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**Research article**

## New Einstein-Randers metrics on certain homogeneous manifolds arising from the generalized Wallach spaces

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**Abstract:** In this article, we find several new non-Riemannian Einstein-Randers metrics on some homogeneous manifolds arising from the generalized Wallach spaces. We first prove the existence of Riemannian Einstein metrics on these homogeneous manifolds. Based on these metrics, we prove that there exist non-Riemannian Einstein-Randers metrics on these homogeneous manifolds.

**Keywords:** Riemannian Einstein metrics; Einstein-Randers metrics; Finsler manifolds; homogeneous manifolds; the generalized Wallach spaces

**Mathematics Subject Classification:** 53C25, 53C30

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### 1. Introduction

Randers metrics were introduced by Randers in the context of general relativity, and later named after him by Ingarden. A Randers metric  $F$  is built from a Riemannian metric and a 1-form, i.e.,

$$F = \alpha + \beta$$

where  $\alpha$  is a Riemannian metric and  $\beta$  is a 1-form and the length of  $\beta$  corresponding to the Riemannian metric  $\alpha$  is less than 1 everywhere.

Then,  $F$  is Riemannian if and only if  $F(x, y) = F(x, -y)$ . But sometimes it is convenient to use the following definition of a Randers metric in [2], i.e.,

$$F(x, y) = \frac{\sqrt{[h(W, y)]^2 + h(y, y)\lambda}}{\lambda} - \frac{h(W, y)}{\lambda}. \quad (1.1)$$

Here,  $\lambda = 1 - h(W, W) > 0$ , and  $F$  is Riemannian if and only if  $W = 0$ . The pair  $(h, W)$  is called the navigation data of the corresponding Randers metric  $F$ .

The Ricci scalar  $Ric(x, y)$  of a Finsler metric is defined by the sum of those  $n - 1$  flag curvatures  $K(x, y, e_v)$ , where  $\{e_v : v = 1, 2, \dots, n - 1\}$  is any collection of  $n - 1$  orthonormal transverse edges perpendicular to the flagpole, i.e.,

$$Ric(x, y) = \sum_{v=1}^{n-1} R_{vv}. \quad (1.2)$$

The Ricci tensor is defined by

$$Ric_{ij} = (\frac{1}{2} F^2 Ric)_{y^i y^j}. \quad (1.3)$$

The Ricci scalar depends on the position  $x$  and the flagpole  $y$  (see [2, 3]), but does not depend on the specific  $n - 1$  flags with transverse edges orthogonal to  $y$ . It is well known that the Ricci scalar in Riemannian geometry only depends on  $x$ . Thus, it is quite interesting to study a Finsler manifold whose Ricci scalar does not depend on the flagpole  $y$ . Generally, a Finsler metric with such a property is called an Einstein metric, i.e.,

$$Ric(x, y) = (n - 1)K(x) \quad (1.4)$$

for some function  $K(x)$  on  $M$ . In particular, for a Randers manifold  $(M, F)$  with  $\dim M \geq 3$ ,  $F$  is an Einstein metric if and only if  $K(x)$  is a constant on  $M$  (see [2]). The following lemma is an important result on Einstein-Randers metrics.

**Lemma 1.1** ([2]). *Suppose that  $(M, F)$  is a Randers space with the navigation data  $(h, W)$ . Then,  $(M, F)$  is Einstein with Ricci scalar  $Ric(x) = (n - 1)K(x)$  if and only if there exists a constant  $\sigma$  satisfying the following conditions:*

- 1)  $h$  is Einstein with Ricci scalar  $(n - 1)K(x) + \frac{1}{16}\sigma^2$ , and
- 2)  $W$  is an infinitesimal homothety of  $h$ , i.e.,  $\Omega_W h = -\sigma h$ .

Furthermore,  $\sigma$  must be zero whenever  $h$  is not Ricci-flat.

It is well known that  $K(x)$  is a constant if  $(M, F)$  is a homogeneous Einstein Finsler manifold. Here, a Finsler manifold  $(M, F)$  is called homogeneous if its full group of isometry acts transitively on  $M$ . Based on Lemma 1.1, Deng-Hou obtained a characterization of homogeneous Einstein-Randers metrics.

**Lemma 1.2** ([8]). *Let  $G$  be a connected Lie group and  $H$  a closed subgroup of  $G$  such that  $G/H$  is a reductive homogeneous space with a decomposition  $\mathfrak{g} = \mathfrak{h} + \mathfrak{m}$ . Suppose that  $h$  is a  $G$ -invariant Riemannian metric on  $G/H$  and  $W \in \mathfrak{m}$  is invariant under  $H$  with  $h(W, W) < 1$ . Let  $\tilde{W}$  be the corresponding  $G$ -invariant vector field on  $G/H$  with  $\tilde{W}|_o = W$ . Then, the Randers metric  $F$  with the navigation data  $(h, \tilde{W})$  is Einstein with the Ricci constant  $K$  if and only if  $h$  is Einstein with the Ricci constant  $K$  and  $W$  satisfies*

$$\langle [W, X]_{\mathfrak{m}}, Y \rangle + \langle X, [W, Y]_{\mathfrak{m}} \rangle = 0, \quad \forall X, Y \in \mathfrak{m}, \quad (1.5)$$

where  $\langle , \rangle$  is the restriction of  $h$  on  $T_o(G/H) \simeq \mathfrak{m}$ . In this case,  $\tilde{W}$  is a Killing vector field with respect to the Riemannian metric  $h$ .

Just as in the Riemannian case, it is a fundamental problem to classify homogeneous Einstein Finsler spaces. In particular, it is very important to know if a homogeneous manifold admits invariant Einstein

Finsler metrics. In fact, there are many studies on homogenous Einstein Randers manifolds [5, 6, 9–11, 14–17] and a little progress on homogeneous Einstein  $(\alpha, \beta)$ -metrics [18].

The main goal of this paper is to find non-Riemannian Einstein-Randers metrics on homogeneous manifolds  $G/H$  for  $G = E_7$  and  $G = E_6$ , which means we prove the following theorem.

**Theorem 1.3.** *1) There are at least six families of  $G$ -invariant non-Riemannian Einstein-Randers metrics on the homogeneous manifold  $E_7/A_1$ .*

- 2) There are at least four families of  $G$ -invariant non-Riemannian Einstein-Randers metrics on the homogeneous manifold  $E_7/A_5$ .*
- 3) There are at least two families of  $G$ -invariant non-Riemannian Einstein-Randers metrics on the homogeneous manifold  $E_7/(A_1 \times A_5)$ .*
- 4) There are at least ten families of  $G$ -invariant non-Riemannian Einstein-Randers metrics on the homogeneous manifold  $E_6/A_1$ .*
- 5) There are at least four families of  $G$ -invariant non-Riemannian Einstein-Randers metrics on the homogeneous manifold  $E_6/A_3$ .*
- 6) There are at least four families of  $G$ -invariant non-Riemannian Einstein-Randers metrics on the homogeneous manifold  $E_6/(A_1 \times A_3)$ .*
- 7) There are at least two families of  $G$ -invariant non-Riemannian Einstein-Randers metrics on the homogeneous manifold  $E_6/(A_1 \times A_1)$ .*
- 8) There is at least one family of  $G$ -invariant non-Riemannian Einstein-Randers metric on the homogeneous manifold  $E_6/(A_1 \times A_1 \times A_3)$ .*

## 2. The Ricci tensor for reductive homogeneous spaces

Let  $G$  be a compact semi-simple Lie group with Lie algebra  $\mathfrak{g}$ ,  $K$  a connected closed subgroup of  $G$  with Lie algebra  $\mathfrak{k}$ . Through this paper, we denote by  $B$  the negative of the Killing form of  $\mathfrak{g}$ , which is positive definite because of the compactness of  $G$ . As a result,  $B$  can be treated as an inner product on  $\mathfrak{g}$ . Let  $\mathfrak{g} = \mathfrak{k} \oplus \mathfrak{m}$  be the reductive decomposition with respect to  $B$  such that  $[\mathfrak{k}, \mathfrak{m}] \subset \mathfrak{m}$ , where  $\mathfrak{m}$  is the tangent space of  $G/K$ . We assume that  $\mathfrak{m}$  can be decomposed into mutually non-equivalent irreducible  $\text{Ad}(K)$ -modules as follows:

$$\mathfrak{m} = \mathfrak{m}_1 \oplus \cdots \oplus \mathfrak{m}_q.$$

We will write  $\mathfrak{k} = \mathfrak{k}_0 \oplus \mathfrak{k}_1 \oplus \cdots \oplus \mathfrak{k}_p$ , where  $\mathfrak{k}_0 = Z(\mathfrak{k})$  is the center of  $\mathfrak{k}$  and  $\mathfrak{k}_i$  is the simple ideal for  $i = 1, \dots, p$ . Let  $G \times K$  act on  $G$  by  $(g_1, g_2)g = g_1gg_2^{-1}$ . Then,  $G \times K$  acts almost effectively on  $G$  with isotropy group  $\Delta(K) = \{(k, k) | k \in K\}$ . As a result,  $G$  can be treated as the coset space  $(G \times K)/\Delta(K)$  and we have  $\mathfrak{g} \oplus \mathfrak{k} = \Delta(\mathfrak{k}) \oplus \Omega$ , where  $\Omega \cong T_0((G \times K)/\Delta(K)) \cong \mathfrak{g}$  via the linear map  $(X, Y) \rightarrow (X - Y) \in \mathfrak{g}$ ,  $(X, Y) \in \Omega$ .

A direct conclusion is that there exists a 1-1 corresponding between all  $G$ -invariant metrics on the reductive homogeneous space  $G/K$  and  $\text{Ad}_G(K)$ -invariant inner products on  $\mathfrak{m}$ . Now, we have a orthogonal decomposition of  $\mathfrak{g}$  with respect to the Killing form of  $\mathfrak{g}$ :  $\mathfrak{g} = \mathfrak{k}_0 \oplus \mathfrak{k}_1 \oplus \cdots \oplus \mathfrak{k}_p \oplus \mathfrak{m}_1 \oplus \cdots \oplus \mathfrak{m}_q = (\mathfrak{k}_0 \oplus \mathfrak{k}_1 \oplus \cdots \oplus \mathfrak{k}_p) \oplus (\mathfrak{k}_{p+1} \oplus \cdots \oplus \mathfrak{k}_{p+q})$ , with  $\mathfrak{m}_1 \oplus \cdots \oplus \mathfrak{m}_q = \mathfrak{k}_{p+1} \oplus \cdots \oplus \mathfrak{k}_{p+q}$ . In addition, we assume that  $\dim_{\mathbb{R}} \mathfrak{k}_0 \leq 1$  and the ideals  $\mathfrak{k}_i$  are mutually non-isomorphic for  $i = 1, \dots, p$ . Then, we consider all  $\text{Ad}(K)$ -invariant metrics on  $G$  of the form:

$$\langle \cdot, \cdot \rangle = x_0 \cdot B|_{\mathfrak{k}_0} + x_1 \cdot B|_{\mathfrak{k}_1} + \cdots + x_{p+q} \cdot B|_{\mathfrak{k}_{p+q}}, \quad (2.1)$$

where  $x_i \in \mathbb{R}^+$  for  $i = 0, 1, \dots, p+q$ . And all  $\text{Ad}(K)$ -invariant metrics on  $G/K$  of the form:

$$(,) = x_{p+1} \cdot B|_{\mathfrak{k}_{p+1}} + \dots + x_{p+q} \cdot B|_{\mathfrak{k}_{p+q}}, \quad (2.2)$$

where  $x_i \in \mathbb{R}^+$  for  $i = p+1, \dots, p+q$ .

Set from now on  $d_i = \dim_{\mathbb{R}} \mathfrak{k}_i$  and  $\{e_\alpha^i\}_{\alpha=1}^{d_i}$  be a  $B$ -orthonormal basis adapted to the decomposition of  $\mathfrak{g}$  which means  $e_\alpha^i \in \mathfrak{k}_i$  and  $\alpha$  is the number of basis in  $\mathfrak{k}_i$ . Then, we consider the numbers  $A_{\alpha,\beta}^\gamma = B([e_\alpha^i, e_\beta^j], e_\gamma^k)$  such that  $[e_\alpha^i, e_\beta^j] = \sum_\gamma A_{\alpha,\beta}^\gamma e_\gamma^k$ , and set

$$(ijk) := \left[ \begin{array}{c} i \\ j \ k \end{array} \right] = \sum (A_{\alpha,\beta}^\gamma)^2,$$

where the sum is taken over all indices  $\alpha, \beta, \gamma$  with  $e_\alpha^i \in \mathfrak{k}_i, e_\beta^j \in \mathfrak{k}_j, e_\gamma^k \in \mathfrak{k}_k$ . Then,  $(ijk)$  is independent of the choice for the  $B$ -orthonormal basis of  $\mathfrak{k}_i, \mathfrak{k}_j, \mathfrak{k}_k$ , and symmetric for all three indices which means  $(ijk) = (jik) = (jki)$ .

In [1] and [13], the authors obtained the formulas for the components of the Ricci tensor with respect to the left-invariant metric given by (2.1), which can be described by the following lemma.

**Lemma 2.1.** *Let  $G$  be a compact connected semi-simple Lie group endowed with the left-invariant metric  $\langle , \rangle$  given by (2.1). Then, the components  $r_0, r_1, \dots, r_{p+q}$  of the Ricci tensor associated to  $\langle , \rangle$  are expressed as follows:*

$$r_k = \frac{1}{2x_k} + \frac{1}{4d_k} \sum_{j,i} \frac{x_k}{x_j x_i} \left[ \begin{array}{c} k \\ j \ i \end{array} \right] - \frac{1}{2d_k} \sum_{j,i} \frac{x_j}{x_k x_i} \left[ \begin{array}{c} j \\ k \ i \end{array} \right], \quad (k = 0, 1, \dots, p+q).$$

Here, the sums are taken over all  $i = 0, 1, \dots, p+q$ . In particular, for each  $k$  it holds that

$$\sum_{i,j}^{p+q} \left[ \begin{array}{c} j \\ k \ i \end{array} \right] = \sum_{ij}^{p+q} (kij) = d_k.$$

For  $k = p+1, \dots, p+q$  and by considering the sums appearing in the expression of  $r_k$  only for  $i, j$  with  $p+1 \leq i, j \leq p+q$ , one obtains the components  $\hat{r}_{p+1}, \dots, \hat{r}_{p+q}$  of the Ricci tensor  $\hat{r}$  of the  $G$ -invariant metric  $(,)$  on  $G/K$  defined by (2.2).

### 3. Decompositions arising from generalized Wallach spaces

We recall the definition of generalized Wallach spaces. Let  $G/K$  be a reductive homogeneous space, where  $G$  is a semi-simple compact connected Lie group,  $K$  is a connected closed subgroup of  $G$ ,  $\mathfrak{g}$  and  $\mathfrak{k}$  are the corresponding Lie algebras, respectively. If  $\mathfrak{m}$ , the tangent space of  $G/K$  at  $o = \pi(e)$ , can be decomposed into three  $\text{ad}(\mathfrak{k})$ -invariant irreducible summands pairwise orthogonal with respect to  $B$  as:

$$\mathfrak{m} = \mathfrak{m}_1 \oplus \mathfrak{m}_2 \oplus \mathfrak{m}_3,$$

satisfying  $[\mathfrak{m}_i, \mathfrak{m}_i] \in \mathfrak{k}$  for  $i \in \{1, 2, 3\}$  and  $[\mathfrak{m}_i, \mathfrak{m}_j] \subset \mathfrak{m}_k$  for  $\{i, j, k\} = \{1, 2, 3\}$ , then we call  $G/K$  a generalized Wallach space.

In 2014, classification for generalized Wallach spaces arising from a compact simple Lie group has been obtained by Nikonorov [12] and Chen, Kang and Liang [7], in particular, Nikonorov investigated the semi-simple case and gave the classification in [12]. For convenience, in this article, we use the notations in [7].

According to [7], each kind of the classification corresponds to two commutative involutive automorphisms of  $\mathfrak{g}$ , the Lie algebra of the compact simple Lie group  $G$ . In [4], the authors calculated all the coefficients  $(ijk)$  in the expression for the components of Ricci tensor with respect to the metric of form (2.1). In [7],  $E_6\text{-II}$  and  $E_7\text{-II}$  have the following decomposition in the view of Lie algebra:

$$\begin{aligned} E_6 &= T \oplus A_1^1 \oplus A_1^2 \oplus A_3 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \\ E_7 &= T \oplus A_1 \oplus A_5 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3. \end{aligned}$$

By the notations in the above section, we give the following two lemmas in [4] for later use.

**Lemma 3.1.** *In the case of  $E_7\text{-II}$ , the possible non-zero coefficients in the expression for the components of Ricci tensor with respect to the metric of the form (2.1) are as follows:*

$$\begin{aligned} (033) &= \frac{4}{9}, (044) = \frac{5}{9}, (055) = 0, (345) = \frac{20}{3}, \\ (111) &= \frac{1}{3}, (133) = 1, (144) = 0, (155) = \frac{5}{3}, \\ (222) &= \frac{35}{3}, (233) = \frac{35}{9}, (244) = \frac{70}{9}, (255) = \frac{35}{3}. \end{aligned}$$

**Lemma 3.2.** *In the case of  $E_6\text{-II}$ , the possible non-zero coefficients in the expression for the components of Ricci tensor with respect to the metric of the form (2.1) are as follows:*

$$\begin{aligned} (044) &= 1/2, (055) = 1/2, (066) = 0, (456) = 4, \\ (111) &= 1/2, (144) = 0, (155) = 1, (166) = 3/2, \\ (222) &= 1/2, (244) = 1, (255) = 0, (266) = 3/2, \\ (333) &= 5, (344) = 5/2, (355) = 5/2, (366) = 5. \end{aligned}$$

#### 4. Einstein metrics on homogenous manifolds

In this section, we will consider the Einstein metrics on some homogeneous manifolds which can be obtained from the generalized Wallach spaces, namely  $E_7/A_1$ ,  $E_7/A_5$ ,  $E_7/(A_1 \times A_5)$ ,  $E_6/A_1$ ,  $E_6/A_3$ ,  $E_6/(A_1 \times A_3)$ ,  $E_6/(A_1 \times A_1)$  and  $E_6/(A_1 \times A_1 \times A_3)$ .

**Case of  $E_7/A_1$ .** Consider the homogeneous manifold  $E_7/A_1$  with the decomposition

$$\mathfrak{m} = T \oplus A_5 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.1)$$

and  $\text{Ad}(A_1)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1 \oplus A_5)$ -invariant on  $E_7/A_1$  defined by

$$\langle , \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{A_5} + x_2 \cdot B|_{\mathfrak{p}_1} + x_3 \cdot B|_{\mathfrak{p}_2} + x_4 \cdot B|_{\mathfrak{p}_3}, \quad (4.2)$$

where  $u_0, x_1, x_2, x_3, x_4 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.1, the components of Ricci tensor with respect to the metric (4.2) are as follows:

$$\begin{cases} r_T = \frac{1}{4} \left( \frac{4u_0}{9x_2^2} + \frac{5u_0}{9x_3^2} \right), \\ r_{A_5} = \frac{1}{140} \left( \frac{35}{3x_1} + \frac{35x_1}{9x_2^2} + \frac{70x_1}{9x_3^2} + \frac{35x_1}{3x_4^2} \right), \\ r_{p_1} = \frac{1}{2x_2} + \frac{5}{36} \left( \frac{x_2}{x_3x_4} - \frac{x_3}{x_4x_2} - \frac{x_4}{x_2x_3} \right) - \frac{1}{48} \left( \frac{4u_0}{9x_2^2} + \frac{35x_1}{9x_2^2} \right), \\ r_{p_2} = \frac{1}{2x_3} + \frac{1}{9} \left( \frac{x_3}{x_4x_2} - \frac{x_4}{x_2x_3} - \frac{x_2}{x_3x_4} \right) - \frac{1}{60} \left( \frac{5u_0}{9x_3^2} + \frac{70x_1}{9x_3^2} \right), \\ r_{p_3} = \frac{1}{2x_4} + \frac{1}{12} \left( \frac{x_4}{x_2x_3} - \frac{x_2}{x_3x_4} - \frac{x_3}{x_4x_2} \right) - \frac{1}{80} \frac{35x_1}{3x_4^2}. \end{cases}$$

We consider the homogeneous Einstein equation as follows:

$$\{r_T - r_{A_5} = 0, r_{A_5} - r_{p_1} = 0, r_{p_1} - r_{p_2} = 0, r_{p_2} - r_{p_3} = 0\}$$

and it turns out to be equivalent to the following system of equations (we normalize the metric by setting  $u_0 = 1$ ).

$$\begin{cases} g_0 = -3x_1^2x_2^2x_3^2 - 2x_1^2x_2^2x_4^2 - x_1^2x_3^2x_4^2 - 3x_2^2x_3^2x_4^2 + 5x_1x_2^2x_4^2 + 4x_1x_3^2x_4^2 = 0, \\ g_1 = 36x_1^2x_2^2x_3^2 + 24x_1^2x_2^2x_4^2 + 47x_1^2x_3^2x_4^2 - 60x_1x_2^3x_3x_4 + 60x_1x_2x_3^3x_4 - 216x_1x_2x_3^2x_4^2 \\ + 60x_1x_2x_3x_4^3 + 36x_2^2x_3^2x_4^2 + 4x_1x_3^2x_4^2 = 0, \\ g_2 = 56x_1x_2^2x_4 - 35x_1x_3^2x_4 + 108x_2^3x_3 - 216x_2^2x_3x_4 - 108x_2x_3^3 + 216x_2x_3^2x_4 - 12x_2x_3x_4^2 \\ + 4x_2^2x_4 - 4x_3^2x_4 = 0, \\ g_3 = 63x_1x_2x_3^2 - 56x_1x_2x_4^2 - 12x_2^2x_3x_4 - 216x_2x_3^2x_4 + 216x_2x_3x_4^2 + 84x_3^3x_4 - 84x_3x_4^3 \\ - 4x_2x_4^2 = 0. \end{cases}$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3, x_4]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, g_3, zx_1x_2x_3x_4 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3 > x_4$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial of  $x_4$  given by  $(20x_4^2 - 72x_4 + 49) \cdot h(x_4)$ , where  $h(x_4)$  is a polynomial of degree 30 given by

$$\begin{aligned}
h(x_4) = & 147263121465141201705316711269370210741339344914420477996120017678192287817728000 x_4^{30} \\
& - 1102576650209541754792174768351559525030817374433243715583773603796209924867276800 x_4^{29} \\
& + 4392608446633715641642277193442303841356600810515312912956157487998587768223467520 x_4^{28} \\
& - 16381030403148760413861847066334150627671304859530037939042221945751621483816924672 x_4^{27} \\
& + 44632927369562126482344261311667463916531381076564814811273737776429261123742947776 x_4^{26} \\
& - 105145519130201191374375907381751124532181350157935176774829297302755654321119817568 x_4^{25} \\
& + 215087055560353068666721180840479776113190447174940731329033435251254151873127864288 x_4^{24} \\
& - 34451837402804811107041226967326977788246179258239908939790910388711063175367161784 x_4^{23} \\
& + 445879403133315681874085042923629294832116921460904234436654259616839387520348170596 x_4^{22} \\
& - 45400315073071285545752682186102790358057291331388888785553759520405925035444447080 x_4^{21} \\
& + 308171762729585306636518011606003014325716616667221520922199158892726288206642411702 x_4^{20} \\
& - 88928520907166973360693993136042781786034406240792078393993953506206913911288935648 x_4^{19} \\
& - 40034430463545533917890194187357087475825295381916586084753660204634626955000430908 x_4^{18} \\
& + 44518815422796823726900600056989666146339579404372825641721394763313021826539183294 x_4^{17} \\
& - 13231125163887960569832230144624600324399606190969212487314466195084966981574676247 x_4^{16} \\
& - 797439538169345403753233795452611566212570589018415600726606800661093425142878880 x_4^{15} \\
& + 1511979559539474891852062236262989481312808910795807182871248144865160620831183632 x_4^{14} \\
& - 379199106951137596458809929302879785793133656913117463400584039552474682406503552 x_4^{13} \\
& + 3779714013977548775186439810978146278678833555681219090436473692966448564191712 x_4^{12} \\
& - 1785948980557995395564396777431636788073937513566295662291234226016378838797312 x_4^{11} \\
& + 361185586014234569587833784293818501106527299574429579442258565335789303734400 x_4^{10} \\
& - 29814163618672798385666919644831689684183641996050223223894494904315006447104 x_4^9 \\
& + 507366481118078491288306527464440790802333871137796008412526195301608704 x_4^8 \\
& + 55049909220845662880254271319404494752671319946031114132616313964384415744 x_4^7 \\
& + 3753669713205988245353403631211241070292590650176001229552989542586368 x_4^6 \\
& - 461850039539407240955273834215202137394483687360075268757733061237596160 x_4^5 \\
& + 14855590404155632214609581796329294787394060775182341751503634575953920 x_4^4 \\
& - 151882382457983307126403603474158205457944387452892184268979018137600 x_4^3 \\
& + 995700410381405891358618695307405787784707293649196430275379200000 x_4^2 \\
& - 13417336677067637864115045789201884045646775900717868981944320000 x_4 \\
& + 568400491719774288687807984496995610757827345058316786073600000.
\end{aligned}$$

Solving  $h(x_4) = 0$  numerically, we see that there exist four positive solutions, which are given approximately by  $x_4 \approx 0.4025066142$ ,  $x_4 \approx 1.032699036$ ,  $x_4 \approx 1.438454385$ ,  $x_4 \approx 4.052215439$ . Further, we remark that  $x_1$ ,  $x_2$  and  $x_3$  can be expressed by polynomials of  $x_4$ , thus the corresponding solutions of the system  $\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, h(x_4) = 0\}$  with  $x_1 x_2 x_3 x_4 \neq 0$  are as follows:

$$\begin{aligned} & \{x_1 \approx 0.4492704820, x_2 \approx 0.7040380158, x_3 \approx 0.7236008036, x_4 \approx 0.4025066142\}, \\ & \{x_1 \approx 0.2625192922, x_2 \approx 1.149055866, x_3 \approx 0.6923507482, x_4 \approx 1.032699036\}, \\ & \{x_1 \approx 0.3393694928, x_2 \approx 0.7028774603, x_3 \approx 1.499308774, x_4 \approx 1.438454385\}, \\ & \{x_1 \approx 1.476432886, x_2 \approx 1.076444066, x_3 \approx 4.066468725, x_4 \approx 4.052215439\}. \end{aligned}$$

For  $20x_4^2 - 72x_4 + 49 = 0$ , we have  $x_4 \approx 0.9111805582$  and  $x_4 \approx 2.688819441$ , whose corresponding solutions of the system  $\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, 20x_4^2 - 72x_4 + 49 = 0\}$  with  $x_1x_2x_3x_4 \neq 0$  are as follows:

$$\begin{aligned} & \{x_1 = x_3 = 1, x_2 = x_4 \approx 0.9111805582\}, \\ & \{x_1 = x_3 = 1, x_2 = x_4 \approx 2.688819441\}. \end{aligned}$$

In conclusion, we find six different  $G$ -invariant Einstein metrics on homogeneous manifold  $E_7/A_1$ .

**Case of  $E_7/A_5$ .** Consider the homogeneous manifold  $E_7/A_5$  with the decomposition

$$\mathfrak{m} = T \oplus A_1 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.3)$$

and  $\text{Ad}(A_5)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1 \oplus A_5)$ -invariant on  $E_7/A_5$  defined by

$$\langle \cdot, \cdot \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{A_1} + x_2 \cdot B|_{\mathfrak{p}_1} + x_3 \cdot B|_{\mathfrak{p}_2} + x_4 \cdot B|_{\mathfrak{p}_3}, \quad (4.4)$$

where  $u_0, x_1, x_2, x_3, x_4 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.1, the components of Ricci tensor with respect to the metric (4.4) are as follows:

$$\begin{cases} r_T = \frac{1}{4} \left( \frac{4u_0}{9x_2^2} + \frac{5u_0}{9x_3^2} \right), \\ r_{A_1} = \frac{1}{12} \left( \frac{1}{3x_1} + \frac{x_1}{x_2^2} + \frac{5x_1}{3x_4^2} \right), \\ r_{\mathfrak{p}_1} = \frac{1}{2x_2} + \frac{5}{36} \left( \frac{x_2}{x_3x_4} - \frac{x_3}{x_4x_2} - \frac{x_4}{x_2x_3} \right) - \frac{1}{48} \left( \frac{4u_0}{9x_2^2} + \frac{x_1}{x_2^2} \right), \\ r_{\mathfrak{p}_2} = \frac{1}{2x_3} + \frac{1}{9} \left( \frac{x_3}{x_4x_2} - \frac{x_4}{x_2x_3} - \frac{x_2}{x_3x_4} \right) - \frac{1}{60} \frac{5u_0}{9x_3^2}, \\ r_{\mathfrak{p}_3} = \frac{1}{2x_4} + \frac{1}{12} \left( \frac{x_4}{x_2x_3} - \frac{x_2}{x_3x_4} - \frac{x_3}{x_4x_2} \right) - \frac{1}{80} \frac{5x_1}{3x_4^2}. \end{cases}$$

Moreover, the homogeneous Einstein equation is given by

$$\{r_T - r_{A_1} = 0, r_{A_1} - r_{\mathfrak{p}_1} = 0, r_{\mathfrak{p}_1} - r_{\mathfrak{p}_2} = 0, r_{\mathfrak{p}_2} - r_{\mathfrak{p}_3} = 0\},$$

which is equivalent to the following system of equations by setting  $u_0 = 1$ :

$$\begin{cases} g_0 = -5x_1^2x_2^2x_3^2 - 3x_1^2x_3^2x_4^2 - x_2^2x_3^2x_4^2 + 5x_1x_2^2x_4^2 + 4x_1x_3^2x_4^2 = 0, \\ g_1 = 60x_1^2x_2^2x_3 + 45x_1^2x_3x_4^2 - 60x_1x_2^3x_4 + 60x_1x_2x_3^2x_4 - 216x_1x_2x_3x_4^2 + 60x_1x_2x_4^3 \\ + 12x_2^2x_3x_4^2 + 4x_1x_3x_4^2 = 0, \\ g_2 = -9x_1x_3^2x_4 + 108x_2^3x_3 - 216x_2^2x_3x_4 - 108x_2x_3^3 + 216x_2x_3^2x_4 - 12x_2x_3x_4^2 + 4x_2^2x_4 \\ - 4x_3^2x_4 = 0, \\ g_3 = 9x_1x_2x_3^2 - 12x_2^2x_3x_4 - 216x_2x_3^2x_4 + 216x_2x_3x_4^2 + 84x_3^3x_4 - 84x_3x_4^3 - 4x_2x_4^2 = 0. \end{cases}$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3, x_4]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, g_3, zx_1x_2x_3x_4 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3 > x_4$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial  $h(x_4)$  of degree 24 given by

$$\begin{aligned}
h(x_4) = & 743926266335759400800157293103178343319505630645046435156250000 x_4^{24} \\
& - 4378144578543039011361358385061081569161674400675383078316875000 x_4^{23} \\
& + 11004903226880209703956063057559301456591525835888793871369802500 x_4^{22} \\
& - 15728343701831201252751181874054831911519725141563733372843754800 x_4^{21} \\
& + 15016244958117620330302995693661960097380760512161595698974470890 x_4^{20} \\
& - 11994308781926707001101666261632390461497787947208181720205033200 x_4^{19} \\
& + 9738633984043261058161384571494874932415716842246643813641906732 x_4^{18} \\
& - 6775088500386513521204014691159118392435727285614732468339912598 x_4^{17} \\
& + 2706243877713720652207013644363017005932182610630242895498110351 x_4^{16} \\
& - 12000153858619368362412920925916507581760380187119703324718112 x_4^{15} \\
& - 470852895612592200583477610033508459972546973038296225090613648 x_4^{14} \\
& + 158657369099271326137588879121559258044753269615019628589825152 x_4^{13} \\
& - 4309426841561170090469256176045397413256930487964567974681056 x_4^{12} \\
& - 6417570511115254453560753629391603456981811382551849470944256 x_4^{11} \\
& + 1526468649104808736601756484180315699890352358674705978019200 x_4^{10} \\
& - 181418867068729727750349886730055784794192965200077070182912 x_4^9 \\
& + 13350908794995611352389703910041292729770053075493309180672 x_4^8 \\
& - 63299614330360037211228864933836702374869844580196724736 x_4^7 \\
& + 18258540029433839516883828010654743829211229647212972032 x_4^6 \\
& - 229558007445965630066366475684636804625527678725234688 x_4^5 \\
& - 2954197894924039316661797095646508811861752372301824 x_4^4 \\
& + 121322059872735831667919068167865971682917554847744 x_4^3 \\
& + 543967044145299748038471088930322206082300116992 x_4^2 \\
& - 69016429739436519614913257303724864151171891200 x_4 \\
& + 849606259908726681491324537499135318226370560.
\end{aligned}$$

By solving  $h(x_4) = 0$  numerically, we get four different positive solutions, which are given approximately by  $x_4 \approx 0.6449885281$ ,  $x_4 \approx 0.7443959787$ ,  $x_4 \approx 1.060416256$ ,  $x_4 \approx 1.501487029$ . Further, we remark that  $x_1, x_2, x_3$  can be expressed by  $x_4$ . Therefore, the real solutions of the system  $\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, h(x_4) = 0\}$  with  $x_1x_2x_3x_4 \neq 0$  are as follows:

$$\begin{aligned} & \{x_1 \approx 1.173286641, x_2 \approx 0.9076818373, x_3 \approx 0.5898460162, x_4 \approx 0.6449885281\}, \\ & \{x_1 \approx 0.05521141092, x_2 \approx 0.9874672655, x_3 \approx 0.5836473658, x_4 \approx 0.7443959787\}, \\ & \{x_1 \approx 0.07379520877, x_2 \approx 0.6077356344, x_3 \approx 1.170652643, x_4 \approx 1.060416256\}, \\ & \{x_1 \approx 1.140144619, x_2 \approx 0.6881276139, x_3 \approx 1.513843081, x_4 \approx 1.501487029\}. \end{aligned}$$

To conclude, we find four different  $G$ -invariant Einstein metrics on homogeneous manifold  $E_7/A_5$ .

**Case of  $E_7/(A_1 \times A_5)$ .** Consider the homogeneous manifold  $E_7/(A_1 \times A_5)$  with the decomposition

$$\mathfrak{m} = T \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.5)$$

and  $\text{Ad}(A_1 \times A_5)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1 \oplus A_5)$ -invariant on  $E_7/(A_1 \times A_5)$  defined by

$$\langle \cdot, \cdot \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{\mathfrak{p}_1} + x_2 \cdot B|_{\mathfrak{p}_2} + x_3 \cdot B|_{\mathfrak{p}_3}, \quad (4.6)$$

where  $u_0, x_1, x_2, x_3 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.1, the components of Ricci tensor with respect to the metric (4.6) are as follows:

$$\begin{cases} r_T = \frac{1}{4} \left( \frac{4u_0}{9x_1^2} + \frac{5u_0}{9x_2^2} \right), \\ r_{\mathfrak{p}_1} = \frac{1}{2x_1} + \frac{5}{36} \left( \frac{x_1}{x_2x_3} - \frac{x_2}{x_3x_1} - \frac{x_3}{x_1x_2} \right) - \frac{1}{48} \frac{4u_0}{9x_1^2}, \\ r_{\mathfrak{p}_2} = \frac{1}{2x_2} + \frac{1}{9} \left( \frac{x_2}{x_3x_1} - \frac{x_3}{x_1x_2} - \frac{x_1}{x_2x_3} \right) - \frac{1}{60} \frac{5u_0}{9x_2^2}, \\ r_{\mathfrak{p}_3} = \frac{1}{2x_3} + \frac{1}{12} \left( \frac{x_3}{x_1x_2} - \frac{x_1}{x_2x_3} - \frac{x_2}{x_3x_1} \right). \end{cases}$$

The homogeneous Einstein equation is equivalent to the following system of equations (we normalize the metric by setting  $u_0 = 1$ ):

$$\begin{cases} g_0 = -15x_1^3x_2 + 15x_1x_2^3 - 54x_1x_2^2x_3 + 15x_1x_2x_3^2 + 15x_1^2x_3 + 13x_2^2x_3 = 0, \\ g_1 = 27x_1^3x_2 - 54x_1^2x_2x_3 - 27x_1x_2^3 + 54x_1x_2^2x_3 - 3x_1x_2x_3^2 + x_1^2x_3 - x_2^2x_3 = 0, \\ g_2 = -3x_1^2x_2 - 54x_1x_2^2 + 54x_1x_2x_3 + 21x_2^3 - 21x_2x_3^2 - x_1x_3 = 0. \end{cases}$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, zx_1x_2x_3 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial  $h(x_3)$  of degree 8 given by

$$\begin{aligned} h(x_3) = & 126500890017116160000x_3^8 - 443159701524243843840x_3^7 \\ & + 626866608280069521864x_3^6 - 439964490275093593980x_3^5 \\ & + 146470811201043075009x_3^4 - 18042589809134073603x_3^3 \\ & + 1095420294616803246x_3^2 - 33677044349391072x_3 \\ & + 438386090662912. \end{aligned}$$

By solving  $h(x_3) = 0$  numerically, there are two positive solutions, which are given approximately by  $x_3 \approx 0.7485101873$  and  $x_3 \approx 1.034311056$ . Moreover, we remark that  $x_1$  and  $x_2$  can be expressed by  $x_3$ . As a result, the real solutions of the system of equations  $\{g_0 = 0, g_1 = 0, g_2 = 0, h(x_3) = 0\}$  are

$$\begin{aligned} & \{x_1 \approx 0.9908972484, x_2 \approx 0.5833878061, x_3 \approx 0.7485101873\}, \\ & \{x_1 \approx 0.6019040936, x_2 \approx 1.154991982, x_3 \approx 1.034311056\}. \end{aligned}$$

In conclusion, we find two different  $G$ -invariant Einstein metrics on the homogeneous manifold  $E_7/(A_1 \times A_5)$ .

**Case of  $E_6/A_1$ .** Consider the homogeneous manifold  $E_6/A_1^1$  with the decomposition

$$\mathfrak{m} = T \oplus A_1^2 \oplus A_3 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.7)$$

and  $\text{Ad}(A_1^1)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1^1 \oplus A_1^2 \oplus A_3)$ -invariant on  $E_6/A_1^1$  defined by

$$\langle \cdot, \cdot \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{A_1^2} + x_2 \cdot B|_{A_3} + x_3 \cdot B|_{\mathfrak{p}_1} + x_4 \cdot B|_{\mathfrak{p}_2} + x_5 \cdot B|_{\mathfrak{p}_3}, \quad (4.8)$$

where  $u_0, x_1, x_2, x_3, x_4, x_5 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.2, the components of Ricci tensor with respect to the metric (4.8) are as follows:

$$\left\{ \begin{array}{l} r_T = \frac{1}{4} \left( \frac{u_0}{2x_3^2} + \frac{u_0}{2x_4^2} \right), \\ r_{A_1^2} = \frac{1}{12} \left( \frac{1}{2x_1} + \frac{x_1}{x_3^2} + \frac{3x_1}{2x_5^2} \right), \\ r_{A_3} = \frac{1}{60} \left( \frac{5}{x_2} + \frac{5x_2}{2x_3^2} + \frac{5x_2}{2x_4^2} + \frac{5x_2}{x_5^2} \right), \\ r_{\mathfrak{p}_1} = \frac{1}{2x_3} + \frac{1}{8} \left( \frac{x_3}{x_4x_5} - \frac{x_4}{x_5x_3} - \frac{x_5}{x_3x_4} \right) - \frac{1}{32} \left( \frac{u_0}{2x_3^2} + \frac{x_1}{x_3^2} + \frac{5x_2}{2x_3^2} \right), \\ r_{\mathfrak{p}_2} = \frac{1}{2x_4} + \frac{1}{8} \left( \frac{x_4}{x_5x_3} - \frac{x_5}{x_3x_4} - \frac{x_3}{x_4x_5} \right) - \frac{1}{32} \left( \frac{u_0}{2x_4^2} + \frac{5x_2}{2x_4^2} \right), \\ r_{\mathfrak{p}_3} = \frac{1}{2x_5} + \frac{1}{12} \left( \frac{x_5}{x_3x_4} - \frac{x_3}{x_4x_5} - \frac{x_4}{x_5x_3} \right) - \frac{1}{48} \left( \frac{3x_1}{2x_5^2} + \frac{5x_2}{x_5^2} \right). \end{array} \right.$$

We study the following homogeneous Einstein equation

$$\{r_T - r_{A_1^2} = 0, r_{A_1^2} - r_{A_3} = 0, r_{A_3} - r_{\mathfrak{p}_1} = 0, r_{\mathfrak{p}_1} - r_{\mathfrak{p}_2} = 0, r_{\mathfrak{p}_2} - r_{\mathfrak{p}_3} = 0\},$$

which is equivalent to the following system of equations by setting  $u_0 = 1$ :

$$\left\{ \begin{array}{l} g_0 = -3x_1^2x_3^2x_4^2 - 2x_1^2x_4^2x_5^2 - x_3^2x_4^2x_5^2 + 3x_1x_3^2x_5^2 + 3x_1x_4^2x_5^2 = 0, \\ g_1 = 3x_1^2x_2x_3^2x_4^2 + 2x_1^2x_2x_4^2x_5^2 - 2x_1x_2^2x_3^2x_4^2 - x_1x_2^2x_3^2x_5^2 - x_1x_2^2x_4^2x_5^2 - 2x_1x_3^2x_4^2x_5^2 \\ + x_2x_3^2x_4^2x_5^2 = 0, \\ g_2 = 6x_1x_2x_4^2x_5^2 + 16x_2^2x_3^2x_4^2 + 8x_2^2x_3^2x_5^2 + 23x_2^2x_4^2x_5^2 - 24x_2x_3^3x_4x_5 + 24x_2x_3x_4^3x_5 \\ - 96x_2x_3x_4^2x_5^2 + 24x_2x_3x_4x_5^3 + 16x_3^2x_4^2x_5^2 + 3x_2x_4^2x_5^2 = 0, \\ g_3 = -2x_1x_4^2x_5 + 5x_2x_3^2x_5 - 5x_2x_4^2x_5 + 16x_3^3x_4 - 32x_3^2x_4x_5 - 16x_3x_4^3 + 32x_3x_4^2x_5 + x_3^2x_5 \\ - x_4^2x_5 = 0, \\ g_4 = 6x_1x_3x_4^2 + 20x_2x_3x_4^2 - 15x_2x_3x_5^2 - 8x_3^2x_4x_5 - 96x_3x_4^2x_5 + 96x_3x_4x_5^2 + 40x_4^3x_5 \\ - 40x_4x_5^3 - 3x_3x_5^2 = 0. \end{array} \right.$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3, x_4, x_5]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, g_3, g_4, zx_1x_2x_3x_4x_5 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3 > x_4 > x_5$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial  $(4x_5^2 - 16x_5 + 11) \cdot h(x_5)$ , where  $h(x_5)$  is a polynomial of degree 74. Since the length of the polynomial  $h(x_5)$  may affect the readers read, we put it in Appendix.

By solving  $h(x_5) = 0$ , there exist eight different positive solutions, which are given approximately by  $x_5 \approx 0.4678383774, x_5 \approx 0.8564644408, x_5 \approx 1.195992694, x_5 \approx 1.331593277, x_5 \approx 2.323059800, x_5 \approx 2.650131702, x_5 \approx 3.078478984, x_5 \approx 3.248554868$ . Moreover, we remark that  $x_1, x_2, x_3, x_4$  can be written into polynomials of  $x_5$ . As a result, we get all solutions of the system  $\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, g_4 = 0, h(x_5) = 0\}$  with  $x_1x_2x_3x_4x_5 \neq 0$ , which are given below:

$$\begin{aligned} &\{x_1 \approx 0.5079142344, x_2 \approx 0.5564563075, x_3 \approx 0.7674520218, x_4 \approx 0.7344343505, x_5 \approx 0.4678383774\}, \\ &\{x_1 \approx 1.391791299, x_2 \approx 0.2653010181, x_3 \approx 1.024348994, x_4 \approx 0.6953615240, x_5 \approx 0.8564644408\}, \\ &\{x_1 \approx 0.1314059305, x_2 \approx 0.2920839602, x_3 \approx 1.311645073, x_4 \approx 0.6903614938, x_5 \approx 1.195992694\}, \\ &\{x_1 \approx 0.1535432666, x_2 \approx 0.3214000505, x_3 \approx 0.7137178306, x_4 \approx 1.420816234, x_5 \approx 1.331593277\}, \\ &\{x_1 \approx 3.641968638, x_2 \approx 1.362098346, x_3 \approx 2.353687875, x_4 \approx 0.9881377660, x_5 \approx 2.323059800\}, \\ &\{x_1 \approx 1.083726408, x_2 \approx 0.7345275471, x_3 \approx 0.9316800862, x_4 \approx 2.652053038, x_5 \approx 2.650131702\}, \\ &\{x_1 \approx 0.3244500133, x_2 \approx 1.435139205, x_3 \approx 3.101828021, x_4 \approx 1.010069757, x_5 \approx 3.078478984\}, \\ &\{x_1 \approx 0.4746512742, x_2 \approx 1.288027846, x_3 \approx 1.023041915, x_4 \approx 3.259446111, x_5 \approx 3.248554868\}. \end{aligned}$$

For  $4x_5^2 - 16x_5 + 11 = 0$ , we have two different positive solutions given approximately by  $x_5 \approx 0.8819660112$  and  $x_5 \approx 3.118033988$ . Thus, the solutions of the system of equations  $\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, g_4 = 0, 4x_5^2 - 16x_5 + 11 = 0\}$  with  $x_1x_2x_3x_4x_5 \neq 0$  are as follows:

$$\begin{aligned} &\{x_1 = x_2 = x_3 = 1, x_4 = x_5 \approx 0.8819660112\}, \\ &\{x_1 = x_2 = x_3 = 1, x_4 = x_5 \approx 3.118033988\}. \end{aligned}$$

In conclusion, we find ten different  $G$ -invariant Einstein metrics on homogeneous manifold  $E_6/A_1$ .

**Remark 4.1.** One can consider the homogeneous manifold  $E_6/A_1^2$ , but the Einstein metrics on which is the same as the above up to isometry.

**Case of  $E_6/A_3$ .** Consider the homogeneous manifold  $E_6/A_3$  with the decomposition

$$\mathfrak{m} = T \oplus A_1^1 \oplus A_1^2 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.9)$$

and  $\text{Ad}(A_3)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1^1 \oplus A_1^2 \oplus A_3)$ -invariant on  $E_6/A_3$  defined by

$$\langle , \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{A_1^1} + x_2 \cdot B|_{A_1^2} + x_3 \cdot B|_{\mathfrak{p}_1} + x_4 \cdot B|_{\mathfrak{p}_2} + x_5 \cdot B|_{\mathfrak{p}_3}, \quad (4.10)$$

where  $u_0, x_1, x_2, x_3, x_4, x_5 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.2, the components of Ricci tensor with respect

to the metric (4.10) are as follows:

$$\begin{cases} r_T = \frac{1}{4} \left( \frac{u_0}{2x_3^2} + \frac{u_0}{2x_4^2} \right), \\ r_{A_1^1} = \frac{1}{12} \left( \frac{1}{2x_1} + \frac{x_1}{x_4^2} + \frac{3x_1}{2x_5^2} \right), \\ r_{A_1^2} = \frac{1}{12} \left( \frac{1}{2x_2} + \frac{x_2}{x_3^2} + \frac{3x_2}{2x_5^2} \right), \\ r_{p_1} = \frac{1}{2x_3} + \frac{1}{8} \left( \frac{x_3}{x_4x_5} - \frac{x_4}{x_5x_3} - \frac{x_5}{x_3x_4} \right) - \frac{1}{32} \left( \frac{u_0}{2x_3^2} + \frac{x_2}{x_3^2} \right), \\ r_{p_2} = \frac{1}{2x_4} + \frac{1}{8} \left( \frac{x_4}{x_5x_3} - \frac{x_5}{x_3x_4} - \frac{x_3}{x_4x_5} \right) - \frac{1}{32} \left( \frac{u_0}{2x_4^2} + \frac{x_1}{x_4^2} \right), \\ r_{p_3} = \frac{1}{2x_5} + \frac{1}{12} \left( \frac{x_5}{x_3x_4} - \frac{x_3}{x_4x_5} - \frac{x_4}{x_5x_3} \right) - \frac{1}{48} \left( \frac{3x_1}{2x_5^2} + \frac{3x_2}{2x_5^2} \right). \end{cases}$$

Further, the homogeneous Einstein equation is

$$\{r_T - r_{A_1^1} = 0, r_{A_1^1} - r_{A_1^2} = 0, r_{A_1^2} - r_{p_1} = 0, r_{p_1} - r_{p_2} = 0, r_{p_2} - r_{p_3} = 0\},$$

which is equivalent to the following system of equations by setting  $u_0 = 1$ :

$$\begin{cases} g_0 = -3x_1^2x_3^2x_4^2 - 2x_1^2x_3^2x_5^2 - x_3^2x_4^2x_5^2 + 3x_1x_3^2x_5^2 + 3x_1x_4^2x_5^2 = 0, \\ g_1 = 3x_1^2x_2x_3^2x_4^2 + 2x_1^2x_2x_3^2x_5^2 - 3x_1x_2^2x_3^2x_4^2 - 2x_1x_2^2x_4^2x_5^2 - x_1x_3^2x_4^2x_5^2 \\ + x_2x_3^2x_4^2x_5^2 = 0, \\ g_2 = 24x_2^2x_3^2x_4 + 22x_2^2x_4x_5^2 - 24x_2x_3^3x_5 + 24x_2x_3x_4^2x_5 - 96x_2x_3x_4x_5^2 + 24x_2x_3x_5^3 \\ + 8x_3^2x_4x_5^2 + 3x_2x_4x_5^2 = 0, \\ g_3 = 2x_1x_3^2x_5 - 2x_2x_4^2x_5 + 16x_3^3x_4 - 32x_3^2x_4x_5 - 16x_3x_4^3 + 32x_3x_4^2x_5 + x_3^2x_5 - x_4^2x_5 = 0, \\ g_4 = 6x_1x_3x_4^2 - 6x_1x_3x_5^2 + 6x_2x_3x_4^2 - 8x_3^2x_4x_5 - 96x_3x_4^2x_5 + 96x_3x_4x_5^2 + 40x_4^3x_5 \\ - 40x_4x_5^3 - 3x_3x_5^2 = 0. \end{cases}$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3, x_4, x_5]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, g_3, g_4, zx_1x_2x_3x_4x_5 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3 > x_4 > x_5$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial  $h(x_5)$  of degree 31 given by

$$\begin{aligned}
h(x_5) = & 1261108165616422801344202981595019832319020116358877329036935168000 x_5^{31} \\
& - 2151788432405784864427357292265755888622754231825843818424015257600 x_5^{30} \\
& - 2774119470594771842109777615792789059591610394561214417063934689280 x_5^{29} \\
& + 527034483192300175170830381675908821659174954446664977785609519104 x_5^{28} \\
& - 267759863524378249157490847814955848518971246169904017730950922240 x_5^{27} \\
& + 3382539429618210066981097554617654004425041655004361517303374807040 x_5^{26} \\
& - 9157476976548260444114456636484244081952190409782190142886449250304 x_5^{25} \\
& + 3265956979395568431463257163604494142186261853585696957494114123776 x_5^{24} \\
& + 1113697025312385709566702871382152926625644868347096267275321556992 x_5^{23} \\
& + 2112333421951376168597459647362124180457453099511088975109229199360 x_5^{22} \\
& - 1631290981926813934052954101478705657981981503958339738878928916480 x_5^{21} \\
& - 1395050115461884401753191844047758010380519323029179088958100982784 x_5^{20} \\
& + 863710166438884070851640864653697900991831475370745010492288656384 x_5^{19} \\
& + 241554806478645820779502711256075967575496587804599790265232872960 x_5^{18} \\
& - 132746306295681893821384103002486474701425792860693217422446582528 x_5^{17} \\
& - 3267602765928478150920615601297457606838015845305743614683411008 x_5^{16} \\
& + 8892989132102003650248145850686458075368692869017300184897905088 x_5^{15} \\
& - 1328201011191036292414037060140630816274620551629810587376210400 x_5^{14} \\
& - 98351649199565867328202679801373812742043192868226187056696624 x_5^{13} \\
& + 57195416139228700916611387885001217660977148332951265434464548 x_5^{12} \\
& - 8584293427926321735849799193107683389353503523382291011168608 x_5^{11} \\
& + 658987123591863795218834332702519196588690183339615796158964 x_5^{10} \\
& - 21514437280227836839597312688258324479163856784281755956140 x_5^9 \\
& - 562803146421214665893608417569171948144549496211000931459 x_5^8 \\
& + 69880133370536372146966764193862531901488645811504750528 x_5^7 \\
& - 642571168743645462407828141773462190939783200613014680 x_5^6 \\
& - 141173037619723104390974578987052013640153011100310880 x_5^5 \\
& + 4922149886224089807286830168142173254755229906374200 x_5^4 \\
& + 92795089337275364876653655416844512848233792832000 x_5^3 \\
& - 6084529180331760636310164863691628563332526660000 x_5^2 \\
& - 51292165983013771433079466375050322032293400000 x_5 \\
& + 4587282600584194356057367005912663107945250000.
\end{aligned}$$

By solving  $h(x_5) = 0$  numerically, we have four positive and one negative solutions, which can be given approximately by  $x_5 \approx 0.7601951682$ ,  $x_5 \approx 0.9753459906$ ,  $x_5 \approx 1.141009718$ ,  $x_5 \approx 1.523722394$ ,  $x_5 \approx -1.755049707$ . With aid of computer we get all the solutions of system of equations

$\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, g_4 = 0, h(x_5) = 0\}$  with  $x_1 x_2 x_3 x_4 x_5 \neq 0$  are as follows:

$$\begin{aligned} & \{x_1 \approx 0.1044773890, x_2 \approx 1.350603440, x_3 \approx 0.9684903881, x_4 \approx 0.6352925758, x_5 \approx 0.7601951682\}, \\ & \{x_1 \approx 1.350603440, x_2 \approx 0.1044773890, x_3 \approx 0.6352925758, x_4 \approx 0.9684903881, x_5 \approx 0.7601951682\}, \\ & \{x_1 \approx 0.1100292802, x_2 \approx 0.1051859906, x_3 \approx 1.141500322, x_4 \approx 0.6242772748, x_5 \approx 0.9753459906\}, \\ & \{x_1 \approx 0.1051859906, x_2 \approx 0.1100292802, x_3 \approx 0.6242772748, x_4 \approx 1.141500322, x_5 \approx 0.9753459906\}, \\ & \{x_1 \approx 1.126287414, x_2 \approx 1.862577839, x_3 \approx 1.181080704, x_4 \approx 0.7489214139, x_5 \approx 1.141009718\}, \\ & \{x_1 \approx 1.862577839, x_2 \approx 1.126287414, x_3 \approx 0.7489214139, x_4 \approx 1.181080704, x_5 \approx 1.141009718\}, \\ & \{x_1 \approx 1.191985212, x_2 \approx 0.1562720393, x_3 \approx 1.546458511, x_4 \approx 0.7400787612, x_5 \approx 1.523722394\}, \\ & \{x_1 \approx 0.1562720393, x_2 \approx 1.191985212, x_3 \approx 0.7400787612, x_4 \approx 1.546458511, x_5 \approx 1.523722394\}. \end{aligned}$$

We remark that among these solutions, the ones with  $x_5$  equal induce the same metrics up to isometry. As a result, there are four different  $G$ -invariant Einstein metrics on homogeneous manifold  $E_6/A_3$ .

**Case of  $E_6/(A_1 \times A_3)$ .** Consider the homogeneous manifold  $E_6/(A_1^1 \times A_3)$  with the decomposition

$$\mathfrak{m} = T \oplus A_1^2 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.11)$$

and  $\text{Ad}(A_1^1 \times A_3)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1^1 \oplus A_1^2 \oplus A_3)$ -invariant on  $E_6/(A_1^1 \times A_3)$  defined by

$$\langle \cdot, \cdot \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{A_1^2} + x_2 \cdot B|_{\mathfrak{p}_1} + x_3 \cdot B|_{\mathfrak{p}_2} + x_4 \cdot B|_{\mathfrak{p}_3}, \quad (4.12)$$

where  $u_0, x_1, x_2, x_3, x_4 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.2, the components of Ricci tensor with respect to the metric (4.12) are as follows:

$$\left\{ \begin{array}{l} r_T = \frac{1}{4} \left( \frac{u_0}{2x_2^2} + \frac{u_0}{2x_3^2} \right), \\ r_{A_1^2} = \frac{1}{12} \left( \frac{1}{2x_1} + \frac{x_1}{x_2^2} + \frac{3x_1}{2x_4^2} \right), \\ r_{\mathfrak{p}_1} = \frac{1}{2x_2} + \frac{1}{8} \left( \frac{x_2}{x_3 x_4} - \frac{x_3}{x_4 x_2} - \frac{x_4}{x_2 x_3} \right) - \frac{1}{32} \left( \frac{u_0}{2x_2^2} + \frac{x_1}{x_2^2} \right), \\ r_{\mathfrak{p}_2} = \frac{1}{2x_3} + \frac{1}{8} \left( \frac{x_3}{x_4 x_2} - \frac{x_4}{x_2 x_3} - \frac{x_2}{x_3 x_4} \right) - \frac{1}{32} \frac{u_0}{2x_3^2}, \\ r_{\mathfrak{p}_3} = \frac{1}{2x_4} + \frac{1}{12} \left( \frac{x_4}{x_2 x_3} - \frac{x_2}{x_3 x_4} - \frac{x_3}{x_4 x_2} \right) - \frac{1}{48} \frac{3x_1}{2x_4^2}. \end{array} \right.$$

Moreover, the homogeneous Einstein equation is

$$\{r_T - r_{A_1^2} = 0, r_{A_1^2} - r_{\mathfrak{p}_1} = 0, r_{\mathfrak{p}_1} - r_{\mathfrak{p}_2} = 0, r_{\mathfrak{p}_2} - r_{\mathfrak{p}_3} = 0\},$$

which is equivalent to the following system of equations by setting  $u_0 = 1$ :

$$\left\{ \begin{array}{l} g_0 = -3x_1^2 x_2^2 x_3^2 - 2x_1^2 x_3^2 x_4^2 - x_2^2 x_3^2 x_4^2 + 3x_1 x_2^2 x_4^2 + 3x_1 x_3^2 x_4^2 = 0, \\ g_1 = 24x_1^2 x_2^2 x_3 + 22x_1^2 x_3 x_4^2 - 24x_1 x_2^3 x_4 + 24x_1 x_2 x_3^2 x_4 - 96x_1 x_2 x_3 x_4^2 + 24x_1 x_2 x_4^3 \\ \quad + 8x_2^2 x_3 x_4^2 + 3x_1 x_3 x_4^2 = 0, \\ g_2 = -2x_1 x_3^2 x_4 + 16x_2^3 x_3 - 32x_2^2 x_3 x_4 - 16x_2 x_3^3 + 32x_2 x_3^2 x_4 + x_2^2 x_4 - x_3^2 x_4 = 0, \\ g_3 = 6x_1 x_2 x_3^2 - 8x_2^2 x_3 x_4 - 96x_2 x_3^2 x_4 + 96x_2 x_3 x_4^2 + 40x_3^3 x_4 - 40x_3 x_4^3 - 3x_2 x_4^2 = 0. \end{array} \right.$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3, x_4]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, g_3, zx_1x_2x_3x_4 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3 > x_4$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial  $h(x_4)$  of degree 24 given by

$$\begin{aligned}
 h(x_4) = & 724760062707181194419174600514595111417722211978444800000 x_4^{24} \\
 & - 4549506617169858886572945697876772369451209831712030720000 x_4^{23} \\
 & + 12407995627074284240831723931001141829288068512693092352000 x_4^{22} \\
 & - 19700816036968099156069552630998662850171899956562598297600 x_4^{21} \\
 & + 21336801748378149912816519240285552502805547195911991459840 x_4^{20} \\
 & - 18881224039130808168709393497734018965438973813620910063616 x_4^{19} \\
 & + 15807728714235458621829548954153222096357328191223379263488 x_4^{18} \\
 & - 11339342221372302981951544924786699818394336801815208656896 x_4^{17} \\
 & + 5112518772844735173801750222596331052099974758094486477824 x_4^{16} \\
 & - 403396001117807932930095402077870112549798991179251580928 x_4^{15} \\
 & - 903609223572125644190343758413680417593338129263631754880 x_4^{14} \\
 & + 461550807018117050045055022022782250907255388226385775616 x_4^{13} \\
 & - 65891887023479142966230863046964022312825636044010288900 x_4^{12} \\
 & - 13675428883180345095301338226487838153532965536692318272 x_4^{11} \\
 & + 7404379070352221664927980007974819413355565861485929092 x_4^{10} \\
 & - 148399736786611237822770061868493553550293718403156128 x_4^9 \\
 & + 175268024365864286996306905494287487314022126666412671 x_4^8 \\
 & - 12983233756373427437475671647726918973459199705907872 x_4^7 \\
 & + 557098428728836786107911315677130125808601844881240 x_4^6 \\
 & - 7614872308049576020573960694434807958692132824768 x_4^5 \\
 & - 444217606400638301563490761267824391591116546648 x_4^4 \\
 & + 16454149288639736436126283769710330401564494976 x_4^3 \\
 & + 517849751430653664520517887668936015103631904 x_4^2 \\
 & - 41649452199426082400520496110816668096290560 x_4 \\
 & + 743025277331758306683208760331078413656560.
 \end{aligned}$$

By solving  $h(x_4) = 0$  numerically, we have four positive solutions, which can be given approximately by  $x_4 \approx 0.7380550950$ ,  $x_4 \approx 0.9342575736$ ,  $x_4 \approx 0.9881689149$  and  $x_4 \approx 1.544605922$ . We remark that  $x_1, x_2, x_3$  can be written into polynomials of  $x_4$ . As a result, we obtain the solutions of the system of equations  $\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, h(x_4) = 0\}$  with  $x_1x_2x_3x_4 \neq 0$  as follows:

$$\begin{aligned} & \{x_1 \approx 1.330255613, x_2 \approx 0.9639605550, x_3 \approx 0.6236782194, x_4 \approx 0.7380550950\}, \\ & \{x_1 \approx 0.1012164462, x_2 \approx 1.121509362, x_3 \approx 0.6122385007, x_4 \approx 0.9342575736\}, \\ & \{x_1 \approx 0.1102891330, x_2 \approx 0.6235977518, x_3 \approx 1.152510591, x_4 \approx 0.9881689149\}, \\ & \{x_1 \approx 1.194390595, x_2 \approx 0.7396962365, x_3 \approx 1.566827089, x_4 \approx 1.544605922\}. \end{aligned}$$

In conclusion, we find four different  $G$ -invariant Einstein metrics on homogeneous manifold  $E_6/(A_1 \times A_3)$ .

**Remark 4.2.** One can consider Einstein metrics on  $E_6/(A_1^2 \times A_3)$ , which is the same as the above results up to isometry.

**Case of  $E_6/(A_1 \times A_1)$ .** Consider the homogeneous manifold  $E_6/(A_1^1 \times A_1^2)$  with the decomposition

$$\mathfrak{m} = T \oplus A_3 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.13)$$

and  $\text{Ad}(A_1^1 \times A_1^2)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1^1 \oplus A_1^2 \oplus A_3)$ -invariant on  $E_6/(A_1^1 \times A_1^2)$  defined by

$$\langle \cdot, \cdot \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{A_3} + x_2 \cdot B|_{\mathfrak{p}_1} + x_3 \cdot B|_{\mathfrak{p}_2} + x_4 \cdot B|_{\mathfrak{p}_3}, \quad (4.14)$$

where  $u_0, x_1, x_2, x_3, x_4 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.2, the components of Ricci tensor with respect to the metric (4.14) are as follows:

$$\left\{ \begin{array}{l} r_T = \frac{1}{4} \left( \frac{u_0}{2x_2^2} + \frac{u_0}{2x_3^2} \right), \\ r_{A_3} = \frac{1}{60} \left( \frac{5}{x_1} + \frac{5x_1}{2x_2^2} + \frac{5x_1}{2x_3^2} + \frac{5x_1}{x_4^2} \right), \\ r_{\mathfrak{p}_1} = \frac{1}{2x_2} + \frac{1}{8} \left( \frac{x_2}{x_3 x_4} - \frac{x_3}{x_4 x_2} - \frac{x_4}{x_2 x_3} \right) - \frac{1}{32} \left( \frac{u_0}{2x_2^2} + \frac{5x_1}{2x_2^2} \right), \\ r_{\mathfrak{p}_2} = \frac{1}{2x_3} + \frac{1}{8} \left( \frac{x_3}{x_4 x_2} - \frac{x_4}{x_2 x_3} - \frac{x_2}{x_3 x_4} \right) - \frac{1}{32} \left( \frac{u_0}{2x_3^2} + \frac{5x_1}{2x_3^2} \right), \\ r_{\mathfrak{p}_3} = \frac{1}{2x_4} + \frac{1}{12} \left( \frac{x_4}{x_2 x_3} - \frac{x_2}{x_3 x_4} - \frac{x_3}{x_4 x_2} \right) - \frac{1}{48} \frac{5x_1}{x_4^2}. \end{array} \right.$$

Moreover, the homogeneous Einstein equation is

$$\{r_T - r_{A_3} = 0, r_{A_3} - r_{\mathfrak{p}_1} = 0, r_{\mathfrak{p}_1} - r_{\mathfrak{p}_2} = 0, r_{\mathfrak{p}_2} - r_{\mathfrak{p}_3} = 0\},$$

which is equivalent to the following system of equations by setting  $u_0 = 1$ :

$$\left\{ \begin{array}{l} g_0 = -2x_1^2 x_2^2 x_3^2 - x_1^2 x_2^2 x_4^2 - x_1^2 x_3^2 x_4^2 - 2x_2^2 x_3^2 x_4^2 + 3x_1 x_2^2 x_4^2 + 3x_1 x_3^2 x_4^2 = 0, \\ g_1 = 16x_1^2 x_2^2 x_3^2 + 8x_1^2 x_2^2 x_4^2 + 23x_1^2 x_3^2 x_4^2 - 24x_1 x_2^3 x_3 x_4 + 24x_1 x_2 x_3^3 x_4 - 96x_1 x_2 x_3^2 x_4^2 \\ \quad + 24x_1 x_2 x_3 x_4^3 + 16x_2^2 x_3^2 x_4^2 + 3x_1 x_3^2 x_4^2 = 0, \\ g_2 = 5x_1 x_2^2 x_4 - 5x_1 x_3^2 x_4 + 16x_2^3 x_3 - 32x_2^2 x_3 x_4 - 16x_2 x_3^3 + 32x_2 x_3^2 x_4 + x_2^2 x_4 - x_3^2 x_4 = 0, \\ g_3 = 20x_1 x_2 x_3^2 - 15x_1 x_2 x_4^2 - 8x_2^2 x_3 x_4 - 96x_2 x_3^2 x_4 + 96x_2 x_3 x_4^2 + 40x_3^3 x_4 - 40x_3 x_4^3 \\ \quad - 3x_2 x_4^2 = 0. \end{array} \right.$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3, x_4]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, g_3, zx_1x_2x_3x_4 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3 > x_4$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial  $h(x_4)$  of degree 17 given by

$$\begin{aligned} h(x_4) = & 110023745349636492444277742390625 x_4^{17} - 633649771052061357670508068621875 x_4^{16} \\ & + 1693219483961674637778856732365000 x_4^{15} - 4316833335694592632675706273589750 x_4^{14} \\ & + 8843142281870361254769815434556400 x_4^{13} - 13636914380079309535418557047150830 x_4^{12} \\ & + 17268089388308771236690983095499516 x_4^{11} - 16542332326263296024484260888236418 x_4^{10} \\ & + 9539252746514719228613145099164214 x_4^9 - 1216363759832722486352208197180776 x_4^8 \\ & - 2038138365512054114906312824023776 x_4^7 + 1194675635182956573455161439625750 x_4^6 \\ & - 183551938256104631960037855371472 x_4^5 - 20803930135660271330029319947786 x_4^4 \\ & + 6531836600529448404797287589404 x_4^3 + 7919351072786394774692666914 x_4^2 \\ & - 64535844819989369786392515511 x_4 + 3728157545978269069973049171. \end{aligned}$$

By solving  $h(x_4) = 0$  numerically, we have two positive and one negative solutions, which can be given approximately by  $x_4 \approx 1.216390732$ ,  $x_4 \approx 3.124923887$  and  $x_4 \approx -0.4548317660$ . With the aid of a computer, we obtain the solutions of the system of equations  $\{g_0 = 0, g_1 = 0, g_2 = 0, g_3 = 0, h(x_4) = 0\}$  with  $x_1x_2x_3x_4 \neq 0$  as follows:

$$\begin{aligned} & \{x_1 \approx 0.2932652315, x_2 \approx 1.330112486, x_3 \approx 0.6899374515, x_4 \approx 1.216390732\}, \\ & \{x_1 \approx 0.2932652315, x_2 \approx 0.6899374515, x_3 \approx 1.330112486, x_4 \approx 1.216390732\}, \\ & \{x_1 \approx 1.438513860, x_2 \approx 3.147955135, x_3 \approx 1.011105585, x_4 \approx 3.124923887\}, \\ & \{x_1 \approx 1.438513860, x_2 \approx 1.011105585, x_3 \approx 3.147955135, x_4 \approx 3.124923887\}. \end{aligned}$$

It is easy to see that the first two solutions induce the same metric up to isometry, and so do the later two solutions. In conclusion, we find two different  $G$ -invariant Einstein metrics on homogeneous manifold  $E_6/(A_1 \times A_1)$ .

**Case of  $E_6/(A_1 \times A_1 \times A_3)$ .** Consider the homogeneous manifold  $E_6/(A_1 \times A_1 \times A_3)$  with the decomposition

$$\mathfrak{m} = T \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3, \quad (4.15)$$

and  $\text{Ad}(A_1 \times A_1 \times A_3)$ -invariant metrics which is also  $\text{Ad}(T \oplus A_1 \oplus A_1 \oplus A_3)$ -invariant on  $E_6/(A_1 \times A_1 \times A_3)$  defined by

$$\langle , \rangle = u_0 \cdot B|_T + x_1 \cdot B|_{\mathfrak{p}_1} + x_2 \cdot B|_{\mathfrak{p}_2} + x_3 \cdot B|_{\mathfrak{p}_3}, \quad (4.16)$$

where  $u_0, x_1, x_2, x_3 \in \mathbb{R}^+$ . By Lemmas 2.1 and 3.2, the components of Ricci tensor with respect to the metric (4.16) are as follows:

$$\begin{cases} r_T = \frac{1}{4} \left( \frac{u_0}{2x_1^2} + \frac{u_0}{2x_2^2} \right), \\ r_{p_1} = \frac{1}{2x_1} + \frac{1}{8} \left( \frac{x_1}{x_2 x_3} - \frac{x_2}{x_3 x_1} - \frac{x_3}{x_1 x_2} \right) - \frac{1}{32} \frac{u_0}{2x_1^2}, \\ r_{p_2} = \frac{1}{2x_2} + \frac{1}{8} \left( \frac{x_2}{x_3 x_1} - \frac{x_3}{x_1 x_2} - \frac{x_1}{x_2 x_3} \right) - \frac{1}{32} \frac{u_0}{2x_2^2}, \\ r_{p_3} = \frac{1}{2x_3} + \frac{1}{12} \left( \frac{x_3}{x_1 x_2} - \frac{x_1}{x_2 x_3} - \frac{x_2}{x_3 x_1} \right). \end{cases}$$

The homogeneous Einstein equation is

$$\{r_T - r_{p_1} = 0, r_{p_1} - r_{p_2} = 0, r_{p_2} - r_{p_3} = 0\},$$

which is equivalent to the following system of equations by setting  $u_0 = 1$ :

$$\begin{cases} g_0 = -8x_1^3x_2 + 8x_1x_2^3 - 32x_1x_2^2x_3 + 8x_1x_2x_3^2 + 8x_1^2x_3 + 9x_2^2x_3 = 0, \\ g_1 = 16x_1^3x_2 - 32x_1^2x_2x_3 - 16x_1x_2^3 + 32x_1x_2^2x_3 + x_1^2x_3 - x_2^2x_3 = 0, \\ g_2 = -8x_1^2x_2 - 96x_1x_2^2 + 96x_1x_2x_3 + 40x_2^3 - 40x_2x_3^2 - 3x_1x_3 = 0. \end{cases}$$

We consider the polynomial ring  $R = \mathbb{Q}[z, x_1, x_2, x_3]$  and an ideal  $I$ , generated by polynomials  $\{g_0, g_1, g_2, zx_1x_2x_3 - 1\}$ . We take a lexicographic ordering  $>$  with  $z > x_1 > x_2 > x_3$  for a monomial ordering on  $R$ . Then, with the aid of computer software MAPLE 15, we use the Gröbner basis for the ideal  $I$  to contain a polynomial  $h(x_3)$  of degree 5 given by

$$h(x_3) = 2359296000x_3^5 - 6217728000x_3^4 + 6426104320x_3^3 - 2832337088x_3^2 + 316436504x_3 - 12778713.$$

By solving  $h(x_3) = 0$  numerically, we have only one solution, which can be given approximately by  $x_3 \approx 0.9460130230$ . With the aid of a computer, we obtain the solutions of the system of equations  $\{g_0 = 0, g_1 = 0, g_2 = 0, h(x_3) = 0\}$  with  $x_1x_2x_3 \neq 0$  as follows:

$$\begin{cases} x_1 \approx 1.131543385, x_2 \approx 0.6115481021, x_3 \approx 0.9460130230, \\ x_1 \approx 0.6115481021, x_2 \approx 1.131543385, x_3 \approx 0.9460130230. \end{cases}$$

It is easy to see that these two solutions induce the same metric up to isometry. In conclusion, we find one  $G$ -invariant Einstein metric on homogeneous manifold  $E_6/(A_1 \times A_1 \times A_3)$ .

To summarize the above conclusions, we write the results in the table below:

**Table 1.** Number of G-invariant Einstein metrics on homogeneous manifolds arising from  $E_7$  and  $E_6$ .

$G$	$K_1$	Decomposition	No.
$E_7$	$A_1$	$\mathfrak{m} = \mathbf{T} \oplus A_5 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3$	6
	$A_5$	$\mathfrak{m} = \mathbf{T} \oplus A_1 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3$	4
	$A_1 \times A_5$	$\mathfrak{m} = \mathbf{T} \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3$	2
$E_6$	$A_1$	$\mathfrak{m} = \mathbf{T} \oplus A_1 \oplus A_3 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3$	10
	$A_3$	$\mathfrak{m} = \mathbf{T} \oplus A_1 \oplus A_1 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3$	4
	$A_1 \times A_3$	$\mathfrak{m} = \mathbf{T} \oplus A_1 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3$	4
	$A_1 \times A_1$	$\mathfrak{m} = \mathbf{T} \oplus A_3 \oplus \mathfrak{p}_1 \oplus \mathfrak{p}_2 \oplus \mathfrak{p}_3$	2
$A_1 \times A_1 \times A_3$			1

## 5. Einstein-Randers metrics on some homogeneous manifolds

Consider  $\text{Ad}(K)$ -invariant Einstein metrics on homogenous spaces  $G/K_1$  given in the above section. By the equivalence of the adjoint representation and the isotropy representation of  $K_1$  on  $\mathfrak{m}$ , the vector field

$$\tilde{W}|_{gK} = d(\tau(g))|_o(W), \quad \forall g \in G, W \in T$$

is well-defined, and it is  $G$ -invariant (see [2]). Furthermore, one can easily verify the equation

$$\langle [W, X]_{\mathfrak{m}}, Y \rangle + \langle X, [W, Y]_{\mathfrak{m}} \rangle = 0$$

holds for any  $W \in T$  and  $X, Y \in \mathfrak{m}$ , using the facts that  $\mathfrak{k}_0 \subset \mathfrak{k}$  and the metric is  $\text{Ad}(K)$ -invariant. Then, by Lemma 1.2 the homogeneous metric

$$F(x, y) = \frac{\sqrt{[\langle W, y \rangle]^2 + \langle y, y \rangle \lambda}}{\lambda} - \frac{\langle W, y \rangle}{\lambda} \quad (5.1)$$

is a  $G$ -invariant Einstein-Randers metric on  $G/K_1$  when  $\langle W, W \rangle < 1$ , and  $F$  is Riemannian if and only if  $W = 0$ . We proved Theorem 1.3.

## Use of AI tools declaration

The authors declare they have not used Artificial Intelligence (AI) tools in the creation of this article.

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## Conflict of interest

The author declares no conflict of interest.

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## Appendix

$$\begin{aligned}
 h(x_5) = & 286735791008485748810470695743810840507818986495807245432092113216647296626784564977861699042109209392129510863897 \\
 & 755418767603186789102115589033086156800000000000 x_5^{74} + \\
 & 279670236049102351139904666053663705638247826260265755362699196643628199738907627007066161079220950613040640661232590966817 \\
 & 902565200918959678967922032640000000000 x_5^{73} + \\
 & 12918753690124428672118010758800670733815974802135577966174265832412429912149815923677622309916577899904523621218634221177559 \\
 & 60294218661159638345628988211200000000 x_5^{72} - \\
 & 49838015377022282279251130264743821517205032160928761049108400013910754342024841575308473266703195440910209939015481829183503 \\
 & 4348773445712068688851215843328000000 x_5^{71} + \\
 & 17216521539856304098571693296595022260329753415833178877106770896256620873221854427717781655120802190456608654524330568510755 \\
 & 680052854967978691662703182741504000000 x_5^{70} - \\
 & 471349069036435310436426801217578014813456647080821064209339538143937211369877977981311276706031387865683496832214887458777511 \\
 & 31762149626365994897585193825075200000 x_5^{69} + \\
 & 116934384679320454309225137495212457425159336561312377838378290690228344729489573155785414789854881376586557619295688061793323 \\
 & 497128548662965095067567582525521920000 x_5^{68} - \\
 & 264392771507680828916864397643085926644639623160755866399798071947508938847235949523041390219916943660436339948477102597136164 \\
 & 02285842606666390096740104458888064000 x_5^{67} + \\
 & 49677072155366370750737655477619921898303204840487137543837707729792145944422386661996698177374552393068153603786173411816173 \\
 & 304359568897591106702610425943005593600 x_5^{66} - \\
 & 876881536358166894240901459194279748931911750922647333848413653721634575821067743560251000955857738838089131633694758391503095 \\
 & 357870203791549339956838253084603842560 x_5^{65} + \\
 & 135399639159938273567114861989120277265582360688906709350902897860154217628275160858347333908834852685714134105690508565510775 \\
 & 2360025105043068189219915919928372232192 x_5^{64} - \\
 & 167910854747543718566056821299464516464377246172589112315858096636340529300100687437065073238315519564563445679042775874831989 \\
 & 0859958076575934933122883870804729135104 x_5^{63} + \\
 & 202437718362660677567851005383141871912224534242567647820422844160212270819252885954634506182997088687923451222169720107494855 \\
 & 9569821991449141717206562113613539049472 x_5^{62} - \\
 & 164326692586213509111883099081963780146254471533547494408285540555936217576125189245558518673618523487285622368175284024806170 \\
 & 3226588747598615321276681754310451658752 x_5^{61} + \\
 & 7355439130809836845486384767020365332424998276903652724643631547557827905497452088087366692297202113530243335823980747406533 \\
 & 440909093909000356081813697305665077248 x_5^{60} - \\
 & 297791290377203296261214349908028990551101553039219372633701366731253478580889525082759150165161055422721934862929291790810936 \\
 & 331477893446142941155917287639786455040 x_5^{59} - \\
 & 950025399861139248442891965419593987991259727120577273513681749959840433597209971153601429658295145115074236743486315152756903 \\
 & 019807916795950137388253917216617005056 x_5^{58} - \\
 & 208736051313171339580776808799742219258300209889846233049189615508113925804802328655335137117314201609467888045473168246895879 \\
 & 745840848534376859628982488734569070592 x_5^{57} +
 \end{aligned}$$

$$\begin{aligned}
& 357820038513677601517671405104585167085413667252626307765376472526401324862847003444795713388616879476776345376853922923052932 \\
& 1482266896205869758729238704226207531008 x_5^{56} - \\
& 663158529960431853656971893878302396064221608535144330731009251904451372008535280554899539322377095054343549855001605098015181 \\
& 6502058416666705061851236622042973339648 x_5^{55} + \\
& 129242530385167128481743588157231297614638175401572620766575234711521316359776854211452524730076304278514688819944292487745435 \\
& 93665149693022263863172649649982524096512 x_5^{54} - \\
& 172429464393321153901442571547631261975664984790646982622831580576997137931578313268799977159882950866560188283114344290363347 \\
& 43236130880952569726910644887583385976832 x_5^{53} + \\
& 165327076363460249103451029975828727406508825987543171322369652086644826683112450572072351036337887941442239573091514025440074 \\
& 12263682604541971778050717414585237766144 x_5^{52} - \\
& 157641322489197222974791760894139639287167405524049995112204311728514369131747209549168434857906361754329989939844856317929648 \\
& 21883295906387428243023437708866212593664 x_5^{51} + \\
& 111515718774216562668698476201485733792099203307133355990409551170489102120550685358628480599783405172754004888814621548232269 \\
& 99430198671154426690003806953657725616128 x_5^{50} - \\
& 2512841271050755700149051067367972282016567293040229541457812367105034995137127259645201522723579066884653570142422691495345 \\
& 9316250316689749388328150148744883470336 x_5^{49} - \\
& 500677076164051058426443385400134248987615816137846857500109109415745747151204297709666521875803117629403380212254167517987363 \\
& 160583621714473963960185791024061284352 x_5^{48} + \\
& 163101638875856546751259251885942683965746990016893375296240776310201730956981584069173297908931810073096408358782797532258040 \\
& 787102890744575765617679983015712260096 x_5^{47} - \\
& 421370843551487087941615573592545372772562641380654830378663535697334040971945605801195876506191929102342776852116782159153 \\
& 6500852703732781275448529779371677417472 x_5^{46} + \\
& 123834109038089934671294472137246930069523060031714250844706619429291169662133105896189791403231667938090032771493056909574046 \\
& 4706259501322146740447122160587271831552 x_5^{45} + \\
& 182020970274658793719058657381437633418879215571814116696924003282242176870141717350730229351112909089925971481867309468099964 \\
& 2196023948751131158452814694886969024512 x_5^{44} + \\
& 4007430504130089296344353908004460283088429242443631522884050976661843576629096044196085928404617097658704094576717598567410 \\
& 209841081365781438099257427726663221248 x_5^{43} - \\
& 42497825537001218362123773281724249357301195971726919570061852259609413271867163640427683432675518196930740559924270812831375 \\
& 563611884253337782941503899872012306432 x_5^{42} - \\
& 173095529689168945586540432091991868245292417145936250123555793382238071295393560910168427565624671717695384495055692485081314 \\
& 3755029162793643288152147445470405148672 x_5^{41} + \\
& 2813734388990484811052711343479211996610364188442552400568780705195161794258825128540138823297990664396921295940767223114649187 \\
& 77229254515877108902112388369782621696 x_5^{40} + \\
& 1478609812747959659776694877072531915021901854211298456892857835323485394858375580779587143734302060866166288995133976482398827 \\
& 757332314718512078892601607805397073920 x_5^{39} - \\
& 3530188807681883402087239068967648238062453085843469031573584103736822143928164648434968805617272052263983328298021721557176348 \\
& 61200763986215679045184328846763065024 x_5^{38} - \\
& 643733485008734767675328103272134976486121021695730058350986749622659224208821096802945223125730240168450387063503985689835364 \\
& 35159839070545812607121789160773867520 x_5^{37} + \\
& 2151585051628674174000873972276961698344627042149673451283893206863233108712310944766380976142565846440254189190380024194891099 \\
& 47723723046530048350802362337272172240 x_5^{36} + \\
& 1571392063478504403026681171500497608982888562929034843238917707762759456375200266902776859110521001010925495789223225311709229
\end{aligned}$$

$$\begin{aligned}
& 66640564949119170839940770175165369472 x_5^{35} - \\
& 6970087198338142892631419637725297682026478060976181883733568435716207363295338978929029041063837804373983101238056995847333791 \\
& 4855629458083803826009955084443825108 x_5^{34} - \\
& 2103953688392417375895946982748248064266470958465030128742828892428745081849383049981237613442379577045563903355672090403995503 \\
& 8834734820067394966023586825729448032 x_5^{33} + \\
& 131965103985659180399532159668209381342109038813694711607670737040719693247898472372338041113690183608893390155266387788851976 \\
& 5501542665365346335641906308193200195 x_5^{32} + \\
& 1201592723878650296985784908558376748726215478216571469977936360539641682241751683146521785538307944026986624023707333956026772 \\
& 434025973387112237446151013198603456 x_5^{31} - \\
& 1519072737952788288006422057910041941150694561452711227709754820142796061741821657930622193887605344329805371911870410187618742 \\
& 230354699710983368659613714243803548 x_5^{30} + \\
& 460150232801847338211144269544669020679298764358699406370892846094853216531336647270118195869886998840760288680189469787171307 \\
& 1441713839568058060774361029056640 x_5^{29} + \\
& 1072037308240519825398550954264078804114922488356381264733151234321243441642332015631365450737480419776126969139779220810709779 \\
& 01975014726759594214544359567844806 x_5^{28} - \\
& 1205007803208173899426678825454693999818863628379750061080029954331824730067438861128836451853733649699841007524332138999666791 \\
& 0512463244554688058984458230168128 x_5^{27} - \\
& 4574919195058858290500536373895931594305239443236153188094986352126908810537249334041955121003101046214104899351816127954000561 \\
& 321104434332847082016135629402216 x_5^{26} + \\
& 8467680055730501728498309795897327385736639557175721083875689127464511872957399436042824245131992208237009234326907187912147656 \\
& 22874112877423147714146807357216 x_5^{25} + \\
& 1118640233931113183860301041539440981114511620785798164764844603070211596622333247489744825642715135285115555446441878653615739 \\
& 07701098119634401155804176907365 x_5^{24} - \\
& 3230498793185293789740278257476629071739752957529095918297569364916959184652288022010351967916670095428738578674403630634796101 \\
& 1129382412439897351900300621440 x_5^{23} - \\
& 1132240307350821444955677467092009730095622147040085979652899973596155315813896118355596399204113581813473133533316374687012948 \\
& 382756870759752289337073156512 x_5^{22} + \\
& 7355743315375505556115230260932539680131456320426724557924376294900041025577364865159839635749294733965542390559879730142707162 \\
& 10030945337432183710523664640 x_5^{21} - \\
& 149995323362431404603425073872321250334869473786891372145385127714484252979740789199336829459667591503835752731730078484262891 \\
& 9315680630275508211504721620 x_5^{20} - \\
& 9379018779009283536647325410855017782509384280662718053972870134250000102254129843800530437703777545808897413872956521512630844 \\
& 012585339723324254084166656 x_5^{19} + \\
& 6152476558283094934553773971593228318618928370617652990425197508023708464026609098695948230413297837398726487448889169039432858 \\
& 51813601630195936736284060 x_5^{18} + \\
& 343436406707028212544215602819335146759901690389617008752172075152010721753553481361238393526468981180152586813158898672271705 \\
& 9036401414884124685763040 x_5^{17} - \\
& 5827838497126701986932835496666133870165981309385028011446993984529163044221367881757842128127005639952539867306999166204887917 \\
& 430059952129732325025811 x_5^{16} + \\
& 8398463073914106680331510770725921157556270719724560279729964962557979324838557623250672554661690384775755381583431190064741138 \\
& 17546450758626666989504 x_5^{15} - \\
& 5356541941694292107005896505026353847285129093142980419524124319758038315540273131024137208018589444482275106372773469600678846 \\
& 3098071336194295739100 x_5^{14} -
\end{aligned}$$

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131837782815027777127336891641150250427350777659622675682615021476138295395231637196585132855588210523219410313968075431468830  
 7874888360639240646784  $x_5^{13}$  +  
 1902077237100881325382190450900412331681260283403499858270746211313757599203601019544182069933187598481473443984343691055578661  
 667174220889910792374  $x_5^{12}$  +  
 4511046190804774740008184941508686054929624817963028231582886217622724959404025149854302776138821203068788152179522042646665949  
 1107595548644366784  $x_5^{11}$  -  
 199182986952917825259863931754226777282549769238367714945744943637238303240545879220782478839550814018029874450605636578192547  
 5758421301210424256  $x_5^{10}$  +  
 5525463302414569174625187350837831177167198568108845685272504608607381315336088049551088675517171587469973529053983408338810511  
 35329509909803360  $x_5^9$  +  
 926220684852009570179257703195743746202785342339468538701102654638957726165560723071303596803573478877546605001672549466872282  
 3336970481722075  $x_5^8$  -  
 5672574475164801873135086854241321214182423482250684291875985672216895556293522790017523397206426241300113918610058494279352311  
 055797710041600  $x_5^7$  -  
 1381908406226132452950875948252498427689518360441836766458835743894878227061669556619556797528744757532621857543694410597505098  
 45805856559400  $x_5^6$  +  
 152653905271550480059785492979754784690651844341401771106162829507870153113840725219008300417233514689735638232300926938731937  
 1263010656000  $x_5^5$  +  
 1867003566556869050458512424751716750175295547371196312368667351875534823826528942209107253593189949525377797225073796702198937  
 19635025000  $x_5^4$  -  
 259004066049704490143711257052597650563737993192612085149492257367139346125068573373223873355302574449971615120713521428248841  
 2067200000  $x_5^3$  -  
 28825647102342460384593078549690884780588475770063139272365052063071456382179858558092016390155095956155727995353578656824734  
 18500000  $x_5^2$  +  
 2224563290664978215283948523395639332992581815936524068816279095684157347469924868471458760599935821864409765646204754501435500  
 0000000  $x_5$  +  
 5321955484562939896106776126356300860859539556887924777311505619930765156630074075061463851472474466548206633328265696065429687  
 50000.



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