Mathematics

## Research article

## Hamiltonian elliptic system involving nonlinearities with supercritical exponential growth

Yony Raúl Santaria Leuyacc*

Universidad Nacional Mayor de San Marcos, Lima, Perú

* Correspondence: Email: ysantarial@unmsm.edu.pe.


#### Abstract

In this paper, we deal with the existence of nontrivial solutions to the following class of strongly coupled Hamiltonian systems: $$
\left\{\begin{array}{rl} -\operatorname{div}(w(x) \nabla u)=g(x, v), & x \in B_{1}(0), \\ -\operatorname{div}(w(x) \nabla v)=f(x, u), & x \in B_{1}(0), \\ u=v=0 & x \in \partial B_{1}(0), \end{array}\right.
$$ where $w(x)=(\log 1 /|x|)^{\gamma}, 0 \leq \gamma<1$, and the nonlinearities $f$ and $g$ possess exponential growth ranges above the exponential critical hyperbola. Our approach is based on Trudinger-Moser type inequalities for weighted Sobolev spaces and variational methods.


Keywords: Hamiltonian system; Trudinger-Moser inequality; supercritical exponential growth; variational methods; linking theorem
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## 1. Introduction

In the literature, the existence of nontrivial solutions for strongly coupled Hamiltonian systems has been extensively studied by many authors [3,6,15-18, 20, 25, 29, 34, 38-41]. A Hamiltonian system is a mathematical expression of the following form:

$$
\begin{cases}-\Delta u=H_{v}(x, u, v), & x \in \Omega,  \tag{1.1}\\ -\Delta v=H_{u}(x, u, v), & x \in \Omega,\end{cases}
$$

where $\Omega$ is a smooth domain in $\mathbb{R}^{N}, N \geq 2$, and $H(x, u, v)$ is a nonlinear function.
In the case $N \geq 3$. A classical model for $H$ is given by $H(x, u, v)=|u|^{p+1} /(p+1)+|v|^{q+1} /(q+1)$
and the maximal growth for the exponents $p$ and $q$ are related to the curve [17,20,29]:

$$
\begin{equation*}
\frac{1}{p+1}+\frac{1}{q+1}=\frac{N-2}{N} \tag{1.2}
\end{equation*}
$$

If the couple $(p, q)$ lies on (1.2), some features of noncompactness arises, this motivates one to name (1.2) as the critical hyperbola, and we say that the nonlinearities $H_{v}=|v|^{q-1} v$ and $H_{u}=|u|^{p-1} u$ possess critical growth; alternatively, if the couple ( $p, q$ ) is below (1.2) the growth of the nonlinearities are denominated subcritical. We want to point out that the critical hyperbola results from the borderline between existence and nonexistence of solutions for (1.1) (see [6]).

In the case when $N=2$, the critical hyperbola is not defined. Notice that, if $\Omega$ is a bounded domain in $\mathbb{R}^{N}$, the Sobolev embeddings state $W_{0}^{1,2}(\Omega) \subset L^{q}(\Omega)$ for all $1 \leq q \leq 2^{*}=2 N /(N-2)$ for $N \geq 3$. In dimension $N=2$, one has $2^{*}=+\infty$ and $W_{0}^{1,2}(\Omega) \not \subset L^{\infty}(\Omega)$. Therefore, $H_{u}$ and $H_{v}$ may have any arbitrary polynomial growth. It was shown independently by Yudovich [44], Pohožaev [32], and Trudinger [43] that the growth is of exponential type. More precisely, $e^{\alpha u^{2}} \in L^{1}(\Omega)$ for all $u \in H_{0}^{1}(\Omega)$ and $\alpha>0$. Furthermore, Moser [30] proved the existence of a positive constant $C=C(\alpha, \Omega)$ such that

$$
\sup _{\substack{u \in H_{0}^{1}(\Omega)  \tag{1.3}\\\|\nabla u\|_{2} \leq 1}} \int_{\Omega} e^{\alpha u^{2}} d x \begin{cases}\leq C, & \alpha \leq 4 \pi \\ =+\infty, & \alpha>4 \pi\end{cases}
$$

From now on, the estimate of the type (1.3) will be referred to as the Trudinger-Moser inequality. These inequalities have been extended in many directions (see [3, 10, 12, 14, 23, 25, 28, 31, 36, 42] among others). The above results motivate us to say that the function $f$ has subcritical exponential growth if

$$
\lim _{s \rightarrow+\infty} \frac{f(s)}{e^{\alpha s^{2}}}=0, \quad \text { for all } \alpha>0
$$

and critical exponential growth if there exists $\alpha_{0}>0$ such that

$$
\lim _{s \rightarrow+\infty} \frac{f(s)}{e^{\alpha s^{2}}}= \begin{cases}0, & \alpha>\alpha_{0}  \tag{1.4}\\ +\infty, & \alpha<\alpha_{0}\end{cases}
$$

Nonlinear equations considering nonlinearities involving subcritical and critical exponential growth were treated by Adimurthi [1], Adimurthi-Yadava [2], de Figueiredo, Miyagaki, and Ruf [19] (see also $[10,12,24,31,35])$. We recall that a nonlinear equation in a domain $\Omega \subset \mathbb{R}^{N}$ with $N \geq 3$ a classical assumption on the nonlinearity is given by $|f(s)| \leq c\left(1+|s|^{q-1}\right)$, with $1<q \leq 2^{*}=2 N /(N-2)$ (see $[5,7,8,11,26,27]$ among others). If there exist positive constants $k$ and $s_{0}$ such that $g_{1}(s) \leq g_{2}(k s)$ for $s \geq s_{0}$, we shall write $g_{1}(s)<g_{2}(s)$. Additionally, we shall say that $g_{1}$ and $g_{2}$ are equivalent and write $g_{1}(s) \sim g_{2}(s)$ if $g_{1}(s)<g_{2}(s)$ and $g_{2}(s) \prec g_{1}(s)$. Therefore, $f$ possesses critical exponential growth if and only if $f(s) \sim e^{|s|^{2}}$.

The existence of a nontrivial solution of the system (1.1) under $H_{v} \sim e^{\nu^{2}}$ and $H_{u} \sim e^{u^{2}}$ and considering $H_{0}^{1}(\Omega) \times H_{0}^{1}(\Omega)$ as the setting space was proved by de Figueiredo, do Ó, and Ruf [18].

Now, we recall some facts about Lorentz-Sobolev spaces. Let $1<r<+\infty, 1 \leq s<+\infty$ and $\Omega$ subset of $\mathbb{R}^{N}$, the Lorentz space $L^{r, s}(\Omega)$ is the collection of all measurable and finite almost everywhere functions on $\Omega$ such that $\|\phi\|_{r, s}<+\infty$, where

$$
\|\phi\|_{r, s}=\left(\int_{0}^{+\infty}\left[\phi^{*}(t) t^{1 / r}\right]^{s} \frac{d t}{t}\right)^{1 / s}
$$

where $\phi *$ denotes the spherically symmetric decreasing rearrangement of $\phi$. In addition, if $\Omega$ is an open bounded domain in $\mathbb{R}^{N}$, the Lorentz-Sobolev space $W_{0}^{1} L^{r, s}(\Omega)$ is defined to be the closure of the compactly supported smooth functions on $\Omega$, with respect to the quasinorm

$$
\|u\|_{W_{0}^{1} L^{\prime, s}}:=\|\nabla u\|_{r, s} .
$$

Brezis and Wainger [9] proposed the following Trudinger-Moser inequality version on LorentzSobolev spaces: If $\Omega$ be a bounded domain in $\mathbb{R}^{2}$ and $s>1$, then $e^{\left.\alpha|u|\right|^{s-1}}$ belongs to $L^{1}(\Omega)$ for all $u \in W_{0}^{1} L^{2, s}(\Omega)$ and $\alpha>0$. Furthermore, Alvino [4] proved the following refinement of (1.3), there exists a positive constant $C=C(\Omega, s, \alpha)$ such that

$$
\sup _{\substack{u \in W_{W^{1}}^{1} L^{2, s}(\Omega)  \tag{1.5}\\\|\nabla u\| 2, s \leq 1}} \int_{\Omega} e^{\alpha|u|^{s-1}} d x \begin{cases}\leq C, & \alpha \leq(4 \pi)^{s /(s-1)}, \\ =+\infty, & \alpha>(4 \pi)^{s /(s-1)}\end{cases}
$$

Ruf [34] showed that, if the setting space of the system (1.1) is given by the product space $W_{0}^{1} L^{2, q}(\Omega) \times$ $W_{0}^{1} L^{2, p}(\Omega)$, the maximal growth of the nonlinearities can be considered like $H_{u} \sim e^{|u|^{p}}$ and $H_{v} \sim e^{|v|^{q}}$ with $p, q>1$ satisfying

$$
\begin{equation*}
\frac{1}{p}+\frac{1}{q}=1 \tag{1.6}
\end{equation*}
$$

In analogy to (1.2), the curve (1.6) is called exponential critical hyperbola. The existence of solutions of the system (1.8) for $p=q=2$ has been treated in many works [ $3,18,38,40,41$ ] among others, and the case where $(p, q)$ lies on the exponential critical hyperbola given by (1.6) was studied in [15,25,39].

Trudinger-Moser type inequalities for radial Sobolev spaces with logarithmic weights were considered by Calanchi and Ruf [12]. Denote by $H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$, the subspace of the radially symmetric functions in the closure of $C_{0}^{\infty}\left(B_{1}\right)$ with respect to the norm

$$
\|u\|:=\left(\int_{B_{1}}|\nabla u|^{2} w(x) d x\right)^{\frac{1}{2}},
$$

where

$$
\begin{equation*}
w(x)=\left[\log \left(\frac{1}{|x|}\right)\right]^{\gamma}, \quad 0 \leq \gamma<1 . \tag{1.7}
\end{equation*}
$$

Calanchi and Ruf [12] found that

$$
\int_{B_{1}} e^{\left.\alpha|u|\right|^{\frac{2}{-\gamma}}} d x<+\infty, \quad \text { for all } \quad u \in H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right) \text { and } \alpha>0
$$

Furthermore, if $\alpha \leq \alpha_{\gamma}^{*}=2[2 \pi(1-\gamma)]^{\frac{2}{2-\gamma}}$, there exists a positive constant $C$ such that

$$
\sup _{\|u\| \leq 1} \int_{B_{1}} e^{\alpha|u| 1-\frac{2}{1-\gamma}} d x \leq C .
$$

The above results represent an increase in the maximal growth of the exponential type. For $\lambda=1$, the weight given by (1.7) allows us to consider double exponential growth, see [12, 13, 37] for more
details. In this paper, we deal with the existence of solutions to the following Hamiltonian system:

$$
\left\{\begin{align*}
-\operatorname{div}(w(x) \nabla u)=g(x, v), & x \in B_{1},  \tag{1.8}\\
-\operatorname{div}(w(x) \nabla v) & =f(x, u), \\
& x \in B_{1}, \\
u=v \equiv 0, & x \in \partial B_{1},
\end{align*}\right.
$$

where $w$ is given by (1.7) and $B_{1}$ denotes the unit open ball center at the origin in $\mathbb{R}^{2}$. In order to use variational methods, we consider an associated functional defined on the space $H_{0, \text { rad }}^{1}\left(B_{1}, w\right) \times$ $H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$, which allows us to have nonlinearities of the form $f(u) \sim e^{|u|^{2 /(1-\gamma)}}$ and $g(v) \sim e^{|\nu|^{2 /(1-\gamma)}}$.

We assume the following conditions on the nonlinerities $f$ and $g$ :
$\left(H_{1}\right) f, g \in C\left(\bar{B}_{1} \times \mathbb{R}\right)$ and $f(x, s)=g(x, s)=o(s)$ as $s \rightarrow 0^{+}$and $f(x, s)=g(x, s)=0$ for all $x \in B_{1}$ and $s \leq 0$.
$\left(H_{2}\right)$ There exist constants $\mu>2, v>2$ and $s_{0}>0$ such that

$$
0<\mu F(x, s) \leq s f(x, s) \quad \text { and } \quad 0<v G(x, s) \leq \operatorname{sg}(x, s), \quad \text { for all } \quad x \in B_{1} \text { and } s>s_{0}
$$

where $F(x, s)=\int_{0}^{s} f(x, t) d t$ and $G(x, s)=\int_{0}^{s} g(x, t) d t$.
$\left(H_{3}\right)$ There exist constants $M>0$ and $s_{1}>0$ such that

$$
0<F(x, s) \leq M f(x, s) \quad \text { and } \quad 0<G(x, s) \leq M g(x, s), \quad \text { for all } \quad x \in B_{1} \text { and } s>s_{1} .
$$

$\left(H_{4}\right)$ There exist constants $\alpha_{0}>0$ and $\beta_{0}>0$ such that

$$
\lim _{s \rightarrow \infty} \frac{f(x, s)}{e^{\alpha s^{2} /(1-\gamma)}}=\left\{\begin{array}{ll}
0, & \alpha>\alpha_{0}, \\
+\infty, & \alpha<\alpha_{0},
\end{array} \quad \text { and } \quad \lim _{s \rightarrow \infty} \frac{g(x, s)}{e^{\beta s^{2 /(1-\gamma)}}}= \begin{cases}0, & \beta>\beta_{0} \\
+\infty, & \beta<\beta_{0}\end{cases}\right.
$$

uniformly on $x \in B_{1}$.
$\left(H_{5}\right)$ There exist constants $p>2$ and $C_{p}>0$ such that

$$
f(x, s) \geq C_{p} s^{p-1} \quad \text { and } \quad g(s) \geq C_{p} s^{p-1}, \quad \text { for all } \quad s \geq 0
$$

where

$$
C_{p}>\frac{(p-2)^{(p-2) / 2}\left(\max \left\{\alpha_{0}, \beta_{0}\right\}\right)^{(p-2)(1-\gamma) / 2} S_{p}^{p}}{p^{(p-2) / 2}\left(\alpha_{\gamma}^{*}\right)^{(p-2)(1-\gamma) / 2}}
$$

and

$$
S_{p}:=\inf _{0 \neq u \in H_{0, r a d}^{1}} \frac{\left(\int_{\left.B_{1}, w\right)} w(x)|\nabla u|^{2} d x\right)^{1 / 2}}{\left(\int_{B_{1}}|u|^{p} d x\right)^{1 / p}}
$$

Setting the product space

$$
E=H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right) \times H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right),
$$

which is a Hilbert space endowed with the inner product

$$
\langle(u, v),(\phi, \psi)\rangle_{E}=\int_{B_{1}} w(x)(\nabla u \nabla \phi+\nabla v \nabla \psi) d x
$$

for all $(u, v),(\phi, \psi) \in E$, to which corresponds the norm

$$
\|(u, v)\|=\|(u, v)\|_{E}:=\left(\|u\|^{2}+\|v\|^{2}\right)^{1 / 2} .
$$

Additionally, we denote the dual space of $E$ with its usual norm by $E^{*}$. We say that $(u, v) \in E$ is a weak solution of (1.8) if

$$
\begin{equation*}
\int_{B_{1}} w(x)(\nabla u \nabla \psi+\nabla v \nabla \phi) d x=\int_{B_{1}}(f(x, u) \phi+g(x, v) \psi) d x, \quad \text { for all } \quad(\phi, \psi) \in E . \tag{1.9}
\end{equation*}
$$

Under the assumption on $f$ and $g$, we establish the Euler-Lagrange functional $J: E \rightarrow \mathbb{R}$ defined by

$$
J(u, v)=\int_{B_{1}} w(x) \nabla u \nabla v d x-\int_{B_{1}} F(x, u) d x-\int_{B_{1}} G(x, v) d x,
$$

for all $(u, v) \in E$. Furthermore, using standard arguments [21], $J \in C^{1}(E, \mathbb{R})$ and, for all $(u, v)$, $(\phi, \psi) \in E$, it holds

$$
J^{\prime}(u, v)(\phi, \psi)=\int_{B_{1}} w(x)(\nabla u \nabla \psi+\nabla v \nabla \phi) d x-\int_{B_{1}} f(x, u) \phi d x-\int_{B_{1}} g(x, v) \psi d x .
$$

In particular, $(u, v) \in E$ is a nontrivial weak solution of the system (1.8) if only if $(u, v) \in E$ is a nontrivial critical point of the functional $J$. Next, we present our existence result for the system (1.8).

Theorem 1.1. Suppose that $f$ and $g$ satisfy $\left(H_{1}\right)-\left(H_{5}\right)$. Then, the Hamiltonian system (1.8) has a nontrivial weak solution.

First, observe that if $\gamma=0$, then the system (1.8) is reduced to $-\Delta u=g(x, v)$ and $-\Delta v=f(x, u)$ with Dirichlet conditions, and the growth of the functions are given by $f(x, u) \sim e^{|u|^{2}}$ and $g(x, v) \sim e^{|v|^{2}}$ uniformly on $x \in B_{1}$, whose existence of nontrivial weak solutions was found in [18]. In the case for $\gamma>0$, the nonlinearities under the assumption $\left(H_{4}\right)$ behaves like $f(x, u) \sim e^{|u|^{p}}$ and $g(x, v) \sim e^{|v|^{q}}$ uniformly on $x \in B_{1}$, where $1 / p+1 / q<1$ and $p=q$, that is, the pair $(p, q)$ lies in the diagonal direction above the exponential critical hyperbola. Therefore, our result treats the Hamiltonian system (1.8) involving nonlinearities with supercritical exponential growth. Consequently, our result complements the works which study nonlinearities $f(x, u) \sim e^{|u|^{p}}$ and $g(x, v) \sim e^{|v|^{q}}$ for values where $(p, q)$ lies under and on the the curve (1.6), that is, for nonlinerities that possess subcritical and critical exponential growth, respectively [15, 18, 25, 38-41].

The paper is organized as follows: Section 2 contains some preliminaries results and properties our setting space. In Section 3, we show that the Euler-Lagrange energy functional possesses the geometry of the linking theorem. In Section 4, it is established the finite-dimensional approximation and estimated the minimax level of the functional. Finally, in Section 5, we present the proof of Theorem 1.1.

## 2. Preliminaries

Let $H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$ be the subspace of the radially symmetric functions in the closure of $C_{0}^{\infty}\left(B_{1}\right)$ with respect to the norm

$$
\|u\|=\|u\|_{H_{0, \text { rad }}^{1}\left(B_{1}, w\right)}:=\left(\int_{B_{1}}|\nabla u|^{2} w(x) d x\right)^{1 / 2}
$$

where $w(x)=(\log 1 /|x|)^{\gamma}$ and $0 \leq \gamma<1$.
Proposition 2.1. The Sobolev weighted space $H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right)$ is a separable Banach space.
Proof. Let $0 \leq \gamma<1$ be fixed and consider $s(t)=(\log 1 /(1-t))^{\gamma}$. Then, $s$ is a continuous positive function defined for the interval $(0,1)$ and satisfies $\lim _{t \rightarrow 0} s(t)=0$. Moreover, if $\Omega=B_{1}(0)$, we have $d(x)=\operatorname{dist}(x, \partial \Omega)=1-|x|$. Thus, $s \circ d=w$. Therefore, we get that $W^{1,2}(\Omega, s \circ d)$ and its restrictions to the radial symmetric functions are separable Banach spaces (see Theorem 3.9 in [22]).

We note that $H_{0, \text { rad }}^{1}(B, w)$ is a Hilbert space endowed with inner product

$$
\langle u, v\rangle:=\int_{B_{1}} w(x) \nabla u \nabla v d x, \quad \text { for all } u, v \in H_{0, \text { rad }}^{1}(B, w) .
$$

Now, we state a compactness result.
Lemma 2.2. The embedding $H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right) \hookrightarrow L^{p}\left(B_{1}\right)$ is continuous and compact for $1 \leq p<\infty$.
Proof. It follows from the Cauchy-Schwarz inequality that

$$
\int_{B_{1}}|\nabla u| d x \leq\left(\int_{B_{1}}|\nabla u|^{2} w(x) d x\right)^{1 / 2} \cdot\left(\int_{B_{1}} w(x)^{-1} d x\right)^{1 / 2} .
$$

Using the change of variable $|x|=e^{-s}$, we obtain

$$
\frac{1}{2 \pi} \int_{B_{1}} w(x)^{-1} d x=\int_{0}^{+\infty} e^{-2 s} s^{-\gamma} d s=\int_{0}^{1} e^{-2 s} s^{-\gamma} d s+\int_{1}^{+\infty} e^{-2 s} s^{-\gamma} d s
$$

where

$$
\int_{0}^{1} e^{-2 s} s^{-\gamma} d s \leq \int_{0}^{1} s^{-\gamma} d s=\frac{1}{1-\gamma}
$$

and

$$
\int_{1}^{+\infty} e^{-2 s} s^{-\gamma} d s \leq \int_{1}^{+\infty} e^{-2 s} d s=\frac{e^{-2}}{2}
$$

Therefore, there exists $C>0$ such that

$$
\|\nabla u\|_{1} \leq C\left(\int_{B_{1}}|\nabla u|^{2} w(x) d x\right)^{1 / 2}
$$

Thus, $H_{0, \text { rad }}^{1}\left(B_{1}, w\right) \hookrightarrow W_{0}^{1,1}\left(B_{1}\right)$ continuosly, which implies the continuous and compact embedding

$$
H_{0, \text { rad }}^{1}\left(B_{1}, w\right) \hookrightarrow L^{p}\left(B_{1}\right), \quad \text { for all } \quad p \geq 1
$$

Proposition 2.3. (See [12]) Let w be the weight given by (1.7). Then,

$$
\begin{equation*}
\int_{B_{1}} e^{\left.\alpha| | u\right|^{\frac{2}{1-\gamma}}} d x<+\infty, \quad \text { for all } u \in H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right) \text { and } \alpha>0 \tag{2.1}
\end{equation*}
$$

Furthermore, if $\alpha \leq \alpha_{\gamma}^{*}=2[2 \pi(1-\gamma)]^{\frac{1}{2-\gamma}}$, there exists a positive constant $C$ such that

$$
\begin{equation*}
\sup _{\|u\| \leq 1} \int_{B_{1}} e^{\alpha|u| \left\lvert\,-\frac{2}{1-\gamma}\right.} d x \leq C . \tag{2.2}
\end{equation*}
$$

Lemma 2.4. Let $\left(u_{n}\right)$ be a sequence in $H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$ such that $u_{n} \rightharpoonup 0$ in $H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$ and $\left\|u_{n}\right\|=1$ for every $n \in \mathbb{N}$. Then, for every $0<\alpha<\alpha_{\gamma}^{*}$, there exists a subsequence, still denoted by ( $u_{n}$ ) such that

$$
\int_{B_{1}}\left(e^{\left.\alpha\left|u_{n}\right|\right|^{\frac{2}{-\gamma}}}-1\right) d x=0 \rightarrow 0, \quad \text { as } \quad n \rightarrow+\infty
$$

Proof. Choosing $\epsilon>0$ such that $\alpha+\epsilon<\alpha_{\gamma}^{*}$. We have the following limits:

$$
\lim _{|t| \rightarrow 0} \frac{e^{\alpha|t|^{\frac{2}{1-\gamma}}}-1}{|t|}=0 \quad \text { and } \quad \lim _{|t| \rightarrow \infty} \frac{e^{\left.\alpha|t|\right|^{2}--\gamma}}{|t| e^{(\alpha+\epsilon)|t|^{\frac{2}{1-\gamma}}}}=0 .
$$

Thus, we can find $C>0$ such that

$$
e^{\alpha|t|^{\frac{2}{1-\gamma}}}-1 \leq C|t|+C|t| e^{\left.(\alpha+\epsilon) t\right|^{\frac{2}{1-\gamma}}}, \quad \text { for all } \quad t \in \mathbb{R} .
$$

Using the Hölder inequality with $r>1$ such that $r(\alpha+\epsilon)<\alpha_{\gamma}^{*}$ and Proposition 2.3, we get

$$
\int_{B_{1}}\left(e^{\alpha\left|u_{n}\right|^{\frac{2}{-\gamma}}}-1\right) d x \leq C\left\|u_{n}\right\|_{1}+C\left\|u_{n}\right\|_{r^{\prime}} \int_{B_{1}} e^{r(\alpha+\epsilon)\left|u_{n}\right| \frac{2}{1-\gamma}} d x \leq C\left\|u_{n}\right\|_{1}+C\left\|u_{n}\right\| r_{r^{\prime}} .
$$

By the weakly convergence of $u_{n} \rightharpoonup 0$ in $H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$ and Lemma 2.2 for a subsequence, we have

$$
\int_{B_{1}}\left(e^{\alpha\left|u_{n}\right|^{\frac{2}{-\gamma}}}-1\right) d x \rightarrow 0 .
$$

## 3. Linking geometry

This section is devoted to prove that the functional $J$ possesses the geometry of the linking theorem. We start setting the following subspaces:

$$
E^{+}=\left\{(u, u): u \in H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right)\right\} \quad \text { and } \quad E^{-}=\left\{(u,-u): u \in H_{0, \mathrm{rad}}^{1}\left(B_{1}, w\right)\right\},
$$

In particular, we have

$$
E=E^{+} \oplus E^{-}
$$

Lemma 3.1. Assume that $\left(H_{1}\right),\left(H_{4}\right)$, and $\left(H_{5}\right)$, are hold. Then, there exist positive constants $\sigma$ and $\rho$ such that $J(z) \geq \sigma$ for all $z \in \partial B_{\rho} \cap E^{+}$.
Proof. Given $\epsilon>0$ and $q>2$, it follows from $\left(H_{1}\right)$ and $\left(H_{4}\right)$, the existence of some $c>0$ such that

$$
|G(x, s)|,|F(x, s)| \leq \epsilon|s|^{2}+c|s|^{q} e^{2 \alpha_{0}|s|^{\frac{2}{2-\gamma}}}, \quad \text { for all } \quad(x, s) \in B_{1} \times \mathbb{R} .
$$

Using the Cauchy-Schwarz inequality and (2.2), we obtain

$$
\int_{B_{1}} F(x, u) d x \leq \epsilon\|u\|_{2}^{2}+c\|u\|_{2 q}^{q} \int e^{4 \alpha_{0}|u|^{\frac{2}{2-\gamma}}} d x \leq \epsilon\|u\|_{2}^{2}+c\|u\|_{2 q}^{q},
$$

provided that $\|u\| \leq \rho_{0}$ for some $\rho_{0}>0$ such that $4 \alpha_{0} \rho_{0}^{\frac{2}{2-\gamma}}<\alpha_{\gamma}^{*}$. A similar inequality holds for the function $G$. By Lemma 2.2, we obtain $c_{1}>0$ and $c_{2}>0$ such that
$J(u, u) \geq\|u\|^{2}-\int_{B_{1}} F(x, u) d x-\int_{B_{1}} G(x, u) d x \geq\|u\|^{2}-2 \epsilon\|u\|_{2}^{2}-2 c\|u\|_{2 q}^{q} \geq\left(1-2 \epsilon c_{1}\right)\|u\|^{2}-c_{2}\|u\|^{q}$.
Therefore, taking $\epsilon>0$ sufficiently small, we can choose $\rho>0$ sufficiently small and $\sigma>0$ such that $J(z) \geq \sigma$, for all $z \in \partial B_{\rho} \cap E^{+}$.

Taking $e \in H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$ such that $\|e\|=1$, we define

$$
Q_{e}=\left\{r(e, e)+(\omega,-\omega):\|(\omega,-\omega)\| \leq R_{0}, 0 \leq r \leq R_{1}\right\},
$$

where $R_{0}>$ and $R_{1}>0$ will be determined by the following result.
Lemma 3.2. Assume that $\left(H_{1}\right),\left(H_{2}\right)$ and $\left(H_{5}\right)$, are hold. Then, there exist positive constants $R_{0}$ and $R_{1}$ such that

$$
J(z) \leq 0, \quad \text { for all } \quad z \in \partial Q_{e} .
$$

Proof. We observe that the boundary $\partial Q_{e}$ of the set $Q_{e}$ consists of three parts that corresponds the following cases:
(i) Let $z=(\omega,-\omega) \in \partial Q_{e} \cap E^{-}$. Therefore,

$$
J(z)=-\|\omega\|^{2}-\int_{B_{1}} F(x, \omega) d x-\int_{B_{1}} G(x,-\omega) d x \leq 0 .
$$

(ii) Let $z=r(e, e)+(\omega,-\omega)=(r e+\omega, r e-\omega) \in \partial Q_{e}$, with $\|(\omega,-\omega)\|=R_{0}$ and $0 \leq r \leq R_{1}$, then

$$
J(z)=r^{2}\|e\|^{2}-\|\omega\|^{2}-\int_{B_{1}} F(x, r e+\omega) d x-\int_{B_{1}} G(x, r e-\omega) d x \leq R_{1}^{2}\|e\|^{2}-\|\omega\|^{2}=R_{1}^{2}-\frac{R_{0}^{2}}{2} .
$$

Thus, $J(z) \leq 0$ if $R_{0}=\sqrt{2} R_{1}$.
(iii) Let $z=R_{1}(e, e)+(\omega,-\omega) \in \partial Q_{e}$, with $\|(\omega,-\omega)\| \leq R_{0}$, it follows from $\left(H_{5}\right)$ that

$$
F(x, s), G(x, s) \geq \frac{C_{p}}{p}|s|^{p} \quad \text { for all } x \in B_{1} \text { and } s \geq 0
$$

Therefore,

$$
\begin{aligned}
J(z) & =R_{1}^{2}\|e\|^{2}-\|\omega\|^{2}-\int_{B_{1}} F\left(x, R_{1} e+\omega\right) d x-\int_{B_{1}} G\left(x, R_{1} e-\omega\right) d x \\
& \leq R_{1}^{2}-\int_{B_{1}} F\left(x, R_{1} e+\omega\right) d x-\int_{B_{1}} G\left(x, R_{1} e-\omega\right) d x \\
& \leq R_{1}^{2}-\frac{C_{p}}{p} \int_{B_{1}}\left(\left|R_{1} e+\omega\right|^{p}+\left|R_{1} e-\omega\right|^{p}\right) d x \\
& \leq R_{1}^{2}-\frac{2 C_{p} R_{1}^{p}}{p} \int_{B_{1}}|e|^{p} d x .
\end{aligned}
$$

Since $p>2$, we can choose $R_{1}>0$ such that $J(z) \leq 0$.

## 4. Finite-dimensional approximation

Observe that the leading part of the functional $J$ is strongly indefinite, that is, $J$ can assume positives and negatives values on infinite-dimensional subspaces of $E$. Therefore, we can not use the linking theorem. To deal with this inconvenience, we follow the arguments developed by de Figueiredo, do Ó, and Ruf [16], that is, we use a finite dimensional approximation.

Let $e=u_{p} \in H_{0, \text { rad }}^{1}\left(B_{1}, w\right)$ be a nonnengative function with $\left\|u_{p}\right\|_{p}=1$ where $S_{p}$ is attained. We consider $\left\{e_{i}\right\}_{i \in \mathbb{N}}$ a Hilbert basis of $\langle e\rangle^{\perp}$ and setting

$$
E_{n}^{+}=\operatorname{Span}\left\{\left(e_{i}, e_{i}\right): i=1,2, \ldots, n\right\}
$$

and

$$
\begin{gathered}
E_{n}^{-}=\operatorname{Span}\left\{\left(e_{i},-e_{i}\right): i=1,2 \ldots, n\right\}, \\
E_{n}=E_{n}^{+} \oplus E_{n}^{-} .
\end{gathered}
$$

we denote by

$$
H_{n}=\mathbb{R}(e, e) \oplus E_{n}, \quad H_{n}^{+}=\mathbb{R}(e, e) \oplus E_{n}^{+} \quad \text { and } \quad H_{n}^{-}=\mathbb{R}(e, e) \oplus E_{n}^{-}
$$

Setting the following class of functions:

$$
\Gamma_{n}=\left\{\gamma \in C\left(Q_{n}, H_{n}\right): \gamma(z)=z, \forall z \in \partial Q_{n}\right\},
$$

where $Q_{n}=Q_{e} \cap H_{n}$, and set

$$
\begin{equation*}
c_{n}=\inf _{\gamma \in \Gamma_{n}} \max _{z \in Q_{n}} J(\gamma(z)) . \tag{4.1}
\end{equation*}
$$

Now, let $J_{n}$ be the restriction of $J$ to the finite-dimensional space $H_{n}$. Moreover, Lemmas 3.1 and 3.2 are still valid for $J_{n}$. Additionally, it follows from [16] that

$$
\begin{equation*}
\gamma\left(Q_{n}\right) \cap\left(\partial B_{\rho} \cap H_{n}^{+}\right) \neq \emptyset, \quad \text { for all } \quad \gamma \in \Gamma_{n}, \tag{4.2}
\end{equation*}
$$

for $\rho$ given by Lemma 3.1. Moreover, Lemma 3.1 and (4.2), implies that

$$
c_{n} \geq \sigma>0, \quad \text { for all } \quad n \geq 1
$$

Using the fact that the identity map $I_{n}: Q_{n} \rightarrow H_{n}$ belongs to $\Gamma_{n}$ and the fact that $F$ and $G$ are nonnegative functions, we obtain

$$
\begin{equation*}
J(z)=r^{2}\|e\|^{2}-\|u\|^{2}-\int_{B_{1}} F(x, r e+u) d x-\int_{B_{1}} G(x, r e-u) d x \leq R_{1}^{2}, \tag{4.3}
\end{equation*}
$$

for each $z=r(e, e)+(u,-u) \in Q_{n}$. Hence,

$$
\begin{equation*}
c_{n} \leq R_{1}^{2}, \quad \text { for all } n \geq 1 \tag{4.4}
\end{equation*}
$$

Next, this proposition follows from the linking theorem for $J_{n}$ (see [33]).

Proposition 4.1. Assume that $f$ and $g$ satisfy $\left(H_{1}\right)-\left(H_{4}\right)$. Then, the functional $J_{n}$ possesses a critical point $z_{n}=\left(u_{n}, v_{n}\right) \in H_{n}$ at level $c_{n}$ for all $n \in \mathbb{N}$, satisfying

$$
\begin{equation*}
J\left(z_{n}\right)=c_{n} \in\left[\sigma, R_{1}^{2}\right], \tag{4.5}
\end{equation*}
$$

where $\sigma$ and $R_{1}>0$ are given by Lemmas 3.1 and 3.2, respectively, and

$$
\begin{equation*}
J_{n}^{\prime}\left(z_{n}\right)(\phi, \psi)=0, \quad \text { for all }(\phi, \psi) \in H_{n}, \tag{4.6}
\end{equation*}
$$

that is

$$
\begin{equation*}
\int_{B_{1}} w(x) \nabla u_{n} \nabla \psi d x=\int_{B_{1}} g\left(x, v_{n}\right) \psi d x \quad \text { and } \quad \int_{B_{1}} w(x) \nabla \phi \nabla v_{n} d x=\int_{B_{1}} f\left(x, u_{n}\right) \phi d x, \tag{4.7}
\end{equation*}
$$

for each $(\phi, \psi) \in H_{n}$.
Lemma 4.2. (See [15, Lemma 10]) If $r, r^{\prime}>1$ satisfy $1 / r+1 / r^{\prime}=1$, then

$$
s t \leq \begin{cases}\left(e^{t^{r}}-1\right)+s(\ln s)^{1 / r}, & \text { for all } t \geq 0 \text { and } s \geq e^{1 / r^{\prime}} \\ \left(e^{t^{r}}-1\right)+\frac{s^{r^{\prime}}}{r^{\prime}}, & \text { for all } t \geq 0 \text { and } 0 \leq s \leq e^{1 / r^{\prime}}\end{cases}
$$

Lemma 4.3. Let $\left(u_{n}, v_{n}\right)$ be a sequence in $E$ satisfying $\left|J\left(u_{n}, v_{n}\right)\right| \leq d$ and

$$
\begin{equation*}
\left|J^{\prime}\left(u_{n}, v_{n}\right)(\phi, \psi)\right| \leq \epsilon_{n}\|(\phi, \psi)\|, \text { for all } \phi, \psi \in\left\{0, u_{n}, v_{n}\right\}, \text { where } \epsilon_{n} \rightarrow 0 \text {. } \tag{4.8}
\end{equation*}
$$

Then, the sequence $\left(u_{n}, v_{n}\right)$ is bounded in $E$.
Proof. Taking $(\phi, \psi)=\left(u_{n}, v_{n}\right)$ in (4.8), we have

$$
\int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x+\int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x \leq\left|2 \int_{B_{1}} w(x) \nabla u_{n} \nabla v_{n} d x\right|+\epsilon_{n}\left\|\left(u_{n}, v_{n}\right)\right\| .
$$

Since

$$
\int_{B_{1}} w(x) \nabla u_{n} \nabla v_{n} d x=J\left(u_{n}, v_{n}\right)+\int_{B_{1}} F\left(x, u_{n}\right) d x+\int_{B_{1}} G\left(x, v_{n}\right) d x,
$$

combined with the fact that $\left|J\left(u_{n}, v_{n}\right)\right| \leq d$, we obtain

$$
\begin{equation*}
\int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x+\int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x \leq 2 d+2 \int_{B_{1}} F\left(x, u_{n}\right) d x+2 \int_{B_{1}} G\left(x, v_{n}\right) d x+\epsilon_{n}\left\|\left(u_{n}, v_{n}\right)\right\| . \tag{4.9}
\end{equation*}
$$

From $\left(H_{2}\right)$, we get

$$
\begin{aligned}
\int_{B_{1}} F\left(x, u_{n}\right) d x & =\int_{\left\{x \in B_{1}:\left|u_{n}(x)\right| \leq s_{0}\right\}} F\left(x, u_{n}\right) d x+\int_{\left\{x \in B_{1}:\left|u_{n}(x)\right|>s_{0}\right\}} F\left(x, u_{n}\right) d x \\
& \leq \int_{\left\{x \in B_{1}:\left|u_{n}(x)\right| \leq s_{0}\right\}} F\left(x, u_{n}\right) d x+\frac{1}{\mu} \int_{\left\{x \in B_{1}:\left|u_{n}(x)\right|>s_{0}\right\}} f\left(x, u_{n}\right) u_{n} d x \\
& \left.=\int_{\left\{x \in B_{1}:\left|u_{n}(x)\right| \leq s_{0}\right\}} F\left(x, u_{n}\right)-\frac{1}{\mu} f\left(x, u_{n}\right) u_{n}\right) d x+\frac{1}{\mu} \int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x
\end{aligned}
$$

$$
\begin{equation*}
\leq M_{f}\left|B_{1}\right|+\frac{1}{\mu} \int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x \tag{4.10}
\end{equation*}
$$

where

$$
M_{f}=\max _{(x, s) \in \overline{\bar{B}_{1} \times\left[0, s_{0}\right]}}\left(|F(x, s)|+\frac{1}{\mu}|f(x, s) s|\right) .
$$

Similarly for the function $g$, there exists $M_{g}>0$ such that

$$
\begin{equation*}
\int_{B_{1}} G\left(x, v_{n}\right) d x \leq M_{g}\left|B_{1}\right|+\frac{1}{v} \int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x \tag{4.11}
\end{equation*}
$$

From (4.10) and (4.11) in (4.9), we obtain

$$
\begin{equation*}
\left(1-\frac{2}{\mu}\right) \int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x+\left(1-\frac{2}{v}\right) \int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x \leq 2 d+2\left(M_{f}+M_{g}\right)\left|B_{1}\right|+\epsilon_{n}\left\|\left(u_{n}, v_{n}\right)\right\| . \tag{4.12}
\end{equation*}
$$

Taking $(\phi, \psi)=\left(v_{n}, 0\right)$ and $(\phi, \psi)=\left(0, u_{n}\right)$ in (4.8), we get

$$
\left\|v_{n}\right\|^{2}=\int_{B_{1}} w(x) \nabla v_{n} \nabla v_{n} d x \leq \int_{B_{1}} f\left(x, u_{n}\right) v_{n} d x+\epsilon_{n}\left\|\left(v_{n}, 0\right)\right\|
$$

and

$$
\left\|u_{n}\right\|^{2}=\int_{B_{1}} w(x) \nabla u_{n} \nabla u_{n} d x \leq \int_{B_{1}} g\left(x, v_{n}\right) u_{n} d x+\epsilon_{n}\left\|\left(0, u_{n}\right)\right\| .
$$

We define

$$
V_{n}=\frac{v_{n}}{\left\|v_{n}\right\|} \quad \text { and } \quad U_{n}=\frac{u_{n}}{\left\|u_{n}\right\|}
$$

Thus,

$$
\begin{equation*}
\left\|v_{n}\right\| \leq \int_{B_{1}} f\left(x, u_{n}\right) V_{n} d x+\epsilon_{n}\left\|v_{n}\right\| \tag{4.13}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|u_{n}\right\| \leq \int_{B_{1}} g\left(x, v_{n}\right) U_{n} d x+\epsilon_{n}\left\|u_{n}\right\| . \tag{4.14}
\end{equation*}
$$

Let $\alpha_{1}=\alpha_{0}+\xi$. By assumption $\left(H_{4}\right)$, there exists $\lambda>0$ such that

$$
\begin{equation*}
|f(x, s)| \leq \lambda e^{\alpha_{1}|s|^{\frac{2}{1-\gamma}}}, \quad \text { for all } \quad(x, s) \in B_{1} \times \mathbb{R} \tag{4.15}
\end{equation*}
$$

Set $\alpha_{2}=\alpha_{\gamma}^{*}-\xi$, using (4.13), we can write

$$
\left\|v_{n}\right\| \leq \frac{\lambda}{\alpha_{2}^{\frac{1-\gamma}{2}}} \int_{B_{1}} \frac{\left|f\left(x, u_{n}(x)\right)\right|}{\lambda} \alpha_{2}^{\frac{1-\gamma}{2}}\left|V_{n}(x)\right| d x .
$$

From Lemma 4.2 with $s=\left|f\left(x, u_{n}(x)\right)\right| / \lambda, t=\alpha_{2}^{\frac{1-\gamma}{2}}\left|V_{n}(x)\right|, r=2 /(1-\gamma)$ and $r^{\prime}=2 /(1+\gamma)$, we obtain

$$
\left\|v_{n}\right\| \leq \frac{\lambda}{\alpha_{2}^{\frac{1-\gamma}{2}}}\left[\int_{B_{1}}\left(e^{\left.\alpha_{2}\left|V_{n}\right|\right|^{\frac{2}{-\gamma}}}-1\right) d x+\frac{1+\gamma}{2 \lambda^{\frac{2}{1+\gamma}}} \int_{\left\{x \in B_{1}: \frac{\left|\frac{\mid\left(x, u_{n}(x)\right.}{n}\right|}{} e^{\left.\frac{1-\gamma}{2} \frac{2}{1+\gamma}\right\}}\right\}}\left|f\left(x, u_{n}\right)\right|^{\frac{2}{1+\gamma}} d x\right.
$$

$$
\begin{equation*}
\left.+\int_{\left\{x \in B_{1}: \frac{\mid f\left(x, u_{n}(x) \mid\right.}{1} \geq e^{\frac{1-y}{2} \frac{2}{1+\gamma}}\right\}} \frac{\left|f\left(x, u_{n}\right)\right|}{\lambda}\left(\ln \frac{\left|f\left(x, u_{n}\right)\right|}{\lambda}\right)^{\frac{1-\gamma}{2}} d x\right]+\epsilon_{n}\left\|v_{n}\right\| . \tag{4.16}
\end{equation*}
$$

By (4.15), we obtain

$$
\begin{equation*}
\int_{\left\{x \in B_{1}: \frac{\mid f\left(x, u_{n}(x) \mid\right.}{\lambda} \geq e^{\left.\frac{1-\gamma}{2} \frac{2}{1+\gamma}\right\}}\right.} \frac{\left|f\left(x, u_{n}\right)\right|}{\lambda}\left(\ln \frac{\left|f\left(x, u_{n}\right)\right|}{\lambda}\right)^{\frac{1-\gamma}{2}} d x \leq \frac{\alpha_{1}^{\frac{1-\gamma}{2}}}{\lambda} \int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x . \tag{4.17}
\end{equation*}
$$

Observe that

$$
\begin{equation*}
\left.\left.\int_{\left\{x \in B_{1}: \frac{\mid f\left(x, u_{n}(x) \mid\right.}{1} \leq e^{\left.\frac{1-\gamma}{2} \frac{2}{1+\gamma}\right\}}\right.} \right\rvert\, f\left(x, u_{n}\right)\right)^{\frac{2}{1+\gamma}} d x \leq\left(\lambda e^{\frac{1-\gamma}{2} \frac{2}{1+\gamma}}\right)^{\frac{2}{1+\gamma}}\left|B_{1}\right| . \tag{4.18}
\end{equation*}
$$

From Proposition 2.3, we have

$$
\begin{equation*}
\int_{B_{1}}\left(e^{\alpha_{2}\left|V_{n}\right|^{\frac{2}{1-\gamma}}}-1\right) d x \leq C . \tag{4.19}
\end{equation*}
$$

By replacing (4.17)-(4.19) in (4.16), we get $c_{1}>0$ and $c_{2}>0$ such that

$$
\begin{equation*}
\left\|v_{n}\right\| \leq c_{1} \int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x+c_{2}+\epsilon_{n}\left\|v_{n}\right\| \tag{4.20}
\end{equation*}
$$

Similarly, we get

$$
\begin{equation*}
\left\|u_{n}\right\| \leq c_{1} \int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x+c_{2}+\epsilon_{n}\left\|u_{n}\right\| \tag{4.21}
\end{equation*}
$$

Using (4.12), (4.20) and (4.21), we can find $c>0$ such that

$$
\left\|v_{n}\right\| \leq c+\epsilon_{n}\left\|\left(u_{n}, v_{n}\right)\right\|+\epsilon_{n}\left\|v_{n}\right\| \quad \text { and } \quad\left\|u_{n}\right\| \leq c+\epsilon_{n}\left\|\left(u_{n}, v_{n}\right)\right\|+\epsilon_{n}\left\|u_{n}\right\| .
$$

We finally obtain

$$
\|\left(u_{n}, v_{n}\left\|\leq c+\epsilon_{n}\right\|\left(u_{n}, v_{n}\right) \|\right.
$$

which implies that $\left\|\left(u_{n}, v_{n}\right)\right\| \leq c$, for every $n \in \mathbb{N}$, for some positive constant $c$.
Lemma 4.4. Assuming the conditions $\left(H_{1}\right)-\left(H_{5}\right)$, are hold. Let $\left(u_{n}, v_{n}\right)$ be a sequence in $E$ and $(u, v) \in$ $E$ such that $\left(u_{n}, v_{n}\right) \rightharpoonup(u, v)$ weakly in $E, J\left(u_{n}, v_{n}\right) \rightarrow c$ and $\left\|J^{\prime}\left(u_{n}, v_{n}\right)\right\|_{E^{*}} \rightarrow 0$. Then,
(i) $f\left(x, u_{n}\right) \rightarrow f(x, u)$ in $L^{1}\left(B_{1}\right)$ and $g\left(x, u_{n}\right) \rightarrow g(x, u)$ in $L^{1}\left(B_{1}\right)$,
(ii) $F\left(x, u_{n}\right) \rightarrow F(x, u)$ in $L^{1}\left(B_{1}\right)$ and $G\left(x, u_{n}\right) \rightarrow G(x, u)$ in $L^{1}\left(B_{1}\right)$.

Proof. From Lemma 2.2, we can suppose that $u_{n}$ converges to $u$ in $L^{1}\left(B_{1}\right)$. By Proposition 2.3, and the assumptions $\left(H_{1}\right)$ and $\left(H_{4}\right)$, we imply that $f\left(x, u_{n}\right) \in L^{1}\left(B_{1}\right)$. Moreover, using $J^{\prime}\left(u_{n}, v_{n}\right)\left(u_{n}, v_{n}\right)=o_{n}(1)$, we can find $c>0$ such that

$$
\int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x+\int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x \leq c .
$$

According to [19, Lemma 2.10], we obtain the limit (i). On the other hand, from (i), we obtain

$$
\int_{B_{1}} f\left(x, u_{n}\right) d x \rightarrow \int_{B_{1}} f(x, u) d x
$$

Therefore, there exists $p \in L^{1}\left(B_{1}\right)$ such that

$$
\begin{equation*}
f\left(x, u_{n}\right) \leq p(x) \text { almost everywhere in } B_{1} . \tag{4.22}
\end{equation*}
$$

From $\left(H_{1}\right)$ and $\left(H_{3}\right)$, we obtain

$$
\begin{equation*}
F(x, t) \leq \max _{(x, t) \in \bar{B}_{1} \times\left[0, s_{0}\right]} F(x, t)+M f(x, t), \quad \text { for all } \quad(x, t) \in B_{1} \times \mathbb{R} . \tag{4.23}
\end{equation*}
$$

Using (4.22) and (4.23), we have

$$
\begin{equation*}
F\left(x, u_{n}\right) \leq \max _{(x, t) \in \bar{B}_{1} \times\left[0, s_{0}\right]} F(x, t)+M p(x), \quad \text { for all } \quad x \in B_{1} . \tag{4.24}
\end{equation*}
$$

Therefore, $F\left(x, u_{n}\right) \rightarrow F(x, u)$ in $L^{1}\left(B_{1}\right)$, which follows from Lebesgue's dominated convergence theorem.

Let recall that for $p>2, u_{p} \in E$ denotes the nonnegative function such that $\left\|u_{p}\right\|_{p}=1$ and

$$
\begin{equation*}
S_{p}=\inf _{0 \neq u \in E} \frac{\left(\int_{B_{1}} w(x)|\nabla u|^{2} d x\right)^{1 / 2}}{\left(\int_{B_{1}}|u|^{p} d x\right)^{1 / p}}=\left\|u_{p}\right\| . \tag{4.25}
\end{equation*}
$$

Lemma 4.5. Suppose that $f$ and $g$ satisfy $\left(H_{1}\right)-\left(H_{5}\right)$. Then, the following inequality holds:

$$
\sup _{z \in \mathbb{R}^{+}\left(u_{p}, u_{p}\right)+E^{-}} J(z)<\left(\frac{\alpha_{\gamma}^{*}}{\max \left\{\alpha_{0}, \beta_{0}\right\}}\right)^{1-\gamma} .
$$

Proof. Let $z=t\left(u_{p}, u_{p}\right)+(v,-v)$ with $t \geq 0, v \in E$ and $u_{p}$ is given by (4.25). Then,

$$
J(z)=t^{2}\left\|u_{p}\right\|^{2}-\|v\|^{2}-\int_{B_{1}} F\left(x, t u_{p}+v\right) d x-\int_{B_{1}} G\left(x, t u_{p}-v\right) d x .
$$

Using condition $\left(H_{5}\right)$, we have

$$
\left.J(z) \leq t^{2}\left\|u_{p}\right\|^{2}-\left.\frac{C_{p}}{p} \int_{B_{1}}\left(\left|t u_{p}+v\right|^{p}+\mid t u_{p}-v\right)\right|^{p}\right) d x \leq t^{2}\left\|u_{p}\right\|^{2}-\frac{2 C_{p} t^{p}}{p} \int_{B_{1}}\left|u_{p}\right|^{p} d x .
$$

Since $\left\|u_{p}\right\|=S_{p}$ and $\left\|u_{p}\right\|_{p}=1$, we obtain

$$
\sup _{z \in \mathbb{R}^{+}\left(u_{p}, u_{p}\right)+E^{-}} J(z) \leq \max _{t \geq 0}\left\{t^{2} S_{p}^{2}-\frac{2 C_{p} t^{p}}{p}\right\} .
$$

Since the function $\lambda(t)=t^{2} S_{p}^{2}-\frac{2 C_{p} t^{p}}{p}$ achieves its maximum on $t_{0}=\frac{S_{p}^{2 /(p-2)}}{C_{p}^{1 /(p-2)}}$ and using the estimate of $C_{p}$, we have

$$
\sup _{z \in \mathbb{R}^{+}\left(u_{p}, u_{p}\right)+E^{-}} J(z)=\frac{(p-2) S_{p}^{2 p /(p-2)}}{p C_{p}^{2 /(p-2)}}<\left(\frac{\alpha_{\gamma}^{*}}{\max \left\{\alpha_{0}, \beta_{0}\right\}}\right)^{1-\gamma} .
$$

Remark 4.6. By Lemma 4.5, there exists $\delta>0$ such that

$$
c_{n} \leq \max _{Q_{n}} J(z) \leq \sup _{\mathbb{R}^{+}\left(u_{p}, u_{p}\right) \oplus E_{n}^{-}} J(z) \leq \sup _{\mathbb{R}(e, e) \oplus E^{-}} J(z) \leq\left(\frac{\alpha^{*}}{\max \left\{\alpha_{0}, \beta_{0}\right\}}\right)^{1-\gamma}-\delta,
$$

for every $n \in \mathbb{N}$.

## 5. Proof of the Theorem 1.1

Let $\left(u_{n}, v_{n}\right) \in H_{n}$ be the sequence given by Proposition 4.1. From Lemma 4.3, this sequence is bounded in $E$. Thus, up to a subsequence, we can assume that $(u, v) \in E$ such that $\left(u_{n}, v_{n}\right) \rightharpoonup(u, v)$ weakly in $E$, for some $(u, v) \in E$. Taking $(0, \psi)$ and $(\phi, 0)$ in (4.7), where $\phi$ and $\psi$ belongs to $C_{0, \text { rad }}^{\infty}\left(B_{1}\right) \cap$ $H_{n}$. Therefore,

$$
\begin{equation*}
\int_{B_{1}} w(x) \nabla u_{n} \nabla \psi d x=\int_{B_{1}} g\left(x, v_{n}\right) \psi d x \tag{5.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{B_{1}} w(x) \nabla v_{n} \nabla \phi d x=\int_{B_{1}} f\left(x, u_{n}\right) \phi d x . \tag{5.2}
\end{equation*}
$$

Taking the limit in (5.1) and (5.2) as $n \rightarrow \infty$, by Lemma 4.4 and the density $C_{0, \text { rad }}^{\infty}\left(B_{1}\right) \cap\left(\bigcup_{n \in \mathbb{N}} H_{n}\right)$ in $H_{0, \mathrm{rad}}^{1}\left(B_{1}, \omega\right)$, we obtain

$$
\begin{equation*}
\int_{B_{1}} w(x) \nabla u \nabla \psi d x=\int_{B_{1}} g(x, v) \psi d x, \text { for all } \psi \in H_{0, \text { rad }}^{1}\left(B_{1}, \omega\right) \tag{5.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{B_{1}} w(x) \nabla v \nabla \phi d x=\int_{B_{1}} f(x, u) \phi d x, \text { for all } \phi \in H_{0, \text { rad }}^{1}\left(B_{1}, \omega\right) . \tag{5.4}
\end{equation*}
$$

Therefore, $(u, v) \in E$ is a critical point of $J$. Now, we prove that $(u, v)$ is nontrivial. Since the system (1.8) is strongly coupled, if we assume that $u \equiv 0$ we get that $v \equiv 0$. Therefore, by Lemma 2.2, up to a subsequence, we have

$$
\begin{equation*}
u_{n} \rightarrow 0 \text { and } v_{n} \rightarrow 0 \text { in } L^{p}\left(B_{1}\right), \text { for all } p \geq 1 \tag{5.5}
\end{equation*}
$$

and

$$
u_{n} \rightarrow 0 \text { and } v_{n} \rightarrow 0 \text { almost everywhere in } \mathbb{R}^{2} .
$$

If we suppose that $\left\|u_{n}\right\|$ is not bounded below by a positive constant, we can get a subsequence of $\left(u_{n}\right)$ such that $\left\|u_{n}\right\| \rightarrow 0$. Therefore,

$$
\begin{equation*}
\int_{B_{1}} w(x) \nabla u_{n} \nabla v_{n} d x \rightarrow 0 \tag{5.6}
\end{equation*}
$$

Considering the pairs of functions $(\phi, \psi)=\left(u_{n}, 0\right)$ and $(\phi, \psi)=\left(0, v_{n}\right)$ in (4.7), we have

$$
\begin{equation*}
\int_{B_{1}} w(x) \nabla u_{n} \nabla v_{n} d x=\int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x=\int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x . \tag{5.7}
\end{equation*}
$$

Using Lemma 4.4 and (5.5), we have

$$
\begin{equation*}
\int_{B_{1}} F\left(x, u_{n}\right) d x \rightarrow 0 \text { and } \int_{B_{1}} G\left(x, v_{n}\right) d x \rightarrow 0, \quad \text { as } \quad n \rightarrow+\infty \tag{5.8}
\end{equation*}
$$

Thus, by the above limits, we get that $J\left(u_{n}, v_{n}\right)$ tends to zero which contradicts (4.5); consequently, $\left\|u_{n}\right\|$ is bounded below by a positive constant, in particular, we can assume that $\left\|u_{n}\right\| \neq 0$ for all $n \in \mathbb{N}$. Now, taking $(\phi, \psi)=\left(0, u_{n}\right)$ in (4.7), we get

$$
\begin{equation*}
\left\|u_{n}\right\|^{2}=\int_{B_{1}} g\left(x, v_{n}\right) u_{n} d x \tag{5.9}
\end{equation*}
$$

Thus,

$$
\begin{equation*}
\left\|u_{n}\right\| \leq \int_{B_{R}} g\left(x, v_{n}\right) \frac{u_{n}}{\left\|u_{n}\right\|} d x . \tag{5.10}
\end{equation*}
$$

We assume that $\max \left\{\alpha_{0}, \beta_{0}\right\}=\alpha_{0}$. Then, we can write

$$
\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left\|u_{n}\right\| \leq \int_{B_{1}}\left|g\left(x, v_{n}\right) \| \bar{u}_{n}\right| d x
$$

where

$$
\begin{equation*}
\bar{u}_{n}=\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}} \frac{u_{n}}{\left\|u_{n}\right\|} . \tag{5.11}
\end{equation*}
$$

Applying Lemma 4.2 with

$$
s=\frac{\left|g\left(x, v_{n}(x)\right)\right|}{\alpha_{0}^{\frac{1-\gamma}{2}}}, \quad t=\alpha_{0}^{\frac{1-\gamma}{2}}\left|\bar{u}_{n}(x)\right|, \quad r=\frac{2}{1-\gamma} \quad \text { and } \quad r^{\prime}=\frac{2}{1+\gamma},
$$

we have

$$
\begin{align*}
& \left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left\|u_{n}\right\| \leq\left[\int_{B_{1}}\left(e^{\left.\alpha_{0} \mid \bar{u}_{n}\right)^{\frac{2}{1-\gamma}}}-1\right) d x+\frac{1+\gamma}{2} \int_{\left\{x \in B_{1}: \frac{\mid\left(x, v_{n}(x)\right.}{\frac{1}{1-\gamma}} \leq e^{\left.\frac{1-\gamma}{2+\gamma}\right\}}\right.} \frac{\left|g\left(x, v_{n}\right)\right|^{\frac{2}{1+\gamma}}}{\alpha_{0}^{\frac{1-\gamma}{1+\gamma}}} d x\right. \\
& \left.+\int_{\left\{x \in B_{1} \frac{\mid \underline{g}\left(x, v_{n}(x)\right.}{\frac{1-y}{2}} \alpha_{0}^{2}\right.} e^{\left.\frac{1-\gamma}{1+\gamma}\right\}} \frac{\left|g\left(x, v_{n}\right)\right|}{\alpha_{0}^{\frac{1-\gamma}{2}}}\left(\ln \frac{\left|g\left(x, v_{n}\right)\right|}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right)^{\frac{1-y}{2}} d x\right] . \tag{5.12}
\end{align*}
$$

By Lemma 2.4 and (5.11) the first integral tends to zero, using dominated dominated theorem and the fact that $v_{n} \rightarrow 0$ almost everywhere in $B_{1}$ the second integral tends to zero. Hence,

$$
\begin{equation*}
\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left\|u_{n}\right\| \leq \int_{B_{1}} \frac{\left|g\left(x, v_{n}\right)\right|}{\alpha_{0}^{\frac{1-\gamma}{2}}}\left(\ln \frac{\left|g\left(x, v_{n}\right)\right|}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right)^{\frac{1-\gamma}{2}} d x+o_{n}(1) \tag{5.13}
\end{equation*}
$$

Given $\epsilon \in\left(0, \frac{\alpha_{0} \delta}{4\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)}\right)$, where $\delta>0$ is given by Remark 4.6. By $\left(H_{4}\right)$ and the assumption $\alpha_{0} \geq \beta_{0}$, we can find $C_{\epsilon}>0$ such that

$$
|g(x, s)| \leq C_{\epsilon} e^{\left(\alpha_{0}+\epsilon\right)|s|^{\frac{2}{1-\gamma}}}, \quad \text { for all } \quad(x, s) \in B_{1} \times \mathbb{R}
$$

Replacing the above inequality in (5.13), we get

$$
\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left\|u_{n}\right\| \leq \frac{1}{\alpha_{0}^{\frac{1-\gamma}{2}}} \int_{B_{1}}\left|g\left(x, v_{n}\right)\right|\left(\ln \frac{C_{\epsilon} e^{\left(\alpha_{0}+\epsilon\left|v_{n}\right|^{\frac{2}{-\gamma}}\right.}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right)^{\frac{1-\gamma}{2}} d x+o_{n}(1) .
$$

Thus,

$$
\begin{equation*}
\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left\|u_{n}\right\| \leq \frac{1}{\alpha_{0}^{\frac{1-\gamma}{2}}} \int_{B_{1}}\left|g\left(x, v_{n}\right)\right|\left[\ln \frac{1-\gamma}{2}\left(\frac{C_{\epsilon}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right)+\left(\alpha_{0}+\epsilon\right)^{\frac{1-\gamma}{2}}\left|v_{n}\right|\right]+o_{n}(1) . \tag{5.14}
\end{equation*}
$$

Let $I_{n}=\int_{B_{1}}\left|g\left(x, v_{n}\right)\right|\left[\ln \ln ^{\frac{1-\gamma}{2}}\left(\frac{C_{\epsilon}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right)+\left(\alpha_{0}+\epsilon\right)^{\frac{1-\gamma}{2}}\left|v_{n}\right|\right]$ and set

$$
Y_{n}:=\left\{x \in B_{1}: \ln \frac{1-\gamma}{2}\left(\frac{C_{\epsilon}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right) \leq\left(\left(\alpha_{0}+2 \epsilon\right)^{\frac{1-\gamma}{2}}-\left(\alpha_{0}+\epsilon\right)^{\frac{1-\gamma}{2}}\right)\left|v_{n}\right|\right\} .
$$

Hence,

$$
\begin{aligned}
I_{n}= & \ln ^{\frac{1-\gamma}{2}}\left(\frac{C_{\epsilon}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right) \int_{B_{1} \backslash Y_{n}}\left|g\left(x, v_{n}\right)\right| d x+\left(\alpha_{0}+\epsilon\right)^{\frac{1-\gamma}{2}} \int_{B_{1} \backslash Y_{n}} g\left(x, v_{n}\right) v_{n} d x \\
& +\int_{Y_{n}}\left|g\left(x, v_{n}\right)\right|\left[\ln \frac{\frac{1-\gamma}{2}}{}\left(\frac{C_{\epsilon}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right)+\left(\alpha_{0}+\epsilon\right)^{\frac{1-\gamma}{2}}\left|v_{n}\right|\right] d x \\
\leq & \ln ^{\frac{1-\gamma}{2}}\left(\frac{C_{\epsilon}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right) \int_{B_{1} \backslash Y_{n}}\left|g\left(x, v_{n}\right)\right| d x+\left(\alpha_{0}+\epsilon\right)^{\frac{1-\gamma}{2}} \int_{B_{1} \backslash Y_{n}} g\left(x, v_{n}\right) v_{n} d x+\left(\alpha_{0}+2 \epsilon\right)^{\frac{1-\gamma}{2}} \int_{Y_{n}} g\left(x, v_{n}\right) v_{n} d x .
\end{aligned}
$$

Then,

$$
\begin{equation*}
I_{n} \leq \ln ^{\frac{1-\gamma}{2}}\left(\frac{C_{\epsilon}}{\alpha_{0}^{\frac{1-\gamma}{2}}}\right) \int_{B_{\backslash} \backslash Y_{n}}\left|g\left(x, v_{n}\right)\right| d x+\left(\alpha_{0}+2 \epsilon\right)^{\frac{1-y}{2}} \int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x \tag{5.15}
\end{equation*}
$$

Since $v_{n} \rightarrow 0$ almost everywhere in $B_{1}$ and $g$ is bounded in $B_{1} \backslash Y_{n}$ for all $n \in \mathbb{N}$ (being independent of $n$ ), by the dominated convergence theorem, we get

$$
\begin{equation*}
\int_{B_{1} \backslash Y_{n}}\left|g\left(x, v_{n}\right)\right| d x=o_{n}(1) \tag{5.16}
\end{equation*}
$$

Using (5.14)-(5.16), we have

$$
\begin{equation*}
\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left\|u_{n}\right\| \leq\left(1+\frac{2 \epsilon}{\alpha_{0}}\right)^{\frac{1-\gamma}{2}} \int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x+o_{n}(1) \tag{5.17}
\end{equation*}
$$

Arguing similarly, we get

$$
\begin{equation*}
\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left\|v_{n}\right\| \leq\left(1+\frac{2 \epsilon}{\alpha_{0}}\right)^{\frac{1-\gamma}{2}} \int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x+o_{n}(1) \tag{5.18}
\end{equation*}
$$

On the other hand, using Proposition 4.1, Remark 4.6 and (5.8), we obtain

$$
\begin{equation*}
\left|\int_{B_{1}} w(x) \nabla u_{n} \nabla v_{n} d x\right| \leq o_{n}(1)+\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{1-\gamma} . \tag{5.19}
\end{equation*}
$$

Since $J_{n}^{\prime}\left(u_{n}, v_{n}\right)\left(u_{n}, v_{n}\right)=0$, we get

$$
\begin{equation*}
\int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x+\int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x=2\left|\int_{B_{1}} w(x) \nabla u_{n} \nabla v_{n} d x\right| . \tag{5.20}
\end{equation*}
$$

By (5.19) and (5.20), we find

$$
\begin{equation*}
\int_{B_{R}} f\left(x, u_{n}\right) u_{n} d x+\int_{B_{R}} g\left(x, v_{n}\right) v_{n} d x \leq 2\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{1-\gamma}+o_{n}(1) . \tag{5.21}
\end{equation*}
$$

Combining (5.17), (5.18) and (5.21), we obtain

$$
\begin{aligned}
\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}\left(\left\|u_{n}\right\|+\left\|v_{n}\right\|\right) & \leq\left(1+\frac{2 \epsilon}{\alpha_{0}}\right)^{\frac{1-\gamma}{2}}\left(\int_{B_{R}} f\left(x, u_{n}\right) u_{n} d x+\int_{B_{1}} g\left(x, v_{n}\right) v_{n} d x\right)+o_{n}(1) \\
& \leq 2\left(1+\frac{2 \epsilon}{\alpha_{0}}\right)^{\frac{1-\gamma}{2}}\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{1-\gamma}+o_{n}(1) .
\end{aligned}
$$

According to the election of $\epsilon$, for every $n \in \mathbb{N}$, we obtain

$$
\left\|u_{n}\right\|+\left\|v_{n}\right\| \leq 2\left(1+\frac{2 \epsilon}{\alpha_{0}}\right)^{\frac{1-\gamma}{2}}\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\delta\right)^{\frac{1-\gamma}{2}}+o_{n}(1) \leq 2\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\frac{\delta}{2}\right)^{\frac{1-\gamma}{2}}+o_{n}(1) .
$$

Thus, there exists $n_{0} \in \mathbb{N}$ such that

$$
\left\|u_{n}\right\|+\left\|v_{n}\right\| \leq 2\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\frac{\delta}{4}\right)^{\frac{1-\gamma}{2}}, \quad \text { for all } n \geq n_{0}
$$

Moreover, we can suppose that

$$
\left\|u_{n}\right\| \leq\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\frac{\delta}{4}\right)^{\frac{1-\gamma}{2}}, \quad \text { for all } n \geq n_{0}
$$

Taking $\xi>0$ such that $\left(\alpha_{0}+\xi\right)\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\frac{\delta}{4}\right)<\alpha_{\gamma}^{*}$, by $\left(H_{4}\right)$ there exists $c>0$ such that

$$
|f(x, s)| \leq c e^{\left(\alpha_{0}+\xi\right)|s|^{1-\gamma}}, \quad \text { for all } \quad(x, s) \in B_{1} \times \mathbb{R} .
$$

Let $p>1$ close enough to 1 satisfying $p\left(\alpha_{0}+\xi\right)\left(\frac{\alpha_{\gamma}^{*}}{\alpha_{0}}-\frac{\delta}{4}\right)<\alpha_{\gamma}^{*}$. By Proposition 2.3 and the Hölder inequality, we have

$$
\begin{aligned}
\int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x & \leq c \int_{B_{1}}\left|u_{n}\right| e^{\left(\alpha_{0}+\xi\right)\left|u_{n}\right| 1-\frac{2}{1-\gamma}} d x \\
& \leq c\left\|u_{n}\right\|_{p^{\prime}}\left(\int_{B_{1}} e^{\left.p\left(\alpha_{0}+\xi\right)\left|u_{n}\right|\right|^{\frac{2}{1-\gamma}}} d x\right)^{1 / p} \\
& \leq c\left\|u_{n}\right\|_{p^{\prime}}\left(\int_{B_{1}} e^{\left.p\left(\alpha_{0}+\xi\right)\left(\frac{\alpha_{\bar{\gamma}}^{*}}{\alpha_{0}}-\frac{\delta}{4}\right)\left(\frac{\left|u_{n}\right|}{\left|u_{n}\right|}\right)\right)^{\frac{2}{1-\gamma}}} d x\right)^{1 / p} \\
& \leq c\left\|u_{n}\right\|_{p^{\prime}} .
\end{aligned}
$$

Applying (5.5), we obtain

$$
\int_{B_{1}} f\left(x, u_{n}\right) u_{n} d x \rightarrow 0, \quad \text { as } n \rightarrow+\infty .
$$

Therefore, using (5.7) and (5.8), one has

$$
\lim _{n \rightarrow+\infty} J\left(u_{n}, v_{n}\right)=0
$$

which represents a contradiction with the fact that $J\left(z_{n}\right) \geq \sigma$ for all $n \geq 1$. Therefore, $(u, v)$ is a nontrivial weak solution. This complete the proof.

## 6. Conclusions

In this work, we apply variational methods to find a nontrivial solution for a Hamiltonian systems where the nonlinearities possess maximal growth related to Trudinger-Moser type inequalities. To the best of our knowledge, this is the first result to demonstrate the existence of nontrivial solutions for a Hamiltonian involving supercritical exponential growth in the sense of the exponential critical hyperbola in the literature. According to our definition of logarithmic weight, we restricted the domain to the unit ball. It is of interest to further our results to solutions for Hamiltonian systems involving supercritical exponential growth on the whole space $\mathbb{R}^{2}$.

## Use of AI tools declaration

The authors declare they have not used Artificial Intelligence (AI) tools in the creation of this article.

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## Conflict of interest

The author declares no conflicts of interest.

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