



Research article

Two-grid finite element method with an H2N2 interpolation for two-dimensional nonlinear fractional multi-term mixed sub-diffusion and diffusion wave equation

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Abstract: In this paper, we studied the two-grid method (TGM) for two-dimensional nonlinear time fractional multi-term mixed sub-diffusion and diffusion wave equation. A fully discrete scheme with the quadratic Hermite and Newton interpolation (H2N2) method was considered in the temporal direction and the expanded finite element method is used to approximate the spatial direction. In order to reduce computational time, a dual grid method based on Newton iteration was constructed with order $\alpha \in (0, 1)$ and $\beta \in (1, 2)$. The global convergence order of the two-grid scheme reaches $O(\tau^{3-\beta} + h^{r+1} + H^{2r+2})$, where τ , H and h are the time step size, coarse grid mesh size and fine grid mesh size, respectively. The error estimation and stability of the fully discrete scheme were derived. Theoretical analysis shows that the two grid algorithms maintain asymptotic optimal accuracy while saving computational costs. In addition, numerical experiments further confirmed the theoretical results.

Keywords: nonlinear fractional multi-term mixed sub-diffusion and diffusion wave equations; two-grid method (TGM); finite element method; H2N2 interpolation; stability and convergence

Mathematics Subject Classification: 26A33, 65M60, 65N30

1. Introduction

In recent decades, fractional partial differential equations (FPDEs) have become popular for modeling anomalous transport processes. In virtue of the difficulty for looking for the exact solutions of FPDEs, the construction of numerical methods for FPDEs have been studied extensively by many scholars which mainly cover the finite difference method, finite element method, spectral method and so on. Unlike general PDEs, it has been found that fractional derivatives are nonlocal, with long memory and weak singular kernels, and it takes more time to solve fractional differential equations when using some well-known L-type discretization formulas. In order to increase computational

efficiency, some fast algorithms have been presented, such as the sum-of-exponentials technique [1, 2] and the two-grid technique [3–9].

There have been many studies on the approximation of time fractional order equations. Schumer [10] first developed the fractional-order mobile/immobile model in which the time drift term $\partial u/\partial t$ is added to describe the motion time, thus helps to distinguish the status of particles conveniently. Zhang and his coauthors [11] proposed a difference scheme combining the compact difference approach for spatial discretization and the alternating direction implicit (ADI) method in the time stepping for the two-dimensional time fractional diffusion-wave equation. In [12], the authors established a fully-discrete approximate scheme for the 2D multi-term time-fractional mixed diffusion and diffusion-wave equations with a spatial variable coefficient by using the linear triangle finite element method in space and classical L1 time-stepping method combined with the Crank–Nicolson scheme in time. Sun et al. [13] gave an L-type difference scheme for time-fractional mixed sub-diffusion and diffusion wave equations, the new analytical technique with a $\min\{2 - \alpha, 3 - \beta\}$ order accuracy in the discrete H^1 norm and second order accuracy in space, where $\alpha \in (0, 1)$, $\beta \in (1, 2)$. In recent years, Shen et al. [14] derived the H2N2 method to develop a known numerical differential formula of the Caputo fractional derivative to obtain a $(3 - \alpha)$ order temporal convergence rate with $\alpha \in (1, 2)$. Moreover, Shen and his coauthors [15] also gave the finite difference methods based on an H2N2 interpolation for two-dimensional time fractional mixed sub-diffusion and diffusion-wave equations with a $(3 - \beta)$ order temporal convergence rate and second order accuracy in space.

With further research, it is found that nonlinear partial differential equations can describe some model problems more effectively compared to linear fractional differential equations. As we know, numerical methods for nonlinear parabolic fractional partial differential problems have been extensively studied, especially for the two-grid approximation of nonlinear reaction terms. In [4, 5], Xu utilized the two-grid method to discretize asymmetric, indefinite, and nonlinear partial differential equations. Afterward, many authors devoted themselves to the research of the two-grid method. For example, Chen and Chen [6] investigated a scheme for nonlinear reaction-diffusion equations by using the mixed finite element methods, in which the two-grid method (TGM) studied provided a new approach to take the advantage of some nice properties hidden in a complex problem. Based on the finite difference approximation in the time direction and the finite element method in the spatial direction, a class of time fractional fourth order reaction diffusion problems with nonlinear reaction terms was solved by Liu et al. [16]. Almost simultaneously, Liu et al. [17] considered a nonlinear fractional Cable equation by a two-grid algorithm combined arriving at a second-order convergence rate independent of fractional parameters α ($0 < \alpha < 1$) and β ($0 < \beta < 1$) with finite element method. In addition, Bi et al. [18] established a discontinuous Galerkin finite element scheme for the second-order nonlinear elliptic problem, using piecewise polynomials of degree $r \geq 2$ based on pointwise error estimation, and provided an error estimate of the energy norm of a two-grid mesh algorithm. Li and Rui [8] introduced and analyzed a two-grid block-centered finite difference scheme for the nonlinear time-fractional parabolic equation with the Neumann condition, and the error estimates established on a non-uniform rectangular grid with the discrete $L^\infty(L^2)$ and $L^2(H^1)$ errors were $O(t^{2-\alpha} + h^2 + H^3)$. Jin et al. [19] demonstrated a general criterion for showing the fractional discrete Grönwall inequality and verified it for the L1 scheme and convolution quadrature generated by backward difference formulas for time-fractional nonlinear parabolic partial differential equations. In 2019, Li et al. [20] presented and developed a Newton linearized Galerkin finite element method to solve nonlinear time fractional

parabolic problems with non-smooth solutions in time direction. Recently, Chen et al. [7] constructed and studied a two-grid finite element method for 2D nonlinear time fractional two-term mixed sub-diffusion and diffusion wave equations and got $\min(2-\alpha, 3-\beta)$ order in the time direction where $\alpha \in (0, 1)$, $\beta \in (1, 2)$.

However, to our knowledge, there is no two-grid finite element method with an H2N2 interpolation for two-dimensional nonlinear fractional multi-term mixed sub-diffusion and diffusion wave equations, and our aim is to present the corresponding algorithm in this paper. We consider the numerical solution of the following nonlinear time-fractional multi-term mixed sub-diffusion and diffusion wave equations:

$$\begin{cases} \partial_t u(\mathbf{x}, t) + \sum_{i=1}^{M_1} p_i \partial_t^{\alpha_i} u(\mathbf{x}, t) + \sum_{j=1}^{M_2} q_j \partial_t^{\beta_j} u(\mathbf{x}, t) - \Delta u(\mathbf{x}, t) = g(\mathbf{u}), (\mathbf{x}, t) \in \Omega \times (0, T] \\ u(\mathbf{x}, 0) = u_0(\mathbf{x}), u_t(\mathbf{x}, 0) = u_t^0(\mathbf{x}), \mathbf{x} \in \Omega, \end{cases} \quad (1.1)$$

where M_1, M_2 are positive integers, p_i, q_j are nonnegative constants, $\Omega \subset \mathbb{R}^2$ is a bounded convex polygonal region with boundary $\partial\Omega$, $\mathbf{x} = [x, y]$ and $g(\cdot)$ is twice continuously differentiable. Δ is a Laplace operator, and $u_0(\mathbf{x}), u_t^0(\mathbf{x})$ are given as sufficiently smooth functions. The Caputo fractional derivative $\partial_t^{\alpha_i} u(\mathbf{x}, t), \partial_t^{\beta_j} u(\mathbf{x}, t)$ is defined by

$$\partial_t^{\alpha_i} u(\mathbf{x}, t) = \frac{1}{\Gamma(1-\alpha_i)} \int_0^t \frac{\partial u(\mathbf{x}, s)}{\partial s} \frac{1}{(t-s)^{\alpha_i}} ds, \quad (1.2)$$

$$\partial_t^{\beta_j} u(\mathbf{x}, t) = \frac{1}{\Gamma(2-\beta_j)} \int_0^t \frac{\partial^2 u(\mathbf{x}, s)}{\partial s^2} \frac{1}{(t-s)^{\beta_j-1}} ds, \quad (1.3)$$

with $0 < \alpha_0 < \dots < \alpha_i < \dots < \alpha_{M_1} < 1$ ($i = 1, 2, \dots, M_1$) and $1 < \beta_0 < \dots < \beta_j < \dots < \beta_{M_2} < 2$ ($j = 1, 2, \dots, M_2$). For convenience, we define $\alpha_{M_1} = \alpha, \beta_{M_2} = \beta$. Throughout this paper, the notation C denotes a generic constant, which may vary at different occurrences, but it is always independent of the coarse grid mesh size H , fine grid mesh size h and time step size τ .

The remaining outline of the article is constructed as follows. In section two, we propose some preliminary knowledge of fractional derivatives and some lemmas on time approximations. The unconditional stability and the error estimates of the two-grid finite element method (FEM) are discussed in section three. In section four, numerical examples are presented to demonstrate the efficiency of the theoretical results. Finally, a brief conclusion of the article was provided.

2. Construction of fully-discrete scheme

Let $W_p^m(\Omega)$ be the standard Sobolev space with a norm $\|\cdot\|_{m,p}$ given by $\|v\|_{m,p}^p = \sum_{|\alpha| \leq m} \|D^\alpha v\|_{L^p(\Omega)}^p$. For $p = 2$, we denote $H^m(\Omega) = W_2^m(\Omega)$ and $H_0^1(\Omega)$ to be the subspace of $H^1(\Omega)$ consisting of functions with a vanishing trace on $\partial\Omega$. $\|\cdot\|_m = \|\cdot\|_{m,2}$ and $\|\cdot\| = \|\cdot\|_{0,2}$. Take an integer N and denote $\tau = \frac{T}{N}$, $t_n = n\tau$, $t_{n-\frac{1}{2}} = \frac{1}{2}(t_n + t_{n-1})$. For a sequence of smooth functions $\{u(t)\}_{n=0}^N$ on $[0, T]$, we denote

$$u^n = u(t_n), u_t(\mathbf{x}, 0) = u_t^0(\mathbf{x}), u^{n-\frac{1}{2}} = \frac{u^n + u^{n-1}}{2}, \delta_t u^{k-\frac{1}{2}} = \frac{u^k - u^{k-1}}{2} (1 \leq k \leq N), \delta_t u^0 = 0.$$

For any function u defined on the interval $[0, t_1]$, we consider the Hermite quadratic interpolation polynomial $H_{2,0}(t)$. On the interval $[t_{k-1}, t_{k+1}]$ ($1 \leq k \leq N-1$), the Newton quadratic interpolation

polynomial $N_{2,k}(t)$ will be used. With the Caputo-fractional derivatives (1.1) on the time grid points of the form $\{t_{\frac{1}{2}}, t_{\frac{3}{2}}, \dots, t_{N-\frac{1}{2}}\}$ [14, 15], we can obtain

$$\begin{aligned} \sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} u(t_{n-\frac{1}{2}}) &\approx \sum_{i=1}^{M_1} \frac{p_i}{\Gamma(1-\alpha_i)} \left[\int_{t_0}^{t_{\frac{1}{2}}} H'_{2,0}(t)(t_{n-\frac{1}{2}}-t)^{-\alpha_i} dt + \sum_{k=1}^{n-1} \int_{t_{k-\frac{1}{2}}}^{t_{k+\frac{1}{2}}} N'_{2,k}(t)(t_{n-\frac{1}{2}}-t)^{-\alpha_i} dt \right] \\ &= \sum_{i=1}^{M_1} \frac{p_i}{\Gamma(1-\alpha_i)} \left\{ (t_{n-\frac{1}{2}}-t_0)^{1-\alpha_i} u_t^0 - (t_{n-\frac{1}{2}}-t_{\frac{1}{2}})^{1-\alpha_i} \delta_t u^{\frac{1}{2}} \right. \\ &\quad + \int_{t_0}^{t_{\frac{1}{2}}} H''_{2,0}(t)(t_{n-\frac{1}{2}}-t)^{1-\alpha_i} dt + \sum_{k=1}^{n-1} \left[(t_{n-\frac{1}{2}}-t_{k-\frac{1}{2}})^{1-\alpha_i} \delta_t u^{k-\frac{1}{2}} \right. \\ &\quad \left. \left. - (t_{n-\frac{1}{2}}-t_{k+\frac{1}{2}})^{1-\alpha_i} \delta_t u^{k+\frac{1}{2}} + \int_{t_{k-\frac{1}{2}}}^{t_{k+\frac{1}{2}}} N''_{2,k}(t)(t_{n-\frac{1}{2}}-t)^{1-\alpha_i} dt \right] \right\} \\ &= \sum_{i=1}^{M_1} \frac{p_i}{\Gamma(1-\alpha_i)} \left[a_{n-1}^{(n,\alpha_i)} \delta_t u^{n-\frac{1}{2}} + \sum_{k=1}^{n-1} (a_{k-1}^{(n,\alpha_i)} - a_k^{(n,\alpha_i)}) \delta_t u^{k-\frac{1}{2}} + (A_0^{(n,\alpha_i)} - a_0^{(n,\alpha_i)}) u_t^0 \right], \end{aligned} \tag{2.1}$$

where

$$A_0^{(n,\alpha_i)} = (t_{n-\frac{1}{2}} - t_0)^{1-\alpha_i} = t_{n-\frac{1}{2}}^{1-\alpha_i}, \tag{2.2}$$

$$a_{n-1}^{(n,\alpha_i)} = \frac{2}{\tau} \int_{t_0}^{t_{\frac{1}{2}}} (t_{n-\frac{1}{2}} - t)^{1-\alpha_i} dt = \frac{2\tau^{1-\alpha_i}}{2-\alpha_i} \left[\left(n - \frac{1}{2}\right)^{2-\alpha_i} - (n-1)^{2-\alpha_i} \right], \tag{2.3}$$

$$a_{n-k-1}^{(n,\alpha_i)} = \frac{1}{\tau} \int_{t_{k-\frac{1}{2}}}^{t_{k+\frac{1}{2}}} (t_{n-\frac{1}{2}} - t)^{1-\alpha_i} dt = \frac{\tau^{1-\alpha_i}}{2-\alpha_i} \left[(n-k)^{2-\alpha_i} - (n-1-k)^{2-\alpha_i} \right], 1 \leq k \leq n-1. \tag{2.4}$$

The H2N2 interpolation for the Caputo derivative $\mathcal{D}_t^{\beta_j} u(t_{n-\frac{1}{2}})$ can be written as follows:

$$\begin{aligned} \sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} u(t_{n-\frac{1}{2}}) &\approx \sum_{j=1}^{M_2} \frac{q_j}{\Gamma(2-\beta_j)} \left[\int_{t_0}^{t_{\frac{1}{2}}} H''_{2,0}(t)(t_{n-\frac{1}{2}}-t)^{1-\beta_j} dt + \sum_{k=1}^{n-1} \int_{t_{k-\frac{1}{2}}}^{t_{k+\frac{1}{2}}} N''_{2,k}(t)(t_{n-\frac{1}{2}}-t)^{1-\beta_j} dt \right] \\ &= \sum_{j=1}^{M_2} \frac{q_j}{\Gamma(2-\beta_j)} \left[b_0^{(n,\beta_j)} \delta_t u^{n-\frac{1}{2}} - \sum_{k=1}^{n-1} (b_{n-1-k}^{(n,\beta_j)} - b_{n-k}^{(n,\beta_j)}) \delta_t u^{k-\frac{1}{2}} - b_{n-1}^{(n,\beta_j)} u_t^0 \right]. \end{aligned} \tag{2.5}$$

Here

$$b_{n-1}^{(n,\beta_j)} = \frac{2}{\tau} \int_{t_0}^{t_{\frac{1}{2}}} (t_{n-\frac{1}{2}} - t)^{1-\beta_j} dt = \frac{2\tau^{1-\beta_j}}{2-\beta_j} \left[\left(n - \frac{1}{2}\right)^{2-\beta_j} - (n-1)^{2-\beta_j} \right], \tag{2.6}$$

$$b_{n-k-1}^{(n,\beta_j)} = \frac{1}{\tau} \int_{t_{k-\frac{1}{2}}}^{t_{k+\frac{1}{2}}} (t_{n-\frac{1}{2}} - t)^{1-\beta_j} dt = \frac{\tau^{1-\beta_j}}{2-\beta_j} \left[(n-k)^{2-\beta_j} - (n-1-k)^{2-\beta_j} \right], 1 \leq k \leq n-1. \tag{2.7}$$

Lemma 2.1. For $A_0^{(n,\alpha_i)}, a_k^{(n,\alpha_i)}, b_k^{(n,\beta_j)}, 0 \leq k \leq n-1$ defined by (2.2)–(2.4), (2.6) and (2.7), we gain

$$\begin{aligned} 0 &< a_0^{(n,\alpha_i)} < a_1^{(n,\alpha_i)} < \dots < a_{n-1}^{(n,\alpha_i)} < A_0^{(n,\alpha_i)}, \\ 0 &< t_{n-\frac{1}{2}}^{1-\beta_j} < b_{n-1}^{(n,\beta_j)} < b_{n-2}^{(n,\beta_j)} < \dots < b_0^{(n,\beta_j)} = \frac{\tau^{1-\beta_j}}{2-\beta_j}. \end{aligned} \tag{2.8}$$

Lemma 2.2. [14, 15] Suppose $u(t) \in C^3[t_0, t_n]$, and denote $R_{\alpha_i}^{n-\frac{1}{2}} = \sum_{i=1}^{M_1} p_i \partial_t^{\alpha_i} u(t_{n-\frac{1}{2}}) - \sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} u(t_{n-\frac{1}{2}})$ then we can get

$$\|R_{\alpha_i}^{n-\frac{1}{2}}\| \leq C_0 \max_{t_0 \leq t \leq t_{n-\frac{1}{2}}} |u'''(t)| \tau^{3-\alpha_i}.$$

If $u(t) \in C^3[t_0, t_n]$, denote $R_{\beta_j}^{n-\frac{1}{2}} = \sum_{j=1}^{M_2} q_j \partial_t^{\beta_j} u(t_{n-\frac{1}{2}}) - \sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} u(t_{n-\frac{1}{2}})$, then we have

$$\|R_{\beta_j}^{n-\frac{1}{2}}\| \leq C_1 \max_{t_0 \leq t \leq t_{n-\frac{1}{2}}} |u''(t)| \tau^{3-\beta_j}.$$

The basic mechanism in this algorithm is the construction of two shape-regular triangulations of Ω , which we denote by \mathcal{T}_H and \mathcal{T}_h , with different mesh sizes H and h ($h \ll H$). This can be accomplished by successively refining the triangulation \mathcal{T}_H to obtain the fine mesh \mathcal{T}_h . The corresponding finite element spaces are V_H and V_h , which will be called coarse and fine space, respectively. We note that by such a construction we have $V_H \subset V_h$. Let V_h be the two-dimensional subspace of $H_0^1(\Omega)$, which consists of continuous piecewise polynomials of degree r ($r \geq 1$) on \mathcal{T}_h and $V_h^0 = \{v \in V_h, v|_{\partial\Omega} = 0\}$, $V_H^0 = \{v \in V_H, v|_{\partial\Omega} = 0\}$. Note that the corresponding weak formulation of (1.1) is to find $u(\mathbf{x}, t) : (0, T] \rightarrow H_0^1(\Omega)$, then

$$\begin{cases} (\delta_t u(\mathbf{x}, t), v) + \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} u(\mathbf{x}, t), v) + \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} u(\mathbf{x}, t), v) + (\nabla u(\mathbf{x}, t), \nabla v) = (g(u), v), \\ u(\mathbf{x}, 0) = u_0(\mathbf{x}), u_t(\mathbf{x}, 0) = u_t^0(\mathbf{x}), \mathbf{x} \in \Omega. \end{cases} \quad (2.9)$$

First, we derive the two-grid fully discrete scheme for problem (1.1). The process is divided into two steps. Step 1: On the coarse grid \mathcal{T}_H , for any $v_H \in V_H^0$, find $u_H^n \in V_H^n$ for the following nonlinear system, such that

$$\begin{cases} (\delta_t u_H^{n-\frac{1}{2}}, v_H) + \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} u_H^{n-\frac{1}{2}}, v_H) + \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} u_H^{n-\frac{1}{2}}, v_H) + (\nabla u_H^{n-\frac{1}{2}}, \nabla v_H) = (g(u_H^{n-\frac{1}{2}}), v_H), \\ u_H^0 = R_H u_0(\mathbf{x}), u_{Ht}^0 = R_H u_t^0(\mathbf{x}), \mathbf{x} \in \Omega. \end{cases} \quad (2.10)$$

Step 2: On the fine grid \mathcal{T}_h , find $u_h^n \in V_h^n$ for the following linear system, then

$$\begin{cases} (\delta_t u_h^{n-\frac{1}{2}}, v_h) + \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} u_h^{n-\frac{1}{2}}, v_h) + \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} u_h^{n-\frac{1}{2}}, v_h) + (\nabla u_h^{n-\frac{1}{2}}, \nabla v_h), \\ = (g(u_H^{n-\frac{1}{2}}) + g'(u_H^{n-\frac{1}{2}})(u_h^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}), v_h), \forall v_h \in V_h^0, \\ u_h^0 = R_h u_0(\mathbf{x}), u_{ht}^0 = R_h u_t^0(\mathbf{x}), \mathbf{x} \in \Omega. \end{cases} \quad (2.11)$$

Next, the two-grid fully discrete scheme based on the Newton iteration idea will be analyzed. In the two-grid algorithm, we solve the nonlinear fractional equation on the coarse grid \mathcal{T}_H to produce a rough approximation, and then use the rough approximation as the initial guess to solve the linearized equation on the fine grid \mathcal{T}_h .

3. Unconditional stability and error estimate

Below, the stability and convergence of the two-grid fully discrete scheme will be derived. For the demand of analysis, we introduce some lemmas as follows:

Lemma 3.1. [21] Let $R_{\tilde{h}} : H_0^1(\Omega) \rightarrow V_{\tilde{h}}^0$, be the Rize-projection operator satisfying

$$(\nabla(u - R_{\tilde{h}}u), \nabla v) = 0, \quad \forall v \in V_{\tilde{h}}^0. \tag{3.1}$$

For $u \in H_0^1(\Omega) \cap H^{r+1}(\Omega)$, it holds that

$$\|u - R_{\tilde{h}}u\| + \tilde{h} \|u - R_{\tilde{h}}u\|_1 \leq C\tilde{h}^{r+1} \|u\|_{r+1}, \tag{3.2}$$

where \tilde{h} is coarse grid step length H or fine grid size h .

Lemma 3.2. [22] If A_m, B_s are nonnegative real sequences and the sequence Z_m satisfies

$$Z_m \leq A_m + \sum_{s=0}^{m-1} B_s Z_s, \quad m \geq 0,$$

where $A_0 \geq 0$, then the sequence Z_m satisfies

$$Z_m \leq A_m \exp\left(\sum_{s=0}^{m-1} B_s\right), \quad m \geq 0.$$

3.1. Stability

In this subsection, we shall present the main algorithm of the paper. The energy method is applied to establish the stability of the two-grid fully discrete scheme. First, we derive the stability of the coarse grid system.

Theorem 3.1. For $u_H^n \in V_H^0$, the two-grid finite element analysis system (2.10) can obtain the following inequality

$$\|u_H^n\|^2 \leq C\left(\|\nabla u_H^0\|^2 + \sum_{j=1}^{M_2} \frac{t_n^{2-\beta_j} q_j}{\Gamma(3-\beta_j)} \|u_{H_t}^0\|^2\right).$$

Proof. Suppose $\{u_H^{n-\frac{1}{2}}\}$ is the solution of a semi discrete scheme (2.10), the $v_H = \delta_t u_H^{n-\frac{1}{2}} \in V_H^0$, then the first item of (2.10) holds that

$$(\delta_t u_H^{n-\frac{1}{2}}, \delta_t u_H^{n-\frac{1}{2}}) = \|\delta_t u_H^{n-\frac{1}{2}}\|^2. \tag{3.3}$$

Using Eq (2.1), the second item of (2.10) follows that

$$\begin{aligned} & \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} u_H^{n-\frac{1}{2}}, \delta_t u_H^{n-\frac{1}{2}}) \\ &= \sum_{i=1}^{M_1} \frac{p_i}{\Gamma(2-\alpha_i)} \left(a_0^{(n,\alpha_i)} \delta_t u_H^{n-\frac{1}{2}} - \sum_{k=1}^{n-1} (a_{n-k-1}^{(n,\alpha_i)} - a_{n-k}^{(n,\alpha_i)}) \delta_t u_H^{k-\frac{1}{2}} + (A_0^{(n,\alpha_i)} - a_{n-1}^{(n,\alpha_i)}) u_{H_t}^0, \delta_t u_H^{n-\frac{1}{2}} \right) \\ &= \sum_{i=1}^{M_1} \frac{p_i \tau^{1-\alpha_i}}{\Gamma(3-\alpha_i)} (\delta_t u_H^{n-\frac{1}{2}}, \delta_t u_H^{n-\frac{1}{2}}) - \sum_{i=1}^{M_1} \frac{p_i}{\Gamma(2-\alpha_i)} \sum_{k=1}^{n-1} (a_{n-k-1}^{(n,\alpha_i)} - a_{n-k}^{(n,\alpha_i)}) (\delta_t u_H^{k-\frac{1}{2}}, \delta_t u_H^{n-\frac{1}{2}}) \\ & \quad + \sum_{i=1}^{M_1} \frac{p_i}{\Gamma(2-\alpha_i)} (A_0^{(n,\alpha_i)} - a_{n-1}^{(n,\alpha_i)}) (u_{H_t}^0, \delta_t u_H^{n-\frac{1}{2}}). \end{aligned} \tag{3.4}$$

By substituting (2.5) into (2.10), we can organize and obtain the third item of (2.10)

$$\begin{aligned}
 & \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} u_H^{n-\frac{1}{2}}, \delta_t u_H^{n-\frac{1}{2}}) \\
 &= \sum_{j=1}^{M_2} \frac{q_j}{\Gamma(2-\beta_j)} (b_0^{(n,\beta_j)} \delta_t u_H^{n-\frac{1}{2}} - \sum_{k=1}^{n-1} (b_{n-1-k}^{(n,\beta_j)} - b_{n-k}^{(n,\beta_j)}) \delta_t u_H^{k-\frac{1}{2}} - b_{n-1}^{(n,\beta_j)} u_{Ht}^0, \delta_t u_H^{n-\frac{1}{2}}) \\
 &= \sum_{j=1}^{M_2} \frac{q_j \tau^{1-\beta_j}}{\Gamma(3-\beta_j)} (\delta_t u_H^{n-\frac{1}{2}}, \delta_t u_H^{n-\frac{1}{2}}) - \sum_{j=1}^{M_2} \frac{q_j}{\Gamma(2-\beta_j)} \sum_{k=1}^{n-1} (b_{n-1-k}^{(n,\beta_j)} - b_{n-k}^{(n,\beta_j)}) (\delta_t u_H^{k-\frac{1}{2}}, \delta_t u_H^{n-\frac{1}{2}}) \\
 &\quad - \sum_{j=1}^{M_2} \frac{q_j b_{n-1}^{(n,\beta_j)}}{\Gamma(2-\beta_j)} (u_{Ht}^0, \delta_t u_H^{n-\frac{1}{2}}).
 \end{aligned} \tag{3.5}$$

Substituting (3.3)–(3.5) into (2.10), and multiplying 2τ on both sides of this equality, we obtain

$$\begin{aligned}
 & 2\tau \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2 + \left\| \nabla u_H^n \right\|^2 - \left\| \nabla u_H^{n-1} \right\|^2 + \sum_{i=1}^{M_1} \frac{2\tau^{2-\alpha_i} p_i}{\Gamma(3-\alpha_i)} \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2 + \sum_{j=1}^{M_2} \frac{2\tau^{2-\beta_j} q_j}{\Gamma(3-\beta_j)} \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2 \\
 & \leq \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} \sum_{k=1}^{n-1} (a_{n-k-1}^{(n,\alpha_i)} - a_{n-k}^{(n,\alpha_i)}) (\left\| \delta_t u_H^{k-\frac{1}{2}} \right\|^2 + \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2) + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{n-1}^{(n,\alpha_i)} - A_0^{(n,\alpha_i)}) \\
 & \quad (\left\| u_{Ht}^0 \right\|^2 + \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2) + \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2-\beta_j)} \sum_{k=1}^{n-1} (b_{n-1-k}^{(n,\beta_j)} - b_{n-k}^{(n,\beta_j)}) (\left\| \delta_t u_H^{k-\frac{1}{2}} \right\|^2 + \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2) \\
 & \quad + \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2-\beta_j)} b_{n-1}^{(n,\beta_j)} (\left\| u_{Ht}^0 \right\|^2 + \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2) + 2\tau (g(u_H^{n-\frac{1}{2}}), \delta_t u_H^{n-\frac{1}{2}}).
 \end{aligned} \tag{3.6}$$

Denoting $E^0 = \left\| \nabla u_H^0 \right\|^2$ and $E^n = \left\| \nabla u_H^n \right\|^2 + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} \sum_{k=1}^n a_{n-k}^{(n,\alpha_i)} \left\| \delta_t u_H^{k-\frac{1}{2}} \right\|^2 + \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2-\beta_j)} \sum_{k=1}^n b_{n-k}^{(n,\beta_j)} \left\| \delta_t u_H^{k-\frac{1}{2}} \right\|^2$, we can calculate that

$$\begin{aligned}
 E^n & \leq E^{n-1} - 2\tau \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2 + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{n-1}^{(n,\alpha_i)} - A_0^{(n,\alpha_i)}) \left\| u_{Ht}^0 \right\|^2 - \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} A_0^{(n,\alpha_i)} \left\| \delta_t u_H^{n-\frac{1}{2}} \right\|^2 \\
 & \quad + \sum_{j=1}^{M_2} \frac{\tau p_j}{\Gamma(2-\beta_j)} b_{n-1}^{(n,\beta_j)} \left\| u_{Ht}^0 \right\|^2 + 2\tau \left| (g(u_H^{n-\frac{1}{2}}), \delta_t u_H^{n-\frac{1}{2}}) \right|,
 \end{aligned}$$

which, by induction, gives

$$\begin{aligned}
 E^n & \leq E^0 - 2\tau \sum_{l=1}^n \left\| \delta_t u_H^{l-\frac{1}{2}} \right\|^2 + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} \sum_{l=1}^n (a_{l-1}^{(l,\alpha_i)} - A_0^{(l,\alpha_i)}) \left\| u_{Ht}^0 \right\|^2 - \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} \\
 & \quad \sum_{l=1}^n A_0^{(l,\alpha_i)} \left\| \delta_t u_H^{l-\frac{1}{2}} \right\|^2 + \sum_{j=1}^{M_2} \frac{\tau p_j}{\Gamma(2-\beta_j)} \sum_{l=1}^n b_{l-1}^{(l,\beta_j)} \left\| u_{Ht}^0 \right\|^2 + 2\tau \sum_{l=1}^n \left| (g(u_H^{l-\frac{1}{2}}), \delta_t u_H^{l-\frac{1}{2}}) \right|.
 \end{aligned} \tag{3.7}$$

Using the Cauchy-Schwarz inequality, the Young's inequality and the nonlinear property of $g(\cdot)$, it is shown that

$$\begin{aligned}
 2\tau \sum_{l=1}^n \left| \left(g(u_H^{l-\frac{1}{2}}), \delta_t u_H^{l-\frac{1}{2}} \right) \right| &\leq \tau \sum_{l=1}^n \left\{ \frac{1}{\Xi} \left\| g(u_H^{l-\frac{1}{2}}) \right\|^2 + \Xi \left\| \delta_t u_H^{l-\frac{1}{2}} \right\|^2 \right\} \\
 &\leq \sum_{l=1}^n \frac{\tau}{\Xi} C_0^2 \left\| u_H^{l-\frac{1}{2}} \right\|^2 + \sum_{l=1}^n \tau \Xi \left\| \delta_t u_H^{l-\frac{1}{2}} \right\|^2,
 \end{aligned}
 \tag{3.8}$$

where

$$\Xi = 2 + \sum_{i=1}^{M_1} \frac{P_i}{\Gamma(2 - \alpha_i)} (a_{n-l}^{(n,\alpha_i)} + A_0^{(n,\alpha_i)}) + \sum_{j=1}^{M_2} \frac{q_j}{\Gamma(2 - \beta_j)} b_{n-l}^{(n,\beta_j)}.$$

Combing with the Lemma 2.1, above analysis uses

$$0 < a_{l-1}^{(l,\alpha_i)} < A_0^{(l,\alpha_i)}, \tag{3.9}$$

$$\sum_{l=1}^n \frac{1}{\Xi} \leq \frac{1}{\sum_{j=1}^{M_2} \frac{q_j}{\Gamma(2-\beta_j)} \sum_{l=1}^n b_{n-l}^{(n,\beta_j)}} \leq \sum_{j=1}^{M_2} \frac{\Gamma(2 - \beta_j) t_n^{\beta_j-1}}{q_j t_n^{n-\frac{1}{2}}}, \tag{3.10}$$

$$\begin{aligned}
 \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2 - \beta_j)} \sum_{l=1}^n b_{l-1}^{(l,\beta_j)} \left\| u_{Hl}^0 \right\|^2 &\leq \sum_{j=1}^{M_2} \frac{\tau^{2-\beta_j} q_j}{\Gamma(3 - \beta_j)} \left\| u_{Hl}^0 \right\|^2 \sum_{l=1}^n \left[(n - l + 1)^{2-\beta_j} - (n - l)^{2-\beta_j} \right] \\
 &\leq \sum_{j=1}^{M_2} \frac{t_n^{2-\beta_j} q_j}{\Gamma(3 - \beta_j)} \left\| u_{Hl}^0 \right\|^2.
 \end{aligned}
 \tag{3.11}$$

Substituting the estimates (3.8)–(3.10) into (3.5), and using the Poincaré inequality, we have

$$\left\| u_H^n \right\|^2 \leq \left\| \nabla u_H^0 \right\|^2 + \sum_{j=1}^{M_2} \frac{t_n^{2-\beta_j} q_j}{\Gamma(3 - \beta_j)} \left\| u_{Hl}^0 \right\|^2 + \sum_{j=1}^{M_2} \frac{\Gamma(2 - \beta_j)}{q_j} t_n^{\beta_j-1} \sum_{l=1}^n C_0^2 \tau \left\| u_H^{l-\frac{1}{2}} \right\|^2.$$

By the Lemma 3.2, we have the desired result.

Theorem 3.2. For the fine grid system (2.11), the following stable inequality for $u_h^n \in V_h^0$

$$\left\| \nabla u_h^n \right\|^2 \leq C \left(\left\| \nabla u_h^0 \right\|^2 + \sum_{j=1}^{M_2} \frac{t_n^{2-\beta_j} q_j}{\Gamma(3 - \beta_j)} \left\| u_{ht}^0 \right\|^2 + \sum_{j=1}^{M_2} \frac{\Gamma(2 - \beta_j)}{q_j} T^{\beta_j} C_2^2 \max_{0 \leq t \leq T} \left\| u_H(x, t) \right\|^2 \right).$$

Proof. Considering (2.11), we can deduce that

$$\begin{aligned}
 &\sum_{i=1}^{M_1} \frac{P_i}{\Gamma(2 - \alpha_i)} \left(\frac{\tau^{1-\alpha_i}}{2 - \alpha_i} \delta_t u_h^{n-\frac{1}{2}} - \sum_{k=1}^{n-1} (a_{n-k-1}^{(n,\alpha_i)} - a_{n-k}^{(n,\alpha_i)}) \delta_t u_h^{k-\frac{1}{2}} - (a_{n-1}^{(n,\alpha_i)} - A_0^{(n,\alpha_i)}) u_{ht}^0, v_h \right) \\
 &+ \sum_{i=1}^{M_1} \frac{q_j}{\Gamma(2 - \beta_j)} \left(\frac{\tau^{1-\beta_j}}{2 - \beta_j} \delta_t u_h^{n-\frac{1}{2}} - \sum_{k=1}^{n-1} (b_{n-1-k}^{(n,\beta_j)} - b_{n-k}^{(n,\beta_j)}) \delta_t u_h^{k-\frac{1}{2}} - b_{n-1}^{(n,\beta_j)} u_{ht}^0, v_h \right) \\
 &= \left(g(u_H^{n-\frac{1}{2}}) + g'(u_H^{n-\frac{1}{2}})(u_h^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}), v_h \right) \\
 &\leq \left(C(u_H^{n-\frac{1}{2}}) + C_1(u_h^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}), v_h \right) \\
 &= \left(C_2(u_H^{n-\frac{1}{2}}) + C_1(u_h^{n-\frac{1}{2}}), v_h \right),
 \end{aligned}
 \tag{3.12}$$

where C, C_1, C_2 are positive constants.

Let $v_h = \delta_t u_h^{n-\frac{1}{2}}$, $F^n = \|\nabla u_h^n\|^2 + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} \sum_{k=1}^n a_{n-k}^{(n,\alpha_i)} \|\delta_t u_h^{k-\frac{1}{2}}\|^2 + \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2-\beta_j)} \sum_{k=1}^n b_{n-k}^{(n,\beta_j)} \|\delta_t u_h^{k-\frac{1}{2}}\|^2$.
Using the recurrence relation, we have

$$\begin{aligned} F^n &\leq F^{n-1} - 2\tau \|\delta_t u_h^{n-\frac{1}{2}}\|^2 + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{n-1}^{(n,\alpha_i)} - A_0^{(n,\alpha_i)}) \|u_{ht}^0\|^2 - \sum_{j=1}^{M_2} \frac{\tau p_j}{\Gamma(2-\alpha_j)} A_0^{(n,\alpha_j)} \|\delta_t u_h^{n-\frac{1}{2}}\|^2 \\ &\quad + \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2-\beta_j)} b_{n-1}^{(n,\beta_j)} \|u_{ht}^0\|^2 + 2\tau \left(C_2(u_H^{n-\frac{1}{2}}) + C_1(u_h^{n-\frac{1}{2}}), \delta_t u_h^{n-\frac{1}{2}} \right) \\ &\leq F^0 - 2\tau \sum_{l=1}^n \|\delta_t u_h^{l-\frac{1}{2}}\|^2 + \sum_{i=1}^{M_1} \sum_{l=1}^n \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{l-1}^{(l,\alpha_i)} - A_0^{(l,\alpha_i)}) \|u_{ht}^0\|^2 - \sum_{i=1}^{M_1} \sum_{l=1}^n \frac{\tau p_i A_0^{(l,\alpha_i)}}{\Gamma(2-\alpha_i)} \|\delta_t u_h^{l-\frac{1}{2}}\|^2 \\ &\quad + \sum_{j=1}^{M_2} \sum_{l=1}^n \frac{\tau b_{l-1}^{(l,\beta_j)} q_j}{\Gamma(2-\beta_j)} \|u_{ht}^0\|^2 + \sum_{l=1}^n \frac{4\tau C_2^2 \|u_H^{l-\frac{1}{2}}\|^2}{\Xi} + \sum_{l=1}^n \frac{4\tau C_1^2 \|u_h^{l-\frac{1}{2}}\|^2}{\Xi} + \sum_{l=1}^n \tau \Xi \|\delta_t u_h^{l-\frac{1}{2}}\|^2. \end{aligned} \quad (3.13)$$

It follows from (3.10) that

$$\sum_{l=1}^n \frac{4\tau C_2^2 \|u_H^{l-\frac{1}{2}}\|^2}{\Xi} \leq \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} C_2^2 \max_{0 \leq t \leq T} \|u_H(\mathbf{x}, t)\|^2, \quad (3.14)$$

then combining with (3.8) and substituting (3.14) into (3.13) yields

$$\|\nabla u_h^n\|^2 \leq \|\nabla u_h^0\|^2 + \sum_{j=1}^{M_2} \frac{t_n^{2-\beta_j} q_j}{\Gamma(3-\beta_j)} \|u_{ht}^0\|^2 + \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} C_2^2 \max_{0 \leq t \leq T} \|u_H(\mathbf{x}, t)\|^2 + \sum_{l=1}^n \frac{4\tau C_1^2 \|u_h^{l-\frac{1}{2}}\|^2}{\Xi}.$$

Using the Poincaré inequality and the Lemma 3.2, we obtain the proof.

3.2. Convergence

The convergence of the FEM system is considered by the energy method. First, we shall derive the convergence of the coarse grid system as follows:

Theorem 3.3. Let $u(t_n) \in H_0^1(\Omega) \cap H^{r+1}(\Omega)$, $u_t \in H^2(\Omega)$, $u_{tt} \in L^2(\Omega)$, $u_{ttt} \in L^2(\Omega)$, $u_H^n \in V_H^n$ and $u_H^0 = R_h u_0(x)$, then we have

$$\|u(t_n) - u_H^n\| \leq O(H^{r+1} + \tau^{3-\beta}). \quad (3.15)$$

Proof. We combine (1.1) with (2.10) to get

$$\begin{aligned} &(\delta_t u_h^{n-\frac{1}{2}}, v_h) + \left(\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} (u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}), v_h \right) + \left(\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} (u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}), v_h \right) + (\nabla(u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}), \nabla v_h) \\ &= (g(u^{n-\frac{1}{2}}) - g(u_H^{n-\frac{1}{2}}), v_h) - (R_{\alpha_i}^{n-\frac{1}{2}}, v_h) - (R_{\beta_j}^{n-\frac{1}{2}}, v_h). \end{aligned} \quad (3.16)$$

Denoting $\eta^n = u^n - R_h u^n$, $\theta^n = R_h u^n - u_H^n$ and choosing $v_h = \delta_t \theta^{n-\frac{1}{2}}$ in (3.15), we obtain

$$\begin{aligned} & (\delta_t \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) + \left(\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}} \right) + \left(\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}} \right) + (\nabla \theta^{n-\frac{1}{2}}, \nabla \delta_t \theta^{n-\frac{1}{2}}) \\ &= \left(g(u^{n-\frac{1}{2}}) - g(u_H^{n-\frac{1}{2}}), \delta_t \theta^{n-\frac{1}{2}} \right) - (R_{\alpha_i}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - (R_{\beta_j}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - (\delta_t \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) \\ & \quad - \left(\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}} \right) - \left(\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}} \right) - (\nabla \eta^{n-\frac{1}{2}}, \nabla \delta_t \theta^{n-\frac{1}{2}}). \end{aligned} \tag{3.17}$$

From the definitions of $\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} u_H^{n-\frac{1}{2}}$, $\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} u_H^{n-\frac{1}{2}}$ and $\delta_t u_H^{n-\frac{1}{2}}$, we can have that

$$\begin{aligned} \left(\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}} \right) &= \sum_{i=1}^{M_1} \frac{\tau^{1-\alpha_i} p_i}{\Gamma(3-\alpha_i)} \|\delta_t \theta^{n-\frac{1}{2}}\|^2 - \sum_{i=1}^{M_1} \sum_{k=1}^{n-1} \frac{p_i}{\Gamma(2-\alpha_i)} (a_{n-k-1}^{(n,\alpha_i)} - a_{n-k}^{(n,\alpha_i)}) \\ & \quad (\delta_t \theta^{k-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - \sum_{i=1}^{M_1} \frac{p_i}{\Gamma(2-\alpha_i)} (a_{n-1}^{(n,\alpha_i)} - A_0^{(n,\alpha_i)}) (\delta_t \theta^0, \delta_t \theta^{n-\frac{1}{2}}), \end{aligned} \tag{3.18}$$

$$\begin{aligned} \left(\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}} \right) &= \sum_{j=1}^{M_2} \frac{\tau^{1-\beta_j} q_j}{\Gamma(3-\beta_j)} \|\delta_t \theta^{n-\frac{1}{2}}\|^2 - \sum_{j=1}^{M_2} \sum_{k=1}^{n-1} \frac{q_j}{\Gamma(2-\beta_j)} (b_{n-1-k}^{(n,\beta_j)} - b_{n-k}^{(n,\beta_j)}) \\ & \quad (\delta_t \theta^{k-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - \sum_{j=1}^{M_2} \frac{q_j b_{n-1}^{(n,\beta_j)}}{\Gamma(2-\beta_j)} (\delta_t \theta^0, \delta_t \theta^{n-\frac{1}{2}}), \end{aligned} \tag{3.19}$$

and

$$(\nabla \theta^{n-\frac{1}{2}}, \nabla \delta_t \theta^{n-\frac{1}{2}}) = \frac{1}{2\tau} (\|\nabla \theta^n\|^2 - \|\nabla \theta^{n-1}\|^2). \tag{3.20}$$

Since $g(\cdot)$ is twice continuously differentiable, we have

$$\left(g(u^{n-\frac{1}{2}}) - g(u_H^{n-\frac{1}{2}}), \delta_t \theta^{n-\frac{1}{2}} \right) \leq \|u(t_n) - u_H^n\| \|\delta_t \theta^{n-\frac{1}{2}}\| \leq C (\|\eta^{n-\frac{1}{2}}\| + \|\theta^{n-\frac{1}{2}}\|) \|\delta_t \theta^{n-\frac{1}{2}}\|. \tag{3.21}$$

Substituting the above results into (3.16) and multiplying 2τ on both sides of the resulting identity, it holds that

$$\begin{aligned} & 2\tau \|\delta_t \theta^{n-\frac{1}{2}}\|^2 + \|\nabla \theta^n\|^2 - \|\nabla \theta^{n-1}\|^2 + \sum_{i=1}^{M_1} \frac{2\tau^{2-\alpha_i} p_i}{\Gamma(3-\alpha_i)} \|\delta_t \theta^{n-\frac{1}{2}}\|^2 + \sum_{j=1}^{M_2} \frac{2\tau^{2-\beta_j} q_j}{\Gamma(3-\beta_j)} \|\delta_t \theta^{n-\frac{1}{2}}\|^2 \\ &= \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} \sum_{k=1}^{n-1} (a_{n-k-1}^{(n,\alpha_i)} - a_{n-k}^{(n,\alpha_i)}) (\|\delta_t \theta^{k-\frac{1}{2}}\|^2 + \|\delta_t \theta^{n-\frac{1}{2}}\|^2) + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{n-1}^{(n,\alpha_i)} - A_0^{(n,\alpha_i)}) \\ & \quad (\|\delta_t \theta^0\|^2 + \|\delta_t \theta^{n-\frac{1}{2}}\|^2) + \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2-\beta_j)} \sum_{k=1}^{n-1} (b_{n-1-k}^{(n,\beta_j)} - b_{n-k}^{(n,\beta_j)}) (\|\delta_t \theta^{k-\frac{1}{2}}\|^2 + \|\delta_t \theta^{n-\frac{1}{2}}\|^2) \\ & \quad + \frac{\tau b_{n-1}^{(n,\beta_j)}}{\Gamma(2-\beta_j)} (\|\delta_t \theta^0\|^2 + \|\delta_t \theta^{n-\frac{1}{2}}\|^2) - 2\tau (R_{\alpha_i}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - 2\tau (R_{\beta_j}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - 2\tau (\delta_t \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) \\ & \quad - 2\tau \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - 2\tau \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) + 2\tau C (\|\eta^{n-\frac{1}{2}}\| + \|\theta^{n-\frac{1}{2}}\|) \|\delta_t \theta^{n-\frac{1}{2}}\|. \end{aligned} \tag{3.22}$$

Denoting $G^0 = \|\nabla\theta^0\|^2$ and $G^n = \|\nabla\theta^n\|^2 + \sum_{i=1}^{M_1} \sum_{k=1}^n \frac{\tau p_i a_{n-k}^{(n,\alpha_i)}}{\Gamma(2-\alpha_i)} \|\delta_t \theta^{k-\frac{1}{2}}\|^2 + \sum_{j=1}^{M_2} \sum_{k=1}^n \frac{\tau q_j b_{n-k}^{(n,\beta_j)}}{\Gamma(2-\beta_j)} \|\delta_t \theta^{k-\frac{1}{2}}\|^2$, we have that

$$\begin{aligned} G^n &\leq G^{n-1} - 2\tau \|\delta_t \theta^{n-\frac{1}{2}}\|^2 + \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{n-1}^{(n,\alpha_i)} - A_0^{(n,\alpha_i)}) \|\delta_t \theta^0\|^2 - \sum_{i=1}^{M_1} \frac{\tau p_i}{\Gamma(2-\alpha_i)} A_0^{(n,\alpha_i)} \|\delta_t \theta^{n-\frac{1}{2}}\|^2 \\ &+ \sum_{j=1}^{M_2} \frac{\tau q_j}{\Gamma(2-\beta_j)} b_{n-1}^{(n,\beta_j)} \|\delta_t \theta^0\|^2 - 2\tau (R_{\alpha_i}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - 2\tau (R_{\beta_j}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - 2\tau \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) \\ &- 2\tau \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - 2\tau (\delta_t \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) + 2\tau C \|\eta^{n-\frac{1}{2}}\| \|\delta_t \theta^{n-\frac{1}{2}}\| + 2\tau C \|\theta^{n-\frac{1}{2}}\| \|\delta_t \theta^{n-\frac{1}{2}}\|, \end{aligned}$$

which, by induction, gives

$$\begin{aligned} G^n &\leq G^0 - 2\tau \sum_{l=1}^n \|\delta_t \theta^{l-\frac{1}{2}}\|^2 + \sum_{i=1}^{M_1} \sum_{l=1}^n \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{l-1}^{(l,\alpha_i)} - A_0^{(l,\alpha_i)}) \|\delta_t \theta^0\|^2 \\ &- \sum_{i=1}^{M_1} \sum_{l=1}^n \frac{\tau p_i A_0^{(l,\alpha_i)}}{\Gamma(2-\alpha_i)} \|\delta_t \theta^{l-\frac{1}{2}}\|^2 + \sum_{j=1}^{M_2} \sum_{l=1}^n \frac{\tau q_j b_{l-1}^{(l,\beta_j)}}{\Gamma(2-\beta_j)} \|\delta_t \theta^0\|^2 - 2\tau \sum_{l=1}^n (R_{\alpha_i}^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \\ &- 2\tau \sum_{l=1}^n (R_{\beta_j}^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) - 2\tau \sum_{i=1}^{M_1} \sum_{l=1}^n p_i (\mathcal{D}_t^{\alpha_i} \eta^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) - 2\tau \sum_{j=1}^{M_2} \sum_{l=1}^n q_j (\mathcal{D}_t^{\beta_j} \eta^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \\ &- 2\tau \sum_{l=1}^n (\delta_t \eta^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) + 2\tau C \sum_{l=1}^n \|\eta^{l-\frac{1}{2}}\| \|\delta_t \theta^{l-\frac{1}{2}}\| + 2\tau C \sum_{l=1}^n \|\theta^{l-\frac{1}{2}}\| \|\delta_t \theta^{l-\frac{1}{2}}\|. \end{aligned} \tag{3.23}$$

By the Cauchy-Schwarz inequality, the Young's inequality and Lemma 2.2, we obtain

$$2\tau \sum_{l=1}^n (R_{\alpha_i}^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \leq \sum_{l=1}^n \frac{C\tau}{\Xi} \max_{0 \leq t \leq T} \|u_{ttt}(\mathbf{x}, t)\|^2 \max_{0 \leq i \leq M_1} \tau^{6-2\alpha_i} + \sum_{l=1}^n \frac{\tau \Xi}{7} \|\delta_t \theta^{l-\frac{1}{2}}\|^2, \tag{3.24}$$

$$2\tau \sum_{l=1}^n (R_{\beta_j}^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \leq \sum_{l=1}^n \frac{C\tau}{\Xi} \max_{0 \leq t \leq T} \|u_{ttt}(\mathbf{x}, t)\|^2 \max_{0 \leq j \leq M_2} \tau^{6-2\beta_j} + \sum_{l=1}^n \frac{\tau \Xi}{7} \|\delta_t \theta^{l-\frac{1}{2}}\|^2, \tag{3.25}$$

$$2\tau \sum_{l=1}^n (\delta_t \eta^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \leq \sum_{l=1}^n \frac{7\tau \|\delta_t \eta\|^2}{\Xi} + \sum_{l=1}^n \frac{\tau \Xi}{7} \|\delta_t \theta^{l-\frac{1}{2}}\|^2, \tag{3.26}$$

$$2\tau \sum_{i=1}^{M_1} \sum_{l=1}^n p_i (\mathcal{D}_t^{\alpha_i} \eta^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \leq \sum_{l=1}^n \frac{7\tau \|\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} \eta^{l-\frac{1}{2}}\|^2}{\Xi} + \sum_{l=1}^n \frac{\tau \Xi}{7} \|\delta_t \theta^{l-\frac{1}{2}}\|^2, \tag{3.27}$$

$$2\tau \sum_{j=1}^{M_2} \sum_{l=1}^n q_j (\mathcal{D}_t^{\beta_j} \eta^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \leq \sum_{l=1}^n \frac{7\tau \|\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} \eta^{l-\frac{1}{2}}\|^2}{\Xi} + \sum_{l=1}^n \frac{\tau \Xi}{7} \|\delta_t \theta^{l-\frac{1}{2}}\|^2, \tag{3.28}$$

$$2\tau C \sum_{l=1}^n \|\eta^{l-\frac{1}{2}}\| \|\delta_t \theta^{l-\frac{1}{2}}\| \leq \sum_{l=1}^n \frac{7\tau C^2}{\Xi} \|\eta^{l-\frac{1}{2}}\|^2 + \sum_{l=1}^n \frac{\tau \Xi}{7} \|\delta_t \theta^{l-\frac{1}{2}}\|^2, \tag{3.29}$$

$$2\tau C \sum_{l=1}^n \|\theta^{l-\frac{1}{2}}\| \|\delta_l \theta^{l-\frac{1}{2}}\| \leq \sum_{l=1}^n \frac{7\tau C^2}{\Xi} \|\theta^{l-\frac{1}{2}}\|^2 + \sum_{l=1}^n \frac{\tau \Xi}{7} \|\delta_l \theta^{l-\frac{1}{2}}\|^2. \tag{3.30}$$

The definition of Ξ is the same as (3.9). By Lemma 2.1, we have

$$0 < a_{n-l}^{(n,\alpha_i)} < A_0^{(n,\alpha_i)}. \tag{3.31}$$

Based on the above estimates and Lemma 3.1, we have

$$\sum_{l=1}^n \frac{\tau}{\Xi} \max_{1 \leq l \leq n} \|u_{ttt}(\mathbf{x}, t)\|^2 \max_{0 \leq i \leq M_1} \tau^{6-2\alpha_i} \leq \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} \max_{1 \leq l \leq n} \|u_{ttt}(\mathbf{x}, t)\|^2 \tau^{6-2\alpha}, \tag{3.32}$$

$$\sum_{l=1}^n \frac{\tau}{\Xi} \max_{1 \leq l \leq n} \|u_{ttt}(\mathbf{x}, t)\|^2 \max_{0 \leq j \leq M_2} \tau^{6-2\beta_j} \leq \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} \max_{1 \leq l \leq n} \|u_{ttt}(\mathbf{x}, t)\|^2 \tau^{6-2\beta}, \tag{3.33}$$

$$\sum_{l=1}^n \frac{\tau \|\delta_l \eta\|^2}{\Xi} \leq \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} H^{2r+2} \max_{1 \leq l \leq n} \|\delta_l u^{l-\frac{1}{2}}\|_{r+1}^2, \tag{3.34}$$

$$\begin{aligned} \sum_{l=1}^n \frac{7\tau \left\| \sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} \eta^{l-\frac{1}{2}} \right\|^2}{\Xi} &\leq \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} \max_{1 \leq l \leq n} \left\| \sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} \eta^{l-\frac{1}{2}} \right\|^2 \\ &\leq C \max_{1 \leq l \leq n} \left\| \sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} u^{l-\frac{1}{2}} \right\|_{r+1}^2 H^{2r+2} + C \max_{1 \leq l \leq n} \|u_{ttt}(\mathbf{x}, t)\|^2 \tau^{6-2\alpha}, \end{aligned} \tag{3.35}$$

$$\sum_{l=1}^n \frac{7\tau \left\| \sum_{q=0}^{M_2} q_j \mathcal{D}_t^{\beta_j} \eta^{l-\frac{1}{2}} \right\|^2}{\Xi} \leq C \max_{1 \leq l \leq n} \left\| \sum_{q=0}^{M_2} q_j \mathcal{D}_t^{\beta_j} u^{l-\frac{1}{2}} \right\|_{r+1}^2 H^{2r+2} + C \max_{1 \leq l \leq n} \|u_{ttt}(\mathbf{x}, t)\|^2 \tau^{6-2\beta}, \tag{3.36}$$

$$\sum_{l=1}^n \frac{7\tau}{\Xi} \|\eta^{l-\frac{1}{2}}\|^2 \leq \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} \|\eta^{l-\frac{1}{2}}\|^2 \leq \sum_{j=1}^{M_2} \frac{\Gamma(2-\beta_j)}{q_j} T^{\beta_j} \|u^{l-\frac{1}{2}}\|_{r+1}^2 H^{2r+2}. \tag{3.37}$$

Combining (3.22)–(3.37) can give

$$\|\nabla \theta^n\|^2 \leq C \left(\sum_{l=1}^n \frac{7\tau}{\Xi} \|\theta^{l-\frac{1}{2}}\|^2 + H^{2r+2} + \tau^{\min\{6-2\alpha, 6-2\beta\}} \right).$$

By the Poincaré inequality, Lemma 3.1 and the Lemma 3.2, it follows that

$$\|\theta^n\|^2 \leq C(H^{2r+2} + \tau^{6-2\beta}). \tag{3.38}$$

The proof is complete.

Theorem 3.4. Let $u(t_n) \in H_0^1(\Omega) \cap H^{r+1}(\Omega)$, $u_t \in H^2(\Omega)$, $u_{tt} \in L^2(\Omega)$, $u_{ttt} \in L^2(\Omega)$, $u_h^n \in V_h^0$, and $u_h^0 = R_h u_0(\mathbf{x})$, then we arrive at the following

$$\|u(t_n) - u_h^n\| \leq C(\tau^{3-\beta} + h^{r+1} + H^{2r+2}). \tag{3.39}$$

Proof. Subtracting (1.1) from (2.10), it yields

$$\begin{aligned} & (\delta_t(u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), v_h) + \left(\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} (u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), v_h \right) \\ & + \left(\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} (u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), v_h \right) + (\nabla(u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), \nabla v_h) \\ & = (g(u^{n-\frac{1}{2}}) - g(u_H^{n-\frac{1}{2}}) - g'(u_H^{n-\frac{1}{2}})(u_h^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}), v_h) - (R_{\alpha_i}^{n-\frac{1}{2}}, v_h) - (R_{\beta_j}^{n-\frac{1}{2}}, v_h). \end{aligned} \quad (3.40)$$

By the Taylor expansion, we can have the estimate that

$$g(u^{n-\frac{1}{2}}) = g(u_H^{n-\frac{1}{2}}) + g'(u_H^{n-\frac{1}{2}})(u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}}) + \frac{g''(u_H^{n-\frac{1}{2}})}{2}(u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}})^2. \quad (3.41)$$

then it follows from (3.40) that

$$\begin{aligned} & (\delta_t(u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), v_h) + \left(\sum_{i=1}^{M_1} p_i \mathcal{D}_t^{\alpha_i} (u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), v_h \right) \\ & + \left(\sum_{j=1}^{M_2} q_j \mathcal{D}_t^{\beta_j} (u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), v_h \right) + (\nabla(u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}), \nabla v_h) \\ & = (g'(u_H^{n-\frac{1}{2}})(u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}) + \frac{g''(u_H^{n-\frac{1}{2}})}{2}(u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}})^2, v_h) - (R_{\alpha_i}^{n-\frac{1}{2}}, v_h) - (R_{\beta_j}^{n-\frac{1}{2}}, v_h) \\ & \leq (C(u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}) + \frac{C_1}{2}(u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}})^2, v_h). \end{aligned} \quad (3.42)$$

Let $u^{n-\frac{1}{2}} - R_h u^{n-\frac{1}{2}} = \eta^{n-\frac{1}{2}}$, $R_h u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}} = \theta^{n-\frac{1}{2}}$, and $v_h = \delta_t \theta^{n-\frac{1}{2}}$. Hence, (3.42) becomes

$$\begin{aligned} & (\delta_t \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) + \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) + \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} \theta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) + (\nabla \theta^{n-\frac{1}{2}}, \nabla v_h) \\ & \leq -(\delta_t \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - \sum_{i=1}^{M_1} p_i (\mathcal{D}_t^{\alpha_i} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - \sum_{j=1}^{M_2} q_j (\mathcal{D}_t^{\beta_j} \eta^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) - (R_{\alpha_i}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) \\ & \quad - (R_{\beta_j}^{n-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) + (C(u^{n-\frac{1}{2}} - u_h^{n-\frac{1}{2}}) + \frac{C_3}{2}(u^{n-\frac{1}{2}} - u_H^{n-\frac{1}{2}})^2, \delta_t \theta^{n-\frac{1}{2}}). \end{aligned} \quad (3.43)$$

Defining $H^0 = \nabla \theta^0$, $H^n = \|\nabla \theta^n\|^2 + \sum_{i=1}^{M_1} \sum_{k=1}^n \frac{\tau p_i}{\Gamma(2-\alpha_i)} a_{n-k}^{(n,\alpha_i)} \|\delta_t \theta^{k-\frac{1}{2}}\|^2 + \sum_{j=1}^{M_2} \sum_{k=1}^n \frac{\tau q_j}{\Gamma(2-\beta_j)} b_{n-k}^{(n,\beta_j)} \|\delta_t \theta^{k-\frac{1}{2}}\|^2$, we have following results:

$$\begin{aligned} H^n & \leq H^0 - \sum_{l=1}^n 2\tau \|\delta_t \theta^{l-\frac{1}{2}}\|^2 + \sum_{i=1}^{M_1} \sum_{l=1}^n \frac{\tau p_i}{\Gamma(2-\alpha_i)} (a_{l-1}^{(l,\alpha_i)} - A_0^{(l,\alpha_i)}) \|\delta_t \theta^0\|^2 \\ & \quad - \sum_{i=1}^{M_1} \sum_{l=1}^n \frac{\tau p_i A_0^{(l,\alpha_i)}}{\Gamma(2-\alpha_i)} \|\delta_t \theta^{l-\frac{1}{2}}\|^2 + \sum_{j=1}^{M_2} \sum_{l=1}^n \frac{\tau q_j b_{l-1}^{(l,\beta_j)}}{\Gamma(2-\beta_j)} \|\delta_t \theta^0\|^2 - 2\tau \sum_{i=1}^{M_1} \sum_{l=1}^n p_i (\mathcal{D}_t^{\alpha_i} \eta^{l-\frac{1}{2}}, \delta_t \theta^{n-\frac{1}{2}}) \end{aligned}$$

$$\begin{aligned}
& - 2\tau \sum_{j=1}^{M_2} \sum_{l=1}^n q_j (\mathcal{D}_t^{\beta_j} \eta^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) - 2\tau \sum_{l=1}^n (R_{\alpha_i}^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) - 2\tau \sum_{l=1}^n (R_{\beta_j}^{l-\frac{1}{2}}, \delta_t \theta^{l-\frac{1}{2}}) \\
& + 2\tau \sum_{l=1}^n \left(C(u^{l-\frac{1}{2}} - u_h^{l-\frac{1}{2}}) + \frac{C_3}{2}(u^{l-\frac{1}{2}} - u_H^{l-\frac{1}{2}}), \delta_t \theta^{l-\frac{1}{2}} \right).
\end{aligned}$$

We use the similar process of proof to the estimate (3.37). By the Poincaré inequality, Lemma 3.1, Lemma 3.2 and the estimate of (3.37), it follows that

$$\begin{aligned}
\|\theta^n\|^2 & \leq C \left[\|\theta^0\|^2 + \tau^{6-2\beta} + \sum_{l=1}^n \left\| (u^{l-\frac{1}{2}} - u_H^{l-\frac{1}{2}})^2 \right\|^2 + h^{2r+2} \right] \\
& \leq C \left(\|\theta^0\|^2 + \tau^{6-2\beta} + H^{4r+4} + h^{2r+2} \right),
\end{aligned}$$

which leads to

$$\|u - u_h\| \leq C(\tau^{3-\beta} + H^{2r+2} + h^{r+1}). \quad (3.44)$$

This concludes the proof.

Remark 4.1. In the estimate (3.44), we observe that the TGM algorithm can achieve the convergence rate h^{r+1} as long as the mesh sizes satisfy $H = O(h^{\frac{1}{2}})$.

4. Numerical examples

In this section, we present a numerical example to demonstrate the theoretical analysis and illustrate the efficiency of the algorithm discussed in section three. To investigate the spatial and temporal convergence order, we use a bilinear finite element approximation and the computation is performed by using Matlab.

Example 4.1. The following equation has an exact solution $u(x, y, t) = t^{3+\alpha+\beta} \sin \pi x \sin \pi y$:

$$\begin{cases} \partial_t u(x, y, t) + \partial_t^\alpha u(x, y, t) + \partial_t^\beta u(x, y, t) - \Delta u(x, y, t) = -u^2 + g(x, y, t), & (x, y, t) \in \Omega \times (0, T], \\ u(x, y, 0) = u_0(x, y), \quad u_t(x, y, 0) = u_t^0(x, y), & x, y \in \Omega, \end{cases}$$

where Ω is the unit square $(0, 1) \times (0, 1)$, $T = 1$ and g is a known function. The domain Ω is divided into families \mathcal{T}_H and \mathcal{T}_h of quadrilaterals, and $V_H, V_h \subset H_0^1(\Omega)$ are linear spaces of piecewise continuous bilinear functions defined on \mathcal{T}_H and \mathcal{T}_h , respectively.

Let $H_x = H_y = H$, $h_x = h_y = h$ and $h = H^2$. Tables 1–3 show that the spatial convergence rates in the L^2 -norm of FEM and Algorithm 3.1 are both equivalent to two. The convergence results are in good agreement with the results $O(h^{r+1})$ of the theoretical analysis. Moreover, combining Tables 2–4, it can be seen that the TGM scheme will save more time than the general FEM scheme as M increases.

Table 1. Comparison of the spatial convergence order and elapsed CPU(s) time of TGM and FEM for Example 4.1 with different (α, β) , $\tau = 1000$.

H	h	$\alpha = 0.3, \beta = 1.7$			$\alpha = 0.5, \beta = 1.5$			$\alpha = 0.9, \beta = 1.3$		
		$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU
$\frac{1}{2}$	$\frac{1}{4}$	4.815e-2	*	14.39	4.780e-2	*	13.49	4.780e-2	*	13.92
$\frac{1}{3}$	$\frac{1}{9}$	8.811e-3	2.09	29.54	8.720e-3	2.10	29.28	8.631e-3	2.10	29.85
$\frac{1}{4}$	$\frac{1}{16}$	2.635e-3	2.10	67.18	2.614e-3	2.09	67.62	2.590e-3	2.09	69.60
$\frac{1}{5}$	$\frac{1}{25}$	1.058e-3	2.04	175.41	1.063e-3	2.02	173.26	1.058e-3	2.01	174.18
h	$\alpha = 0.3, \beta = 1.7$			$\alpha = 0.5, \beta = 1.5$			$\alpha = 0.9, \beta = 1.3$			
	$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU	
$\frac{1}{4}$	4.816e-2	*	23.66	4.780e-2	*	24.70	4.780e-2	*	25.74	
$\frac{1}{9}$	8.811e-3	2.09	41.52	8.720e-3	2.10	53.05	8.631e-3	2.10	53.44	
$\frac{1}{16}$	2.636e-3	2.10	77.83	2.614e-3	2.09	118.43	2.590e-3	2.09	120.49	
$\frac{1}{25}$	1.059e-3	2.04	326.60	1.063e-3	2.02	296.73	1.058e-3	2.01	298.67	

Table 2. L^2 -errors and temporal convergence rate for TGM and FEM with $\alpha=0.3, \beta = 1.7$, $\alpha=0.5, \beta = 1.5$ and $\alpha=0.7, \beta = 1.7$.

τ	$\alpha = 0.3, \beta = 1.7$			$\alpha = 0.5, \beta = 1.5$			$\alpha = 0.7, \beta = 1.7$		
	$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU
$\frac{1}{12}$	1.397e-2	*	30.34	6.519e-3	*	41.16	1.706e-2	*	30.31
$\frac{1}{14}$	1.145e-2	1.29	35.40	5.169e-3	1.51	51.90	1.396e-2	1.30	35.44
$\frac{1}{16}$	9.633e-3	1.29	41.20	4.231e-3	1.50	58.28	1.174e-2	1.30	40.65
$\frac{1}{18}$	8.274e-3	1.29	47.23	3.549e-3	1.49	62.54	1.008e-2	1.30	45.44
τ	$\alpha = 0.3, \beta = 1.7$			$\alpha = 0.5, \beta = 1.5$			$\alpha = 0.7, \beta = 1.7$		
	$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU
$\frac{1}{12}$	1.395e-2	*	132.68	6.501e-3	*	130.38	1.704e-2	*	133.69
$\frac{1}{14}$	1.143e-2	1.29	153.62	5.152e-3	1.51	152.17	1.395e-2	1.30	152.98
$\frac{1}{16}$	9.623e-3	1.29	186.10	4.215e-3	1.50	175.50	1.173e-2	1.30	175.26
$\frac{1}{18}$	8.265e-3	1.29	267.50	3.533e-3	1.50	198.20	1.007e-2	1.30	200.97

Example 4.2. The following equation has the exact solution $u(x, y, t) = (1 + t^{3+\alpha+\beta})\sin\pi x \sin\pi y$:

$$\begin{cases} \partial_t u(x, y, t) + \sum_{i=1}^2 \partial_t^{\alpha_i} u(x, y, t) + \sum_{j=1}^2 \partial_t^{\beta_j} u(x, y, t) - \Delta u(x, y, t) = -u^2 + g(x, y, t), & (x, y, t) \in \Omega \times (0, T], \\ u(x, y, 0) = u_0(x, y), \quad u_t(x, y, 0) = u_t^0(x, y), & x, y \in \Omega, \end{cases}$$

where Ω is the unit square $(0, 1) \times (0, 1)$, $T = 1$ and g is a known function.

Table 3. Comparison of the spatial convergence order and elapsed CPU(s) time of TGM and FEM for Example 4.2 with different (α_i, β_j) , $\tau = 1000$.

H	h	$\alpha_i = 0.2, 0.3, \beta_j = 1.6, 1.7$			$\alpha_i = 0.4, 0.5, \beta_j = 1.4, 1.5$			$\alpha_i = 0.8, 0.9, \beta_j = 1.2, 1.3$		
		$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU
$\frac{1}{2}$	$\frac{1}{4}$	4.417e-2	*	16.71	4.597e-2	*	17.13	4.598e-2	*	15.87
	$\frac{1}{3}$	8.036e-3	2.10	31.46	8.421e-3	2.10	34.52	8.426e-3	2.10	33.49
	$\frac{1}{4}$	2.413e-3	2.09	72.02	2.539e-3	2.09	70.44	2.543e-3	2.08	74.08
	$\frac{1}{5}$	9.730e-4	2.04	188.77	1.028e-3	2.02	181.33	1.034e-3	2.02	173.04
h		$\alpha_i = 0.2, 0.3, \beta_j = 1.6, 1.7$			$\alpha_i = 0.4, 0.5, \beta_j = 1.4, 1.5$			$\alpha_i = 0.8, 0.9, \beta_j = 1.2, 1.3$		
		$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU
$\frac{1}{4}$	$\frac{1}{4}$	4.417e-2	*	26.53	4.597e-2	*	25.27	4.326e-2	*	25.69
	$\frac{1}{9}$	8.036e-3	2.10	41.09	8.421e-3	2.10	42.90	7.878e-3	2.10	41.97
	$\frac{1}{16}$	2.413e-3	2.09	98.70	2.539e-3	2.08	97.41	2.3703e-3	2.09	98.14
	$\frac{1}{25}$	9.730e-4	2.04	367.81	1.028e-3	2.02	320.95	9.570e-4	2.03	367.51

Table 4. L^2 -errors and temporal convergence rate and elapsed CPU(s) time of TGM and FEM with different (α_i, β_j) , $m = 121$.

τ	$\alpha_i = 0.2, 0.3, \beta_j = 1.6, 1.7$			$\alpha_i = 0.4, 0.5, \beta_j = 1.4, 1.5$			$\alpha_i = 0.8, 0.9, \beta_j = 1.6, 1.7$		
	$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU	$\ u^n - u_h^n\ $	$Rate_h$	CPU
$\frac{1}{12}$	1.893e-2	*	36.93	8.676e-3	*	32.44	2.184e-2	*	30.83
$\frac{1}{14}$	1.544e-2	1.32	35.26	6.845e-3	1.54	35.76	1.777e-2	1.34	35.92
$\frac{1}{16}$	1.294e-2	1.32	42.71	5.577e-3	1.53	40.98	1.487e-2	1.34	40.89
$\frac{1}{18}$	1.107e-2	1.32	45.54	4.657e-3	1.53	46.42	1.271e-2	1.33	46.27
τ	$\alpha_i = 0.2, 0.3, \beta_j = 1.6, 1.7$			$\alpha_i = 0.4, 0.5, \beta_j = 1.4, 1.5$			$\alpha_i = 0.8, 0.9, \beta_j = 1.6, 1.7$		
	$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU	$\ u^n - U^n\ $	$Rate_h$	CPU
$\frac{1}{12}$	6.474e-3	*	137.84	4.772e-3	*	135.30	1.373e-2	*	135.73
$\frac{1}{14}$	5.273e-3	1.33	153.89	3.767e-3	1.53	154.05	1.117e-2	1.33	155.12
$\frac{1}{16}$	4.417e-3	1.33	176.16	3.072e-3	1.53	178.36	9.351e-3	1.33	196.50
$\frac{1}{18}$	3.779e-3	1.32	198.87	2.568e-3	1.52	200.17	7.993e-3	1.33	325.33

5. Conclusions

In this paper, we proposed a new H2N2 formula to approximate the multi-term fractional derivative $\sum_{i=1}^{M_1} p_i \partial_t^{\alpha_i} u(x, t)$, $\sum_{j=1}^{M_2} q_j \partial_t^{\beta_j} u(x, t)$, $\alpha_i \in (0, 1)$, $\beta_j \in (1, 2)$. Based on the H2N2 approximation in time and the finite element method for the spatial discretization, we have presented a fully discrete TGM scheme for two-dimensional nonlinear time fractional multi-term mixed sub-diffusion and diffusion wave equations and proved they are of second-order convergence in space and can reach the optimal convergence order $3 - \beta$ in time, which is not related to α . In future work, we will consider the L2-1 $_{\sigma}$

formulation, which can reach second-order accuracy in time.

Use of AI tools declaration

The authors declare they have not used Artificial Intelligence (AI) tools in the creation of this article.

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Conflict of interest

There is no conflicts of interest between all authors.

References

1. S. Jiang, J. Zhang, Q. Zhang, Z. Zhang, Fast evaluation of the Caputo fractional derivative and its applications to fractional diffusion equations, *Commun. Comput. Phys.*, **21** (2017), 650–678. <https://doi.org/10.4208/cicp.OA-2016-0136>
2. Y. Yu, P. Perdikaris, G. E. Karniadaki, Fractional modeling of viscoelasticity in 3D cerebral arteries and aneurysms, *J. Comput. Phys.*, **323** (2016), 219–242. <https://doi.org/10.1016/j.jcp.2016.06.038>
3. Q. Li, Y. Chen, Y. Huang, Y. Wang, Two-grid methods for semilinear time fractional reaction diffusion equations by expanded mixed finite element method, *Commun. Comput. Phys.*, **157** (2020), 38–54. <https://doi.org/10.1016/j.apnum.2020.05.024>
4. J. Xu, A novel two-grid method for semilinear equations, *SIAM. J. Sci. Comput.*, **15** (1994), 231–237. <https://doi.org/10.1137/0915016>
5. J. Xu, Two-grid discretization techniques for linear and non-linear PDEs, *SIAM. J. Numer. Anal.*, **33** (1996), 1759–1777. <https://doi.org/10.1137/S0036142992232949>
6. L. Chen, Y. Chen, Two-grid method for nonlinear reaction-diffusion equations by mixed finite element methods, *J. Sci. Comput.*, **49** (2011), 383–401. <https://doi.org/10.1007/s10915-011-9469-3>
7. Y. Chen, Q. Gu, Q. Li, Y. Huang, A two-grid finite element approximation for nonlinear time fractional two-term mixed sub-diffusion and diffusion wave equations, *J. Comput. Math.*, **40** (2022), 938–956. <https://doi.org/10.4208/jcm.2104-m2021-0332>
8. X. Li, H. Rui, A two-grid block-centered finite difference method for the nonlinear time-fractional parabolic equation, *J. Sci. Comput.*, **72** (2017), 863–891. <https://doi.org/10.1007/s10915-017-0380-4>
9. Y. Tang, A characteristic mixed finite element method for bilinear convection-diffusion optimal control problems, *J. Nonlinear Funct. Anal.*, **2022** (2022), 39. <https://doi.org/10.23952/jnfa.2022.39>

10. R. Schumer, D. A. Benson, M. M. Meerschaert, B. Baeumer, Fractal mobile/immobile solute transport, *Water Resour. Res.*, **39** (2003), 1296–1308. <https://doi.org/10.1029/2003WR002141>
11. Y. Zhang, Z. Sun, X. Zhao, Compact alternating direction implicit scheme for the two-dimensional fractional diffusion-wave equation, *SIAM. J. Numer. Anal.*, **50** (2012), 1535–1555. <https://doi.org/10.1137/110840959>
12. Y. Zhao, F. Wang, X. Hu, Anisotropic linear triangle finite element approximation for multi-term time-fractional mixed diffusion and diffusion-wave equations with variable coefficient on 2D bounded domain, *Comput. Math. Appl.*, **78** (2019), 1705–1719. <https://doi.org/10.1016/j.camwa.2018.11.028>
13. Z. Sun, C. Ji, R. Du, A new analytical technique of the L-type difference schemes for time fractional mixed sub-diffusion and diffusion-wave equations, *Appl. Math. Lett.*, **102** (2020), 106–115. <https://doi.org/10.1016/j.aml.2019.106115>
14. J. Shen, C. Li, Z. Sun, An H2N2 interpolation for Caputo derivative with order in (1, 2) and its application to time fractional hyperbolic equation in more than one space dimension, *J. Sci. Comput.*, **83** (2020), 38–67. <https://doi.org/10.1007/s10915-020-01219-8>
15. J. Shen, X. M. Gu, Two finite difference methods based on an H2N2 interpolation for two-dimensional time fractional mixed diffusion and diffusion-wave equations, *Discrete. Cont. Dyn. Syst. B*, **27** (2022), 1179–1207. <https://doi.org/10.3934/dcdsb.2021086>
16. Y. Liu, Y. Du, H. Li, S. He, W. Gao, Finite difference/finite element method for a nonlinear time-fractional fourth-order reaction-diffusion problem, *Comput. Math. Appl.*, **70** (2015), 573–591. <https://doi.org/10.1016/j.camwa.2015.05.015>
17. Y. Liu, Y. Du, H. Li, J. Wang, A two-grid finite element approximation for a nonlinear time-fractional Cable equation, *Nonlinear Dyn.*, **85** (2016), 2535–2548. <https://doi.org/10.1007/s11071-016-2843-9>
18. C. Bi, C. Wang, Y. Lin, Pointwise error estimates and two-grid algorithms of discontinuous Galerkin method for strongly nonlinear elliptic problems, *J. Sci. Comput.*, **67** (2016), 153–175. <https://doi.org/10.1007/s10915-015-0072-x>
19. B. Jin, B. Li, Z. Zhou, Numerical analysis of nonlinear subdiffusion equations, *SIAM. J. Numer. Anal.*, **56** (2018), 1–23. <https://doi.org/10.1137/16M1089320>
20. D. Li, C. Wu, Z. Zhang, Linearized Galerkin FEMs for nonlinear time fractional parabolic problems with non-smooth solutions in time direction, *J. Sci. Comput.*, **80** (2019), 403–419. <https://doi.org/10.1007/s10915-019-00943-0>
21. P. Ciarlet, *The finite element method for elliptic problems*, New York: North-Hollan, 1978. <https://doi.org/10.1115/1.3424474>
22. I. H. Sloan, V. Thomée, Time discretization of an integro-differential equation of parabolic type, *SIAM. J. Numer. Anal.*, **23** (1986), 1052–1061. <https://doi.org/10.1137/0723073>