



Research article

A novel method for Mannheim curves in the Galilean 3–space G_3

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Abstract: This research presents a novel method for Mannheim curves in three-dimensional Galilean space G_3 . Using this method, the necessary and sufficient conditions, along with the established results, must be satisfied for a curve in G_3 to qualify as a Mannheim curve. Furthermore, relevant examples and graphs are provided to demonstrate how Mannheim curves and their partners can correspond to Salkowski and anti-Salkowski curves. Finally, in G_3 , the Mannheim partner curves are described.

Keywords: Mannheim curves; slant helices; Salkowski curves; Curves of Anti-Salkowski; Galilean Space

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1. Introduction

In the Galilean 3-space G_3 , Ogreenmis et al. [11] proposed that if both curvatures, κ and τ are positive constants along the curve β , then the curve β is considered a circular helix with respect to the Frenet frame $\mathbf{t}(s)$, $\mathbf{n}(s)$, and $\mathbf{b}(s)$.

Furthermore, concerning the Frenet frame, a curve β is characterized as a general helix if the ratio $\frac{\kappa}{\tau}$ remains constant, and the converse is also true.

Additionally, the concept of a slant helix was introduced by Karacan and Tuncer [10]. If a constant vector field W exists in G_3 and the function $\langle W, n(s) \rangle_{G_3}$ is constant, then a curve β is called a slant helix. Moreover, if one of the two functions,

$$\pm \frac{\kappa^2}{\tau^3} \frac{d}{ds} \left(\frac{\tau}{\kappa} \right)$$

is constant everywhere and τ does not vanish, then the curve is identified as a slant helix in G_3 , and the converse is also true.

Bertrand explored curves in Euclidean 3-space whose principal normal is the principal normal of another curve in the classical differential geometry of curves. He demonstrated that a linear relationship with constant coefficients between the first and second curvatures of the original curve is a necessary and sufficient condition for the existence of such a second curve [3], i.e., κ and τ satisfy the equation $c_1\kappa + c_2\tau = 1$, where c_1 and c_2 belong to \mathbb{R} . A related curve is called a Mannheim curve, in which the binormal vector field of one curve is the principal normal vector field of another. Liu and Wang [17] examined Mannheim partner curves in both Minkowski and Euclidean 3-space. Since these works, many studies on Mannheim curves in Euclidean 3-space, Minkowski 3-space, dual 3-space, and Galilean spaces have been published [2, 4, 7–9, 12–15].

Let $\gamma(s)$ be an admissible Mannheim curve, and take $\widetilde{\gamma}(\overline{s})$ as the admissible Mannheim companion curve of $\gamma(s)$ through the Galilean 3-space G_3 . In scholarly works, $\widetilde{\gamma}(\overline{s})$ is expressed as

$$\widetilde{\gamma}(\overline{s}) = \widetilde{\gamma}(g(s)) = \gamma(s) + \lambda(s) \mathbf{n}(s), \quad (1.1)$$

in which case the function $g(s)$ is differentiable, and $\mathbf{n}(s)$ is the major normal line of $\gamma(s)$ [16]. In relation (1.1), the vector $\overrightarrow{\widetilde{\gamma} \gamma}$ does not have to be linearly dependent with the $\mathbf{n}(s)$ vector.

Hence, the Mannheim mate curve, $\widetilde{\gamma}$, is assumed to be produced by

$$\widetilde{\gamma}(\overline{s}) = \widetilde{\gamma}(g(s)) = \gamma(s) + \lambda_1(s) \mathbf{t}(s) + \lambda_2(s) \mathbf{n}(s) + \lambda_3(s) \mathbf{b}(s), \quad (1.2)$$

where the Frenet frame of $\gamma(s)$ is represented by $\mathbf{t}(s)$, $\mathbf{n}(s)$, and $\mathbf{b}(s)$. During this situation, if we accept $\lambda_1(s) = \lambda_3(s) = 0$, we receive the instance that is referenced inside the written works. Therefore, in our article we took $\lambda_1 \neq 0$, and $\lambda_3 \neq 0$ to provide the Mannheim generalized curve in G_3 .

Throughout this work, we establish the circumstances that must be satisfied for a G_3 curve to qualify as a Mannheim curve and give instances that correspond to this new way of defining Mannheim curves. Finally, we describe Mannheim mate curves within G_3 .

2. Galilean 3-space curves

We will provide a few definitions in this section that will be used in our paper. For more fundamental concepts, see [1, 5, 6].

Galilean space in three dimensions, G_3 , is defined as the Cayley-Klein space, where the characteristic protective metric has the signature $(0, 0, +, +)$. In Galilean space, the absolute is represented by a triple (V, E, J) . V indicates the ideal plane, E defines a line in V , and J represents an elliptic point of involution $(0 : 0 : r_2 : r_3) \rightarrow (0 : 0 : r_3 : -r_2)$.

A plane is called Euclidean if it contains E , otherwise it is called isotropic. A vector $r = (r_1, r_2, r_3)$ is described as not being isotropic if $r_1 \neq 0$. The form of all unit non-isotropic vectors is $r = (1, r_2, r_3)$. Regarding the isotropic vectors, r_1 is equal to zero.

Allow $\overrightarrow{\eta} = (\eta_1, \eta_2, \eta_3)$ and $\overrightarrow{\xi} = (\xi_1, \xi_2, \xi_3)$ to be two vectors in Galilean 3-space G_3 . In G_3 , a dot product is described as

$$\langle \overrightarrow{\eta}, \overrightarrow{\xi} \rangle_{G_3} = \begin{cases} \eta_1 \xi_1 & \text{if } \eta_1 \neq 0 \text{ or } \xi_1 \neq 0; \\ \eta_2 \xi_2 + \eta_3 \xi_3 & \text{if } \eta_1 = 0 \text{ and } \xi_1 = 0. \end{cases}$$

The vector $\vec{\zeta} = (\zeta_1, \zeta_2, \zeta_3)$, in G_3 , has the following norm: $\|\vec{\zeta}\| = \sqrt{\langle \vec{\zeta}, \vec{\zeta} \rangle}$.
The Galilean vector product can be defined as

$$\vec{\eta} \times \vec{\xi} = \begin{vmatrix} 0 & e_2 & e_3 \\ \eta_1 & \eta_2 & \eta_3 \\ \xi_1 & \xi_2 & \xi_3 \end{vmatrix}.$$

In a coordinate form, assume that $\beta : I \rightarrow G_3$ is a curve in Galilean space G_3 given by

$$\beta(t) = (u(t), v(t), w(t)),$$

such that $u(t), v(t)$, and $w(t) \in C^3$, $t \in I$. In the case of $u'(t) \neq 0$, then β is known as an admissible curve.

Assuming that the curve β is admissible in G_3 , and that the parameter s is the arc length, which is derived by

$$\beta(s) = (s, v(s), w(s)),$$

$\kappa(s)$, and $\tau(s)$, the first and second curvature functions, respectively, are given by

$$\kappa(s) = \|\beta''(s)\| = \sqrt{v''^2(s) + w''^2(s)},$$

and

$$\tau(s) = \frac{\det(\beta'(s), \beta''(s), \beta'''(s))}{\kappa^2(s)},$$

respectively. The Frenet frame associated with $\beta(s)$ is written as

$$\begin{aligned} \mathbf{t}(s) &= \beta'(s) = (1, v'(s), w'(s)), \\ \mathbf{n}(s) &= \frac{\beta''(s)}{\kappa(s)} = \frac{1}{\kappa(s)}(0, v''(s), w''(s)), \\ \mathbf{b}(s) &= \frac{1}{\kappa(s)}(0, -w''(s), v''(s)), \end{aligned}$$

so that the velocity vector, principal normal vector, and binormal vector of the curve β are represented, respectively, by the symbols $\mathbf{t}(s)$, $\mathbf{n}(s)$, and $\mathbf{b}(s)$.

Regarding $\beta(s)$, the Frenet equations are written as

$$\frac{d}{ds} \begin{bmatrix} \mathbf{t}(s) \\ \mathbf{n}(s) \\ \mathbf{b}(s) \end{bmatrix} = \begin{bmatrix} 0 & \kappa & 0 \\ 0 & 0 & \tau \\ 0 & -\tau & 0 \end{bmatrix} \begin{bmatrix} \mathbf{t}(s) \\ \mathbf{n}(s) \\ \mathbf{b}(s) \end{bmatrix}.$$

3. A new method for Mannheim curves within Galilean space G_3

Definition 3.1. An admissible curve $\gamma : I \subset \mathbb{R} \rightarrow G_3$ with non-vanishing curvatures is claimed to be a Mannheim curve if there exists a curve $\tilde{\gamma} : \tilde{I} \subset \mathbb{R} \rightarrow G_3$ in such a way that the principle normal vector field of $\gamma(s)$ coincides with the bi-normal vector field of $\tilde{\gamma}(\tilde{s})$ at $s \in I$ and $\tilde{s} \in \tilde{I}$. $\gamma(s)$ will be referred to as a Mannheim curve, $\tilde{\gamma}(\tilde{s})$ as a Mannheim partner, and the pair $(\gamma, \tilde{\gamma})$ as a pair of Mannheim's.

Let $\gamma : I \rightarrow G_3$, $I \subset \mathbb{R}$ be an admissible Mannheim curve in G_3 with the Frenet frame $\{\mathbf{t}(s), \mathbf{n}(s), \mathbf{b}(s)\}$ and non-vanishing curvatures $\kappa(s)$ and $\tau(s)$. Furthermore, assuming the Frenet frame $\{\tilde{\mathbf{t}}, \tilde{\mathbf{n}}, \tilde{\mathbf{b}}\}$, let $\tilde{\gamma} : \tilde{I} \rightarrow G_3$, $\tilde{I} \subset \mathbb{R}$ be an admissible Mannheim partner curve for γ using non-vanishing curvatures $\tilde{\kappa}(s)$ and $\tilde{\tau}(s)$.

Then, $\tilde{\gamma}$ can be formulated as

$$\tilde{\gamma}(\tilde{s}) = \tilde{\gamma}(g(s)) = \gamma(s) + \lambda_1(s) \mathbf{t}(s) + \lambda_2(s) \mathbf{n}(s) + \lambda_3(s) \mathbf{b}(s),$$

where the functions $\lambda_1(s)$, $\lambda_2(s)$, and $\lambda_3(s)$ exhibit differentiability on I .

Theorem 3.1. Assume that $\gamma : I \rightarrow G_3$, $I \subset \mathbb{R}$ is an admissible Mannheim curve within G_3 whose curvatures $\kappa(s)$, $\tau(s)$ do not vanish. If there are differentiable functions λ_1 , λ_2 , and λ_3 that satisfy the given conditions:

$$\lambda'_2 + \kappa\lambda_1 = \tau\lambda_3, \quad (1 + \lambda'_1)\kappa = (\lambda'_3 + \tau\lambda_2)\tau, \quad \lambda'_3 + \tau\lambda_2 \neq 0. \quad (3.1)$$

Hence, curve γ is a Mannheim curve with $\tilde{\gamma}$ as its Mannheim partner; in addition, the converse is also true.

Proof. Consider that γ is an admissible Mannheim curve with non-zero curvatures κ , τ , using arc length s as the parameter, and $\tilde{\gamma}$ represents the Mannheim partner of γ , where the parameter is the arc length \tilde{s} . Then, the curve $\tilde{\gamma}$ can be represented as

$$\tilde{\gamma}(\tilde{s}) = \tilde{\gamma}(g(s)) = \gamma(s) + \lambda_1(s) \mathbf{t}(s) + \lambda_2(s) \mathbf{n}(s) + \lambda_3(s) \mathbf{b}(s), \quad (3.2)$$

for $s \in I$, such that $\lambda_1(s)$, $\lambda_2(s)$, and $\lambda_3(s)$ are differentiable functions on I . By using s to differentiate Eq (3.2), we arrive at

$$g'(s) \tilde{t}(\tilde{s}) = (1 + \lambda'_1(s)) \mathbf{t}(s) + (\kappa(s)\lambda_1(s) - \tau(s)\lambda_3(s) + \lambda'_2(s)) \mathbf{n}(s) + (\tau(s)\lambda_2(s) + \lambda'_3(s)) \mathbf{b}(s). \quad (3.3)$$

Taking the scalar product of Eq (3.3) with $\mathbf{n}(s)$, we get

$$\kappa(s)\lambda_1(s) - \tau(s)\lambda_3(s) + \lambda'_2(s) = 0. \quad (3.4)$$

Substituting from (3.4) into (3.3), implies

$$g'(s) \tilde{t}(\tilde{s}) = (1 + \lambda'_1(s)) \mathbf{t}(s) + (\tau(s)\lambda_2(s) + \lambda'_3(s)) \mathbf{b}(s). \quad (3.5)$$

Taking the scalar product of Eq (3.5) with itself yields

$$(g'(s))^2 = (1 + \lambda'_1(s))^2 + (\tau(s)\lambda_2(s) + \lambda'_3(s))^2. \quad (3.6)$$

If we take

$$\alpha(s) = \frac{1 + \lambda'_1(s)}{g'(s)}, \quad \text{and} \quad \theta(s) = \frac{\tau(s)\lambda_2(s) + \lambda'_3(s)}{g'(s)},$$

we have

$$\tilde{t}(\tilde{s}) = \alpha(s) \mathbf{t}(s) + \theta(s) \mathbf{b}(s). \quad (3.7)$$

When we differentiate Eq (3.7) in relation to s , we obtain

$$\tilde{\kappa}(\tilde{s}) \tilde{n}(\tilde{s}) g'(s) = \alpha'(s) \mathbf{t}(s) + (\kappa(s) \alpha(s) - \tau(s) \theta(s)) \mathbf{n}(s) + \theta'(s) \mathbf{b}(s). \quad (3.8)$$

Using the scalar product of Eq (3.8) and $\mathbf{n}(s)$, we get $\kappa(s) \alpha(s) - \tau(s) \theta(s) = 0$, and therefore

$$(1 + \lambda_1'(s))\kappa(s) = (\tau(s)\lambda_2(s) + \lambda_3'(s))\tau(s), \quad (3.9)$$

such that $\tau(s)\lambda_2(s) + \lambda_3'(s) \neq 0$.

In contrast, consider γ to be a curve with the non-vanishing curvatures κ and τ and conditions (3.1) hold for differentiable functions λ_1 , λ_2 , and λ_3 .

Then, we can create an additional curve, $\tilde{\gamma}$, as follows

$$\tilde{\gamma}(s) = \tilde{\gamma}(g(s)) = \gamma(s) + \lambda_1(s) \mathbf{t}(s) + \lambda_2(s) \mathbf{n}(s) + \lambda_3(s) \mathbf{b}(s). \quad (3.10)$$

Differentiating Eq (3.10) with respect to s yields

$$g'(s) \tilde{t}(s) = (1 + \lambda_1'(s)) \mathbf{t}(s) + (\tau(s)\lambda_2(s) + \lambda_3'(s)) \mathbf{b}(s), \quad (3.11)$$

which gives that

$$g'(s) = \left(\langle \tilde{t}(s), \tilde{t}(s) \rangle_{G_3} \right)^{\frac{1}{2}} = \frac{q_1}{\kappa} (\tau(s)\lambda_2(s) + \lambda_3'(s)) \sqrt{\kappa^2 + \tau^2}, \quad (3.12)$$

where $q_1 = \text{sgn}(\tau(s)\lambda_2(s) + \lambda_3'(s))$.

Substituting into Eq (3.11), we have

$$\tilde{t}(s) = \frac{q_1}{\sqrt{\kappa^2 + \tau^2}} (\tau(s) \mathbf{t}(s) + \kappa(s) \mathbf{b}(s)), \quad (3.13)$$

and hence $\langle \tilde{t}(s), \tilde{t}(s) \rangle_{G_3} = 1$.

Taking

$$\alpha_1(s) = \frac{q_1 \tau(s)}{\sqrt{\kappa^2 + \tau^2}}, \quad \text{and} \quad \alpha_2(s) = \frac{q_1 \kappa(s)}{\sqrt{\kappa^2 + \tau^2}},$$

then

$$\tilde{t}(s) = \alpha_1(s) \mathbf{t}(s) + \alpha_2(s) \mathbf{b}(s). \quad (3.14)$$

Equation (3.14) can be differentiated with respect to s to yield

$$\frac{d\tilde{t}}{ds} = \frac{1}{g'(s)} (\alpha_1' \mathbf{t}(s) + \alpha_2' \mathbf{b}(s)),$$

and then

$$\tilde{\kappa} = \sqrt{\left\langle \frac{d\tilde{t}}{ds}, \frac{d\tilde{t}}{ds} \right\rangle} = \frac{\sqrt{\alpha_1'^2(s) + \alpha_2'^2(s)}}{g'(s)} = \frac{q_2 (\kappa\tau' - \tau\kappa')}{g'(s)(\kappa^2 + \tau^2)} = \frac{q_2 \kappa^2 \left(\frac{\tau}{\kappa}\right)'}{g'(s)(\kappa^2 + \tau^2)}, \quad (3.15)$$

such that $q_2 = \text{sgn}(\kappa\tau' - \tau\kappa')$. Then, we may determine $\tilde{n}(s)$ as follows:

$$\tilde{n}(s) = \frac{q_1}{q_2 \sqrt{\kappa^2 + \tau^2}} (\kappa \mathbf{t}(s) - \tau \mathbf{b}(s)), \quad (3.16)$$

and $\langle \tilde{n}(\tilde{s}), \tilde{n}(\tilde{s}) \rangle_{G_3} = 1$. Also, $\tilde{b}(\tilde{s})$ can be defined as

$$\tilde{b}(\tilde{s}) = \tilde{t}(\tilde{s}) \times \tilde{n}(\tilde{s}) = \frac{1}{q_2} \mathbf{n}(s), \quad (3.17)$$

and $\langle \tilde{b}(\tilde{s}), \tilde{b}(\tilde{s}) \rangle_{G_3} = 1$. Finally, we get

$$\tilde{\tau}(\tilde{s}) = - \langle \tilde{b}'(\tilde{s}), \tilde{n}(\tilde{s}) \rangle = \frac{q_1 \tau^2}{\sqrt{\kappa^2 + \tau^2}} \neq 0. \quad (3.18)$$

Therefore, $\tilde{\gamma}$ is a Mannheim partner curve of γ . Then γ is a Mannheim curve. \square

In Theorem (3.1), if we set $\lambda_1(s) = \lambda_3(s) = 0$, using the Mannheim mate curve $\tilde{\gamma}$, which is expressed as

$$\tilde{\gamma}(s) = \gamma(s) + \lambda_2(s) \mathbf{n}(s),$$

we may derive the criteria of classical Mannheim curves seen in the literature.

Corollary 3.2. *Assume that the admissible curve $\gamma : I \rightarrow G_3$, $I \subset \mathbb{R}$ has non-vanishing curvatures κ and τ . After that, the curve γ is a Mannheim curve, and its Mannheim partner $\tilde{\gamma}$ is defined as*

$$\tilde{\gamma}(s) = \gamma(s) + \lambda_2(s) \mathbf{n}(s),$$

if $\lambda_2(s)$ exists and it fulfills $\kappa = \lambda_2 \tau^2$, and the opposite is again valid.

Corollary 3.3. *Let $\gamma(s)$ constitute an admissible Mannheim curve through G_3 using s as the parameter for arc length, and let $\tilde{\gamma}(\tilde{s})$ be an admissible Mannheim partner curve of γ . Assume $\gamma(s)$ is a generalized helix, then $\tilde{\gamma}(\tilde{s})$ represents a straight line.*

Proof. Assume $\gamma : I \rightarrow G_3$, $I \subset \mathbb{R}$ is an admissible Mannheim general helix that has non-vanishing curvatures κ and τ . Following that, the ratio $\frac{\tau}{\kappa}$ is constant, which gives from Eq (3.15) that $\tilde{\kappa} = 0$. Then, $\tilde{\gamma}(\tilde{s})$ defines a line. \square

Corollary 3.4. *Take $\gamma : I \subset \mathbb{R} \rightarrow G_3$ to be an admissible Mannheim curve with non-vanishing curvatures κ and τ , and let curve $\tilde{\gamma}$ be the Mannheim partner curve of γ having $\tilde{\kappa}$ and $\tilde{\tau}$ non-zero curvatures. Thus, given that $g'(s)$ has a non-zero value, $\tilde{\gamma}$ indicates a general helix if the slant helix is represented by γ . In addition, the opposite also holds.*

Proof. Consider $\gamma : I \subset \mathbb{R} \rightarrow G_3$ as an admissible Mannheim curve with κ ; and τ non-vanishing curvatures and let $\tilde{\gamma}$ be the Mannheim partner curve of γ having non-zero curvatures $\tilde{\kappa}$ and $\tilde{\tau}$. And, following that, from Eqs (3.15) and, (3.18), we obtain

$$\frac{\tilde{\kappa}}{\tilde{\tau}} = \frac{q_2 \kappa^2 \left(\frac{\tau}{\kappa}\right)'}{q_1 g'(s) \tau^3 \sqrt{\left(\frac{\tau}{\kappa}\right)^2 + 1}}.$$

If $g'(s)$ is a non-zero constant, then $\tilde{\gamma}$ can only be considered a general helix in the event that γ is a slant helix. \square

In the next part, we introduce special cases for the Mannheim curve in G_3 and its Mannheim partner curve.

Case 3.5. Assuming that $\gamma : I \rightarrow G_3$, $I \subset \mathbb{R}$ is an admissible Mannheim curve with non-zero curvatures $\kappa(s)$ and $\tau(s)$, then the requirements of Theorem (3.1) have been met. Suppose that $\lambda_2 = \lambda \in \mathbb{R}$. Next, we are able to acquire

$$\kappa(s)\lambda_1(s) = \tau(s)\lambda_3(s) \quad \text{and} \quad (1 + \lambda'_1(s))\kappa(s) = (\lambda'_3(s) + \lambda\tau(s))\tau(s),$$

which gives that

$$\lambda_3(s) = \frac{\lambda\tau^2(s) - \kappa(s)}{\kappa(s)\left(\frac{\tau}{\kappa}\right)'} \quad \text{and} \quad \lambda_1(s) = \frac{\tau(s)[\lambda\tau^2(s) - \kappa(s)]}{\kappa^2(s)\left(\frac{\tau}{\kappa}\right)'}$$

Therefore, we get the Mannheim partner curve $\tilde{\gamma}$ as

$$\tilde{\gamma}(\tilde{s}) = \gamma(s) + \frac{\tau(s)[\lambda\tau^2(s) - \kappa(s)]}{\kappa^2(s)\left(\frac{\tau}{\kappa}\right)'} \mathbf{t}(s) + \lambda \mathbf{n}(s) + \frac{\lambda\tau^2(s) - \kappa(s)}{\kappa(s)\left(\frac{\tau}{\kappa}\right)'} \mathbf{b}(s).$$

The following illustration displays the Mannheim curve as an example of the Salkowski curve.

Example 3.6. Consider the Salkowski curve in G_3 given by

$$\gamma(s) = \left(s, \frac{1}{4} (3 - 4s)\cos(2\sqrt{s}) + 6\sqrt{s} \sin(2\sqrt{s}), \frac{1}{4} ((3 - 4s)\sin(2\sqrt{s}) - 6\sqrt{s} \cos(2\sqrt{s})) \right),$$

along with curvatures $\kappa(s) = 1$ and $\tau(s) = \frac{1}{\sqrt{s}}$ and the Frenet frame

$$\begin{aligned} \mathbf{t}(s) &= \left(1, \sqrt{s}\sin(2\sqrt{s}) + \frac{1}{2}(\cos(2\sqrt{s})), -\sqrt{s}\cos(2\sqrt{s}) + \frac{1}{2}(\sin(2\sqrt{s})) \right), \\ \mathbf{n}(s) &= (0, \cos(2\sqrt{s}), \sin(2\sqrt{s})), \\ \mathbf{b}(s) &= (0, -\sin(2\sqrt{s}), \cos(2\sqrt{s})). \end{aligned}$$

If we set $\lambda = 0$ in case (3.5), the Mannheim partner curve $\tilde{\gamma}(\tilde{s})$ may be derived as

$$\tilde{\gamma}(\tilde{s}) = \left(3s, \frac{3}{4}(\cos(2\sqrt{s}) + 2\sqrt{s}\sin(2\sqrt{s})), \frac{3}{4}(\sin(2\sqrt{s}) - 2\sqrt{s}\cos(2\sqrt{s})) \right), \quad (3.19)$$

along with curvatures

$$\tilde{\kappa} = \frac{\sqrt{s(s+1)}}{6s(s+1)^2}, \quad \text{and} \quad \tilde{\tau} = \frac{\sqrt{s(s+1)}}{s(s+1)},$$

and it is possible to acquire the Frenet frame as follows

$$\begin{aligned} \tilde{\mathbf{t}}(\tilde{s}) &= \left(\frac{1}{\sqrt{s+1}}, \frac{1}{2\sqrt{s+1}}\cos(2\sqrt{s}), \frac{1}{2\sqrt{s+1}}\sin(2\sqrt{s}) \right), \\ \tilde{\mathbf{n}}(\tilde{s}) &= \left(\frac{-\sqrt{s}}{\sqrt{s+1}}, -\sqrt{s+1}\sin(2\sqrt{s}) - \frac{\sqrt{s}}{2\sqrt{s+1}}\cos(2\sqrt{s}), \sqrt{s+1}\cos(2\sqrt{s}) - \frac{\sqrt{s}}{2\sqrt{s+1}}\sin(2\sqrt{s}) \right), \end{aligned}$$

$$\tilde{b}(\bar{s}) = (0, -\cos(2\sqrt{s}), -\sin(2\sqrt{s})).$$

We can easily obtain $n(s) = -\tilde{b}(\bar{s})$, implying that the Mannheim curve $\gamma(s)$ possesses a Mannheim partner curve $\tilde{\gamma}(\bar{s})$. And, $\tilde{\gamma}(\bar{s})$ is not a general helix (Figure 1).

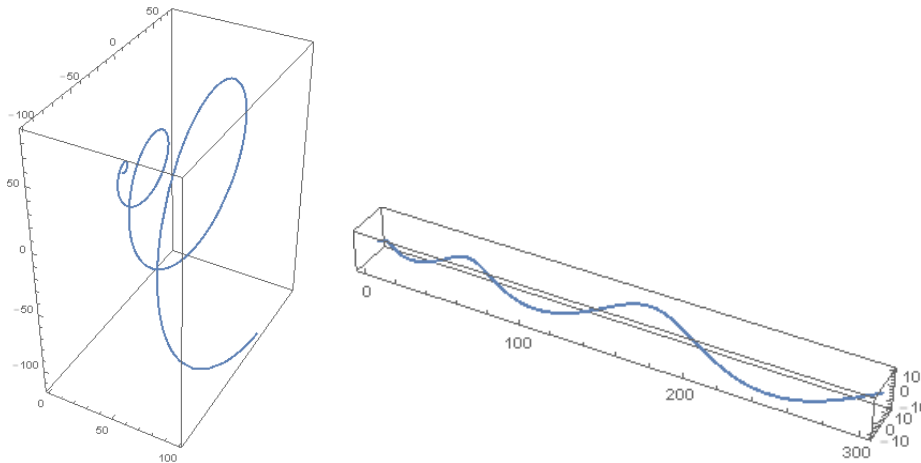


Figure 1. Curves $\gamma(s)$ and $\tilde{\gamma}(\bar{s})$ diagram of Example 3.6.

In the next part, we present a Mannheim Curve as an illustration of the anti-Salkowski curve inside G_3 .

Example 3.7. Define $\gamma(s)$ to be an anti-Salkowski curve in G_3 , which is given by

$$\gamma(s) = \left(s, \frac{16}{289} \left[8 \sin(s) \sinh\left(\frac{s}{4}\right) - 15 \cos(s) \cosh\left(\frac{s}{4}\right) \right], -\frac{16}{289} \left[8 \cos(s) \sinh\left(\frac{s}{4}\right) + 15 \sin(s) \cosh\left(\frac{s}{4}\right) \right] \right)$$

with curvatures $\kappa(s) = \cosh\left(\frac{s}{4}\right)$ and $\tau(s) = 1$, and the Frenet frame in the manner described below:

$$\mathbf{t}(s) = \left(1, \frac{16}{17} \left[\sin(s) \cosh\left(\frac{s}{4}\right) + \frac{1}{4} \cos(s) \sinh\left(\frac{s}{4}\right) \right], -\frac{16}{17} \left[\cos(s) \cosh\left(\frac{s}{4}\right) - \frac{1}{4} \sin(s) \sinh\left(\frac{s}{4}\right) \right] \right),$$

$$\mathbf{n}(s) = (0, \cos(s), \sin(s)),$$

$$\mathbf{b}(s) = (0, -\sin(s), \cos(s)).$$

Taking $\lambda = 0$ in case (3.5), the Mannheim partner curve $\tilde{\gamma}(\bar{s})$ is then obtained in the following manner:

$$\begin{aligned} \tilde{\gamma}(\bar{s}) = & \left(s + 4 \coth\left(\frac{s}{4}\right), \frac{128}{289} \sin(s) \sinh\left(\frac{s}{4}\right) + \frac{32}{289} \cos(s) \cosh\left(\frac{s}{4}\right) - \frac{4}{17} \sin(s) \cosh\left(\frac{s}{4}\right) \coth\left(\frac{s}{4}\right), \right. \\ & \left. -\frac{128}{289} \cos(s) \sinh\left(\frac{s}{4}\right) + \frac{32}{289} \sin(s) \cosh\left(\frac{s}{4}\right) + \frac{4}{17} \cos(s) \cosh\left(\frac{s}{4}\right) \coth\left(\frac{s}{4}\right) \right), \end{aligned}$$

with curvatures $\tilde{\kappa} = \frac{-\operatorname{sgn}\left(\frac{1}{4} \sinh\left(\frac{s}{4}\right)\right) \tanh\left(\frac{s}{4}\right)}{4 \left(1 - \operatorname{sech}^2\left(\frac{s}{4}\right)\right) \left(1 + \cosh^2\left(\frac{s}{4}\right)\right)^{\frac{3}{2}}}$ and $\tilde{\tau} = \frac{\operatorname{sgn}\left(1 - \operatorname{cosech}^2\left(\frac{s}{4}\right)\right)}{\sqrt{1 + \cosh^2\left(\frac{s}{4}\right)}}$, and the Frenet frame obtained

as

$$\tilde{\mathbf{t}}(\bar{s}) = \frac{\operatorname{sgn}\left(1 - \operatorname{cosech}^2\left(\frac{s}{4}\right)\right)}{\sqrt{1 + \cosh^2\left(\frac{s}{4}\right)}} \left(1, \frac{1}{17} \left[-\sin(s) \cosh\left(\frac{s}{4}\right) + 4 \cos(s) \sinh\left(\frac{s}{4}\right) \right], \frac{1}{17} \left[\cos(s) \cosh\left(\frac{s}{4}\right) + 4 \sin(s) \sinh\left(\frac{s}{4}\right) \right] \right),$$

$$\begin{aligned} \bar{n}(s) &= \frac{\operatorname{sgn}\left(1 - \operatorname{cosech}^2\left(\frac{s}{4}\right)\right)}{\operatorname{sgn}\left(\frac{1}{4} \sinh\left(\frac{s}{4}\right)\right) \sqrt{1 + \cosh^2\left(\frac{s}{4}\right)}} \left(\cosh\left(\frac{s}{4}\right), \cosh\left(\frac{s}{4}\right) \left[\frac{16}{17} \sin(s) \cosh\left(\frac{s}{4}\right) + \frac{4}{17} \cos(s) \sinh\left(\frac{s}{4}\right) \right] + \sin(s), \right. \\ &\quad \left. \cosh\left(\frac{s}{4}\right) \left[-\frac{16}{17} \cos(s) \cosh\left(\frac{s}{4}\right) + \frac{4}{17} \sin(s) \sinh\left(\frac{s}{4}\right) \right] - \cos(s) \right), \\ \bar{b}(s) &= \frac{1}{\operatorname{sgn}\left(\frac{1}{4} \sinh\left(\frac{s}{4}\right)\right)} (0, \cos(s), \sin(s)). \end{aligned}$$

It is simple to obtain $n(s) = \operatorname{sgn}\left(\frac{1}{4} \sinh\left(\frac{s}{4}\right)\right) \bar{b}(s)$, indicating that $\gamma(s)$ is a Mannheim curve and that $\bar{\gamma}(s)$ is its Mannheim partner curve. Again, $\bar{\gamma}(s)$ is not a general helix (Figure 2).

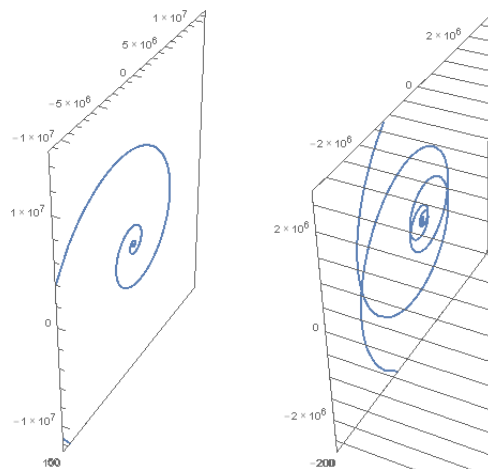


Figure 2. Diagram of curves $\gamma(s)$ and $\bar{\gamma}(s)$ in Example 3.7.

Example 3.8. Consider the curve in G_3 defined by

$\gamma(s) = \left(s, -\sqrt{\pi} s \operatorname{FresnelS}\left(\frac{s}{\sqrt{\pi}}\right) - 2\cos\left(\frac{s^2}{2}\right), \sqrt{\pi} s \operatorname{FresnelC}\left(\frac{s}{\sqrt{\pi}}\right) - 2\sin\left(\frac{s^2}{2}\right) \right)$, with the curvatures $\kappa(s) = s^2$ and $\tau(s) = s$, where $\operatorname{FresnelS}(x) = \int \sin\left(\frac{\pi x^2}{2}\right) dx$ and $\operatorname{FresnelC}(x) = \int \cos\left(\frac{\pi x^2}{2}\right) dx$, and the Frenet frame is obtained as

$$\begin{aligned} \mathbf{t}(s) &= \left(1, s \sin\left(\frac{s^2}{2}\right) - \int \sin\left(\frac{s^2}{2}\right) ds, -s \cos\left(\frac{s^2}{2}\right) + \int \cos\left(\frac{s^2}{2}\right) ds \right), \\ \mathbf{n}(s) &= \left(0, \cos\left(\frac{s^2}{2}\right), \sin\left(\frac{s^2}{2}\right) \right), \\ \mathbf{b}(s) &= \left(0, -\sin\left(\frac{s^2}{2}\right), \cos\left(\frac{s^2}{2}\right) \right). \end{aligned}$$

Taking $\lambda = 0$ in case (3.5), we acquire the Mannheim partner curve $\bar{\gamma}(s)$ according to

$$\bar{\gamma}(s) = \left(2s, -2s \int \sin\left(\frac{s^2}{2}\right) ds - 2\cos\left(\frac{s^2}{2}\right), 2s \int \cos\left(\frac{s^2}{2}\right) ds - 2\sin\left(\frac{s^2}{2}\right) \right),$$

with curvatures $\bar{\kappa} = \frac{1}{2(s^2+1)^{\frac{3}{2}}}$ and $\bar{\tau} = \frac{s}{\sqrt{s^2+1}}$. The Frenet frame is calculated as

$$\begin{aligned}\tilde{t}(\bar{s}) &= \left(\frac{1}{\sqrt{s^2+1}}, \frac{-1}{\sqrt{s^2+1}} \int \sin\left(\frac{s^2}{2}\right) ds, \frac{1}{\sqrt{s^2+1}} \int \cos\left(\frac{s^2}{2}\right) ds \right), \\ \tilde{n}(\bar{s}) &= \left(-\frac{s}{\sqrt{s^2+1}}, -\sqrt{s^2+1} \sin\left(\frac{s^2}{2}\right) + \frac{s}{\sqrt{s^2+1}} \int \sin\left(\frac{s^2}{2}\right) ds, \sqrt{s^2+1} \cos\left(\frac{s^2}{2}\right) - \frac{s}{\sqrt{s^2+1}} \int \cos\left(\frac{s^2}{2}\right) ds \right), \\ \tilde{b}(\bar{s}) &= \left(0, -\cos\left(\frac{s^2}{2}\right), -\sin\left(\frac{s^2}{2}\right) \right).\end{aligned}$$

Then, we can deduce that $n(s) = -\tilde{b}(\bar{s})$, which means that $\gamma(s)$ is a Mannheim curve, along with $\tilde{\gamma}(\bar{s})$ serving as its Mannheim partner curve. Also, $\tilde{\gamma}(\bar{s})$ is not a general helix (Figure 3).

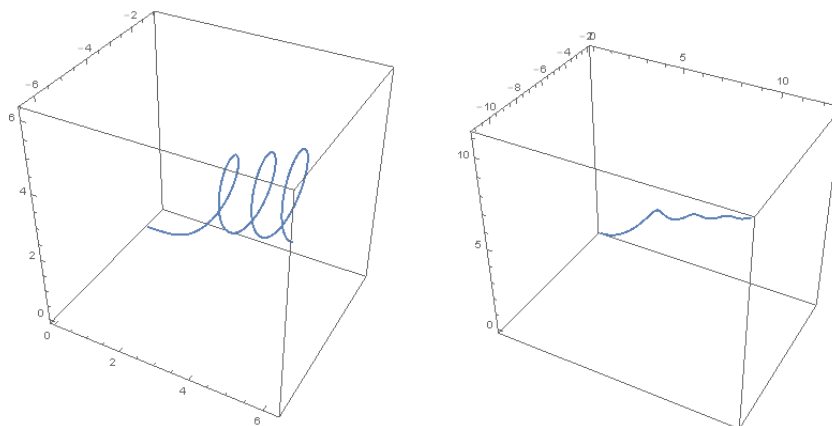


Figure 3. Curves $\gamma(s)$ and $\tilde{\gamma}(\bar{s})$ of Example 3.8.

Theorem 3.9. Let $\tilde{\gamma} : \tilde{I} \rightarrow G_3$, $\tilde{I} \subset \mathbb{R}$ be an admissible curve utilizing the $\{\tilde{t}, \tilde{n}, \tilde{b}\}$ Frenet frame as well as the non-zero curvatures $\tilde{\kappa}$ and $\tilde{\tau}$. For each given Mannheim curve, if its Mannheim partner curve is represented by $\tilde{\gamma}$, then there exist differentiable functions μ_1, μ_2, μ_3 , and δ satisfying the following two cases:

(1) If $\dot{B} \neq 0$, then we have

$$\tilde{\tau}\mu_2(s) + \dot{\mu}_3(\bar{s}) = 0, \quad \frac{1 + \dot{\mu}_1(\bar{s})}{\dot{f}} \neq 0, \quad \dot{\delta} = -\tilde{\kappa}, \quad \mu_2(\bar{s}) + \tilde{\kappa}\mu_1(\bar{s}) - \tilde{\tau}\mu_3(\bar{s}) = \delta(1 + \mu_1(\bar{s})). \quad (3.20)$$

(2) If $\dot{B} = 0$, then we have

$$\tilde{\tau}\mu_2(\bar{s}) + \mu_3(\bar{s}) = 0, \quad \dot{\mu}_2(\bar{s}) - \tilde{\tau}\mu_3(\bar{s}) + \tilde{\kappa}(-\bar{s} + c_o) = d\dot{f}, \quad (3.21)$$

where c_o and d are non-zero constants and $B = \frac{\mu_2(\bar{s}) + \tilde{\kappa}\mu_1(\bar{s}) - \tilde{\tau}\mu_3(\bar{s})}{\dot{f}}$. In this case, “ $\dot{\cdot}$ ” signifies the derivative in relation to \bar{s} .

Proof. Suppose that γ is an admissible Mannheim curve with non-zero κ and τ curvatures, specified by the arc length s , and the Mannheim partner curve of curve γ is represented by curve $\tilde{\gamma}$, where $\tilde{\gamma}$ is parameterized by the arc length \bar{s} . Thus, we may define the curve γ as

$$\gamma(s) = \gamma(f(\bar{s})) = \tilde{\gamma}(\bar{s}) + \mu_1(\bar{s})\tilde{t}(\bar{s}) + \mu_2(\bar{s})\tilde{n}(\bar{s}) + \mu_3(\bar{s})\tilde{b}(\bar{s}), \quad (3.22)$$

with every $\bar{s} \in \bar{I}$, such that $\mu_1(\bar{s})$, $\mu_2(\bar{s})$, and $\mu_3(\bar{s})$ are differentiable functions on \bar{I} .

When we differentiate Eq (3.22) considering \bar{s} , we arrive at

$$\mathbf{t}(s)\dot{f} = \left(1 + \dot{\mu}_1(\bar{s})\right)\bar{t}(\bar{s}) + \left(\dot{\mu}_2(\bar{s}) + \bar{\kappa}\mu_1(\bar{s}) - \bar{\tau}\mu_3(\bar{s})\right)\bar{n}(\bar{s}) + \left(\bar{\tau}\mu_2(\bar{s}) + \dot{\mu}_3(\bar{s})\right)\bar{b}(\bar{s}). \quad (3.23)$$

By taking the dot product of (3.23) with $\bar{b}(\bar{s})$, we get

$$\bar{\tau}\mu_2(\bar{s}) + \dot{\mu}_3(\bar{s}) = 0. \quad (3.24)$$

Inserting (3.24) in (3.23), we have

$$\mathbf{t}(s)\dot{f} = \left(1 + \dot{\mu}_1(\bar{s})\right)\bar{t}(\bar{s}) + \left(\dot{\mu}_2(\bar{s}) + \bar{\kappa}\mu_1(\bar{s}) - \bar{\tau}\mu_3(\bar{s})\right)\bar{n}(\bar{s}). \quad (3.25)$$

By taking the scalar product of (3.25) with itself, we deduce

$$(\dot{f})^2 = \left(1 + \dot{\mu}_1(\bar{s})\right)^2 + \left(\dot{\mu}_2(\bar{s}) + \bar{\kappa}\mu_1(\bar{s}) - \bar{\tau}\mu_3(\bar{s})\right)^2. \quad (3.26)$$

If we set

$$A = \frac{1 + \dot{\mu}_1(\bar{s})}{\dot{f}}, \quad \text{and} \quad B = \frac{\dot{\mu}_2(\bar{s}) + \bar{\kappa}\mu_1(\bar{s}) - \bar{\tau}\mu_3(\bar{s})}{\dot{f}}, \quad (3.27)$$

we obtain

$$\mathbf{t}(s) = A\bar{t}(\bar{s}) + B\bar{n}(\bar{s}). \quad (3.28)$$

Differentiating (3.28) with respect to \bar{s} , we obtain

$$\dot{f} \kappa \mathbf{n}(s) = \dot{A}\bar{t}(\bar{s}) + (A\bar{\kappa} + \dot{B})\bar{n}(\bar{s}) + \bar{\tau}B\bar{b}(\bar{s}). \quad (3.29)$$

Since the principle normal vector $\mathbf{n}(s)$ of the curve $\gamma(s)$ and the bi-normal vector $\bar{b}(\bar{s})$ of its Mannheim Partner curve is linearly dependent, we reach

$$\bar{\kappa} = \frac{-\dot{B}}{A}.$$

We have two cases:

- If $\dot{B} \neq 0$, then $A \neq 0$ and $B \neq 0$, which implies

$$\frac{1 + \dot{\mu}_1(\bar{s})}{\dot{f}} \neq 0, \quad \text{and} \quad \dot{\mu}_2(\bar{s}) + \bar{\kappa}\mu_1(\bar{s}) - \bar{\tau}\mu_3(\bar{s}) = \delta(1 + \dot{\mu}_1(\bar{s})),$$

where $\delta = \frac{B}{A}$.

Again, from (3.29), utilizing $\bar{b}(\bar{s})$ to obtain the scalar product, we acquire $\dot{A} = 0$ and $A\bar{\kappa} + \dot{B} = 0$, which gives

$$\dot{\delta} = -\bar{\kappa}.$$

- If $\dot{B} = 0$, then $A = 0$, which implies $\dot{\mu}_1 = -1$ and $\mu_1 = -\bar{s} + c_o$, where c_o is constant. Also, we can deduce that B is non-zero constant (d , say). Now, from Eq (3.27), we have

$$\dot{\mu}_2(\bar{s}) - \bar{\tau}\mu_3(\bar{s}) + \bar{\kappa}(-\bar{s} + c_o) = d\dot{f}$$

□

If we set $\mu_1(\bar{s}) = \mu_2(\bar{s}) = 0$ within Theorem (3.9), we arrive at the requirements of traditional Mannheim partner curves found in the literature by using

$$\gamma(s) = \gamma(f(\bar{s})) = \bar{\gamma}(\bar{s}) + \mu_3(\bar{s})\bar{b}(\bar{s}).$$

Also, we get

$$\dot{\mu}_3(\bar{s}) = 0, \quad \dot{\delta} = -\bar{\kappa}, \quad -\bar{\tau}\mu_3(\bar{s}) = \delta.$$

So, we have

$$\bar{\tau} = \frac{-1}{\mathbf{a}} \int \bar{\kappa} d\bar{s},$$

for some non-zero real number \mathbf{a} . Therefore, this leads us to the next corollary.

Corollary 3.10. *Suppose that $\bar{\gamma} : \bar{I} \rightarrow G_3$, $\bar{I} \subset \mathbb{R}$ is an admissible curve with non-zero curvatures $\bar{\kappa}$ and $\bar{\tau}$, and with the Frenet frame $\{\bar{t}, \bar{n}, \bar{b}\}$. If $\bar{\gamma}$ is a Mannheim partner curve of a certain Mannheim curve given by*

$$\gamma(s) = \gamma(f(\bar{s})) = \bar{\gamma}(\bar{s}) + \mu_3(s)\bar{b}(\bar{s}),$$

then a non-zero real number \mathbf{a} exists such that

$$\bar{\tau} = \frac{-1}{\mathbf{a}} \int \bar{\kappa} d\bar{s}.$$

4. Conclusions

In this study, a novel method for Mannheim curves in three-dimensional Galilean space was presented. The necessary and sufficient conditions for a curve to be a Mannheim curve were obtained. Finally, some examples were introduced.

Author contributions

Mervat Elzawy: conceptualization, methodology, software, reviewing and editing; Safaa Mosa: data curation, writing-original draft preparation, visualization, investigation, reviewing and editing.

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Conflict of interest

All authors declare no conflicts of interest in this paper.

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