Mathematics

## Research article

# Investigating stochastic solutions for fourth order dispersive NLSE with quantic nonlinearity 

Yazid Alhojilan ${ }^{1}$ and Islam Samir ${ }^{2, *}$<br>${ }^{1}$ Department of Mathematics, College of Science, Qassim University, P.O. Box 6644, Buraydah 51452, Saudi Arabia<br>${ }^{2}$ Department of Physics and Mathematics Engineering, Faculty of Engineering, Ain Shams University, Cairo, Egypt<br>* Correspondence: Email: islamsamir@eng.asu.edu.eg.


#### Abstract

In this paper, the stochastic fourth order nonlinear Schrödinger equation with quantic nonlinearity and affected by multiplicative noise is considered. This model is used to mimic the wave propagation through optical fibers. The improved modified extended tanh method is used to extract optical solutions for the investigated model. Various types of stochastic solutions are provided such as bright soliton, dark soliton, singular soliton, singular periodic solution and Weierstrass elliptic solution. Moreover, Matlab software packages are used to introduce the effect of the multiplicative noise on the raised solutions. The noise intensity is varied to show the robust of the extracted solutions against the noise.


Keywords: quantic nonlinearity; nonlinear equations; partial differential equations; stochastic NLSE; multiplicative noise
Mathematics Subject Classification: 35C08, 35C09, 35C07

## 1. Introduction

Nonlinear partial differential equations (NLPDEs) can be used to mimic many nonlinear phenomena that arise in a various scientific fields, including geology,fluid dynamics, biology, chemical physics, Optics, plasma physics and solid state physics [1-9]. Many researchers have recently become interested in studying the dynamics of these models [10-13]. One of the basic models of nonlinear waves is the nonlinear Schrodinger equation (NLSE) [14,15]. It is used in a variety of fields, including electromechanical systems [16], laser beams [17], and the theory of crystals [18] and solids [19].

General random disturbances Spontaneous emissions and Thermal fluctuations can be modeled using stochastic Schrodinger models. Many authors investigated the uniqueness and existence of
stochastic NLSEs with multiplicative or additive noise. The case of additive noise is investigated in [20,21] while the case of multiplicative noise if reported in [22,23]. Both types of noises are investigated in $[24,25]$. In addition, Numerical methods are used to study the stochastic NLSE in [26,27].

In this context, the stochastic Fourth order NLSE with quantic nonlinearity and affected by multiplicative noise. This stochastic model has not been studied before. This model reads as:

$$
\begin{equation*}
i\left(H_{t}+\alpha_{1} H_{x}+\alpha_{3} H_{x x x}\right)+\alpha_{2} H_{x x}+\alpha_{4} H_{x x x x}+\beta_{1}|H|^{2} H+\beta_{2}|H|^{4} H+\sigma H B_{t}=0 \tag{1.1}
\end{equation*}
$$

where $H_{t}$ represents the linear temporal evolution. $H_{x}$ represents the inter-modal dispersion. $H_{x x x}$ and $H_{x x x x}$ represent the third order and fourth order dispersion. $H_{x x}$ represents the group velocity dispersion. $\beta_{1}$ and $\beta_{2}$ represent the coefficients of self phase modulation effect and quantic nonlinearity. $B_{t}$ is the stochastic term which represents the white noise. $B(t)$ and follow the next conditions:
(i) For $t \geq 0, B(t)$ has continuous trajectories.
(ii) For $s<t, B(t)-B(s)$ has independent increments.
(iii) $B(t)-B(s)$ has a normal distribution with variance $=t-s$ and mean $=0$.

This model is used to describe the wave propagation through optical fibers. This model includes the dispersion and nonlinearity effects. Dispersion is also known as wave broadening or wave spreading. Some of the optical power in a propagated wave is delayed at the fiber's output end, which results in dispersion [28]. Dispersion shows how optical power that enters the fiber simultaneously leaves at various moments. Because of dispersion, the optical wave width grows continually across the fiber. The output signal varies from the input signal when dispersion reaches the data rate's maximum value [29]. In addition, The intensity dependence of the medium's refractive index causes nonlinear effects in optical fiber. Optical soliton is a special type of solutions that can propagate to very large distance retaining its shape and speed. This optical wave is generated due to dedicated balance which occurs between the dispersion and nonlinear effects [30,31]. In addition, adding the noise effect to the investigated model is very necessary to show the robust of the extracted solutions against the noise. Research into all these different influences is making a huge contribution to the development of the telecommunications industry [29].

The nonlinear stochastic Chiral Schrödinger equation (CNLSE) in two dimensions was investigated in literature [32] to extract periodic envelopes, explosive, dissipative, symmetric solitons, and blow up waves. It was confirmed that the noise factor is dominant on all the wave conversion, growing and damping of envelopes and shocks. In ref [33], the stochastic NLSE with group velocity dispersion and self modulation nonlinear effect was studied by applying the unified solver approach to extract rational, dissipative, explosive, envelope, periodic, and localized soliton solutions.

In this current work, the improved modified extended tanh function scheme is utilized to provide optical solitons for the stochastic NLSE in (1.1). In addition, this method provide other solutions such as trigonometric, hyperbolic and exponential type solutions. The proposed method is briefly introduced in Section 2. Then, the stochastic solutions are extracted for the investigated model in Section 3. The effect of the multiplicative noise on the obtained solutions is discussed in Section 4. Finally, the work is concluded.

## 2. Revisitation of the technique

This section introduces the improved modified extended tanh scheme [34].

Considering the following NLPDE:

$$
\begin{equation*}
G\left(H, H_{t}, H_{x}, H_{x x}, H_{x t}, \ldots .\right)=0, \tag{2.1}
\end{equation*}
$$

where $G$ is polynomial in $H(x, t)$ and its partial derivatives
Step 1: The following mathematical transformation is applied

$$
\begin{equation*}
H(x, t)=F(\xi), \xi=x-\lambda t . \tag{2.2}
\end{equation*}
$$

Therefore, Eq (2.1) is converted to the following ODE:

$$
\begin{equation*}
F\left(F, F^{\prime}, F^{\prime \prime}, F^{\prime \prime \prime}, \ldots .\right)=0 . \tag{2.3}
\end{equation*}
$$

Step 2: The solutions for Eq (2.3) is supposed in the form:

$$
\begin{equation*}
F(\xi)=a_{0}+\sum_{r=1}^{N} a_{r} A^{r}(\xi)+b_{r} A^{-r}(\xi) \tag{2.4}
\end{equation*}
$$

where $A(\xi)$ holds:

$$
\begin{equation*}
A^{\prime}(\xi)=\sqrt{d_{0}+d_{1} A(\xi)+d_{2} A^{2}(\xi)+d_{3} A^{3}(\xi)+d_{4} A^{4}(\xi)} \tag{2.5}
\end{equation*}
$$

By changing the values of the constants $d_{0}, d_{1}, d_{2}, d_{3}$ and $d_{4}$, one can obtain different general solutions for Eq (2.5).
Step 3: $N$ that raised in Eq (2.4) can be evaluated by applying the balancing rule on Eq (2.3).
Step 4: Substituting by Eq (2.4) with (2.5) into (2.3), a polynomial in $A(\xi)$ is provided. Collecting the terms which have the same power together and equating them to zero, one get a system of nonlinear equations that can be solved using packages of Mathematica to determine $\lambda, a_{k}$ and $b_{k}$.

## 3. Results

In this section, the improved modified extended tanh method is applied to provide stochastic solutions for the proposed model in the following form:

$$
\begin{equation*}
H(x, t)=F(\xi) e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)}, \xi=x-\lambda t . \tag{3.1}
\end{equation*}
$$

Substituting by (3.1) into (1.1) provides

$$
\begin{equation*}
F^{(3)}\left(\alpha_{3}-4 \alpha_{4} k\right)+F^{\prime}\left(\alpha_{1}+4 \alpha_{4} k^{3}-3 \alpha_{3} k^{2}-2 \alpha_{2} k-\lambda\right)=0, \tag{3.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\beta_{2} F^{5}+\alpha_{4} F^{(4)}+\beta_{1} F^{3}+F^{\prime \prime}\left(\alpha_{2}+6 \alpha_{4} k^{2}\right)+F\left(-3 \alpha_{4} k^{4}-\alpha_{2} k^{2}+\alpha_{1} k+\sigma^{2}-w\right)=0 \tag{3.3}
\end{equation*}
$$

where Eq (3.2) represents the imaginary part while Eq (3.3) is the real part.
Equating the coefficients of (3.2) to zero gives

$$
\alpha_{3}=4 \alpha_{4} k,
$$

$$
\begin{equation*}
\lambda=\alpha_{1}-8 \alpha_{4} k^{3}-2 \alpha_{2} k . \tag{3.4}
\end{equation*}
$$

Balancing $F^{(4)}$ with $F^{5}$, we get $N=1$. Then, the solution of $\mathrm{Eq}(3.3)$ can be represented as follow:

$$
\begin{equation*}
F(\xi)=a_{0}+a_{1} A(\xi)+\frac{b_{1}}{A(\xi)} \tag{3.5}
\end{equation*}
$$

By employing Step 4 which discussed in the last section, we can obtain the next results for (1.1) :
Case 1. $d_{0}=d_{1}=d_{3}=0$

## Result (1)

$$
\begin{aligned}
& a_{0}=0, b_{1}=0, d_{4}=\frac{a_{1}^{2} \sqrt{-\beta_{2}}}{2 \sqrt{6} \sqrt{\alpha_{4}}}, \alpha_{1}=\frac{\frac{\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}\left(k^{2}-d_{2}\right)}{\sqrt{\beta_{2}}}+\alpha_{4}\left(-10 d_{2} k^{2}+9 d_{2}^{2}-3 k^{4}\right)-\sigma^{2}+w}{k}, \\
& d_{2}=-\frac{-\frac{\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}}{\sqrt{\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{10 \alpha_{4}} .
\end{aligned}
$$

Then, we obtain the following stochastic bright soliton and singular periodic solutions for (1.1).

$$
\begin{gather*}
H(x, t)=\sqrt[4]{6} \sqrt{\frac{-\frac{\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}}{\sqrt{\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{5 \sqrt{\alpha_{4}} \sqrt{-\beta_{2}}}} \operatorname{sech}\left(\frac{\sqrt{-\frac{-\frac{\sqrt{6} \sqrt{-\sigma_{4} \beta_{1}}+\alpha_{2}+6 \alpha_{4} k^{2}}{\sqrt{\beta_{2}}}}{\alpha_{4}}}(x-\lambda t)}{\sqrt{10}}\right) \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)},  \tag{3.6}\\
H(x, t)=\sqrt[4]{6} \sqrt{\frac{-\frac{\sqrt{6} \sqrt{\alpha_{4} \beta_{1}}}{\sqrt{-\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{5 \sqrt{\alpha_{4}} \sqrt{-\beta_{2}}}} \sec \left(\frac{\sqrt{\frac{-\frac{\sqrt{6} \sqrt{\sigma_{4} \beta_{1}}}{\sqrt{-\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{\alpha_{4}}}(x-\lambda t)}{\sqrt{10}}\right) \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)} . \tag{3.7}
\end{gather*}
$$

Case 2. $d_{1}=d_{3}=0$

## Result (1)

$$
\begin{aligned}
& a_{0}=0, b_{1}=0, d_{4}=\frac{a_{1}^{2} \sqrt{-\beta_{2}}}{2 \sqrt{6} \sqrt{\alpha_{4}}}, d_{2}=-\frac{-\frac{\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}}{\sqrt{\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{10 \alpha_{4}}, d_{0}=\frac{d_{2}^{2}}{4 d_{4}} \\
& \alpha_{1}=\frac{\sqrt{6} \sqrt{\alpha_{4}} \sqrt{-\beta_{2}} \beta_{1}\left(k^{2}-d_{2}\right)+\alpha_{4} \beta_{2}\left(-10 d_{2} k^{2}+6 d_{2}^{2}-3 k^{4}\right)+\beta_{2}\left(w-\sigma^{2}\right)}{\beta_{2} k}
\end{aligned}
$$

Then, we obtain the following stochastic dark soliton and singular periodic solutions for (1.1).

$$
H(x, t)=\sqrt[4]{\frac{3}{2}} \sqrt{\frac{-\frac{\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}}{\sqrt{\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{5 \sqrt{\alpha_{4}} \sqrt{-\beta_{2}}}} \tanh \left(\frac{\sqrt{\frac{-\frac{\sqrt{6} \sqrt{-\alpha_{4} \beta_{1}}}{\sqrt{\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{\alpha_{4}}}(x-\lambda t)}{2 \sqrt{5}}\right) \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)},
$$

$$
\begin{equation*}
H(x, t)=\sqrt[4]{\frac{3}{2}} \sqrt{-\frac{-\frac{\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}}{\sqrt{\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{5 \sqrt{\alpha_{4}} \sqrt{-\beta_{2}}}} \tan \left(\frac{\sqrt{-\frac{-\frac{\sqrt{6} \sqrt{-\sigma_{4}} \beta_{1}}{\sqrt{\beta_{2}}}+\alpha_{2}+6 \alpha_{4} k^{2}}{\alpha_{4}}}(x-\lambda t)}{2 \sqrt{5}}\right) \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)} \tag{3.8}
\end{equation*}
$$

## Result (2)

$$
\begin{aligned}
& a_{0}=0, a_{1}=0, \alpha_{1}=\frac{6 \beta_{1} b_{1}^{2} d_{2}^{2} d_{4}\left(d_{2}-k^{2}\right)+2 \beta_{2} b_{1}^{4} d_{4}^{2}\left(10 d_{2} k^{2}-6 d_{2}^{2}+3 k^{4}\right)+3 d_{2}^{4}\left(w-\sigma^{2}\right)}{3 d_{2}^{4} k}, \\
& d_{4}=\frac{\sqrt{3} \sqrt{\alpha_{4}} d_{2}^{2}}{b_{1}^{2} \sqrt{-2 \beta_{2}}}, d_{2}=\frac{-\alpha_{2} \sqrt{-\beta_{2}}-\sqrt{6} \sqrt{\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{-\beta_{2}} k^{2}}{10 \alpha_{4} \sqrt{-\beta_{2}}}, d_{0}=\frac{d_{2}^{2}}{4 d_{4}}
\end{aligned}
$$

Then, we obtain the following stochastic singular soliton and singular periodic solutions for (1.1).

$$
\begin{align*}
& H(x, t)=\sqrt[4]{\frac{3}{2}} \sqrt{\frac{-\alpha_{2} \sqrt{-\beta_{2}}-\sqrt{6} \sqrt{\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{-\beta_{2}} k^{2}}{5 \sqrt{\alpha_{4}} \beta_{2}}} \operatorname{coth}\left(\frac{\sqrt{-\frac{-\alpha_{2} \sqrt{-\beta_{2}}-\sqrt{6} \sqrt{\sigma_{4} \beta_{1}-6 \alpha_{4} \sqrt{-\beta_{2}} k^{2}}}{\alpha_{4} \sqrt{-\beta_{2}}}} x^{2 \sqrt{5}}(x-\lambda t)}{\times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)}}\right) \\
& H(x, t)=\sqrt[4]{\frac{3}{2}} \sqrt{-\frac{-\alpha_{2} \sqrt{-\beta_{2}}-\sqrt{6} \sqrt{\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{-\beta_{2}} k^{2}}{5 \sqrt{\alpha_{4}} \beta_{2}}} \cot \left(\frac{\sqrt{\frac{-\alpha_{2} \sqrt{-\beta_{2}-\sqrt{6} \sqrt{\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{-\beta_{2}} k^{2}}}{\alpha_{4} \sqrt{-\beta_{2}}}}(x-\lambda t)}{2 \sqrt{5}}\right)  \tag{3.10}\\
& \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)} .
\end{align*}
$$

## Result (3)

$$
\begin{aligned}
& a_{0}=0, d_{4}=\frac{a_{1} d_{2}}{2 b_{1}}, \alpha_{1}=\frac{-\frac{6 a_{1} \beta_{1} b_{1}\left(2 d_{2}+k^{2}\right)}{d_{2}}+\frac{a_{1}^{2} \beta_{2} b_{1}^{2}\left(-20 d_{2} k^{2}-36 d_{2}^{2}+3 k^{4}\right)}{d_{2}^{2}}+6\left(w-\sigma^{2}\right)}{6 k}, \\
& b_{1}=-\frac{\sqrt{\frac{3}{2}}\left(\alpha_{2} \sqrt{-\beta_{2}}+\sqrt{6} \sqrt{\alpha_{4}} \beta_{1}+6 \alpha_{4} \sqrt{-\beta_{2}} k^{2}\right)}{10 a_{1} \sqrt{\alpha_{4}} \beta_{2}}, d_{2}=\frac{a_{1} b_{1} \sqrt{-\beta_{2}}}{\sqrt{6} \sqrt{\alpha_{4}}}, d_{0}=\frac{d_{2}^{2}}{4 d_{4}}
\end{aligned}
$$

Then, we obtain the following stochastic singular soliton and singular periodic solutions for (1.1).

$$
H(x, t)=\left\{\frac{\sqrt[4]{3} \sqrt{\frac{\sqrt{6 \beta_{1}+\frac{\sqrt{-\beta_{2}}\left(a_{2}+6 \sigma_{4} k^{2}\right)}{\sqrt{\alpha_{4}}}}}{a_{1} \beta_{2}}} \tanh \left(\frac{\sqrt{-\frac{\sqrt{6} \beta_{1}}{\sqrt{\alpha_{4}} \sqrt{-\beta_{2}}-\frac{\alpha_{2}}{\alpha_{4}}-6 k^{2}}(x-\lambda t)}}{2 \sqrt{10}}\right)}{2^{3 / 4} \sqrt{5}}\right.
$$

$$
\begin{align*}
& \left.+\sqrt{-a_{1}} \operatorname{coth}\left(\frac{\sqrt{-\frac{\sqrt{6} \beta_{1}}{\sqrt{\alpha_{4}} \sqrt{-\beta_{2}}}-\frac{\alpha_{2}}{\alpha_{4}}-6 k^{2}}(x-\lambda t)}{2 \sqrt{10}}\right)\right\} \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)},  \tag{3.12}\\
& H(x, t)=\left\{\frac{\sqrt[4]{3} \sqrt{-\frac{\sqrt{6} \beta_{1}+\frac{\sqrt{-\beta_{2}}\left(\alpha_{2}+6 \alpha_{4} k^{2}\right)}{a_{1}}}{a_{1} \beta_{2}}} \tan \left(\frac{\sqrt{\frac{\sqrt{6} \beta_{1}}{\sqrt{\sigma_{4}} \sqrt{-\beta_{2}}+\frac{\alpha_{2}}{\alpha_{1}}+6 k^{2}}(x-\lambda t)}}{2 \sqrt{10}}\right)}{2^{3 / 4} \sqrt{5}}\right. \\
& \left.+\sqrt{a_{1}} \cot \left(\frac{\sqrt{\frac{\sqrt{6} \beta_{1}}{\sqrt{\alpha_{4}} \sqrt{-\beta_{2}}}+\frac{\alpha_{2}}{\alpha_{4}}+6 k^{2}}(x-\lambda t)}{2 \sqrt{10}}\right)\right\} \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)} . \tag{3.13}
\end{align*}
$$

Case 3. $d_{0}=d_{1}=0$

## Result (1)

$$
\begin{aligned}
& b_{1}=0, d_{2}=\frac{a_{0} d_{3}}{a_{1}}, \alpha_{1}=\frac{2 a_{0}^{2} a_{1}^{2} \beta_{2} k^{3}}{d_{3}^{2}}-\frac{2 a_{0} a_{1} k\left(5 a_{0}^{2} \beta_{2}+3 \beta_{1}\right)}{3 d_{3}}-\frac{a_{0}^{4} \beta_{2}+a_{0}^{2} \beta_{1}+\sigma^{2}-w}{k}, \\
& \alpha_{2}=\frac{2 a_{0} a_{1}\left(-5 a_{0}^{2} \beta_{2} d_{3}+6 a_{0} a_{1} \beta_{2} k^{2}-3 \beta_{1} d_{3}\right)}{3 d_{3}^{2}}, d_{3}=\frac{\sqrt{2} a_{0} a_{1} \sqrt{-\beta_{2}}}{\sqrt{3 \alpha_{4}}}, d_{4}=\frac{d_{3}^{2}}{4 d_{2}}
\end{aligned}
$$

Then, we obtain the following stochastic dark soliton for (1.1).

$$
\begin{equation*}
H(x, t)=\left\{a_{0}\left(\tanh \left(\frac{a_{0} \sqrt{\frac{\sqrt{-\beta-\beta_{2}}}{\sqrt{a_{4}}}}(x-\lambda t)}{2^{3 / 4} \sqrt[4]{3}}\right)+1\right)+a_{0}\right\} \times e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)} \tag{3.14}
\end{equation*}
$$

Case 4. $d_{3}=d_{4}=0$

## Result (1)

$$
\begin{aligned}
& a_{0}=0, a_{1}=0, d_{0}=-\frac{b_{1}^{2} \sqrt{-\beta_{2}}}{2 \sqrt{6} \sqrt{\alpha_{4}}}, \alpha_{1}=\frac{-\frac{\sqrt{6} \sqrt{\alpha_{4}} \beta_{1}\left(k^{2}-d_{2}\right)}{\sqrt{-\beta_{2}}}+\alpha_{4}\left(-10 d_{2} k^{2}+9 d_{2}^{2}-3 k^{4}\right)-\sigma^{2}+w}{k}, \\
& d_{2}=\frac{-\alpha_{2} \sqrt{\beta_{2}}-\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{\beta_{2}} k^{2}}{10 \alpha_{4} \sqrt{\beta_{2}}}, d_{1}=0 .
\end{aligned}
$$

Then, we obtain the following stochastic singular soliton and singular periodic solutions for Eq (1.1).

$$
H(x, t)=\sqrt[4]{6} \sqrt{\frac{\alpha_{2} \sqrt{\beta_{2}}+\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{\beta_{2}} k^{2}}{5 \sqrt{-\alpha_{4}} \beta_{2}}} \operatorname{csch}\left(\frac{\sqrt{\frac{\alpha_{2} \sqrt{\beta_{2}}+\sqrt{6} \sqrt{-\alpha_{4} \beta_{1}-6 \alpha_{4} \sqrt{\beta_{2}} k^{2}}}{\alpha_{4}}(x-\lambda t)}}{\sqrt{10}}\right)
$$

$$
\begin{align*}
& x e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)}  \tag{3.15}\\
& H(x, t)=\sqrt[4]{6} \sqrt{\frac{-\alpha_{2} \sqrt{\beta_{2}}-\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{\beta_{2}} k^{2}}{5 \sqrt{-\alpha_{4}} \beta_{2}}} \csc \left(\frac{\sqrt{-\frac{-\alpha_{2} \sqrt{\beta_{2}}-\sqrt{6} \sqrt{-\alpha_{4}} \beta_{1}-6 \alpha_{4} \sqrt{\beta_{2}} k^{2}}{\alpha_{4} \sqrt{\beta_{2}}}}(x-\lambda t)}{\sqrt{10}}\right) \\
& x e^{i\left(w t-k x-\sigma^{2} t+\sigma B(t)\right)} . \tag{3.16}
\end{align*}
$$

## Result (2)

$$
\begin{aligned}
& a_{1}=0, d_{2}=\frac{a_{0} d_{1}}{b_{1}}, \alpha_{1}=\frac{2 a_{0}^{2} b_{1}^{2} \beta_{2} k^{3}}{d_{1}^{2}}-\frac{2 a_{0} b_{1} k\left(5 a_{0}^{2} \beta_{2}+3 \beta_{1}\right)}{3 d_{1}}-\frac{a_{0}^{4} \beta_{2}+a_{0}^{2} \beta_{1}+\sigma^{2}-w}{k}, \\
& \alpha_{2}=\frac{2 a_{0} b_{1}\left(6 a_{0} b_{1} \beta_{2} k^{2}-5 a_{0}^{2} \beta_{2} d_{1}-3 \beta_{1} d_{1}\right)}{3 d_{1}^{2}}, d_{1}=\frac{\sqrt{2} a_{0} b_{1} \sqrt{-\beta_{2}}}{\sqrt{3 \alpha_{4}}}, d_{0}=\frac{d_{1}^{2}}{4 d_{2}}
\end{aligned}
$$

Then, we obtain the next stochastic exponential solution for (1.1).

Case 5. $d_{2}=d_{4}=0$

## Result (1)

$$
\begin{aligned}
& a_{1}=0, d_{0}=\frac{b_{1}^{2} \sqrt{-\beta_{2}}}{2 \sqrt{6} \sqrt{\alpha_{4}}}, d_{1}=\frac{\sqrt{\frac{2}{3}} a_{0} b_{1} \sqrt{-\beta_{2}}}{\sqrt{\alpha_{4}}}, d_{3}=-\frac{2 \sqrt{\frac{2}{3}} a_{0}^{3} \sqrt{-\beta_{2}}}{\sqrt{\alpha_{4}} b_{1}}, a_{0}=\frac{\sqrt{-\frac{6 \beta_{1}+\frac{\sqrt{6} \sqrt{-\beta_{2}\left(6 \sigma_{4} 4^{2}-\alpha_{2}\right)}}{\beta_{2}}}{\sqrt{\alpha_{4}}}}}{\sqrt{30}}, \\
& \alpha_{1}=\frac{-11 a_{0}^{4} \beta_{2}+a_{0}^{2}\left(5 \sqrt{6} \sqrt{-\alpha_{4}} \sqrt{\beta_{2}} k^{2}-3 \beta_{1}\right)-\sigma^{2}+w}{k}-3 \alpha_{4} k^{3}-\frac{\sqrt{6} \sqrt{\alpha_{4}} \beta_{1} k}{\sqrt{-\beta_{2}}} .
\end{aligned}
$$

Then, we obtain the next stochastic Weierstrass elliptic solution for (1.1).

where $g_{2}=-4 d_{1} / d_{3}$ and $g_{3}=-4 d_{0} / d_{3}$.

## 4. Influence of the noise on the extracted solutions

In this section, we investigate the effect of the multiplicative noise on the derived solutions. 2 D and 3D simulations using different values for $\sigma$ are presented using Matlab packages. In Figure 1, 3D graphical illustrations of Eq (3.6) are introduced with $\alpha_{4}=-0.13, \alpha_{2}=2, \beta_{2}=0.055, \beta_{1}=$ $-2, k=-2, \lambda=-0.095, w=1$. In Figure 2, 2D graphical illustrations of Eq (3.6) is presented using different noise intensities. In Figure 3, 3D graphical illustrations of Eq (3.8) are provided with $\alpha_{4}=0.49, \alpha_{2}=-2, \beta_{2}=-2, \beta_{1}=-2, k=-2, \lambda=-0.63, w=1$. In Figure 4, 2D graphical illustrations of Eq (3.8) is presented using different noise intensities. It is observed that as the noise intensity increases, the extracted wave begins to degrade. From Figures $1-4$, we also observed that the higher level wave is more robust to noise than the lower level wave. One can notice that the propagated wave of Eq (3.6) is fully distorted when $\sigma=5$ as shown in Figures 1 and 2 while the propagated wave of Eq (3.8) is fully distorted when $\sigma=2$ as shown in Figures 3 and 4. The signal level of the obtained solutions can be controlled to be robust against the noise by adjusting coefficients of the fourth order dispersion, group velocity dispersion, self phase modulation and the quantic nonlinearity.


Figure 1. 3D graphs of solution $H(x, t)$ of Eq (3.6).


Figure 2. 2D graphs of solution $H(x, t)$ of Eq (3.6).


Figure 3. 3D graphs of solution $H(x, t)$ of Eq (3.8).


Figure 4. 2D graphs of solution $H(x, t)$ of Eq (3.8).

## 5. Conclusions

In this work, The improved modified extended tanh function method is implemented successfully to investigate stochastic forth order NLSE with quantic nonlinearity and affected by multiplicative noise. Many exact optical solutions are extracted in terms of hyperbolic, trigonometric, exponential and Weierstrass elliptic solutions. This stochastic model with fourth-order dispersion and quantic nonlinearity has not been studied before. Therefore, all of these extracted solutions are new. These solutions are crucial in understanding some complex physical phenomena. The extracted solutions will be useful in future studies like coastal water motions, industrial studies, quasi particle theory, optical fiber and space plasma applications. Matlab packages are used to simulate the effect of the multiplicative noise on the derived solutions.

## Acknowledgments

The researchers would like to thank the Deanship of Scientific Research, Qassim University, for funding the publication of this project.

## Conflict of interest

The authors declare no conflict of interest.

## References

1. I. Samir, N. Badra, H. M. Ahmed, A. H. Arnous, A. S. Ghanem, Solitary wave solutions and other solutions for Gilson-Pickering equation by using the modified extended mapping method, Results Phys., 36 (2022), 105427. https://doi.org/10.1016/j.rinp.2022.105427
2. I. Samir, N. Badra, A. R. Seadawy, H. M. Ahmed, A. H. Arnous, Computational extracting solutions for the perturbed Gerdjikov-Ivanov equation by using improved modified extended analytical approach, J. Geom. Phys., 176 (2022), 104514. https://doi.org/10.1016/j.geomphys.2022.104514
3. A. R. Alharbi, Traveling-wave and numerical solutions to a Novikov-Veselov system via the modified mathematical methods, AIMS Math., 8 (2023), 1230-1250. https://doi.org/10.3934/math. 2023062
4. M. Sharaf, E. El-Shewy, M. Zahran, Fractional anisotropic diffusion equation in cylindrical brush model, J. Taibah Univ. Sci., 14 (2020), 1416-1420. https://doi.org/10.1080/16583655.2020.1824743
5. I. Samir, N. Badra, H. M. Ahmed, A. H. Arnous, Solitary wave solutions for generalized Boiti-Leon-Manna-Pempinelli equation by using improved simple equation method, Int. J. Appl. Comput. Math., 8 (2022), 1-12. https://doi.org/10.1007/s40819-022-01308-2
6. H. Abdelwahed, Nonlinearity contributions on critical MKP equation, J. Taibah Univ. Sci., 14 (2020), 777-782. https://doi.org/10.1080/16583655.2020.1774136
7. I. Samir, N. Badra, H. M. Ahmed, A. H. Arnous, Solitons in birefringent fibers for CGL equation with Hamiltonian perturbations and Kerr law nonlinearity using modified extended direct algebraic method, Commun. Nonlinear Sci., 102 (2021), 105945. https://doi.org/10.1016/j.cnsns.2021.105945
8. T. A. Nofal, I. Samir, N. Badra, A. Darwish, H. M. Ahmed, A. H. Arnous, Constructing new solitary wave solutions to the strain wave model in micro-structured solids, Alex. Engin. J., 61 (2022), 11879-11888. https://doi.org/10.1016/j.aej.2022.05.050
9. I. Samir, N. Badra, A. R. Seadawy, H. M. Ahmed, A. H. Arnous, Exact wave solutions of the fourth order non-linear partial differential equation of optical fiber pulses by using different methods, Optik, 2021, 166313. https://doi.org/10.1016/j.ijleo.2021.166313
10. H. X. Jia, D. W. Zuo, X. H. Li, X. S. Xiang, Breather, soliton and rogue wave of a two-component derivative nonlinear Schrödinger equation, Phys. Lett. A, 405 (2021), 127426. https://doi.org/10.1016/j.physleta.2021.127426
11. M. A. Abdelrahman, S. Hassan, M. Inc, The coupled nonlinear Schrödinger-type equations, Mod. Phys. Lett. B, 34 (2020), 2050078. https://doi.org/10.1142/S0217984920500785
12. A. R. Alharbi, A study of traveling wave structures and numerical investigation of two-dimensional Riemann problems with their stability and accuracy, CMES-Comp. Model. Eng., 134 (2023), 21932209. https://doi.org/10.32604/cmes.2022.018445
13. S. Frassu, T. Li, G. Viglialoro, Improvements and generalizations of results concerning attraction-repulsion chemotaxis models, Math. Method. Appl. Sci., 45 (2022), 11067-11078. https://doi.org/10.1002/mma. 8437
14. K. L. Geng, D. S. Mou, C. Q. Dai, Nondegenerate solitons of 2-coupled mixed derivative nonlinear Schrödinger equations, Nonlinear Dyn., 111 (2023), 603-617. https://doi.org/10.1007/s11071-022-07833-5
15. W. B. Bo, R. R. Wang, Y. Fang, Y. Y. Wang, C. Q. Dai, Prediction and dynamical evolution of multipole soliton families in fractional Schrödinger equation with the PT-symmetric potential and saturable nonlinearity, Nonlinear Dyn., 111 (2023), 1577-1588. https://doi.org/10.1007/s11071-022-07884-8
16. M. Blencowe, Quantum electromechanical systems, Phys. Rep., 395 (2004), 159-222. https://doi.org/10.1016/j.physrep.2003.12.005
17. P. Kelley, Self-focusing of optical beams, Phys. Rev. Lett., 15 (1965), 1005. https://doi.org/10.1103/PhysRevLett.15.1005
18. H. Chu, Eigen energies and eigen states of conduction electrons in pure bismuth under size and magnetic field quantizations, J. Phys. Chem. Solids, 50 (1989), 319-324. https://doi.org/10.1016/0022-3697(89)90494-0
19. N. Ashcroft, N. Mermin, Solid state physics, New York, Cengage Learning, 1976.
20. G. Falkovich, I. Kolokolov, V. Lebedev, S. Turitsyn, Statistics of soliton-bearing systems with additive noise, Phys. Rev. E, 63 (2001), 025601. https://doi.org/10.1103/PhysRevE.63.025601
21. A. Debussche, C. Odasso, Ergodicity for a weakly damped stochastic non-linear Schrödinger equation, J. Evol. Equ., 5 (2005), 317-356. https://doi.org/10.1007/s00028-005-0195-x
22. M. A. Abdelrahman, W. W. Mohammed, The impact of multiplicative noise on the solution of the chiral nonlinear Schrödinger equation, Phys. Scripta, 95 (2020), 085222. https://doi.org/10.1088/1402-4896/aba3ac
23. S. Albosaily, W. W. Mohammed, M. A. Aiyashi, M. A. Abdelrahman, Exact solutions of the (2+1)-dimensional stochastic chiral nonlinear Schrödinger equation, Symmetry, 12 (2020), 1874. https://doi.org/10.3390/sym12111874
24. K. Cheung, R. Mosincat, Stochastic nonlinear Schrödinger equations on tori, Stoch. Partial Differ., 7 (2019), 169-208. https://doi.org/10.1007/s40072-018-0125-x
25. A. Debussche, L. D. Menza, Numerical simulation of focusing stochastic nonlinear Schrödinger equations, Physica D, 162 (2002), 131-154. https://doi.org/10.1016/S0167-2789(01)00379-7
26. J. Cui, J. Hong, Z. Liu, W. Zhou, Strong convergence rate of splitting schemes for stochastic nonlinear Schrödinger equations, J. Differ. Equ., 266 (2019), 5625-5663. https://doi.org/10.1016/j.jde.2018.10.034
27. J. Cui, J. Hong, Z. Liu, Strong convergence rate of finite difference approximations for stochastic cubic Schrödinger equations, J. Differ. Equ., 263 (2017), 3687-3713. https://doi.org/10.1016/j.jde.2017.05.002
28. Y. Fang, G. Z. Wu, X. K. Wen, Y. Y. Wang, C. Q. Dai, Predicting certain vector optical solitons via the conservation-law deep-learning method, Opt. Laser Technol., 155 (2022), 108428. https://doi.org/10.1016/j.optlastec.2022.108428
29. I. Samir, A. Abd-Elmonem, H. M. Ahmed, General solitons for eighth-order dispersive nonlinear Schrödinger equation with ninth-power law nonlinearity using improved modified extended tanh method, Opt. Quant. Electron., 55 (2023), 470. https://doi.org/10.1007/s11082-023-04753-5
30. R. R. Wang, Y. Y. Wang, C. Q. Dai, Influence of higher-order nonlinear effects on optical solitons of the complex Swift-Hohenberg model in the mode-locked fiber laser, Opt. Laser Technol., 152 (2022), 108103. https://doi.org/10.1016/j.optlastec.2022.108103
31. J. J. Fang, D. S. Mou, H. C. Zhang, Y. Y. Wang, Discrete fractional soliton dynamics of the fractional Ablowitz-Ladik model, Optik, 228 (2021), 166186. https://doi.org/10.1016/j.ijleo.2020.166186
32. Y. F. Alharbi, E. K. El-Shewy, M. A. Abdelrahman, Effects of Brownian noise strength on new chiral solitary structures, J. Low Freq. Noise V. A., 2022.
33. Y. F. Alharbi, E. El-Shewy, M. A. Abdelrahman, New and effective solitary applications in Schrödinger equation via Brownian motion process with physical coefficients of fiber optics, AIMS Math., 8 (2023), 4126-4140. https://doi.org/10.3934/math. 2023205
34. Z. Yang, B. Y. Hon, An improved modified extended tanh-function method, Z. Naturforsch. A, 61 (2006), 103-115. https://doi.org/10.1515/zna-2006-3-401
© 2023 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/4.0)
