approximations (dot lines) of the proposed problem are compared with different parameters $\rho = 1/4, 1/3, 1/2$, and $k = 1/2, 3/4$.

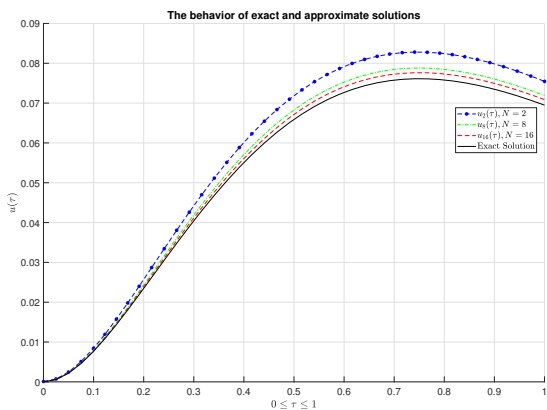
Figure 1 displays the behavior of the exact and approximate solutions under the fixed value $\alpha = 9/10$, and vary values $\rho = 1/4, 1/3, 1/2$, $k = 1/2, 3/4$, $N = 2, 8, 16$ for the proposed problem (5.1). While Figure 2 displays the behavior of the absolute error between the exact and approximate solutions under the fixed values $\alpha = 9/10$, and vary values $\rho = 1/4, 1/3, 1/2$, $k = 1/2, 3/4$, $N = 2, 8, 16$ for the proposed problem (5.1). Figures 1a–1b illustrate that as the value k increases, the approximate solution quickly converges to the exact solution. Moreover, the absolute error values rapidly decrease to zero, as can be seen from Figures 2a–2b. Conversely, Figures 1c–1d indicate that as the value ρ increases, the approximate solution converges to the exact solution more quickly and to greater values. While the absolute error values increase, as can be seen from the Figures 2c–2d.



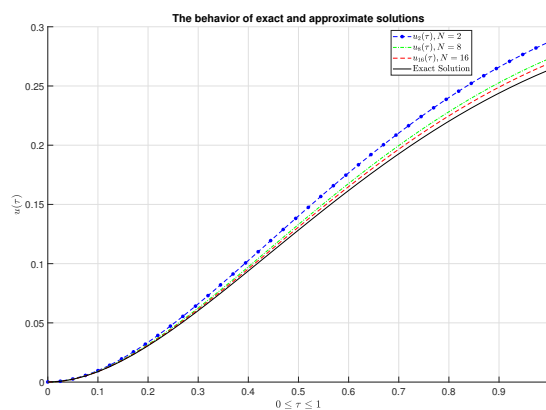
(a) $\rho = 1/4$ and $k = 1/2$.



(b) $\rho = 1/4$ and $k = 3/4$.



(c) $\rho = 1/3$ and $k = 3/4$.



(d) $\rho = 1/2$ and $k = 3/4$.

Figure 1. The behavior of exact and approximate solutions under $\alpha = 9/10$, $\rho = 1/4, 1/3, 1/2$, $k = 1/2, 3/4$, and $N = 2, 8, 16$ for the proposed problem (5.1).

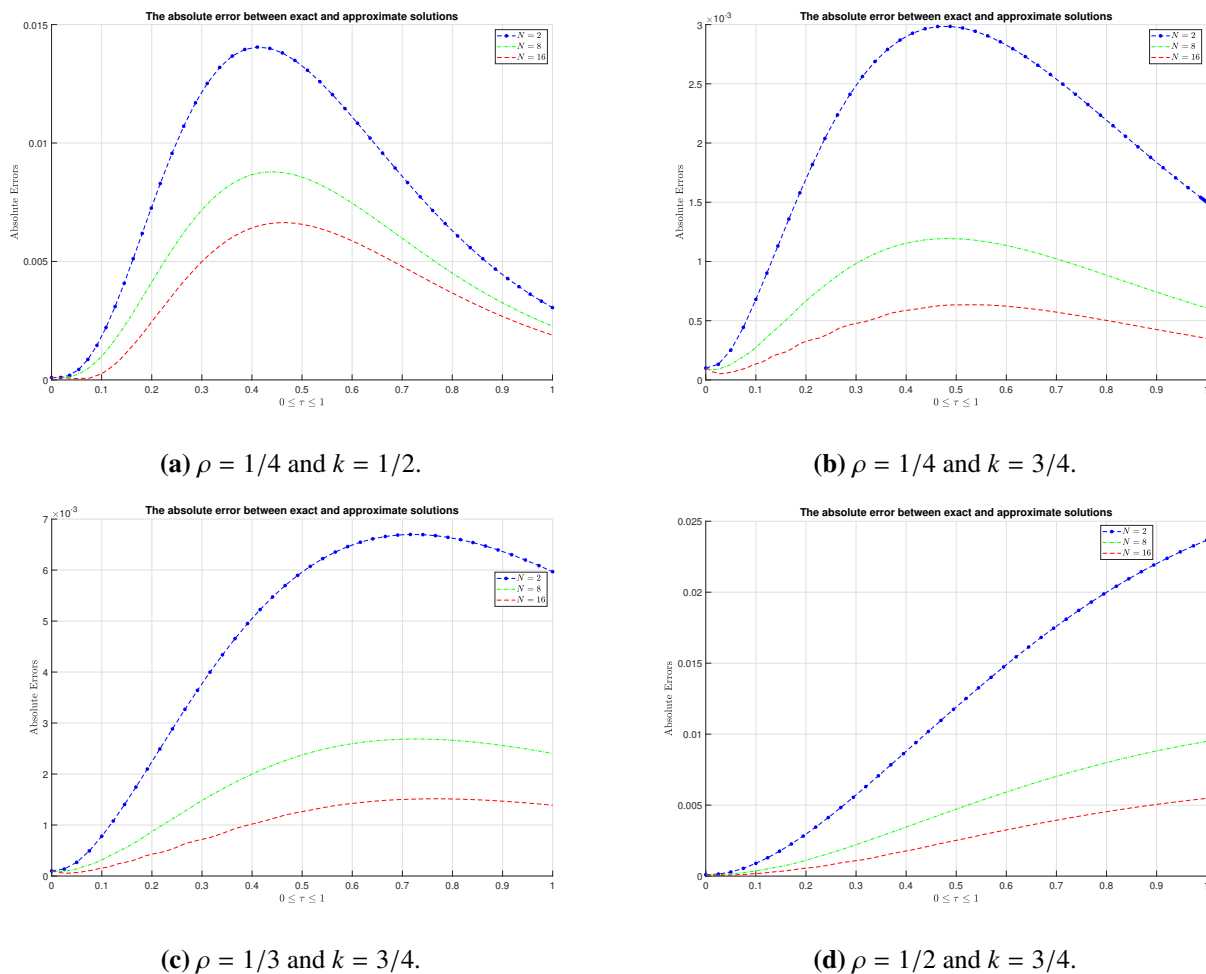


Figure 2. The absolute error between exact and approximate solutions under $\alpha = 9/10$, $\rho = 1/4, 1/3, 1/2, k = 1/2, 3/4$, and $N = 2, 8, 16$ for the proposed problem (5.1).

5.2. An Application to Blood Alcohol Levels (BALs) problem

Now, an application has been developed to support our theoretical results. We study a basic model for determining BALs that is characterized by two differential equations with real-data individuals. The BALs problem is explained by the following Cauchy-type system of the form:

$$\begin{cases} \frac{dS}{d\tau} = -c_1 S(\tau), & S(0) = S_0, \\ \frac{dB}{d\tau} = c_1 S(\tau) - c_2 B(\tau), & B(0) = B_0 = 0, \end{cases} \tag{5.6}$$

where $S(\tau)$ and $B(\tau)$ denote the concentrations (mg/L) of alcohol in the stomach and blood of a human body at time τ (min), respectively. The rate law constants i (min^{-1}) are given by $c_i, i = 1, 2$, and S_0, B_0 represent the subject’s initial alcohol intake in the stomach and blood, respectively. The exact solution for the system (5.6) is fairly simple and may be determined using the Laplace transform approach

$$S(\tau) = S_0 \exp(-c_1 \tau), \tag{5.7}$$

$$\mathcal{B}(\tau) = \frac{c_1 \mathcal{S}_0}{c_1 - c_2} (\exp(-c_2 \tau) - \exp(-c_1 \tau)), \quad c_1 \neq c_2. \quad (5.8)$$

Fractional-order models have been recognized as being more useful in estimating experimental data than integer-order models. Another reason is that various parameters can be adjusted accordingly. Then, by taking the (k, ψ) -Caputo-PFDO into the left-sided of the problem (5.6), the problem (5.6) can be rewritten as:

$$\begin{cases} {}^C_{0,k} \mathfrak{D}^{\alpha,\rho;\psi} \mathcal{S}(\tau) = -c_1 \mathcal{S}(\tau), & \alpha \in (0, 1], \quad \rho \in (0, 1], \quad k > 0, \\ {}^C_{0,k} \mathfrak{D}^{\alpha,\rho;\psi} \mathcal{B}(\tau) = c_1 \mathcal{S}(\tau) - c_2 \mathcal{B}(\tau), & \alpha \in (0, 1], \quad \rho \in (0, 1], \quad k > 0, \\ \mathcal{S}(0) = \mathcal{S}_0, \quad \mathcal{B}(0) = \mathcal{B}_0 = 0. \end{cases} \quad (5.9)$$

Similarly procedure in Example 5.1, we obtain that

$$\left\{ \begin{array}{l} A_N(\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}} \left[(1-\rho)\mathcal{S}(\tau) + k\rho \frac{\mathcal{S}'(\tau)}{\psi'(\tau)} \right] \\ - e^{\frac{\rho-1}{k\rho}\psi(\tau)} \sum_{i=1}^N B_{N,i}(\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}-i} \mathcal{V}_i(\tau) = -c_1 \mathcal{S}(\tau), \\ A_N(\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}} \left[(1-\rho)\mathcal{B}(\tau) + k\rho \frac{\mathcal{B}'(\tau)}{\psi'(\tau)} \right] \\ - e^{\frac{\rho-1}{k\rho}\psi(\tau)} \sum_{i=1}^N B_{N,i}(\psi(\tau) - \psi(0))^{1-\frac{\alpha}{k}-i} \mathcal{U}_i(\tau) = c_1 \mathcal{S}(\tau) - c_2 \mathcal{B}(\tau), \\ \mathcal{V}'_i(\tau) = (\psi(\tau) - \psi(0))^{i-1} \psi'(\tau) e^{\frac{1-\rho}{k\rho}\psi(\tau)} \left[(1-\rho)\mathcal{S}(\tau) + k\rho \frac{\mathcal{S}'(\tau)}{\psi'(\tau)} \right], \quad i = 1, 2, \dots, N, \\ \mathcal{U}'_i(\tau) = (\psi(\tau) - \psi(0))^{i-1} \psi'(\tau) e^{\frac{1-\rho}{k\rho}\psi(\tau)} \left[(1-\rho)\mathcal{B}(\tau) + k\rho \frac{\mathcal{B}'(\tau)}{\psi'(\tau)} \right], \quad i = 1, 2, \dots, N, \\ \mathcal{S}(0) = \mathcal{S}_0, \quad \mathcal{B}(0) = \mathcal{B}_0 = 0, \quad \mathcal{V}_i(0) = 0, \quad \mathcal{U}_i(0) = 0. \end{array} \right. \quad (5.10)$$

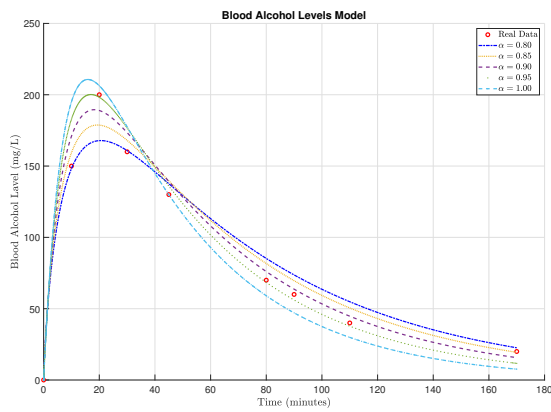
In all graphical simulations, we use the parameters $k_1 = 0.12$, $k_2 = 0.01$, and $\psi(\tau) = \tau$ based on the experimental results for the BALs of a real individual in Table 1. We separated them into three cases, which are shown in Figure 3.

Table 1. Experimental results for the BALs of a real individual.

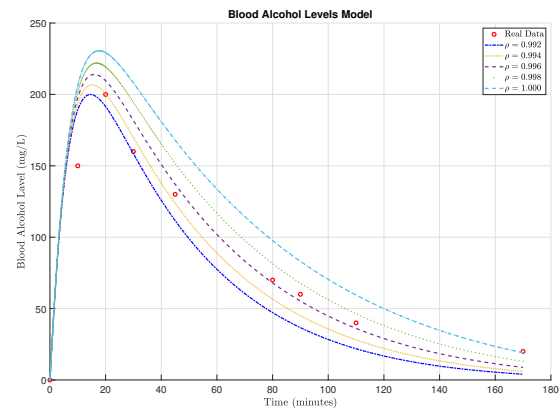
Time (min)	0	10	20	30	45	80	90	110	170
BAL (mg/L)	0	150	200	160	130	170	60	40	20

Case (1): We set the initial conditions $\mathcal{S}_0 = 330$, $\mathcal{B}_0 = 0$, $\rho = 0.988$, and $k = 0.995$ with the varied $\alpha \in \{0.80, 0.85, 0.90, 0.95, 1.00\}$. The graphical simulation of $\mathcal{B}(\tau)$ with the common parameters is shown as in Figure 3a. Case (2): We set the initial conditions $\mathcal{S}_0 = 280$, $\mathcal{B}_0 = 0$, $\alpha = 0.95$, and $k = 0.9$ with the varied $\rho \in \{0.992, 0.994, 0.996, 0.998, 1.000\}$. The graphical simulation of $\mathcal{B}(\tau)$ with the common parameters is shown as in Figure 3b. Case (3): We set the initial conditions $\mathcal{S}_0 = 255$, $\mathcal{B}_0 = 0$, $\alpha = 0.95$, and $\rho = 0.998$ with the varied $k \in \{0.88, 0.90, 0.92, 0.94, 0.96\}$. The graphical

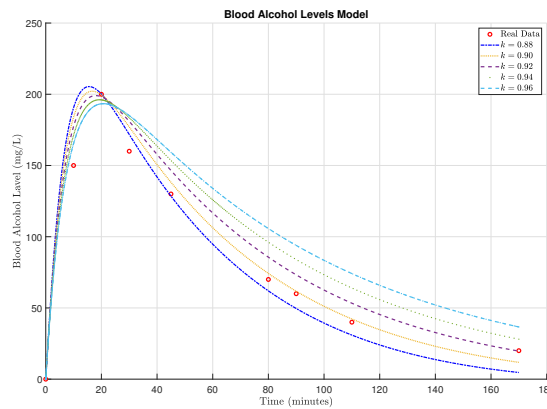
simulation of $B(\tau)$ with the common parameters is shown as in Figure 3c. Case (4): We set the initial conditions $S_0 = 255$, $B_0 = 0$, $\alpha = 0.90$, $\rho = 0.25$, $k = 0.5$, and $N = 2, 5, 20$. The graphical simulation of $B(\tau)$ with the common parameters is shown as in Figure 4.



(a) $\alpha \in \{0.80, 0.85, 0.90, 0.95, 1.00\}$.



(b) $\rho \in \{0.992, 0.994, 0.996, 0.998, 1.000\}$.



(c) $k \in \{0.88, 0.90, 0.92, 0.94, 0.96\}$.

Figure 3. BAL compared with real data in Table 1 under $\alpha \in \{0.80, 0.85, 0.90, 0.95, 1.00\}$, $\rho \in \{0.992, 0.994, 0.996, 0.998, 1.000\}$, and $k \in \{0.88, 0.90, 0.92, 0.94, 0.96\}$ for the BAL model (5.9).

We calculate the relative absolute error, $E_N = |(x_e - x_a)/x_e|$ between the exact solution $x_e(\tau)$ and the approximation solution $x_a(\tau)$ for the concentration of alcohol in the blood of a human body at time τ , $B(\tau)$, as displayed in Table 2. The presented technique is applicable. Also, the results agree with the exact solutions, and the error decreases as N increases.

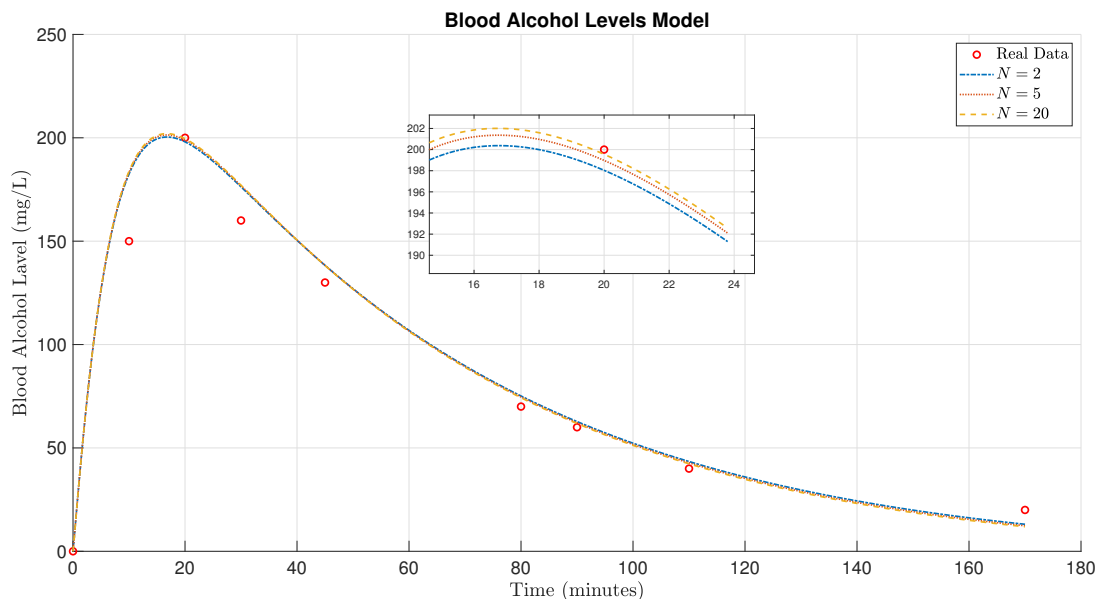


Figure 4. $N \in \{2, 5, 20\}$.

Table 2. The relative absolute error E_N at $\alpha = 0.90$, $\rho = 0.25$, $k = 0.5$ with different value of $N = 2, 5$, and 20 .

Time (min)	BAL (mg/L)	S_2	S_5	S_{20}	E_2	E_5	E_{20}
0	0	0.0000	0.0000	0.0000	-	-	-
10	150	182.4541	183.2365	183.7853	0.2164	0.2216	0.2252
20	200	198.0376	198.9739	199.5539	0.0098	0.0051	0.0022
30	160	176.2629	176.7983	177.0553	0.1016	0.1050	0.1066
45	130	138.3014	138.3158	138.2298	0.0639	0.0640	0.0633
80	70	75.1240	74.6521	74.2932	0.0732	0.0665	0.0613
90	60	62.7600	62.2160	61.8149	0.0460	0.0369	0.0302
110	40	43.4365	42.7890	42.3281	0.0859	0.0697	0.0582
170	20	13.0767	12.3641	11.8921	0.3462	0.3818	0.4054

6. Conclusions

In this work, we analyzed an extended Gronwall inequality in the context of the (k, ψ) -PFOs and proved its properties. A Cauchy-type problem under the (k, ψ) -Caputo-PFDO is the highlight of our presentation. First, the existence and uniqueness results of the proposed problem (1.1) were established by applying Banach's contraction mapping principle and Leray-Schauder's nonlinear alternative. Then, UH-ML stability was studied to guarantee the existing results. Moreover, a numerical technique is demonstrated based on a decomposition formula for the (k, ψ) -Caputo-PFDO. Finally, numerical examples are shown to verify the theoretical results. As seen in Example 5.1, we set the specific type of Cauchy problem and showed numeric calculations to confirm the accuracy of Theorem 3.2, which expresses the unique solution. We also showed the calculation satisfied the

conditions for various Ulam stability to guarantee the results. For the fixed value of order α and with the different parameters ρ and k , we found the result's behavior through the graphic numerical simulation that the approximate solution of the system converges to the exact solution as ρ and k increased. In contrast, the absolute error values decreased to zero for increasing value k but increased for decreasing value ρ . Furthermore, to strengthen our study in a broader domain, we applied the established numerical technique to the BALs problem as seen in Example 5.2. We expressed a numerical simulation of this problem under different parameters. The graphical results illustrated that various values of α, ρ , and k gave the corresponding behavior of the system in the same trend as the data. Consequently, the fractional-order form may approximate the alcohol concentration in a human's blood under a range of circumstances, as proven with data.

In future works, we can discuss applying (k, ψ) -Caputo-PFDO to real-world problems. This larger paradigm greatly contributes to the development of fractional calculus. It is paving the way for exciting future studies in this dynamic and developing discipline.

Author contributions

Weerawat Sudsutad: Problem statement, conceptualization, methodology, investigation, software, writing the original draft, writing, reviewing, and editing; *Jutarat Kongson*: Supervision, problem statement, conceptualization, methodology, investigation, writing the original draft, reviewing, and editing; *Chatthai Thaiprayoon*: Methodology, software, investigation, writing the original draft, reviewing, and editing; *Aphirak Aphithana, Weerapan Sae-dan*: Investigation, software, writing, reviewing, and editing. All the authors read and approved the final manuscript.

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Conflict of interest

The authors declare no conflicts of interest.

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