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Research article

Approximate analytical solution of singularly perturbed boundary value problems in MAPLE

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Abstract: In this paper, based on a new initial value method, an approximate analytical solution of singularly perturbed boundary value problems is obtained using a new and simple Maple program *spivp*. The algorithm is designed for practicing engineers or applied mathematicians who need a practical tool for solving these problems. Some examples are solved to illustrate the implementation of the program and the accuracy of the algorithm. Analytical and numerical results are compared with results in literature. The results confirm that the present method is accurate and offers a simple and easy practical tool for obtaining approximate analytical and numerical solutions of singularly perturbed boundary value problems.

Keywords: singularly perturbed boundary value problems; initial-value algorithms; mechanization **Mathematics Subject Classification:** 34B16, 34E15

1. Introduction

Singularly perturbed boundary value problems (SPBVPs), also known as stiff, are not easily treated analytically or numerically [1–20] and so that having a simple program to obtain analytical and numerical solutions of these problems would be interesting and helpful. The most common analytical method to study SPBVPs is the method of matched asymptotic expansions which involves finding outer and inner solutions of the problem and their matching [1–7]. Recently, Wang [9] present a Maple program and its corresponding algorithm to obtain a rational approximate solution for a class of nonlinear SPBVPs. Such an algorithm is based on the boundary value method presented by Wang [8] and using Taylor series and Pade expansion to approximate the boundary condition at infinity. In fact, the accuracy of the obtained boundary layer solution in [8,9] depends upon how far

2273

the Pade expansion approximates the solution, which consequently depends upon the radius of convergence and how many terms considered in the expansion [21]. Therefore, to get accurate results over the boundary layer region, more higher-order terms may be needed [10]. Moreover, the technique of using Taylor series and Pade expansion for solving boundary layer problems with a boundary condition at infinity often results in nonlinear algebraic equations and hence multiple solutions appear which leads to unsatisfied results. To overcome these drawbacks, we have established a new and simple Maple program *spivp*, based on a new initial value method [11], to obtain approximate analytical and numerical solutions of SPBVPs. The algorithm is designed for practicing engineers or applied mathematicians who need a practical tool for solving these problems. Some examples are solved to illustrate the implementation of the program and the accuracy of the algorithm. Analytical and numerical results are compared with results in literature. The results confirm that the present method is accurate and offers a simple and easy practical tool for obtaining approximate analytical and numerical solutions of SPBVPs.

2. Boundary value method [8]

Let us first recall the basic principles of the boundary value method presented in [8]. Consider the nonlinear SPBVP of the form

$$\varepsilon \frac{d^2 y}{dx^2} + p(x, y) \frac{dy}{dx} + q(x, y) = 0, x \in [0, 1],$$
(2.1)

with conditions

$$y(0) = \alpha, y(1) = \beta,$$
 (2.2)

where $0 < \varepsilon \ll 1$, α and β are given constants, p(x, y) and q(x, y) are assumed to be sufficiently continuously differentiable functions on [0, 1], and $p(x, y) \ge M > 0$ for every $x \in [0, 1]$, where *M* is a positive constant. Under these assumptions, the SPBVP (2.1-2.2) has a unique solution y(x) which in general displays a boundary layer at x = 0 for small values of ε .

Theorem 1. The solution y(x) of the nonlinear SPBVP (2.1-2.2) can be expressed as:

$$y(x) = u(x) + v(t) + O(\varepsilon), \ t = \frac{x}{\varepsilon}$$

where u(x) and v(t) are the solutions of the following IVP and BVP given respectively by:

$$\begin{cases} p(x,u)\frac{du}{dx} + q(x,u) = 0, \\ u(1) = \beta, \end{cases}$$
(2.3)

and

$$\begin{cases} \frac{d^2v}{dt^2} + p(0, u(0) + v(t))\frac{dv}{dt} = 0, \\ v(0) = \alpha - u(0), \lim_{t \to +\infty} v(t) = 0. \end{cases}$$
(2.4)

Proof. See Ref. [8].

In fact, v(t) cannot be always solved from Eq. (2.4) and so that one can assume that $v'(0) = \delta$ and use Taylor series and Pade expansion to determine the value of δ using the condition at infinity and obtain an approximate rational solution for Eq. (2.4) [8,9].

3. Initial value method

In this section, an initial value method is established from the boundary value method in Theorem1 for solving SPBVP (2.1-2.2).

Theorem 2. The solution y(x) of the nonlinear SPBVP (2.1-2.2) can be expressed as:

$$y(x) = u(x) + v(t) + O(\varepsilon), \quad t = \frac{x}{\varepsilon},$$
(3.1)

where u(x) and v(t) are the solutions of the following IVPs:

$$\begin{cases} p(x, u)\frac{du}{dx} + q(x, u) = 0, \\ u(1) = \beta, \end{cases}$$
(3.2)

and

$$\begin{cases} \frac{dv}{dt} + f(0, u(0) + v(t)) = f(0, u(0)) \\ v(0) = \alpha - u(0). \end{cases}$$
(3.3)

where

$$f(0, u(0) + v(t)) = \int p(0, u(0) + v(t)) \, dv$$

Proof. Integrating Eq. (2.4) results in

$$\int_{t}^{\infty} \frac{d^2 v}{ds^2} ds + \int_{t}^{\infty} \left(p(0, u(0) + v(s)) \frac{dv}{ds} \right) ds = 0,$$
(3.4)

$$\left[\frac{dv}{ds} + f(0, u(0) + v(s))\right]_{s=t}^{s \to +\infty} = 0,$$
(3.5)

where $f(0, u(0) + v(s)) = \int p(0, u(0) + v(s)) dv$.

Thus we have

$$\frac{dv}{dt} + f(0, u(0) + v(t)) = \lim_{s \to +\infty} \frac{dv}{ds} + f(0, u(0) + \lim_{s \to +\infty} v(s)),$$
(3.6)

which results in a first order initial value problem given by

$$\frac{dv}{dt} + f(0, u(0) + v(t)) = f(0, u(0)),$$
(3.7)

with the initial condition

$$v(0) = \alpha - u(0).$$

Thus, we have replaced the second order BVP (2.4) by the approximate first order IVP (3.7). Obviously Theorem 2 agrees with the theoretical results in [11].

Remark. The value of δ used in the boundary value method [8,9] is easily obtained from Eq. (3.3) as $\delta = f(0, u(0)) - f(0, \alpha)$

AIMS Mathematics

2275

3.1. Theoretical results for linear problems

Using the initial value method presented in Theorem 2 for linear problems results in some useful theoretical results.

Let's consider the linear SPBVP

$$\varepsilon \frac{d^2 y}{dx^2} + p(x)\frac{dy}{dx} + q(x)y = 0, x \in [0, 1],$$
(3.8)

with conditions

$$y(0) = \alpha, y(1) = \beta, \tag{3.9}$$

where $0 < \varepsilon \ll 1$, α and β are given constants, p(x) and q(x) are assumed to be sufficiently continuously differentiable functions on [0, 1], and $p(x) \ge M > 0$ for every $x \in [0, 1]$, where *M* is a positive constant. The assumption merely implies that the boundary layer will be in the neighborhood of x = 0.

Theorem 3. The solution of the linear SPBVP (3.8-3.9) can be approximated by:

$$y(x) = \beta e^{\int_{1}^{x} \frac{-q(x)}{p(x)} ds} + \left((\alpha - u(0)) e^{\frac{-p(0)x}{\varepsilon}} \right) + O(\varepsilon)$$
(3.10)

Proof. According to the initial value method presented in Theorem 2, we have

$$y(x) = u(x) + v(t) + O(\varepsilon), \quad t = \frac{x}{\varepsilon},$$
(3.11)

where u(x) and v(t) are the solutions of the following IVPs:

$$\begin{cases} p(x)\frac{du}{dx} + q(x)u = 0, \\ u(1) = \beta, \end{cases}$$
(3.12)

and

$$\begin{cases} \frac{dv}{dt} + p(0)v(t) = 0\\ v(0) = \alpha - u(0). \end{cases}$$
(3.13)

The reduced problem (3.12) and the boundary layer corrected problem (3.13) are easily solvable since they are separable and thus (3.11) results in the asymptotic analytical solution given by

$$y(x) = \beta e^{\int_1^x \frac{-q(x)}{p(x)} ds} + \left((\alpha - u(0)) e^{\frac{-p(0)x}{\varepsilon}} \right) + O(\varepsilon),$$

which is the well known asymptotic solution given in [2,11].

Obviously, the value of δ used in the boundary value method in Theorem 1 can be obtained directly from (3.13) as $\delta = v'(0) = -p(0)(\alpha - u(0))$ which agrees with the theoretical results in [8,9].

4. Approximate analytical solution in Maple

Mathematical mechanization is one of the important methods of mathematical studies [9,12-13,22-26]. In Maple, obtaining analytical solution of the nonlinear SPBVP (2.1-2.2) through solving the IVPs (3.2) and (3.3) can be well established a very simple symbolic program that can

AIMS Mathematics

easily calculate the approximate analytical and numerical solutions. Now if we want to solve SPBVP (2.1-2.2) using the initial value method in Theorem 2, everything we have to do is just to input information about the SPBVP, and the program will give out the solution. The program *spivp* is as follows:

spivp := proc (p, q, alpha, beta) local p0, q0, pu, qu, f, us, k, vs, sol, vx, iniu, iniv; p0 := unapply(p, x, y); q0 := unapply(q, x, y); pu := p0(x, u(x)); qu := q0(x, u(x)); f := unapply(int(p, y), x, y); iniu := u(1) = beta; us := rhs(dsolve(pu*(diff(u(x), x))+qu = 0, iniu, u(x))); k := eval(us, x = 0); iniv := v(0) = alpha-k; vs := rhs(dsolve(diff(v(t), t)+f(0, v(t)+k) = f(0, k), iniv, v(t))); $vx := subs(t = x/\epsilon, vs);$ sol := simplify(us+vx, size) end proc

For convenience, in the program *spivp*, we always let *x* represents the independent variable, and set the parameters as following:

p: the coefficient, p(x, y), in the equation which to be solved.

q: the coefficient, q(x, y), in the equation which to be solved.

 α, β : the boundary values.

For example, if we consider the SPBVP:

$$\varepsilon \frac{d^2 y}{dx^2} + e^y \frac{dy}{dx} - \frac{\pi}{2} \sin\left(\frac{\pi x}{2}\right) e^{2y} = 0, \ y(0) = 0, \ y(1) = 0.$$

One just input the following command

spivp(exp(y), -pi/2.sin(pi.x/2).exp(2.y), 0, 0);

5. Examples

In this section, four examples are solved to illustrate the implementation and the accuracy of the method.

Example 1. Consider the nonlinear SPBVP from Bender and Orszag [6]

$$\varepsilon \frac{d^2 y}{dx^2} + 2\frac{dy}{dx} + e^y = 0, \ y(0) = y(1) = 0.$$
(5.1)

AIMS Mathematics

To solve this problem, we input the following command:

spivp(2, exp(y), 0, 0)

The output is:

$$y(x) - \ln\left(\frac{1}{2}x + \frac{1}{2}\right) - \ln(2)e^{-\frac{2x}{\varepsilon}}$$

which is the uniform valid approximate solution given by Bender and Orszag [6]. The solution of Example 1 at different values of ε is presented in Figure 1.



Figure 1. Solution of Example 1 at different values of ε .

Example 2. Consider the nonlinear SPBVP from O'Malley [2]

$$\varepsilon \frac{d^2 y}{dx^2} + e^y \frac{dy}{dx} - \frac{\pi}{2} \sin\left(\frac{\pi x}{2}\right) e^{2y} = 0, \ y(0) = y(1) = 0.$$
(5.2)

Input the following command in Maple

spivp(exp(y), -Pi/2.sin(Pi.x/2).exp(2.y), 0, 0);

The output is

$$y(x) = -\ln\left(\cos(\frac{1}{2}\pi x) + 1\right) + \ln\left(\frac{-2}{e^{-\frac{1}{2}\frac{x}{e}} - 2}\right),$$

which is the uniform valid approximate solution given by O'Malley [2]. The solution of Example 2 at different values of ε is presented in Figure 2.

AIMS Mathematics



Figure 2. Solution of Example 2 at different values of ε .

As shown from the previous two examples, the present algorithm offers a simple and easy tool for obtaining asymptotic analytical and approximate numerical solutions for SPBVPs.

Example 3. Consider the nonlinear SPBVP from Kevorkian and Cole [7]

$$\varepsilon \frac{d^2 y}{dx^2} + y \frac{dy}{dx} - y = 0, \ y(0) = -1, y(1) = 3.9995$$
(5.3)

For comparison purpose, we write the uniform approximation provided by Kevorkian and Cole [7] as

$$y(x) = x + c_1 \tanh[c_1(x/\varepsilon + c_2)/2],$$
 (5.4)

where $c_1 = 2.9995$ and $c_2 = 1/c_1 \ln[(c_1 - 1)/(c_1 + 1)]$. For this problem, Wang [9] has obtained a solution given by:

$$y(x) = x + 2.9995 + \frac{L(t)}{Q(t)},$$
(5.5)

where $t = x/\varepsilon$, and

$$L(t) = -3.9995 + (1.99975 + \delta) t - 0.39995(1 + 2\delta) t^{2} + [0.0333292(1 + 2\delta) -0.25\delta + 0.1\delta(1 + 2\delta) + 0.166667(\delta - \delta^{2})] t^{3},$$

$$O(t) = 1 - 0.5\delta + 0.1(1 + 2\delta)t^{2} - 0.0083333(1 + 2\delta)t^{3},$$
(5.6)

where $\delta = -1.999791999$.

Here we apply the present method by inputting the following command in Maple

AIMS Mathematics

spivp(*y*, -*y*, -1, 3.9995); The output is



Figure 3. Comparison of the numerical solutions of Example 3 at $\varepsilon = 0.001$.

x	Present	Results in[16]	Results in[17]	Results in[7]
0.00	-1.000000	-1.000000	-1.000000	-1.0000000
0.5ε	1.1484593	2.1363630	1.1484592	1.1484592
1.0ε	2.4569397	2.8140590	2.4569396	2.4569396
0.1	3.0994999	3.0995020	3.0962686	3.0994999
0.2	3.1995000	3.1995020	3.1963699	3.1995000
0.3	3.2995000	3.2995020	3.2964650	3.2995000
0.4	3.3995000	3.3995010	3.3965544	3.3995000
0.5	3.4995000	3.4995010	3.4466386	3.4995000
0.6	3.5995000	3.5995000	3.5967184	3.5995000
0.7	3.6995000	3.6995000	3.6967939	3.6995000
0.8	3.7995000	3.7995000	3.7968656	3.7995000
0.9	3.8995000	3.8995000	3.8969335	3.8995000
1.0	3.9995000	3.9995000	3.9969980	3.9995000

Table 1. Comparison of the numerical solutions of Example 3 at $\varepsilon = 10^{-3}$.

	Tuble 2. Solution of Example 5 at american values of 0.				
x	$\varepsilon = 10^{-2}$	$\varepsilon = 10^{-4}$	$\varepsilon = 10^{-5}$	$\varepsilon = 10^{-6}$	
0.00	-1.0000000	-1.0000000	-1.0000000	-1.0000000	
0.5ε	1.1529593	1.1480093	1.1479643	1.1479598	
1.0ε	2.4659397	2.4560397	2.4559497	2.4559407	
0.1	3.0995000	3.0995000	3.0995000	3.0995000	
0.2	3.1995000	3.1995000	3.1995000	3.1995000	
0.3	3.2995000	3.2995000	3.2995000	3.2995000	
0.4	3.3995000	3.3995000	3.3995000	3.3995000	
0.5	3.4995000	3.4995000	3.4995000	3.4995000	
0.6	3.5995000	3.5995000	3.5995000	3.5995000	
0.7	3.6995000	3.6995000	3.6995000	3.6995000	
0.8	3.7995000	3.7995000	3.7995000	3.7995000	
0.9	3.8995000	3.8995000	3.8995000	3.8995000	
1.0	3.9995000	3.9995000	3.9995000	3.9995000	

Table 2. Solution of Example 3 at different values of ε .

In Figure 3, the present solution (5.7) and the obtained solutions by Wang [9], Eqs. (5.5-5.6) are compared with the uniform approximate solution in [7] at $\varepsilon = 0.001$. As shown from Figure 3 the solution (5.5-5.6), solid red line, deviates much from our solution (5.7), dashed blue line, and the uniform approximate solution, dotted blue line, in [7]. In fact, applying the condition $\lim_{t\to+\infty} v(t) = 0$ to (5.6) results in two different values for δ ($\delta = -1.999791999$, $\delta = -0.5$) and each one of them leads to unsatisfied results. The numerical results of Example 3 are compared with the results in literature in Table 1. Moreover, the results for different values of ε are presented in Table 2.

Example 4. Finally, we consider the following linear SPBVP from Bender and Orszag [6]

$$\varepsilon \frac{d^2 y}{dx^2} + \frac{dy}{dx} - y = 0, \ y(0) = y(1) = 1.$$
(5.8)

For this problem, Wang [9] has obtained the following rational solution:

$$y(x) = \frac{e^x}{e} + \frac{L(x)}{Q(x)},$$
(5.9)

where

$$L(x) = \varepsilon \left(-19115.326\varepsilon^4 + 200691.9566x^4 - 19115.32x^3\varepsilon + 8427734.7x^2\varepsilon^2 - 1070458.24x\varepsilon^3 \right),$$

and

$$Q(x) = -30240\varepsilon^5 + 196560x\varepsilon^4 + 102480x^2\varepsilon^3 + 22260x^3\varepsilon^2 + 2490x^4\varepsilon + 125x^5.$$

Using the following command :

spivp(1, -y, 1, 1),

AIMS Mathematics

the obtained solution is

$$y(x) = e^{x-1} + (e-1)e^{-x/\varepsilon - 1}.$$
(5.10)

The numerical results of Example 4 are shown in Tables 3, 4 and Figure 4.

Obviously, the solution in (5.9) needs more higher-order terms in Pade's expansion to match the boundary layer behavior [10].



Figure 4. Solution of Example 4 at different values of ε .

Table 3. Comparison of the numerical solutions of Example 4 at $\varepsilon = 10^{-3}$.

x	Present	Results in[6]	Results in[9]	Results in[17]	Exact solution
0.00	1.0000000	1.0000000	1.0000000	1.0000000	1.0000000
0.01	0.3716054	0.3716050	39.621608	0.3712379	0.3719724
0.02	0.3753111	0.3753109	34.943933	0.3749439	0.3756784
0.03	0.3790831	0.3790832	29.856212	0.3787160	0.3794502
0.04	0.3828929	0.3828929	25.706355	0.3825260	0.3832599
0.05	0.3867410	0.3867409	22.458425	0.3863742	0.3871079
0.10	0.4065697	0.4065696	13.632736	0.4062043	0.4069350
0.30	0.4965853	0.4965853	5.5064473	0.4962382	0.4969324
0.50	0.6065307	0.6065305	3.6923714	0.6062278	0.6068334
0.70	0.7408182	0.7408182	2.9700713	0.7405963	0.7410401
0.90	0.9048374	0.9048373	2.6496651	0.9047471	0.9049277
1.00	1.0000000	1.0000000	2.5738182	1.0000000	1.0000000

AIMS Mathematics

	Tuble 11 Solution of Example 1 at anterent values of 0.				
x	$\varepsilon = 10^{-2}$	$\varepsilon = 10^{-4}$	$\varepsilon = 10^{-5}$	$\varepsilon = 10^{-6}$	
0.00	1.00000000	1.000000000	1.00000000	1.00000000	
0.01	0.60412084	0.37157669	0.37157669	0.37157669	
0.02	0.46085931	0.37531109	0.37531109	0.37531109	
0.03	0.41055446	0.37908303	0.37908303	0.37908303	
0.04	0.39447057	0.38289288	0.38289288	0.38289288	
0.05	0.39100021	0.38674102	0.38674102	0.38674102	
0.10	0.40659835	0.40656965	0.40656965	0.40656965	
0.30	0.49658530	0.49658530	0.49658530	0.49658530	
0.50	0.60653065	0.60653065	0.60653065	0.60653065	
0.70	0.74081822	0.74081822	0.74081822	0.74081822	
0.90	0.90483741	0.90483741	0.90483741	0.90483741	
1.00	1.00000000	1.00000000	1.00000000	1.00000000	

Table 4. Solution of Example 4 at different values of ε .

6. Conclusion

In this paper, we have presented approximate analytical and numerical solutions of SPBVPs using a new and simple Maple program *spivp*. The method is based on an initial value method that replaces the original SPBVP by two asymptotically approximate IVPs. These IVPs are solved analytically and their solutions are combined to approximate the analytical solution of the original SPBVP. The method is very easy to implement on any computer with a minimum problem preparation. We have implemented it in Maple and applied it to four examples and compared the obtained results with the results in literature. The results indicate that the method is accurate and offers a simple and easy practical tool for the practicing engineers or applied mathematicians who need a practical tool for obtaining approximate analytical and numerical solution of SPBVPs.

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Conflict of interest

The author declares no conflict of interest.

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2282

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2283

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