



Research article

Stability for discrete time waveform relaxation methods based on Euler schemes

Junjiang Lai and Zhencheng Fan*

College of Mathematics and Data Science, Minjiang University, Fuzhou 350108, China

* **Correspondence:** Email: fanzhencheng@yeah.net.

Abstract: Stability properties of discrete time waveform relaxation (DWR) methods based on Euler schemes are analyzed by applying them to two dissipative systems. Some sufficient conditions for stability of the considered methods are obtained; at the same time two examples of instability are given. To investigate the influence of the splitting functions and underlying numerical methods on stability of DWR methods, DWR methods based on different splittings and different numerical schemes are considered. The obtained results show that the stabilities of waveform relaxation methods based on an implicit Euler scheme are better than those based on explicit Euler scheme, and that the stabilities of waveform relaxation methods based on the classical splittings such as Gauss-Jacobi and Gauss-Seidel splittings are worse than those based on the eigenvalue splitting presented in this paper. Finally, numerical examples that confirm the theoretical results are presented.

Keywords: stability; discrete time waveform relaxation methods; Euler methods; ordinary differential equations

Mathematics Subject Classification: 34D20, 65L05, 65L20

1. Introduction

Consider the initial problem

$$\dot{x} = f(t, x), \quad t \in [0, T]; \quad x(0) = x_0. \tag{1.1}$$

Its waveform relaxation (WR) methods based on Euler methods can be constructed as follows:

(i) Taking the splitting function $F(t, x, x) = f(t, x)$, construct the iterative scheme

$$\begin{cases} \dot{x}^{(k+1)} = F(t, x^{(k+1)}, x^{(k)}), \quad t \in [0, T], \\ x^{(k+1)}(0) = x_0, \end{cases} \tag{1.2}$$

with $k \in \mathbb{Z}^+$ and $x^{(0)} \equiv x_0$, where \mathbb{Z}^+ denotes all nonnegative integers;

(ii) Discretizing (1.2) by employing the explicit and implicit Euler methods, one arrives at the expected methods

$$\begin{cases} x_{n+1}^{(k+1)} = x_n^{(k+1)} + hF(t_n, x_n^{(k+1)}, x_n^{(k)}), & n = 0, 1, \dots, N, \\ x_0^{(k+1)} = x_0, \end{cases} \quad (1.3)$$

and

$$\begin{cases} x_{n+1}^{(k+1)} = x_n^{(k+1)} + hF(t_{n+1}, x_{n+1}^{(k+1)}, x_{n+1}^{(k)}), & n = 0, 1, \dots, N, \\ x_0^{(k+1)} = x_0, \end{cases} \quad (1.4)$$

where the step size h satisfies the following: $Nh = T$, mesh points $t_n = nh$ and $x_n^{(0)} = x_0$ for any n .

The WR method was introduced for the first time by Lelarasme et al. [1] for the time domain analysis of large-scale nonlinear dynamical systems. Its two advantages having a multirate property and inherent parallelism, can be observed, making it quite competitive with the classical methods based, for examples, on discrete variable methods such as the Runge-Kutta, linear multistep or predictor-correction methods for differential systems [2–4]. The convergence of WR methods has been studied for various types of ordinary differential equations (ODEs). Here we only mention the references [2, 5–10].

However, a convergent WR method may be impractical. To confirm this a convergent and unstable example of WR methods is given as follows:

Applying (1.4) to the equation

$$\dot{x} = -3x, t \in [0, T]; x(0) = x_0,$$

one arrives at the WR method given by

$$\begin{cases} x_{n+1}^{(k+1)} = x_n^{(k+1)} - \alpha h x_{n+1}^{(k+1)} + (\alpha - 3)h x_{n+1}^{(k)}, \\ x_0^{(k+1)} = x_0, x_n^{(0)} \equiv x_0. \end{cases} \quad (1.5)$$

This method is convergent, that is $\lim_{k \rightarrow \infty, N \rightarrow \infty} \max_{0 \leq n \leq T/N} |x_n^{(k)} - x(t_n)| = 0$. For simplicity we assume that all calculations associated with (1.5) are exact and only initial values have tiny perturbations. Let $\{\varepsilon_n^{(k)}\}$ denote these perturbations; then, the resulting perturbed solution $\{\tilde{x}_n^{(k)}\}$ satisfies

$$\begin{cases} \tilde{x}_{n+1}^{(k+1)} = \tilde{x}_n^{(k+1)} - \alpha h \tilde{x}_{n+1}^{(k+1)} + (\alpha - 3)h \tilde{x}_{n+1}^{(k)}, \\ \tilde{x}_0^{(k+1)} = x_0 + \varepsilon_0^{(k+1)}, \tilde{x}_n^{(0)} \equiv x_0 + \varepsilon_n^{(0)}. \end{cases} \quad (1.6)$$

In Example 3.9 of Section 3 we will show that the differences $x_n^{(k)} - \tilde{x}_n^{(k)}$, $k, n \in \mathbb{Z}^+$ may be unbounded for any h when $\alpha = 1$. This means that the tiny perturbations of initial values may lead to a huge change in the solution for the above WR method. A method with a similar property is called unstable and is clearly impractical. Hence, one should study the stability property of WR methods. Two simplified approaches have been used to understand partly the stability of WR methods. One approach studies stability of the numerical method generated by letting the iteration index of a WR method go to infinity [11, 12]. Another approach studies the stability of the time-point relaxation method, a variant of WR in which each window size equals the step size used in numerical integration [11, 13].

In [14] a very different approach is adopted, which investigates whether or not the methods obtained by applying some standard discretized time WR methods to a dissipative system can preserve the contractivity properties of the system.

Now, there exist few studies on the stability properties of WR methods except for the aforementioned references. Thus, further investigation is necessary. In this paper we study the stability of WR methods by using a similar approach as that used in [14]. First, apply the WR methods to some stable test equations. Second, study under what conditions the WR methods can preserve the stability of exact solutions. Last, explore the key factors for determining the stability of WR methods. In fact, this approach has been widely used to study the stability of classical numerical methods of ODEs. Being different from [14] in this paper we will consider more general stability properties than contractivity, use more splittings than the three special splittings used in [14] and derive some new interesting results. Based on these stability results, it is expected that the discrete-time WR methods can be applied to some significant inverse problems of mathematical physics [15–18], and we also refer the reader to [19–25] for more related discussions.

The organization of this paper is as follows. In Section 2, we will present the definition of the stability of WR methods and prove two useful lemmas. Section 3 investigates the stability of the WR methods applied to (2.1) and explores the key factors impacting the stability of WR methods. In Section 4, we study the stability of the WR methods applied to (2.2) and try to provide some stable conditions. At last, Section 5 presents some numerical experiments to illustrate the theories obtained.

2. Preliminary

In this section we provide some stability definitions for the WR methods and two lemmas for the study of stability in the next two sections.

Consider the following two dissipative systems:

$$\dot{x} = Ax, t \geq 0; x(0) = x_0, \quad (2.1)$$

with $A \in \mathbb{R}^{d \times d}$ and $\max_i \Re(\lambda_i(A)) < 0$ ($\Re(\lambda_i(A))$ denoting the real part of eigenvalues λ_i of A), and

$$\dot{x} = f(t, x), t \geq 0; x(0) = x_0, \quad (2.2)$$

where f satisfies the one-sided Lipschitz condition

$$\langle x_1 - x_2, f(t, x_1) - f(t, x_2) \rangle \leq -c \|x_1 - x_2\|^2. \quad (2.3)$$

Here, $x_1, x_2 \in \mathbb{R}^d$ and $c > 0$. Here and hereafter $\|\cdot\|$ denotes, with the exception of Section 4, a vector norm or its induced matrix norm.

Assume that all calculations in the WR methods are exact with the exception of initial values and let $\{\varepsilon_n^{(0)}, \varepsilon_0^{(k)}, n, k \in \mathbb{Z}^+\}$ be the perturbations of initial values. This is a classical assumption for an investigation into the stability of numerical methods for ODEs. Let $\{x_n^{(k)}, n \in \mathbb{Z}^+, k \in \mathbb{Z}^+\}$ denote the approximate solution generated by a WR method with the initial values $\{x_n^{(0)}, x_0^{(k)}, n, k \in \mathbb{Z}^+\}$, and let $\{\tilde{x}_n^{(k)}, n \in \mathbb{Z}^+, k \in \mathbb{Z}^+\}$ denote the perturbed solution generated by the perturbed system obtained by replacing the initial values of the above WR method by $\{x_n^{(0)} + \varepsilon_n^{(0)}, x_0^{(k)} + \varepsilon_0^{(k)}, n, k \in \mathbb{Z}^+\}$.

In general, we say that a numerical method is stable if the differences between the approximate solution and its perturbed solution are controllable when it is applied to a dissipative system.

Definition 2.1. A WR method is called stable if the approximate solution $\{x_n^{(k)}, n \in \mathbb{Z}^+, k \in \mathbb{Z}^+\}$ and its perturbed solution $\{\tilde{x}_n^{(k)}, n \in \mathbb{Z}^+, k \in \mathbb{Z}^+\}$ generated by applying it to the dissipative systems (2.1) or (2.2) satisfy that

$$\exists C > 0, st : \sup_{k,n} \|x_n^{(k)} - \tilde{x}_n^{(k)}\| \leq C \max \left\{ \sup_k \|x_0^{(k)} - \tilde{x}_0^{(k)}\|, \sup_n \|x_n^{(0)} - \tilde{x}_n^{(0)}\| \right\}.$$

The WR method is called contractive if it is stable and $C < 1$. The WR method is called asymptotically stable if it is stable and satisfies

$$\forall \varepsilon > 0, \exists K > 0, N > 0, st : \|x_n^{(k)} - \tilde{x}_n^{(k)}\| < \varepsilon \text{ for all } k > K \text{ and } n > N,$$

that is

$$\lim_{n,k \rightarrow \infty} \|x_n^{(k)} - \tilde{x}_n^{(k)}\| = 0.$$

Lemma 2.2. Suppose that $a \geq 0, b \geq 0, c \geq 0, \rho = a + b + c < 1$ and the sequence of positive numbers $\{u_n^{(k)}, n, k \in \mathbb{Z}^+\}$ satisfy

$$u_{n+1}^{(k+1)} \leq au_n^{(k+1)} + bu_n^{(k)} + cu_{n+1}^{(k)}. \quad (2.4)$$

Then,

$$\sup_{j \geq i} \left\{ \max \left\{ u_j^{(i)}, u_i^{(j)} \right\} \right\} \leq \rho^i \sup_{j \geq 0} \left\{ \max \left\{ u_j^{(0)}, u_0^{(j)} \right\} \right\}, \quad (2.5)$$

for all $i \geq 1$.

Proof. Let M denote $\sup_{j \geq 0} \left\{ \max \left\{ u_j^{(0)}, u_0^{(j)} \right\} \right\}$. We will show that (2.5) holds for all $i \geq 1$ by the induction.

Our proof is divided into two steps.

Step 1: Show that (2.5) holds given that $i = 1$, that is

$$\max \left\{ u_j^{(1)}, u_1^{(j)} \right\} \leq \rho M \quad (2.6)$$

for all $j \geq 1$. For this the mathematical induction is used again. By virtue of (2.4) one can derive easily that (2.6) holds for $j = 1$. Suppose that (2.6) holds for the index j . Then, it is enough to show that (2.6) also holds when j is replaced with $j + 1$. By (2.4) and (2.6), we have

$$u_{j+1}^{(1)} \leq au_j^{(1)} + bu_j^{(0)} + cu_{j+1}^{(0)} \leq \rho M + bM + cM \leq \rho M,$$

and

$$u_1^{(j+1)} \leq au_0^{(j+1)} + bu_0^{(j)} + cu_1^{(j)} \leq aM + bM + c\rho M \leq \rho M.$$

Hence, (2.6) is true when j is replaced with $j + 1$.

Step 2: Assume that (2.5) holds for the index i , and we will show that it still holds if i is replaced by $i + 1$, that is

$$\max \left\{ u_j^{(i+1)}, u_{i+1}^{(j)} \right\} \leq \rho^{i+1} M, \quad (2.7)$$

for all $j \geq i + 1$. By (2.4) and (2.5) one can derive easily that

$$u_{i+1}^{(i+1)} \leq au_i^{(i+1)} + bu_i^{(i)} + cu_{i+1}^{(i)} \leq \rho^{i+1} M.$$

Hence, (2.7) is true for $j = i + 1$. Now assume that (2.7) is true for the index j . Then, it is enough to show that (2.7) holds when j is replaced by $j + 1$ by the induction. By virtue of (2.4), (2.5) and (2.7), we can obtain that

$$u_{j+1}^{(i+1)} \leq au_j^{(i+1)} + bu_j^{(i)} + cu_{j+1}^{(i)} \leq a\rho^{i+1}M + b\rho^iM + c\rho^iM \leq \rho^{i+1}M,$$

and

$$u_{i+1}^{(j+1)} \leq au_i^{(j+1)} + bu_i^{(j)} + cu_{i+1}^{(j)} \leq a\rho^iM + b\rho^iM + c\rho^{i+1}M \leq \rho^{i+1}M,$$

which imply that (2.7) is true when j is replaced with $j + 1$.

Consequently, Steps 1 and 2 show that (2.5) holds for all $i \geq 1$ by the induction. The proof is complete. \square

Lemma 2.3. *Suppose that $a \geq 0, b \geq 0, c \geq 0, \rho = a + b + c \geq 1$ and the sequence of positive numbers $\{u_n^{(k)}, n \in \mathbb{Z}^+, k \in \mathbb{Z}^+\}$ satisfy*

$$u_{n+1}^{(k+1)} = au_n^{(k+1)} + bu_n^{(k)} + cu_{n+1}^{(k)}. \quad (2.8)$$

Then,

$$\inf_{j \geq i} \left\{ \min \left\{ u_j^{(i)}, u_i^{(j)} \right\} \right\} \geq \rho^i \inf_{j \geq 0} \left\{ \min \left\{ u_j^{(0)}, u_0^{(j)} \right\} \right\}, \quad (2.9)$$

for all $i \geq 1$.

Proof. The proof of the lemma is similar to that of Lemma 2.2, so we omit it. \square

3. Stability of WR methods for Eq (2.1)

Choosing the splitting function $F(t, x, x) = A_1x + A_2x = Ax$ and applying the WR methods (1.3) and (1.4) to Eq (2.1) one arrives at

$$\begin{cases} x_{n+1}^{(k+1)} = (I + hA_1)x_n^{(k+1)} + hA_2x_n^{(k)}, \\ x_0^{(k+1)} = x_0, x_n^{(0)} = x_0, n, k \in \mathbb{Z}^+ \end{cases} \quad (3.1)$$

and

$$\begin{cases} \tilde{x}_{n+1}^{(k+1)} = (I - hA_1)^{-1}x_n^{(k+1)} + (I - hA_1)^{-1}hA_2x_n^{(k)}, \\ \tilde{x}_0^{(k+1)} = x_0, \tilde{x}_n^{(0)} = x_0, n, k \in \mathbb{Z}^+. \end{cases} \quad (3.2)$$

It is easy to give the sufficient conditions for the stability of WR methods (3.1) and (3.2) by using Lemma 2.2.

Theorem 3.1. *The WR method (3.1) is contractive and asymptotically stable if $\|I + hA_1\| + \|hA_2\| < 1$.*

Proof. Let $\{x_n^{(k)}\}$ be the approximate solution generated by (3.1) and $\{\varepsilon_n^{(0)}, \varepsilon_0^{(k)}\}$ be the perturbations of the initial values. Then, the perturbation solution $\{\tilde{x}_n^{(k)}\}$ caused by $\{\varepsilon_n^{(0)}, \varepsilon_0^{(k)}\}$ satisfies

$$\begin{cases} \tilde{x}_{n+1}^{(k+1)} = (I + hA_1)\tilde{x}_n^{(k+1)} + hA_2\tilde{x}_n^{(k)}, \\ \tilde{x}_0^{(k+1)} = x_0 + \varepsilon_0^{(k+1)}, \tilde{x}_n^{(0)} = x_0 + \varepsilon_n^{(0)}, n, k \in \mathbb{Z}^+. \end{cases} \quad (3.3)$$

Let $e_n^{(k)}$ denote $\|x_n^{(k)} - \tilde{x}_n^{(k)}\|$. By virtue of (3.1) and (3.3) we have

$$e_{n+1}^{(k+1)} \leq \|I + hA_1\|e_n^{(k+1)} + \|hA_2\|e_n^{(k)}, n, k \in \mathbb{Z}^+.$$

This together with Lemma 2.2, proves the theorem. \square

Theorem 3.2. *The WR method (3.2) is contractive and asymptotically stable if $\|(I - hA_1)^{-1}\| + \|(I - hA_1)^{-1}hA_2\| < 1$.*

Proof. The proof is similar to that of Theorem 3.1, so we omit it. \square

Now we try to derive some interesting results of stability for the WR methods (3.1) and (3.2) by using Theorems 3.1 and 3.2. We call the WR methods (3.1) and (3.2) the block Gauss-Jacobi (BGJ) WR methods if

$$A_1 = \begin{pmatrix} A_{11} & & \\ & \ddots & \\ & & A_{ss} \end{pmatrix} \text{ (Block diagonal matrix),}$$

and the block Gauss-Seidel (BGS) WR methods if

$$A_1 = \begin{pmatrix} A_{11} & & \\ \vdots & \ddots & \\ A_{n1} & \dots & A_{ss} \end{pmatrix} \text{ (Block lower triangular matrix),}$$

where $A_{ii}, i = 1, 2, \dots, s$ denotes the square matrices, and \tilde{a}_{ij} , the i -th line and j -th row of A_1 , is equal to a_{ij} when $\tilde{a}_{ij} \neq 0$. The BGJ and BGS splitting may be the most common splitting. Bellen and his coauthors in [14] have studied the stability of WR methods with the Gauss-Jacobi (GJ) and Gauss-Seidel (GS) splittings, and their study shows that the methods are stable when A is diagonally dominant with negative diagonal elements; however, the techniques developed in [14] may be unsuitable to study the stability of WR methods with the BGJ and BGS splitting. The following two corollaries show that the WR methods are stable when using the splitting matrix satisfying some conditions including the BGJ and BGS splittings.

Lemma 3.3. ([26]) *Let $A = (a_{ij})_{d \times d}$ and $r_i(A) = \sum_{i \neq j} |a_{ij}|$. If A is diagonally dominant, then,*

$$\|A^{-1}\|_{\infty} \leq \frac{1}{\min_i (|a_{ii}| - r_i(A))}.$$

Here, and throughout, $\|\cdot\|_{\infty}$ denotes the maximum norm in \mathbb{C}^d or the maximum row sum matrix norm.

Corollary 3.4. *Let $A = (a_{ij})_{d \times d}$ be diagonally dominant with negative diagonal elements. Let D denote the diagonal matrix $\text{diag}(a_{11}, a_{22}, \dots, a_{dd})$, and let U, V be any $d \times d$ matrix satisfying $U + V = A - D$. Suppose that*

$$\max_i \Re(a_{ii}) + a < 0, \text{ where } a = \|U\|_{\infty} + \|V\|_{\infty}. \quad (3.4)$$

Then, the method (3.1) with the splitting $A_1 = D + U$ is contractive and asymptotically stable if

$$h < \min \left\{ \frac{1}{a}, \min_i \frac{-2(\Re(a_{ii}) + a)}{|a_{ii}|^2 - a^2} \right\}. \quad (3.5)$$

Proof. By (3.4), we have

$$a < -\Re(a_{ii}) < |a_{ii}|, \forall i.$$

This and (3.5) yield

$$2(\Re(a_{ii}) + a) + h(|a_{ii}|^2 - a^2) < 0, \forall i. \quad (3.6)$$

Multiplying h and adding 1 to the two sides of (3.6) we get

$$1 + 2h\Re(a_{ii}) + h^2(\Re(a_{ii}))^2 + h^2(\Im(a_{ii}))^2 < 1 - 2ah + a^2, \forall i.$$

Hence, we have

$$(1 + h\Re(a_{ii}))^2 + (h\Im(a_{ii}))^2 < (1 - ah)^2, \forall i,$$

that is

$$|1 + ha_{ii}|^2 < (1 - ah)^2, \forall i.$$

Noting that (3.5) implies that $1 - ah > 0$, we get the following from the above inequality

$$|1 + ha_{ii}| + ah < 1, \forall i.$$

This means that

$$\|I + hD\|_\infty + h\|U\|_\infty + h\|V\|_\infty < 1.$$

Note that $A_2 = V$ if $A_1 = D + U$ for the method (3.1). Thus, the above inequality yields

$$\|I + hA_1\|_\infty + \|hA_2\|_\infty < 1.$$

This together with Theorem 3.1, completes the proof of Corollary 3.4. \square

Corollary 3.5. *Let A, D, U and V be the same as those in Corollary 3.4 and $D + U$ be diagonally dominant. Suppose that*

$$\|V\|_\infty + r_i(U) + \Re(a_{ii}) < 0, \forall i. \quad (3.7)$$

Then, the method (3.2) with the splitting $A_1 = D + U$ is contractive and asymptotically stable for any step size h .

Proof. Noting that $|\Re(a_{ii})| < |a_{ii}|$ we can obtain from (3.7) that

$$|a_{ii}|^2 - (\|V\|_\infty + r_i(U))^2 > 0, \forall i.$$

The above condition and (3.7) show that for any step size h

$$2h(\|V\|_\infty + r_i(U) + \Re(a_{ii})) \leq (|a_{ii}|^2 - (\|V\|_\infty + r_i(U))^2)h^2, \forall i.$$

This means that

$$1 + 2h(\|V\|_\infty + r_i(U)) + (\|V\|_\infty + r_i(U))^2h^2 \leq 1 - 2h\Re(a_{ii}) + (|a_{ii}|^2)h^2, \forall i,$$

that is

$$(1 + h(\|V\|_\infty + r_i(U)))^2 \leq |1 - ha_{ii}|^2, \forall i.$$

Hence, we get

$$1 + h\|V\|_\infty \leq \min_i (|1 - ha_{ii}| - hr_i(U)).$$

This together with Lemma 3.3, shows that

$$\|(I - h(D + U))^{-1}\|_\infty + \|(I - h(D + U))^{-1}\|_\infty \|V\|_\infty h < 1.$$

Noting that if $A_1 = D + U$ then $A_2 = V$ for the WR method (3.2); the above condition and Theorem 3.2 complete the proof of Corollary 3.5. \square

A severe restriction on the matrix is required to satisfy the condition of diagonal dominance in the presence of negative diagonal elements. The following two corollaries show that the condition that all eigenvalues of A in (2.1) have negative real parts is enough to guarantee the stability of the WR methods (3.1) and (3.2) whenever the splitting is suitable.

Corollary 3.6. *Let A be the matrix in (2.1) and T be a non-singular matrix satisfying*

$$A = T^{-1} \begin{pmatrix} \lambda_1 & \mu_1 & & 0 \\ & \lambda_2 & \ddots & \\ & & \ddots & \mu_{d-1} \\ 0 & & & \lambda_d \end{pmatrix} T,$$

where λ_i is the eigenvalue of A and $\mu_i = 0$ or 1 . Assume that

$$A_1 = T^{-1} \begin{pmatrix} \lambda_1 & & & 0 \\ & \lambda_2 & & \\ & & \ddots & \\ 0 & & & \lambda_d \end{pmatrix} T \text{ and } A_2 = T^{-1} \begin{pmatrix} 0 & \mu_1 & & 0 \\ & 0 & \ddots & \\ & & \ddots & \mu_{d-1} \\ 0 & & & 0 \end{pmatrix} T.$$

Then, the WR method (3.1) is asymptotically stable when the step size h satisfies

$$0 < h < \min_i \frac{-2\Re(\lambda_i)}{|\lambda_i|^2}. \quad (3.8)$$

Proof. Let $y_n^{(k)} = T x_n^{(k)}$, $\forall k, n$. By virtue of (3.1) we have

$$y_{n+1}^{(k+1)} = (I + h\Lambda_1)y_n^{(k+1)} + h\Lambda_2 y_n^{(k)}, \quad (3.9)$$

$$\text{where } \Lambda_1 = \begin{pmatrix} \lambda_1 & & & 0 \\ & \lambda_2 & & \\ & & \ddots & \\ 0 & & & \lambda_d \end{pmatrix}, \Lambda_2 = \begin{pmatrix} 0 & \mu_1 & & 0 \\ & 0 & \ddots & \\ & & \ddots & \mu_{d-1} \\ 0 & & & 0 \end{pmatrix}.$$

Hence, (3.9) is asymptotically stable by Theorem 3.1 if

$$\|I + h\Lambda_1\| + \|h\Lambda_2\| < 1, \quad (3.10)$$

where $\|\cdot\|$ is an operator norm of the matrix.

It is clear that (3.1) is asymptotically stable if (3.9) is asymptotically stable. Thus, it is enough to show that (3.8) implies (3.10).

Taking $S = \text{diag}(1, t, \dots, t^d)$ and $\|A\|_S = \|SAS^{-1}\|_\infty$, we have

$$\begin{aligned} \|I + h\Lambda_1\|_S + \|h\Lambda_2\|_S &= \|S(I + h\Lambda_1)S^{-1}\|_\infty + h\|S\Lambda_2S^{-1}\|_\infty \\ &= \|I + h\Lambda_1\|_\infty + h|t^{-1}|\|\Lambda_2\|_\infty \\ &= \max_i |1 + h\lambda_i| + h|t^{-1}|. \end{aligned} \quad (3.11)$$

By (3.8) we have

$$1 + h^2|\lambda_i|^2 + 2h\Re(\lambda_i) < 1.$$

Noting that

$$\begin{aligned} 1 + h^2|\lambda_i|^2 + 2h\Re(\lambda_i) &= 1 + 2h\Re(\lambda_i) + h^2(\Re(\lambda_i))^2 + h^2(\Im(\lambda_i))^2 \\ &= (1 + h\Re(\lambda_i))^2 + h^2(\Im(\lambda_i))^2 = |1 + h\lambda_i|^2, \end{aligned}$$

we can obtain

$$|1 + h\lambda_i| < 1, \quad \forall i.$$

Hence,

$$\max_i |1 + h\lambda_i| < 1. \quad (3.12)$$

By (3.11) and (3.12) we have the following for a sufficiently large t

$$\|I + h\Lambda_1\|_S + \|h\Lambda_2\|_S < 1.$$

This together with Theorem 3.1, completes the proof of Corollary 3.6. \square

Corollary 3.7. *If A_1 and A_2 are the same as those in Corollary 3.6, then the WR method (3.2) is asymptotically stable for any step size h .*

Proof. Let $y_n^{(k)} = Tx_n^{(k)}$, $\forall k, n$. By virtue of (3.2) we have

$$y_{n+1}^{(k+1)} = (I - h\Lambda_1^{-1})y_n^{(k+1)} + (I - h\Lambda_1^{-1})\Lambda_2hy_n^{(k)}. \quad (3.13)$$

Here, Λ_1 and Λ_2 are the same as those in the proof of Corollary 3.6.

Take $S = \text{diag}(1, t, \dots, t^d)$ and $\|A\|_S = \|SAS^{-1}\|_\infty$. By Theorem 3.2, the method (3.13) is asymptotically stable if

$$\|(I - h\Lambda_1)^{-1}\|_S + \|(I - h\Lambda_1)^{-1}\Lambda_2h\|_S < 1.$$

Clearly, it is enough to prove that

$$\|(I - h\Lambda_1)^{-1}\|_\infty(1 + h\|\Lambda_2\|_S) < 1,$$

that is

$$\frac{1}{|1 - h\lambda_i|}(1 + h\|\Lambda_2\|_S) < 1, \quad \forall i. \quad (3.14)$$

Note that $\Re(\lambda_i) < 0, \forall i$. Hence,

$$|1 - h\lambda_i| > 1 - h\Re(\lambda_i), \forall i.$$

Let t be large enough such that

$$\|\Lambda_2\|_S = |t^{-1}| \|\Lambda_2\|_\infty < \min_i (-\Re(\lambda_i)).$$

Thus, for t large enough, we have

$$1 + h\|\Lambda_2\|_S < 1 - h\Re(\lambda_i) < |1 - h\lambda_i|, \forall i,$$

that is (3.14) holds.

Consequently, the proof is complete by the fact that the asymptotically stable property of (3.2) and (3.13) is the same. \square

Lemma 2.3 shows that if $\|I + hA_1\| + \|hA_2\| > 1$ and $e_{n+1}^{(k+1)} = \|I + hA_1\|e_n^{(k+1)} + \|hA_2\|e_n^{(k)}$, then the WR method (3.1) is unstable.

Example 3.8. For the differential equation

$$x' = \lambda x, \lambda < 0 \tag{3.15}$$

and its WR method

$$x_{n+1}^{(k+1)} = (1 + \lambda_1 h)x_n^{(k+1)} + \lambda_2 h x_n^{(k)} (\lambda_1 + \lambda_2 = \lambda), \tag{3.16}$$

let $\tilde{x}_n^{(k)}$ denote the perturbed value of $x_n^{(k)}$, $n, k \in \mathbb{Z}^+$, which satisfies

$$\tilde{x}_{n+1}^{(k+1)} = (1 + \lambda_1 h)\tilde{x}_n^{(k+1)} + \lambda_2 h \tilde{x}_n^{(k)}.$$

Let $\epsilon_n^{(k)}$ denote $\tilde{x}_n^{(k)} - x_n^{(k)}$. When taking $\epsilon_0^{(k)} > 0$ for $k \in \mathbb{Z}^+$, $\epsilon_{2l}^{(0)} > 0$ and $\epsilon_{2l+1}^{(0)} < 0$ for $l \in \mathbb{Z}^+$, $\lambda_1, \lambda_2 < 0$ and $1 + \lambda_1 h < 0$, we have

$$|\epsilon_{n+1}^{(k+1)}| = |1 + \lambda_1 h| |\epsilon_n^{(k+1)}| + |\lambda_2 h| |\epsilon_n^{(k)}|.$$

Suppose that $\inf_{j \geq 0} \left\{ \min \left\{ |\epsilon_j^{(0)}|, |\epsilon_0^{(j)}| \right\} \right\} > 0$. Then, the WR method (3.16) is unstable when $h > \max \left\{ -\frac{1}{\lambda_1}, -\frac{2}{\lambda} \right\}$.

Similarly, the WR method (3.2) may also be unstable if $\|(I - hA_1)^{-1}\| + \|(I - hA_1)^{-1}hA_2\| > 1$ by Lemma 2.3.

Example 3.9. Applying (3.2) to Eq (3.15) we get the following WR method

$$\begin{aligned} x_{n+1}^{(k+1)} &= x_n^{(k+1)} + \lambda_1 h x_{n+1}^{(k+1)} + \lambda_2 h x_{n+1}^{(k)} \\ &= \frac{1}{1 - \lambda_1 h} x_n^{(k+1)} + \frac{\lambda_2 h}{1 - h\lambda_1} x_{n+1}^{(k)} \quad (\lambda_1 + \lambda_2 = \lambda). \end{aligned} \tag{3.17}$$

Let $\lambda_2 < \lambda_1 < 0$. Let $\tilde{x}_n^{(k)}$ denote the perturbed value of $x_n^{(k)}$ and $\epsilon_n^{(k)}$ denote $\tilde{x}_n^{(k)} - x_n^{(k)}$ for $n, k \in \mathbb{Z}^+$. When taking $\epsilon_n^{(0)} < 0$ for $n \in \mathbb{Z}^+$ and $\epsilon_0^{(2l)} < 0$ and $\epsilon_0^{(2l+1)} > 0$ for $l \in \mathbb{Z}^+$, we have

$$|\epsilon_{n+1}^{(k+1)}| = \left| \frac{1}{1 - \lambda_1 h} \right| |\epsilon_n^{(k+1)}| + \left| \frac{\lambda_2 h}{1 - \lambda_1 h} \right| |\epsilon_{n+1}^{(k)}|.$$

Suppose that $\inf_{j \geq 0} \left\{ \min \left\{ |\epsilon_j^{(0)}|, |\epsilon_0^{(j)}| \right\} \right\} > 0$. Then, the WR method (3.17) is unstable for any h .

4. Stability of WR methods for Eq (2.2)

Let $\langle \cdot, \cdot \rangle$ denote an inner product and $\| \cdot \|$ the corresponding inner product norm. Assume that the splitting function $F(t, x, y)$ satisfies the following conditions:

(A1) there exists $C > 0$ such that

$$\langle x - \tilde{x}, F(t, x, y) - F(t, \tilde{x}, y) \rangle \leq -C\|x - \tilde{x}\|^2, \forall t \in \mathbb{R}, \forall x, \tilde{x}, y \in \mathbb{R}^d,$$

and

(A2) there exists $L_1, L_2 > 0$ such that

$$\|F(t, x, y) - F(t, \tilde{x}, \tilde{y})\|^2 \leq L_1\|x - \tilde{x}\|^2 + L_2\|y - \tilde{y}\|^2, \forall t \in \mathbb{R}, \forall x, y, \tilde{x}, \tilde{y} \in \mathbb{R}^d.$$

Consider the WR methods (1.3) and (1.4) of Eq (2.2). Let $\{\varepsilon_n^{(0)}, \varepsilon_0^{(k)}\}$ be the perturbations of the initial values of (1.3) and (1.4). Then, the perturbed solution $\{\tilde{x}_n^{(k)}\}$ corresponding to (1.3) satisfies the following:

$$\begin{cases} \tilde{x}_{n+1}^{(k+1)} = \tilde{x}_n^{(k+1)} + hF(t_n, \tilde{x}_n^{(k+1)}, \tilde{x}_n^{(k)}), \\ \tilde{x}_0^{(k+1)} = x_0 + \varepsilon_0^{(k+1)}, \tilde{x}_n^{(0)} = x_0 + \varepsilon_n^{(0)}, n, k \in \mathbb{Z}^+, \end{cases} \quad (4.1)$$

and the perturbed solution $\{\tilde{x}_n^{(k)}\}$ corresponding to (1.4) satisfies the following:

$$\begin{cases} \tilde{x}_{n+1}^{(k+1)} = \tilde{x}_n^{(k+1)} + hF(t_{n+1}, \tilde{x}_{n+1}^{(k+1)}, \tilde{x}_{n+1}^{(k)}), \\ \tilde{x}_0^{(k+1)} = x_0 + \varepsilon_0^{(k+1)}, \tilde{x}_n^{(0)} = x_0 + \varepsilon_n^{(0)}, n, k \in \mathbb{Z}^+. \end{cases} \quad (4.2)$$

Theorem 4.1. Suppose that

$$\sqrt{L_2} < C \quad (4.3)$$

and the assumptions (A1) and (A2) hold. Then, the WR method (1.3) is contractive and asymptotically stable if

$$h < \min \left\{ \frac{1}{2C - \sqrt{L_2}}, \frac{2(C - \sqrt{L_2})}{L_1 + L_2} \right\}. \quad (4.4)$$

Proof. Let $\{x_n^{(k)}\}$ and $\{\tilde{x}_n^{(k)}\}$ denote the approximate values respectively generated by (1.3) and (4.1), and let $\delta_n^{(k)}, F_n$ and \tilde{F}_n denote $x_n^{(k)} - \tilde{x}_n^{(k)}, F(t_n, x_n^{(k+1)}, x_n^{(k)})$ and $F(t_n, \tilde{x}_n^{(k+1)}, \tilde{x}_n^{(k)})$, respectively. By (1.3) and (4.1) we have

$$\begin{aligned} \langle \delta_{n+1}^{(k+1)}, \delta_{n+1}^{(k+1)} \rangle &= \langle \delta_n^{(k+1)} + h(F_n - \tilde{F}_n), \delta_n^{(k+1)} + h(F_n - \tilde{F}_n) \rangle \\ &= \langle \delta_n^{(k+1)}, \delta_n^{(k+1)} \rangle + 2h\langle \delta_n^{(k+1)}, F_n - \tilde{F}_n \rangle \\ &\quad + h^2\langle F_n - \tilde{F}_n, F_n - \tilde{F}_n \rangle. \end{aligned} \quad (4.5)$$

Using the assumptions (A1) and (A2), we can obtain

$$\begin{aligned} &\langle \delta_n^{(k+1)}, (F_n - \tilde{F}_n) \rangle \\ &= \langle \delta_n^{(k+1)}, F(t_n, x_n^{(k+1)}, x_n^{(k)}) - F(t_n, \tilde{x}_n^{(k+1)}, \tilde{x}_n^{(k)}) \rangle \\ &\quad + \langle \delta_n^{(k+1)}, F(t_n, \tilde{x}_n^{(k+1)}, x_n^{(k)}) - F(t_n, \tilde{x}_n^{(k+1)}, \tilde{x}_n^{(k)}) \rangle \\ &\leq -C\|\delta_n^{(k+1)}\|^2 + \|\delta_n^{(k+1)}\| \|F(t_n, \tilde{x}_n^{(k+1)}, x_n^{(k)}) - F(t_n, \tilde{x}_n^{(k+1)}, \tilde{x}_n^{(k)})\| \\ &\leq -C\|\delta_n^{(k+1)}\|^2 + \sqrt{L_2}\|\delta_n^{(k+1)}\| \|\delta_n^{(k)}\|, \end{aligned} \quad (4.6)$$

and

$$\langle F_n - \tilde{F}_n, F_n - \tilde{F}_n \rangle = \|F_n - \tilde{F}_n\|^2 \leq L_1 \|\delta_n^{(k+1)}\|^2 + L_2 \|\delta_n^{(k)}\|^2. \quad (4.7)$$

By (4.5), (4.6) and (4.7), we get

$$\begin{aligned} \|\delta_{n+1}^{(k+1)}\|^2 &\leq \|\delta_n^{(k+1)}\|^2 - 2Ch\|\delta_n^{(k+1)}\|^2 + 2\sqrt{L_2}h\|\delta_n^{(k+1)}\|\|\delta_n^{(k)}\| \\ &\quad + L_1h^2\|\delta_n^{(k+1)}\|^2 + L_2h^2\|\delta_n^{(k)}\|^2 \\ &= (1 - 2Ch + \sqrt{L_2}h + L_1h^2)\|\delta_n^{(k+1)}\|^2 + (\sqrt{L_2}h + L_2h^2)\|\delta_n^{(k)}\|^2. \end{aligned}$$

It is not difficult to prove that

$$|1 - 2Ch + \sqrt{L_2}h + L_1h^2| + |\sqrt{L_2}h + L_2h^2| < 1,$$

under the conditions (4.3) and (4.4). This together with Lemma 2.2, completes the proof of Theorem 4.1. \square

Theorem 4.2. *Suppose that the condition (4.3) and the assumptions (A1) and (A2) hold. Then, the WR method (1.4) is contractive and asymptotically stable for any step size h .*

Proof. Let $\{x_n^{(k)}\}$ and $\{\tilde{x}_n^{(k)}\}$ denote the approximate values respectively generated by (1.4) and (4.2), and let $\delta_n^{(k)}$, F_{n+1} and \tilde{F}_{n+1} denote $x_n^{(k)} - \tilde{x}_n^{(k)}$, $F(t_{n+1}, x_{n+1}^{(k+1)}, x_{n+1}^{(k)})$ and $F(t_{n+1}, \tilde{x}_{n+1}^{(k+1)}, \tilde{x}_{n+1}^{(k)})$, respectively. Using the property of the inner product we can obtain

$$\begin{aligned} &\frac{1}{h} (\langle \delta_{n+1}^{(k+1)}, \delta_{n+1}^{(k+1)} \rangle - \langle \delta_n^{(k+1)}, \delta_n^{(k+1)} \rangle) \\ &= \frac{1}{h} (\langle \delta_{n+1}^{(k+1)}, \delta_{n+1}^{(k+1)} \rangle - \langle \delta_n^{(k+1)}, \delta_{n+1}^{(k+1)} \rangle + \langle \delta_n^{(k+1)}, \delta_{n+1}^{(k+1)} \rangle - \langle \delta_n^{(k+1)}, \delta_n^{(k+1)} \rangle) \\ &= \left\langle \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}), \delta_{n+1}^{(k+1)} \right\rangle + \left\langle \delta_n^{(k+1)}, \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}) \right\rangle \\ &= \left\langle \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}), \delta_{n+1}^{(k+1)} \right\rangle + \left\langle \delta_{n+1}^{(k+1)}, \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}) \right\rangle \\ &\quad + \left\langle \delta_n^{(k+1)} - \delta_{n+1}^{(k+1)}, \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}) \right\rangle \\ &\leq 2 \left\langle \delta_{n+1}^{(k+1)}, \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}) \right\rangle. \end{aligned} \quad (4.8)$$

Note that (1.4) and (4.2) imply that

$$\delta_{n+1}^{(k+1)} = \delta_n^{(k+1)} + h(F_{n+1} - \tilde{F}_{n+1}).$$

Thus,

$$\begin{aligned} &\left\langle \delta_{n+1}^{(k+1)}, \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}) \right\rangle = \left\langle \delta_{n+1}^{(k+1)}, F_{n+1} - \tilde{F}_{n+1} \right\rangle \\ &= \left\langle \delta_{n+1}^{(k+1)}, F(t_{n+1}, x_{n+1}^{(k+1)}, x_{n+1}^{(k)}) - F(t_{n+1}, \tilde{x}_{n+1}^{(k+1)}, x_{n+1}^{(k)}) \right\rangle \\ &\quad + \left\langle \delta_{n+1}^{(k+1)}, F(t_{n+1}, \tilde{x}_{n+1}^{(k+1)}, x_{n+1}^{(k)}) - F(t_{n+1}, \tilde{x}_{n+1}^{(k+1)}, \tilde{x}_{n+1}^{(k)}) \right\rangle. \end{aligned} \quad (4.9)$$

By (A1), (A2) and (4.9), we get

$$\begin{aligned} & \left\langle \delta_{n+1}^{(k+1)}, \frac{1}{h} (\delta_{n+1}^{(k+1)} - \delta_n^{(k+1)}) \right\rangle \\ &= -C \|\delta_{n+1}^{(k+1)}\|^2 + \|\delta_{n+1}^{(k+1)}\| \|F(t_{n+1}, \tilde{x}_{n+1}^{(k+1)}, x_{n+1}^{(k)}) - F(t_{n+1}, \tilde{x}_{n+1}^{(k+1)}, \tilde{x}_{n+1}^{(k)})\| \\ &\leq -C \|\delta_{n+1}^{(k+1)}\|^2 + \|\delta_{n+1}^{(k+1)}\| \sqrt{L_2} \|\delta_{n+1}^{(k)}\|. \end{aligned} \quad (4.10)$$

By using (4.8) and (4.10), we can derive that

$$\|\delta_{n+1}^{(k+1)}\|^2 \leq \frac{1}{1 + 2Ch - \sqrt{L_2}h} \|\delta_n^{(k+1)}\|^2 + \frac{\sqrt{L_2}h}{1 + 2Ch - \sqrt{L_2}h} \|\delta_{n+1}^{(k)}\|^2. \quad (4.11)$$

Note that the condition (4.3) implies the following for any h

$$\frac{1 + \sqrt{L_2}h}{|1 + 2Ch - \sqrt{L_2}h|} < 1.$$

This together with (4.11) and Lemma 2.2, completes the proof of Theorem 4.2. \square

5. Numerical experiments

In this section we will present some numerical experiments to verify the theories developed in Section 3. Let $\xi_n^{(k)}, i = 1, 2, \dots, d, k \in \mathbb{Z}^+, n \in \mathbb{Z}^+ (k \cdot n = 0)$ be independent random variables, each uniformly distributed on the interval $[-0.5, 0.5]$. Take perturbations of the initial values $x_n^{(k)}$ as $\varepsilon_n^{(k)} = ({}_{1}\xi_n^{(k)}, {}_{2}\xi_n^{(k)}, \dots, {}_{d}\xi_n^{(k)})^T$, where $k, n \in \mathbb{Z}^+$ and $k \cdot n = 0$. Let $\{x_n^{(k)}\}$ and $\{\tilde{x}_n^{(k)}\}$ denote the numerical solutions respectively generated by (3.1) and (3.3), or (3.2) and its perturbed system given by

$$\begin{cases} \tilde{x}_{n+1}^{(k+1)} = (I - hA_1)^{-1} \tilde{x}_n^{(k+1)} + (I - hA_1)^{-1} hA_2 \tilde{x}_{n+1}^{(k)}, \\ \tilde{x}_0^{(k+1)} = x_0 + \varepsilon_0^{(k+1)}, \tilde{x}_n^{(0)} = x_0 + \varepsilon_n^{(0)}, n, k \in \mathbb{Z}^+, \end{cases}$$

and let $e_n^{(k)}$ denote $\|x_n^{(k)} - \tilde{x}_n^{(k)}\|$.

In our experiments we have investigated the relation between the errors $e_n^{(k)}$ and factors such as the splitting, the step size and the iteration number. The results are presented as figures, where we have plotted $\log(e_n^{(k)})$ versus the time nh . By Definition 2.1, the WR methods (3.1) and (3.2) are stable if $\lim_{n,k \rightarrow \infty} \log(e_n^{(k)}) = -\infty$ and unstable if $\lim_{n,k \rightarrow \infty} \log(e_n^{(k)}) = \infty$. Here and hereafter, the use of log denotes the logarithm with base e .

Let M_1, M_2 and M_3 denote the following matrices:

$$\begin{pmatrix} -5 & -2 & 0 & 0 & -1 \\ -1 & -5 & -1 & -1 & 0 \\ 0 & -2 & -6 & -1 & 0 \\ 0 & -2 & 0 & -5 & -1 \\ -1 & -1 & 0 & -1 & -6 \end{pmatrix}, \begin{pmatrix} -5 & -2 & 0 & 0 & 0 \\ -1 & -5 & -1 & 0 & 0 \\ 0 & -2 & -6 & 0 & 0 \\ 0 & 0 & 0 & -5 & -1 \\ 0 & 0 & 0 & -1 & -6 \end{pmatrix}$$

and

$$\begin{pmatrix} -5 & -2 & 0 & 0 & 0 \\ -1 & -5 & -1 & 0 & 0 \\ 0 & -2 & -6 & 0 & 0 \\ 0 & -2 & 0 & -5 & -1 \\ -1 & -1 & 0 & -1 & -6 \end{pmatrix},$$

respectively. We have verified Corollaries 3.4 and 3.5 by investigating the stability of (3.1) and (3.2) with the BGJ splitting $A_1 = M_2, A_2 = M_1 - M_2$ and the BGS splitting $A_1 = M_3, A_2 = M_1 - M_3$. By Corollaries 3.4 and 3.5, the WR method (3.1) is stable for $h < 0.2$, and (3.2) is stable for any h when using the above BGJ and BGS splittings, which is in agreement with the results of the numerical experiments plotted in Figures 1 and 2.

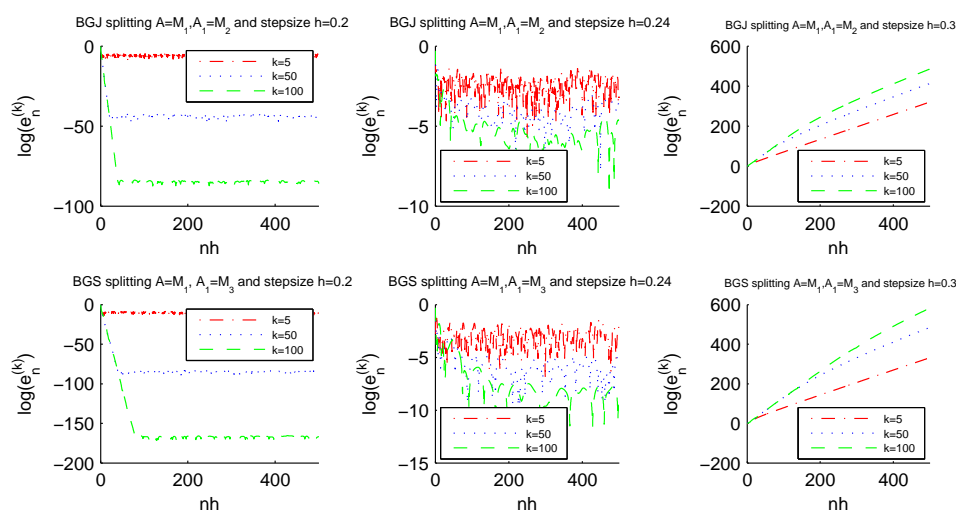


Figure 1. Numerical simulation of WR method (3.1) with BGJ and BGS splitting.

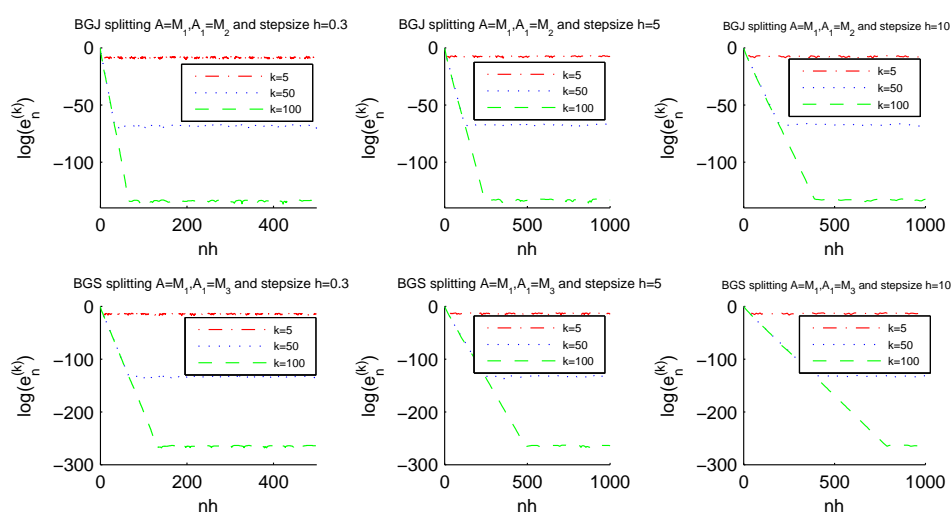


Figure 2. Numerical simulation of WR method (3.2) with BGJ and BGS splitting.

To verify Corollaries 3.6 and 3.7, we take the matrix A in (2.1) to be

$$M_4 = \begin{pmatrix} -0.5 & -1 & 0.5 \\ 10.5 & -8 & 3.5 \\ 13.5 & -7 & 2.5 \end{pmatrix} \text{ and } M_5 = \begin{pmatrix} -2.5 & 1 & -0.5 \\ 2.5 & -2 & 0.5 \\ 9.5 & -5 & 1.5 \end{pmatrix},$$

whose eigenvalues have negative parts, but which do not satisfy the conditions of Corollaries 3.4 and 3.5. We have shown that the WR methods (3.1) and (3.2) in this situation are unstable when using the classical GJ and GS splittings, but stable when using the eigenvalue (EV) splitting given in Corollary 3.6 in this paper.

Let

$$M_6 = \begin{pmatrix} 1 & -2 & 1 \\ 8 & -7 & 3 \\ 4 & -2 & 0 \end{pmatrix}, M_7 = \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, M_8 = \begin{pmatrix} -0.5 & 0 & 0 \\ 0 & -8 & 0 \\ 0 & 0 & 2.5 \end{pmatrix},$$

$$M_9 = \begin{pmatrix} -2.5 & 0 & 0 \\ 0 & -2 & 0 \\ 0 & 0 & 1.5 \end{pmatrix}, M_{10} = \begin{pmatrix} -0.5 & 0 & 0 \\ 10.5 & -8 & 0 \\ 13.5 & -7 & 2.5 \end{pmatrix},$$

$$M_{11} = \begin{pmatrix} -2.5 & 0 & 0 \\ 2.5 & -2 & 0 \\ 9.5 & -5 & 1.5 \end{pmatrix}.$$

Noting that

$$M_4 = \begin{pmatrix} 0 & 1 & 2 \\ 1 & 1 & 4 \\ 2 & -1 & 0 \end{pmatrix} \begin{pmatrix} -1 & 1 & 0 \\ 0 & -2 & 1 \\ 0 & 0 & -3 \end{pmatrix} \begin{pmatrix} 0 & 1 & 2 \\ 1 & 1 & 4 \\ 2 & -1 & 0 \end{pmatrix}^{-1},$$

and

$$M_5 = \begin{pmatrix} 0 & 1 & 2 \\ 1 & 1 & 4 \\ 2 & -1 & 0 \end{pmatrix} \begin{pmatrix} -1 & 1 & 0 \\ 0 & -1 & 1 \\ 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 & 2 \\ 1 & 1 & 4 \\ 2 & -1 & 0 \end{pmatrix}^{-1},$$

we can write the splittings of $A = M_4$ as the EV splitting $A_1 = M_6, A_2 = M_4 - M_6$, the GJ splitting $A_1 = M_8, A_2 = M_4 - M_8$ and the GS splitting $A_1 = M_{10}, A_2 = M_4 - M_{10}$, and the splittings of $A = M_5$ as the EV splitting $A_1 = M_7, A_2 = M_5 - M_7$, the GJ splitting $A_1 = M_9, A_2 = M_5 - M_9$ and the GS splitting $A_1 = M_{11}, A_2 = M_5 - M_{11}$.

By Corollary 3.6 the WR method (3.1) with $A = M_4$ is stable for the step size $h < 2/3$ when using the EV splitting, and that with $A = M_5$ is stable for the step size $h < 2$ when using the EV splitting. These are supported by the results of experiments plotted in Figure 3. By Corollary 3.7 the WR method (3.2) with $A = M_4$ or M_5 is stable for any h when using the EV splitting, which is consistent with the results of experiments plotted in Figure 4.

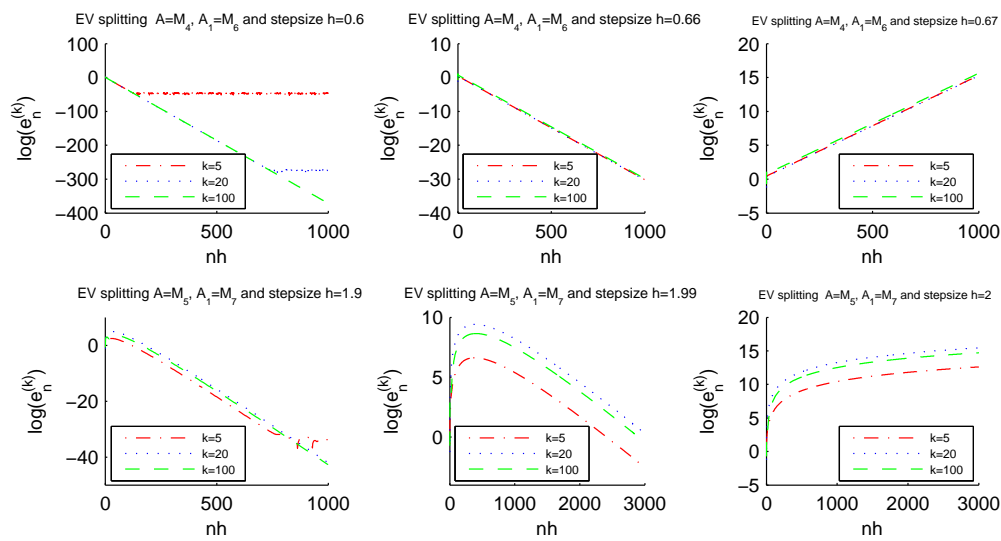


Figure 3. Numerical simulation of WR method (3.1) with the EV splitting.

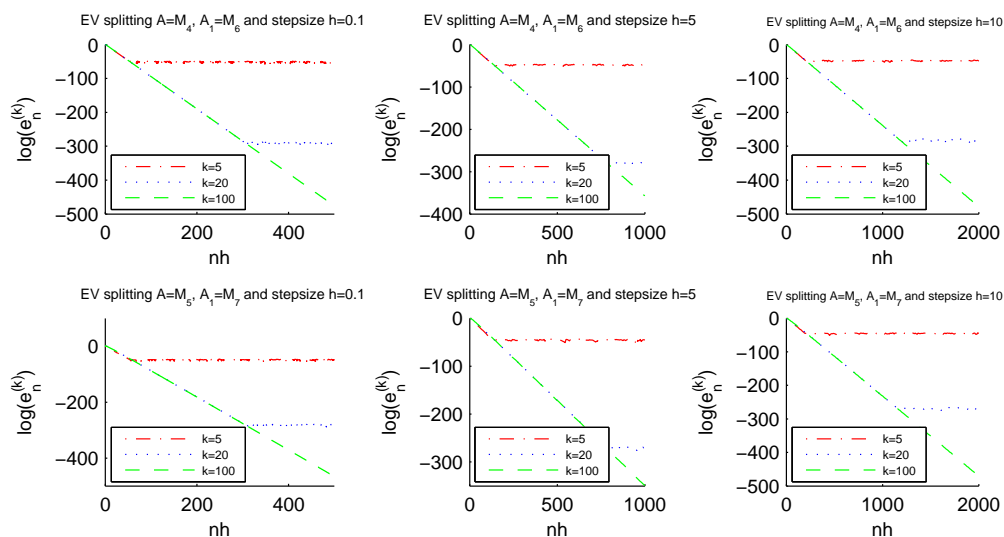


Figure 4. Numerical simulation of WR method (3.2) with the EV splitting.

One may be interested to know about the stability of the WR methods with $A = M_4$ or M_5 when another splitting is used. Here, we provide insight into the problem. Taking the GJ and GS splittings of $A = M_4$ and M_5 , we have obtained two groups of the numerical solutions of WR methods (3.1) and (3.2) displayed in Figures 5 and 6, which show that the WR (3.1) and (3.2) with $A = M_4$ or M_5 are unstable for the GJ and GS splittings.

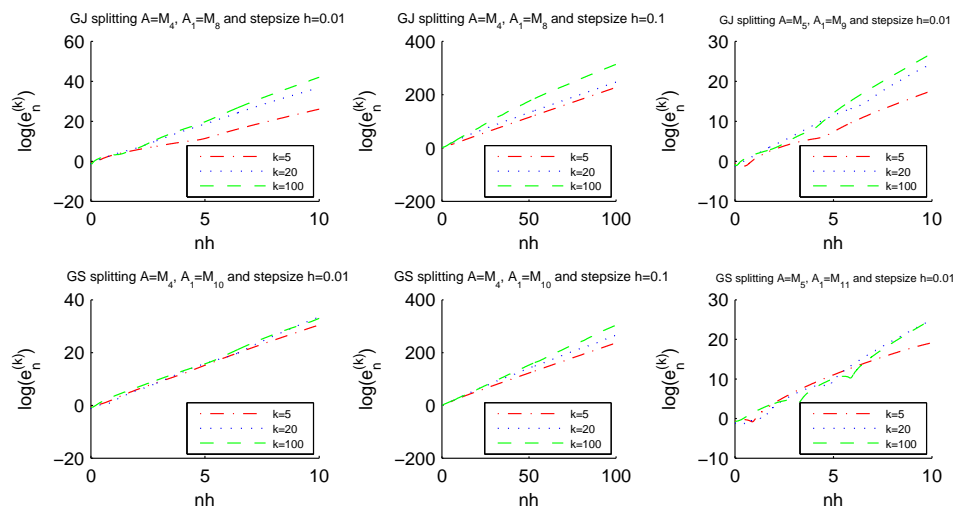


Figure 5. Numerical simulation of WR method (3.1) with GJ and GS splitting.

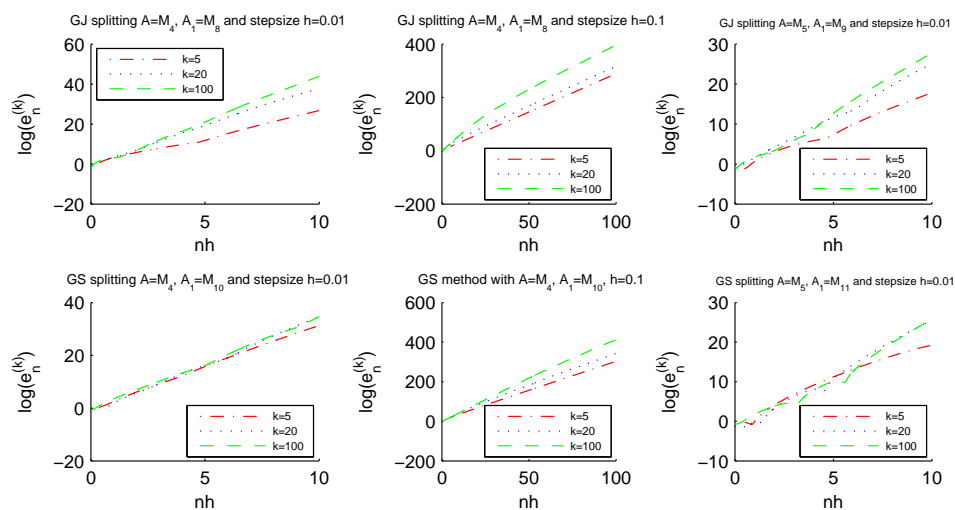


Figure 6. Numerical simulation of WR method (3.2) with GJ and GS splitting.

One may also want to know whether or not the instability of WR methods originates purely from the splitting function or from the underlying explicit Euler method. In Sections 3 and 4 only some sufficient conditions of stability are obtained for the WR methods, which are not enough to deduce the source of instability in WR methods. However, some reasonable conclusions can be reached by performing numerical experiments.

For the explicit Euler method of (2.1) there exists the maximum step size h_{\max} such that the method is stable if and only if the step size $h < h_{\max}$. It seems that there also exists a similar h_{\max} for the WR method (3.1) of (2.1) by Theorem 3.1 and Corollaries 3.4 and 3.6. We have explored whether or not h_{\max} of (3.1) is the same as that of the underlying explicit Euler method by conducting numerical experiments. For comparison in a uniform standard, we take $\varepsilon_0^{(k)} = \varepsilon_n^{(0)} = (1, 1, \dots, 1)^T$ for all $k, n \in \mathbb{Z}^+$.

The approximate values of h_{\max} for the explicit Euler method and the WR method (3.1) as applied to (2.1) have been obtained, which are listed in Tables 1–3, for the following three cases:

Case 1. $A = -3, A_1 = \alpha, A_2 = A - A_1;$

Case 2. $A = M_1, A_1 = \begin{pmatrix} -5 & & & & \\ \alpha & -5 & & & \\ & & -6 & & \\ & & & -5 & \\ & & & & -6 \end{pmatrix}, A_2 = A - A_1;$

Case 3. $A = M_5, A_1 = \begin{pmatrix} -1 & & & & \\ & \alpha & & & \\ & & & & \\ & & & & \\ & & & & -1 \end{pmatrix}, A_2 = A - A_1.$

From Table 1, we see that the h_{\max} of the WR method (3.1) for Case 1 and of the underlying explicit Euler method are the same; particularly, the h_{\max} of (3.1) does not change with the splitting, i.e., according to the value of α . Hence, the instability of (3.1) for Case 1 originates purely from underlying explicit Euler methods and is independent of the splitting. However, from Tables 2 and 3, as well as the comparison of Figures 3 and 5, we can find convincing evidence of a link between the instability of (3.1) and the splitting matrix used. Consequently, the stability of (3.1) for Cases 2 or 3 may change with the splitting.

Table 1. The maximum step size for the stability of numerical methods of (2.1) for Case 1 and the explicit Euler method (EEM).

Meth.	EEM	WR method (3.1) for Case 1					
		$\alpha = -10$	$\alpha = -2$	$\alpha = -1$	$\alpha = 1$	$\alpha = 2$	$\alpha = 10$
h_{\max}	2/3	1.66	1.66	1.66	1.66	1.66	1.66

Table 2. The maximum step size for the stability of numerical methods of (2.1) for Case 2, the EEM and the WR method with BGJ or BGS splitting (WRwBJoBS).

Meth.	EEM	WRwBJoBS	WR method (3.1) for Case 2				
			$\alpha = 0$	$\alpha = 1$	$\alpha = 2$	$\alpha = 3$	$\alpha = 5$
h_{\max}	0.24	0.24	0.24	0.24	0.22	0.21	< 0.01

Table 3. The maximum step size for the stability of numerical methods of (2.1) for Case 3 and the EEM.

Meth.	EEM	WR method (3.1) for Case 3			
		$\alpha = -1$	$\alpha = -0.9$	$\alpha = -0.85$	$\alpha = 0.8$
h_{\max}	1.99	1.99	1.33	1.29	< 0.01

Hence, we reasonably conclude that the instability of WR method (3.1) originates purely from neither the splitting function used nor the underlying explicit Euler method.

6. Conclusions

In this paper, we have discussed the stability properties of DWR methods based on Euler schemes. We pointed out that the DWR methods may be unstable by virtue of a numerical example. There is hence a need to study the stability of DWR methods. In this paper, we have applied DWR methods based on Euler schemes to two dissipative systems and analyzed the properties of the numerical solutions generated. We then obtained some conditions under which the DWR methods can preserve the stability of the exact solutions. These conditions show that underlying methods and splitting ways are two key factors in determining the stability of DWR methods.

Use of AI tools declaration

The authors declare they have not used Artificial Intelligence (AI) tools in the creation of this article.

Acknowledgments

This work was supported by the Natural Science Foundation of Fujian Province, China (grant number 2021J011031) and Research Project of Fashu Foundation (MFK23013).

Conflict of interest

The authors declare no conflict of interest.

References

1. E. Lelarasmee, A. E. Ruehli, L. Sangiovanni-Vincentelli, The waveform relaxation method for time-domain analysis of large scale integrated circuits, *IEEE T. Comput. Aid. D.*, **1** (1982), 131–145. <http://doi.org/10.1109/TCAD.1982.1270004>
2. K. J. in't Hout, On the convergence waveform relaxation methods for stiff nonlinear ordinary differential equations, *Appl. Numer. Math.*, **18** (1995), 175–190. [http://doi.org/10.1016/0168-9274\(95\)00052-v](http://doi.org/10.1016/0168-9274(95)00052-v)
3. Z. Jackiewicz, M. Kwapisz, Convergence of waveform relaxation methods for differential algebraic systems, *SIAM J. Numer. Anal.*, **33** (1996), 2303–2317. <http://doi.org/10.1137/S0036142992233098>
4. R. Jeltsch, B. Pohl, Waveform relaxation with overlapping splittings, *SIAM J. Sci. Comput.*, **16** (1995), 40–49. <http://doi.org/10.1137/0916004>
5. D. Conte, R. D'Ambrosio, B. Paternoster, GPU-acceleration of waveform relaxation methods for large differential systems, *Numer. Algor.*, **71** (2016), 293–310. <http://doi.org/10.1007/s11075-015-9993-6>
6. M. R. Crisci, N. Ferraro, E. Russo, Convergence results for continuous-time waveform methods for Volterra integral equations, *J. Comput. Appl. Math.*, **71** (1996), 33–45. [http://doi.org/10.1016/0377-0427\(95\)00225-1](http://doi.org/10.1016/0377-0427(95)00225-1)

7. Z. Hassanzadeh, D. K. Salkuyeh, Two-stage waveform relaxation method for the initial value problems with non-constant coefficients, *Comp. Appl. Math.*, **33** (2014), 641–654. <http://doi.org/10.1007/s40314-013-0086-7>
8. Y. L. Jiang, Windowing waveform relaxation of initial value problems, *Acta Math. Appl. Sin, Engl. Ser.*, **22** (2006), 575–588. <http://doi.org/10.1007/s10255-006-0331-6>
9. J. Sand, K. Burrage, A Jacobi waveform relaxation method for ODEs, *SIAM J. Sci. Comput.*, **20** (1998), 534–552. <http://doi.org/10.1137/S1064827596306562>
10. X. Yang, On solvability and waveform relaxation methods of linear variable-coefficient differential-algebraic equations, *J. Comp. Math.*, **32** (2014), 696–720. <http://doi.org/10.4208/jcm.1405-m4417>
11. A. Bellen, Z. Jackiewicz, M. Zennaro, Stability analysis of time-point relaxation Heun method, Report “Progetto finalizzato sistemi informatici e calcolo parallelo”, University of Trieste, 1990.
12. J. K. M. Jansen, R. M. M. Mattheij, M. T. M. Penders, W. H. A. Schilders, Stability and efficiency of waveform relaxation methods, *Comput. Math. Appl.*, **28** (1994), 153–166. [http://doi.org/10.1016/0898-1221\(94\)00103-0](http://doi.org/10.1016/0898-1221(94)00103-0)
13. A. Bellen, Z. Jackiewicz, M. Zennaro, Time-point relaxation Runge-Kutta methods for ordinary differential equations, *J. Comput. Appl. Math.*, **45** (1993), 121–137. [http://doi.org/10.1016/0377-0427\(93\)90269-H](http://doi.org/10.1016/0377-0427(93)90269-H)
14. A. Bellen, Z. Jackiewicz, M. Zennaro, Contractivity of waveform relaxation Runge-Kutta iterations and related limit methods for dissipative systems in the maximum norm, *SIAM J. Numer. Anal.*, **31** (1994), 499–523. <http://doi.org/10.1137/0731027>
15. E. Blåsten, H. Liu, Recovering piecewise constant refractive indices by a single far-field pattern, *Inverse Probl.*, **36** (2020), 085005. <http://doi.org/10.1088/1361-6420/ab958f>
16. E. L. K. Blåsten, H. Liu, Scattering by curvatures, radiationless sources, transmission eigenfunctions, and inverse scattering problems, *SIAM J. Math. Anal.*, **53** (2021), 3801–3837. <http://doi.org/10.1137/20M1384002>
17. H. Liu, M. Petrini, L. Rondi, J. Xiao, Stable determination of sound-hard polyhedral scatterers by a minimal number of scattering measurements, *J. Differ. Equations*, **262** (2017), 1631–1670. <http://doi.org/10.1016/J.JDE.2016.10.021>
18. H. Liu, L. Rondi, J. Xiao, Mosco convergence for $H(\text{curl})$ spaces, higher integrability for Maxwell’s equations, and stability in direct and inverse EM scattering problems, *J. Eur. Math. Soc.*, **21** (2019), 2945–2993. <http://doi.org/10.4171/JEMS/895>
19. Y. Chow, Y. Deng, Y. He, H. Liu, X. Wang, Surface-localized transmission eigenstates, super-resolution imaging, and pseudo surface plasmon modes, *SIAM J. Imaging Sci.*, **14** (2021), 946–975. <https://doi.org/10.1137/20M1388498>
20. H. Diao, X. Cao, H. Liu, On the geometric structures of transmission eigenfunctions with a conductive boundary condition and applications, *Commun. Part. Diff. Eq.*, **46** (2021), 630–679. <https://doi.org/10.1080/03605302.2020.1857397>
21. H. Li, J. Z. Li, H. Liu, On quasi-static cloaking due to anomalous localized resonance in \mathbb{R}^3 , *SIAM J. Appl. Math.*, **75** (2015), 1245–1260. <https://doi.org/10.1137/15M1009974>

22. J. Li, H. Liu, Q. Wang, Enhanced multilevel linear sampling methods for inverse scattering problems, *J. Comput. Phys.*, **257** (2014), 554–571. <https://doi.org/10.1016/j.jcp.2013.09.048>
23. J. Li, H. Liu, Y. Wang, Recovering an electromagnetic obstacle by a few phaseless backscattering measurements, *Inverse Probl.*, **33** (2017), 035011. <http://doi.org/10.1088/1361-6420/aa5bf3>
24. H. Liu, A global uniqueness for formally determined inverse electromagnetic obstacle scattering, *Inverse Probl.*, **24** (2008), 035018. <http://doi.org/10.1088/0266-5611/24/3/035018>
25. X. Wang, Y. Guo, S. Bousba, Direct imaging for the moment tensor point sources of elastic waves, *J. Comput. Phys.*, **448** (2022), 110731. <https://doi.org/10.1016/j.jcp.2021.110731>
26. J. M. Varah, A lower bound for the smallest singular value of a matrix, *Linear Algebra Appl.*, **11** (1975), 3–5. [http://doi.org/10.1016/0024-3795\(75\)90112-3](http://doi.org/10.1016/0024-3795(75)90112-3)



AIMS Press

©2023 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (<http://creativecommons.org/licenses/by/4.0>)