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#### Research article

# Radial stationary solutions to a class of wave system as well as their asymptotical behavior

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**Abstract:** For a stationary version to a class of wave system

$$\begin{cases} -\left(a_{1}+b_{1}\int_{\mathbb{R}^{3}}|\nabla u|^{2}dx+c\int_{\mathbb{R}^{3}}|\nabla v|^{2}dx\right)\Delta u+u=\frac{p}{Q}|u|^{p-2}u|v|^{q},\\ -\left(a_{2}+b_{2}\int_{\mathbb{R}^{3}}|\nabla v|^{2}dx+c\int_{\mathbb{R}^{3}}|\nabla u|^{2}dx\right)\Delta v+v=\frac{q}{Q}|u|^{p}|v|^{q-2}v, \end{cases}$$

 $u, v \in H^1_r(\mathbb{R}^3)$ , by establishing a variant variational identity and constraint set, we prove that for  $a_s > 0$ ,  $b_s > 0$ , (s = 1, 2),  $c \ge 0$  and p > 1, q > 1 with  $Q := p + q \in (2, 6)$ , the system admits a positive radially symmetric ground state solution in  $H^1_r(\mathbb{R}^3) \times H^1_r(\mathbb{R}^3)$ . Moreover, for any fixed  $a_1 > 0$  and  $a_2 > 0$ , as  $b_1^2 + b_2^2 + c^2 \to 0$ , this solution converges to a positive radially symmetric solution to

$$-a_1 \Delta u + u = \frac{p}{Q} |u|^{p-2} u |v|^q, \quad -a_2 \Delta v + v = \frac{q}{Q} |u|^p |v|^{q-2} v, \quad u, \ v \in H^1_r(\mathbb{R}^3).$$

**Keywords:** system with Kirchhoff term; radially symmetric solutions; variant variational identity **Mathematics Subject Classification:** 35J20

#### 1. Introduction and main result

Let  $H^1(\mathbb{R}^3) := W^{1,2}(\mathbb{R}^3)$  be the usual Hilbert space. Denoted by  $H^1_r(\mathbb{R}^3)$  the subspace of  $H^1(\mathbb{R}^3)$  which contains all radially symmetric functions in  $H^1(\mathbb{R}^3)$ . In this paper we are concerned with the

existence of positive radial solutions to the following coupled Kirchhoff type system:

$$\begin{cases}
-\left(a_{1}+b_{1}\int_{\mathbb{R}^{3}}|\nabla u|^{2}dx+c\int_{\mathbb{R}^{3}}|\nabla v|^{2}dx\right)\Delta u+u=\frac{p}{Q}|u|^{p-2}u|v|^{q}, \\
-\left(a_{2}+b_{2}\int_{\mathbb{R}^{3}}|\nabla v|^{2}dx+c\int_{\mathbb{R}^{3}}|\nabla u|^{2}dx\right)\Delta v+v=\frac{q}{Q}|u|^{p}|v|^{q-2}v, \\
u,v\in H_{r}^{1}(\mathbb{R}^{3}), \quad u:=u(x), \ v:=v(x), \ x\in\mathbb{R}^{3}
\end{cases}$$
(1.1)

where  $a_1 > 0$ ,  $a_2 > 0$ ;  $b_1 > 0$ ,  $b_2 > 0$ ,  $c \ge 0$  and p > 1, q > 1 with 2 < Q := p + q < 6. The number 6 is the critical exponent of the embedding  $H^1(\mathbb{R}^3) \to L^s(\mathbb{R}^3)$   $(2 \le s \le 6)$ .

Such kind of Kirchhoff system is originated from the study of Kirchhoff string [8] and recent work by Matsuyama and Ruzhansky [20]. In [20], the authors propose the following Kirchhoff systems

$$\begin{cases} \partial_t^2 \varphi_1 - \rho_1 \left( 1 + \int_{\mathbb{R}^3} |\nabla \varphi_1|^2 dx + c \int_{\mathbb{R}^3} |\nabla \varphi_2|^2 dx \right) \Delta \varphi_1 + P_1(x, D_x) \varphi_1 = 0, \\ \partial_t^2 \varphi_2 - \rho_1 \left( 1 + \int_{\mathbb{R}^3} |\nabla \varphi_2|^2 dx + c \int_{\mathbb{R}^3} |\nabla \varphi_1|^2 dx \right) \Delta \varphi_2 + P_2(x, D_x) \varphi_2 = 0 \\ \varphi_1 := \varphi_1(t, x), \quad \varphi_2 := \varphi_2(t, x), \quad t > 0, \ x \in \mathbb{R}^3. \end{cases}$$
(Ksys)

In [20], analytic methods are used to study the Cauchy problem of (Ksys). Stationary version related to (Ksys) with nonlinear perturbation has attracted more and more attentions. In [26], the authors study the existence of solutions to the following system

$$\begin{cases} -\left(a+b\int_{\mathbb{R}^N} |\nabla u|^2 dx\right) \Delta u + V(x)u = \frac{\partial F(x,u,v)}{\partial u}, \\ -\left(c+d\int_{\mathbb{R}^N} |\nabla v|^2 dx\right) \Delta v + V(x)v = \frac{\partial F(x,u,v)}{\partial v}, \\ u(x) \to 0 \quad \text{and} \quad v(x) \to 0 \text{ as } |x| \to \infty. \end{cases}$$

Besides some other conditions, the authors [26] assume that the nonlinear perturbation F(x, u, v) satisfies "4-AR condition" for system in the sense that

$$u\frac{\partial F(x,u,v)}{\partial u}+v\frac{\partial F(x,u,v)}{\partial v}\geq 4F(x,u,v), \quad \forall \ x\in\mathbb{R}^N, \ (u,v)\in\mathbb{R}^2.$$

We point out that a natural and important case of  $F(x, u, v) = |u|^p |v|^q$  with p > 1, q > 1 and p + q < 4 was *not* covered by the results of [26]. In [19], the author considers the following system

$$\begin{cases} -\left(\epsilon^{2}a + \epsilon b \int_{\mathbb{R}^{N}} |\nabla u|^{2} dx\right) \Delta u + V(x)u = \frac{1}{\mu} \frac{\partial F(u, v)}{\partial u}, \\ -\left(\epsilon^{2}a + \epsilon b \int_{\mathbb{R}^{N}} |\nabla v|^{2} dx\right) \Delta v + V(x)v = \frac{1}{\mu} \frac{\partial F(u, v)}{\partial v}, \\ u, \quad v \in H^{1}(\mathbb{R}^{N}). \end{cases}$$

The author assumes that  $\mu > 4$  and for s > 0,  $F(su, sv) = s^{\mu}F(u, v)$  and prove the existence of solutions by variational methods. Again, the case of  $F(u, v) = |u|^p |v|^q$  with p > 1, q > 1 and  $p + q \le 4$  is not studied, either.

The main purpose of the present paper is to study the existence of positive radial solutions to (1.1) for all p > 1, q > 1 with 2 . By a solution to (1.1), we mean a critical point of the following functional

$$I(u,v) = \frac{1}{2} \int \left( a_1 |\nabla u|^2 + a_2 |\nabla v|^2 \right) dx + \frac{1}{2} \int \left( |u|^2 + |v|^2 \right) dx - \frac{1}{Q} \int |u|^p |v|^q dx + \frac{1}{4} \left( b_1 \left( \int |\nabla u|^2 dx \right)^2 + b_2 \left( \int |\nabla v|^2 dx \right)^2 + 2c \int |\nabla u|^2 dx \int |\nabla v|^2 dx \right)$$

defined on  $H_r^1(\mathbb{R}^3) \times H_r^1(\mathbb{R}^3)$ . According to the Sobolev embedding theorem, the functional I is well defined and  $C^1$ . It is easily to see that (0,0) is a solution to (1.1), which is usually called trivial solution. We call (u,v) a semitrivial solution if I'(u,v)=(0,0) and  $u \neq 0, v=0$  or  $u=0, v\neq 0$ . If I'(u,v)=(0,0) and  $u\neq 0, v\neq 0$ , then (u,v) is called a nontrivial solution. If (u,v) is a solution to (1.1) and u>0, v>0, then (u,v) is called a positive solution.

**Definition 1.1.** Denote  $H := H^1_r(\mathbb{R}^3) \times H^1_r(\mathbb{R}^3)$ . A nontrivial solution  $(u, v) \in H$  is called a positive radially symmetric ground state solution to (1.1) if  $I(u, v) \leq I(\bar{u}, \bar{v})$  for any  $(\bar{u}, \bar{v}) \in H \setminus \{(0, 0)\}$  and  $I'(\bar{u}, \bar{v}) = (0, 0)$ .

The first result of the present paper is the following theorem.

**Theorem 1.2.** Assume that  $a_1 > 0$ ,  $a_2 > 0$ ,  $b_1 > 0$ ,  $b_2 > 0$ ,  $c \ge 0$ ; p > 1, q > 1 and  $0 . Then the system (1.1) has a positive radially symmetric ground state solution <math>(u, v) \in H$ .

We emphasize that, as a corollary of Theorem 1.2, when c = 0, Theorem 1.2 generalizes results of [19,26] in the sense that we can get the existence of solutions in the case of  $F(u, v) = |u|^p |v|^q$ , p > 1, q > 1 and 2 .

We also point out that, when  $a_1 = a_2$ ,  $b_1 = b_2$ , c = 0 and u = v, (1.1) becomes a semilinear elliptic equation with Kirchhoff term. In the past ten years, a lot of mathematicians have made contributions to the existence and multiplicity of solutions to

$$-\left(a+b\int_{\mathbb{R}^3}|\nabla u|^2dx\right)\Delta u+V(x)u=f(x,u),\quad x\in\mathbb{R}^3,\ u\in H^1(\mathbb{R}^3),\tag{SK}$$

see e. g. [1,7,9,10,12,14-16,25] and some related background [6,13,17]. In the process of studying (SK), to overcome the difficulties created by the Kirchhoff trem, authors usually need to assume that f(x,u) satisfies either 4–superlinear at infinity in the sense that

$$\lim_{|u| \to +\infty} \frac{\int_0^u f(x, s) ds}{|u|^4} = +\infty \quad \text{uniformly in} \quad x \in \mathbb{R}^3$$

or satisfies '4-Ambrosetti-Rabinowitz condition' with the version

(AR) there is  $\mu > 4$  such that  $0 < \mu \int_0^u f(x, s) ds \le f(x, u)u$  for all  $u \ne 0$ .

Recently, for single equation with a Kirchhoff term, such kind of results has been extended to fractional Kirchhoff type equation or p-Kirchhoff equations, see e. g. [2, 4, 5, 21, 22] as well as the references therein. Therefore a special case (i. e.,  $a_1 = a_2$ ,  $b_1 = b_2$  and u = v) of Theorem 1.2 can be a complement to the results of [11,12,18,24] since we can get solutions for nonlinear growth in the full range of (2, 6).

Noticing that when  $b_1 = 0$ ,  $b_2 = 0$  and c = 0, (1.1) is the usual semilinear elliptic system. While for  $b_1, b_2 > 0$  and c > 0, any solution to (1.1) usually depends on  $b_1, b_2$  and c. A natural and interesting question is: what is the asymptotical behavior of the solution (u, v) as  $|b_1|^2 + |b_2|^2 + |c|^2 \rightarrow 0$ . To answer this question precisely, we denote this (u, v) by  $(u_{b_1,b_2,c}, v_{b_1,b_2,c})$  and the (1.1) by  $(1.1)_{b_1,b_2,c}$ . Then we can prove the following theorem.

**Theorem 1.3.** Let p > 1, q > 1 and  $2 . For any fixed <math>a_1 > 0$  and  $a_2 > 0$ , if sequences  $b_1^{(n)} > 0$ ,  $b_2^{(n)} > 0$ ,  $c^{(n)} \ge 0$ , and  $|b_1^{(n)}|^2 + |b_2^{(n)}|^2 + |c^{(n)}|^2 \to 0$  as  $n \to \infty$ , then the sequence of solutions  $(u_{b_1^{(n)}, b_2^{(n)}, c^{(n)}}, v_{b_1^{(n)}, b_2^{(n)}, c^{(n)}})$  to  $(1.1)_{b_1^{(n)}, b_2^{(n)}, c^{(n)}}$  converges to a positive radially symmetric solution to  $(1.1)_{0,0,0}$ .

In order to prove Theorem 1.2 and Theorem 1.3, we use variational methods. But the methods used in [11, 12, 19, 26] to deal with the case of 4-superlinear or 4-(AR) condition can not be applied here (since our p+q may be less than 4). The main strategy of the present paper is to establish a constrained set  $\mathcal{M}$  and then minimize the functional  $\mathcal{I}$  over  $\mathcal{M}$ , see (2.1) in Section 2. And we manage to prove that the minimum d of  $\mathcal{I}|_{\mathcal{M}}$  can be solved and the minimizer is a solution to (1.1) as required. This idea is inspired from Ruiz [23] where the author studied a class of Schrödinger-Poisson system. But in the present paper, we have to overcome the Kirchhoff term and the coupling between u and v.

This paper is organized as follows. In Section 2, we study the variational structure of (1.1) and establish a set  $\mathcal{M}$ , proving that this set is indeed a manifold and can be a natural constraint of the functional  $\mathcal{I}$ , see Lemma 2.6 and Lemma 2.7. In Section 3, we prove Theorem 1.2. In Section 4, we study the asymptotical behavior of solutions obtained in Theorem 1.2 with respect to  $b_1$ ,  $b_2$  and c, where we finish the proof of Theorem 1.3.

**Notation.** Throughout the paper, all integrals are taken over  $\mathbb{R}^3$  unless specified.  $L^s(\mathbb{R}^3)$   $(1 \le s < +\infty)$  is the usual Lebesgue space with the standard norm  $|u|_s$ . For  $a_1 > 0$ ,  $a_2 > 0$ , we use a norm on H:  $||u||^2 := \int (a_1|\nabla u|^2 + u^2 + a_2|\nabla v|^2 + v^2) dx$  whose inner product is  $\langle (u_1, v_1), (u_2, v_2) \rangle = \int (a_1\nabla u_1\nabla u_2 + u_1u_2 + a_2\nabla v_1\nabla v_2 + v_1v_2) dx$ .

#### 2. Variational structure

In this section, we establish variational framework of (1.1). Since we only assume that p > 1 and q > 1 and 2 < Q := p + q < 6, the standard Nehari type constraint can not be applied here. As we pointed out in the introduction, our strategy is to construct a set  $\mathcal{M}$  which is a manifold and prove that this set can be a natural constraint. Then we define a suitable minimization problem and prove that the minimizer can be achieved. We start with the following Pohozaev identity.

**Lemma 2.1.** Let  $(u, v) \in H$  be a solution to (1.1). Then  $\mathcal{P}(u, v) = 0$ , where

$$\mathcal{P}(u,v) := \int (a_1 |\nabla u|^2 + a_2 |\nabla v|^2) dx + 3 \int (|u|^2 + |v|^2) dx - \frac{6}{Q} \int |u|^p |v|^q dx + \left(b_1 \left(\int |\nabla u|^2 dx\right)^2 + 2c \int |\nabla u|^2 dx \int |\nabla v|^2 dx + b_2 \left(\int |\nabla v|^2 dx\right)^2\right).$$

*Proof.* We only prove it formally. Let  $(u, v) \in H$  be a solution to (1.1). Multiplying the first equation of the system (1.1) by  $x \cdot \nabla u$  and the second equation by  $x \cdot \nabla v$  respectively, and integrating by parts,

we get that

$$\left(a_{1} + b_{1} \int_{\mathbb{R}^{3}} |\nabla u|^{2} dx + c \int_{\mathbb{R}^{3}} |\nabla v|^{2} dx\right) \int \nabla u \nabla(x \cdot \nabla u) dx + \int u(x \cdot \nabla u) dx 
= \frac{p}{Q} \int |u|^{p-2} u|v|^{q} (x \cdot \nabla u) dx; 
\left(a_{2} + b_{2} \int_{\mathbb{R}^{3}} |\nabla v|^{2} dx + c \int_{\mathbb{R}^{3}} |\nabla u|^{2} dx\right) \int \nabla v \nabla(x \cdot \nabla v) dx + \int v(x \cdot \nabla v) dx 
= \frac{q}{Q} \int |u|^{p} |v|^{q-2} v(x \cdot \nabla v) dx.$$

By simple computation, we also have that

$$\int \nabla u \nabla (x \cdot \nabla u) dx = -\frac{1}{2} \int |\nabla u|^2 dx, \qquad \int \nabla v \nabla (x \cdot \nabla v) dx = -\frac{1}{2} \int |\nabla v|^2 dx,$$
$$\int u(x \cdot \nabla u) dx = -\frac{3}{2} \int |u|^2 dx, \qquad \int v(x \cdot \nabla v) dx = -\frac{3}{2} \int |v|^2 dx,$$

and

$$p\int |u|^{p-2}u|v|^q(x\cdot\nabla u)dx=-3\int |u|^p|v|^qdx-q\int |u|^p|v|^{q-2}v(x\cdot\nabla v)dx.$$

Combining the above equalities, we deduce that

$$\int (a_{1}|\nabla u|^{2} + a_{2}|\nabla v|^{2}) dx + 3 \int (|u|^{2} + |v|^{2}) dx - \frac{6}{Q} \int |u|^{p}|v|^{q} dx + \left(b_{1} \left(\int |\nabla u|^{2} dx\right)^{2} + 2c \int |\nabla u|^{2} dx \int |\nabla v|^{2} dx + b_{2} \left(\int |\nabla v|^{2} dx\right)^{2}\right) = 0.$$

The proof is complete.

Next, we define  $\mathcal{G}(u, v) := \frac{1}{4} \langle I'(u, v), (u, v) \rangle + \frac{1}{4} \mathcal{P}(u, v)$ . Then

$$\mathcal{G}(u,v):=\frac{1}{2}\mathcal{A}(u,v)+C(u,v)+\frac{1}{2}\mathcal{B}(u,v)-\frac{Q+6}{4Q}\int|u|^p|v|^qdx,$$

where  $\mathcal{A}(u,v) := \int \left(a_1|\nabla u|^2 + a_2|\nabla v|^2\right)dx$ ,  $C(u,v) := \int \left(|u|^2 + |v|^2\right)dx$  and  $\mathcal{B}(u,v) := b_1\left(\int |\nabla u|^2dx\right)^2 + 2c\int |\nabla u|^2dx\int |\nabla v|^2dx + b_2\left(\int |\nabla v|^2dx\right)^2$ . Define

$$\mathcal{M} := \{(u, v) \in H \setminus \{(0, 0)\} : \mathcal{G}(u, v) = 0\}.$$

Clearly if  $(u, 0) \in \mathcal{M}$ , then u = 0 and if  $(0, v) \in \mathcal{M}$ , then v = 0. Hence

$$\mathcal{M} = \{(u, v) \in H : \mathcal{G}(u, v) = 0 \text{ and } u \neq 0, v \neq 0\}.$$

We define the following minimization problem

$$d := \inf\{ \mathcal{I}(u, v) : (u, v) \in \mathcal{M} \}. \tag{2.1}$$

To study this minimization problem, we firstly characterize the properties of the constrained set  $\mathcal{M}$ .

**Proposition 2.2.** Suppose the conditions of Theorem 1.2 hold. For any  $u, v \in H^1_r(\mathbb{R}^3)\setminus\{0\}$ , there is a unique t := t(u, v) > 0 such that  $(u^t, v^t) \in \mathcal{M}$ , where  $u^t(x) := t^{\frac{1}{4}}u(t^{-\frac{1}{2}}x)$ ,  $v^t(x) := t^{\frac{1}{4}}v(t^{-\frac{1}{2}}x)$ . Particularly, the  $\mathcal{M}$  is not empty.

Before proving Proposition 2.2, we give a lemma.

**Lemma 2.3.** Let  $\alpha, \beta, \gamma, \delta$  be positive constants and  $Q \in (2, 6)$ . For  $t \ge 0$ , we define  $f(t) := \alpha t + \beta t^2 + \gamma t^2 - \delta t^{\frac{Q+6}{4}}$ . Then f has a unique critical point which corresponds to its maximum.

*Proof.* For  $t \ge 0$ , we compute directly that

$$f'(t) = \alpha + 2\beta t + 2\gamma t - \frac{Q+6}{4}\delta t^{\frac{Q+2}{4}},$$
  
$$f''(t) = 2\beta + 2\gamma - \frac{Q+6}{4}\frac{Q+2}{4}\delta t^{\frac{Q-2}{4}}.$$

Since f'' is strictly decreasing and  $f''(0) = 2\beta + 2\gamma > 0$ , there exists  $t_2 > 0$  such that  $f''(t_2) = 0$  and  $f''(t)(t_2 - t) > 0$  for  $t \neq t_2$ .

Since  $f'(0) = \alpha > 0$  and f' is increasing for  $t < t_2$ , f' takes positive values at least for  $t \in [0, t_2]$ . For  $t > t_2$ , f' decreases, tending to  $-\infty$ . Then there exists  $t_0 > t_2$  such that  $f'(t_0) = 0$  and  $f'(t)(t_0 - t) > 0$  for  $t \neq t_0$ .

In conclusion,  $t_0$  is the unique critical point of f and corresponds to its maximum as  $\frac{Q+6}{4} > 2$ .

We are now in a position to prove Proposition 2.2.

**Proof of Proposition 2.2.** For any  $u, v \in H_r^1(\mathbb{R}^3)\setminus\{0\}$  and any t > 0, we define  $u^t(x) := t^{\frac{1}{4}}u(t^{-\frac{1}{2}}x)$  and  $v^t(x) := t^{\frac{1}{4}}v(t^{-\frac{1}{2}}x)$ . Then by direct computation, there hold  $\mathcal{A}(u^t, v^t) = t\mathcal{A}(u, v)$ ,  $C(u^t, v^t) = t^2C(u, v)$ ,  $\mathcal{B}(u^t, v^t) = t^2\mathcal{B}(u, v)$  and  $\int |u^t|^p|v^t|^qdx = t^{\frac{Q+6}{4}}\int |u|^p|v|^qdx$ . Therefore

$$I(u^{t}, v^{t}) = \frac{t}{2}\mathcal{A}(u, v) + \frac{t^{2}}{2}C(u, v) + \frac{t^{2}}{4}\mathcal{B}(u, v) - \frac{1}{Q}t^{\frac{Q+6}{4}}\int |u|^{p}|v|^{q}dx.$$

Denote  $g(t) := I(u^t, v^t)$ . Then g is positive for small t and tends to  $-\infty$  as  $t \to +\infty$  because  $\frac{Q+6}{4} > 2$ . From Lemma 2.3, g(t) has a unique critical point t(u, v) (here and after, t(u, v) means t depends on u and v), corresponding to its maximum. Denoting this t(u, v) by  $t_0$ , then we have that

$$g'(t_0) = \frac{1}{2}\mathcal{A}(u,v) + t_0C(u,v) + \frac{t_0}{2}t_0\mathcal{B}(u,v) - \frac{Q+6}{4Q}t_0^{\frac{Q+2}{4}}\int |u|^p|v|^qdx = 0.$$

Moreover from

$$\mathcal{G}(u^{t_0}, v^{t_0}) = \frac{1}{2} \mathcal{A}(u^{t_0}, v^{t_0}) + C(u^{t_0}, v^{t_0}) + \frac{1}{2} \mathcal{B}(u^{t_0}, v^{t_0}) - \frac{Q+6}{4Q} \int |u^{t_0}|^p |v^{t_0}|^q dx$$
  
=  $t_0 g'(t_0) = 0$ ,

we deduce that  $(u^{t_0}, v^{t_0}) \in \mathcal{M}$ . This proves the proposition and particularly  $\mathcal{M}$  is not empty.  $\square$ 

**Remark 2.4.** From the definition of G(u, v) and Lemma 2.1, we know that if (u, v) is a solution to (1.1), then the unique t(u, v) defined as above satisfies t(u, v) = 1.

**Lemma 2.5.** For any  $u, v \in H_r^1(\mathbb{R}^3) \setminus \{0\}$ , if  $\mathcal{G}(u, v) < 0$ , then  $t(u, v) \in (0, 1)$ , where the t(u, v) is defined as in the proof of Proposition 2.2.

*Proof.* For  $u, v \in H^1_r(\mathbb{R}^3) \setminus \{0\}$ , and let  $t_0 := t(u, v)$  be defined by Proposition 2.2. From

$$\mathcal{G}(u,v) = \frac{1}{2}\mathcal{A}(u,v) + C(u,v) + \frac{1}{2}\mathcal{B}(u,v) - \frac{Q+6}{4Q}\int |u|^p|v|^q dx < 0$$

and

$$\mathcal{G}(u^{t_0}, v^{t_0}) = \frac{t_0}{2} \mathcal{A}(u, v) + t_0^2 C(u, v) + \frac{t_0^2}{2} \mathcal{B}(u, v) - \frac{Q+6}{4Q} t_0^{\frac{Q+6}{4}} \int |u|^p |v|^q dx = 0,$$

we obtain that

$$\frac{1}{2} \left( t_0^{\frac{p+6}{4}} - t_0 \right) \mathcal{A}(u, v) + \left( t_0^{\frac{p+6}{4}} - t_0^2 \right) C(u, v) + \frac{1}{2} \left( t_0^{\frac{p+6}{4}} - t_0^2 \right) \mathcal{B}(u, v) < 0,$$

which implies  $t_0 < 1$ . Therefore  $t_0 := t(u, v) \in (0, 1)$ .

**Lemma 2.6.** Suppose the conditions of Theorem 1.2 hold. Then the  $\mathcal{M}$  is bounded away from zero and  $\mathcal{M}$  is a  $C^1$  manifold.

*Proof.* Firstly, for any  $(u, v) \in \mathcal{M}$ , we deduce from  $\mathcal{G}(u, v) = 0$  and Sobolev inequality that there is  $M_1 > 0$ ,

$$\frac{1}{2}\mathcal{A}(u,v) + C(u,v) \le \frac{1}{2}\mathcal{A}(u,v) + C(u,v) + \frac{1}{2}\mathcal{B}(u,v) 
= \frac{Q+6}{4Q} \int |u|^p |v|^q dx \le M_1 \left(\mathcal{A}(u,v) + C(u,v)\right)^{\frac{Q}{2}}.$$

Hence there is  $M_2 > 0$  such that  $\mathcal{A}(u, v) + C(u, v) \ge M_2$ . This proves that  $\mathcal{M}$  is bounded away from zero

Secondly, we will prove that for any  $(u, v) \in \mathcal{M}$ ,  $\mathcal{G}'(u, v) \neq (0, 0)$ . Arguing by a contradiction, we assume that there is  $(u_0, v_0) \in \mathcal{M}$  such that  $\mathcal{G}'(u_0, v_0) = (0, 0)$ . Then in a weak sense,  $(u_0, v_0)$  satisfies

$$\begin{cases}
-\left(a_{1}+2b_{1}\int|\nabla u_{0}|^{2}dx+2c\int|\nabla v_{0}|^{2}dx\right)\Delta u_{0}+2u_{0} = \frac{p(Q+6)}{4Q}|u_{0}|^{p-2}u_{0}|v_{0}|^{q}, \\
-\left(a_{2}+2b_{2}\int|\nabla v_{0}|^{2}dx+2c\int|\nabla u_{0}|^{2}dx\right)\Delta v_{0}+2v_{0} = \frac{q(Q+6)}{4Q}|u_{0}|^{p}|v_{0}|^{q-2}v_{0},
\end{cases} (2.2)$$

Setting  $h_0 := I(u_0, v_0)$  and  $i := \int (a_1 |\nabla u_0|^2 + a_2 |\nabla v_0|^2) dx$ ,  $j := \int (|u_0|^2 + |v_0|^2) dx$ ,  $k := b_1 (\int |\nabla u_0|^2 dx)^2 + 2c \int |\nabla u_0|^2 dx \int |\nabla v_0|^2 dx + b_2 (\int |\nabla v_0|^2 dx)^2$  and  $e := \int |u_0|^p |v_0|^q dx$ . Then we have from  $h_0 = I(u_0, v_0)$  that

$$\frac{1}{2}i + \frac{1}{2}j + \frac{1}{4}k - \frac{1}{Q}e = h_0.$$
 (2.3)

The  $(u_0, v_0) \in \mathcal{M}$  implies that

$$\frac{1}{2}i + j + \frac{1}{2}k - \frac{Q+6}{4Q}e = 0. {(2.4)}$$

Multiplyling the first equation in the system (2.2) by u and the second equation in (2.2) by v, respectively, and integrating by parts, we obtain that

$$\left(a_{1} + 2b_{1} \int |\nabla u_{0}|^{2} dx + 2c \int |\nabla v_{0}|^{2} dx\right) \int |\nabla u_{0}|^{2} dx + 2 \int |u_{0}|^{2} dx 
= \frac{p(Q+6)}{4Q} \int |u_{0}|^{p} |v_{0}|^{q} dx, 
\left(a_{2} + 2b_{2} \int |\nabla v_{0}|^{2} dx + 2c \int |\nabla u_{0}|^{2} dx\right) \int |\nabla v_{0}|^{2} dx + 2 \int |v_{0}|^{2} dx 
= \frac{q(Q+6)}{4Q} \int |u_{0}|^{p} |v_{0}|^{q},$$

Hence one gets that

$$i + 2j + 2k - \frac{Q+6}{4}e = 0. (2.5)$$

Since  $(u_0, v_0)$  is a weak solution of (2.2), we deduce by a Pohozaev type argument and use the definition of i, j, k and e that

$$\frac{1}{2}i + 3j + k - \frac{3(Q+6)}{4Q}e = 0. {(2.6)}$$

Now solving the equations (2.3), (2.4), (2.5) and (2.6) as the following: multiplying (2.6) by 2 and minus (2.5), one deduces that

$$4j + \left(\frac{Q+6}{4} - \frac{6(Q+6)}{4Q}\right)e = 0. (2.7)$$

Multiplying (2.4) by 4 and minus (2.5), one has that

$$i + 2j = \frac{Q+6}{Q} - \frac{Q+6}{4}e. \tag{2.8}$$

It is now deduced from (2.7) and (2.8) that

$$i = \frac{2 - Q}{8Q}(Q + 6)e < 0, (2.9)$$

which is a contradiction because i > 0, e > 0 and Q > 0. This proves the lemma.

**Lemma 2.7.** Under the conditions of Theorem 1.2, the  $\mathcal{M}$  is a natural constraint in the following sense: if  $(u_0, v_0) \in \mathcal{M}$  is a critical point of  $I|_{\mathcal{M}}$ , then the  $(u_0, v_0) \in \mathcal{M}$  is also a critical point of I on H.

*Proof.* Suppose that  $(u_0, v_0) \in \mathcal{M}$  is a critical point of  $\mathcal{I}|_{\mathcal{M}}$ , then in a weak sense, there is a Lagrange multiplier  $\lambda \in \mathbb{R}$  such that  $\mathcal{I}'(u_0, v_0) = \lambda \mathcal{G}'(u_0, v_0)$ . Therefore in a weak sense, the  $(u_0, v_0)$  satisfies

$$\begin{cases}
-\left(a_{1}+b_{1}\int|\nabla u_{0}|^{2}dx+c\int|\nabla v_{0}|^{2}dx\right)\Delta u_{0}+u_{0}-\frac{p}{Q}|u_{0}|^{p-2}u_{0}|v_{0}|^{q} \\
=\lambda\left(-\left(a_{1}+2b_{1}\int|\nabla u_{0}|^{2}dx+2c\int|\nabla v_{0}|^{2}dx\right)\Delta u_{0}+2u_{0}-\frac{p(Q+6)}{4Q}|u_{0}|^{p-2}u_{0}|v_{0}|^{q}\right), \\
-\left(a_{2}+b_{2}\int|\nabla v_{0}|^{2}dx+c\int|\nabla u_{0}|^{2}dx\right)\Delta v_{0}+v_{0}-\frac{q}{Q}|u_{0}|^{p}|v_{0}|^{q-2}v_{0} \\
=\lambda\left(-\left(a_{2}+2b_{2}\int|\nabla v_{0}|^{2}dx+2c\int|\nabla u_{0}|^{2}dx\right)\Delta v_{0}+2v_{0}-\frac{q(Q+6)}{4Q}|u_{0}|^{p}|v_{0}|^{q-2}v_{0}\right).
\end{cases} (2.10)$$

Which is equivalent to the following system

$$\begin{cases}
-\left((\lambda - 1)a_1 + (2\lambda - 1)b_1 \int |\nabla u_0|^2 dx + (2\lambda - 1)c \int |\nabla v_0|^2 dx\right) \Delta u_0 \\
+ (2\lambda - 1)u_0 &= \left(\frac{p(Q+6)}{4Q}\lambda - \frac{p}{Q}\right) |u_0|^{p-2} u_0 |v_0|^q, \\
-\left((\lambda - 1)a_2 + (2\lambda - 1)b_2 \int |\nabla v_0|^2 dx + (2\lambda - 1)c \int |\nabla u_0|^2 dx\right) \Delta v_0 \\
+ (2\lambda - 1)v_0 &= \left(\frac{q(Q+6)}{4Q}\lambda - \frac{q}{Q}\right) |u_0|^p |v_0|^{q-2} v_0.
\end{cases}$$
(2.11)

Claim:  $\lambda = 0$ .

In order to prove this claim, we denote  $d_0 = \mathcal{I}(u_0, v_0)$ . Then Lemma 2.6 implies that  $d_0 > 0$ . Set

$$i := \int (a_1 |\nabla u_0|^2 + a_2 |\nabla v_0|^2) dx, \qquad j := \int (|u_0|^2 + |v_0|^2) dx,$$
$$k := b_1 \left( \int |\nabla u_0|^2 dx \right)^2 + 2c \int |\nabla u_0|^2 dx \int |\nabla v_0|^2 dx + b_2 \left( \int |\nabla v_0|^2 dx \right)^2$$

and  $e := \int |u_0|^p |v_0|^q dx$ . Firstly, from  $d_0 = \mathcal{I}(u_0, v_0)$  and  $\mathcal{G}(u_0, v_0) = 0$ , we have that

$$\frac{1}{2}i + \frac{1}{2}j + \frac{1}{4}k - \frac{1}{Q}e = d_0 \tag{2.12}$$

and

$$\frac{1}{2}i + j + \frac{1}{2}k - \frac{Q+6}{4Q}e = 0. {(2.13)}$$

Secondly, multiplying the first equation in the system (2.11) by  $u_0$  and integrating by parts, we have that

$$\left( (\lambda - 1)a_1 + (2\lambda - 1)b_1 \int_{\mathbb{R}^3} |\nabla u_0|^2 dx + (2\lambda - 1)c \int_{\mathbb{R}^3} |\nabla v_0|^2 dx \right) \int |\nabla u_0|^2 dx 
+ (2\lambda - 1) \int |u_0|^2 dx = \left( \frac{p(Q+6)}{4Q} \lambda - \frac{p}{Q} \right) \int |u_0|^p |v_0|^q dx;$$

multiplying the second equation in (2.11) by  $v_0$  and integrating by parts, we get that

$$\left( (\lambda - 1)a_2 + (2\lambda - 1)b_2 \int_{\mathbb{R}^3} |\nabla v_0|^2 dx + (2\lambda - 1)c \int_{\mathbb{R}^3} |\nabla u_0|^2 dx \right) \int |\nabla v_0|^2 dx 
+ (2\lambda - 1) \int |v_0|^2 dx = \left( \frac{q(Q+6)}{4Q} \lambda - \frac{q}{Q} \right) \int |u_0|^p |v_0|^q dx.$$

Combining the above two equalities, we deduce that

$$(\lambda - 1)i + (2\lambda - 1)j + (2\lambda - 1)k - \left(\frac{Q+6}{4}\lambda - 1\right)e = 0.$$
 (2.14)

Since  $(u_0, v_0)$  is a weak solution of (2.11), we deduce by a Phozaev type argument that

$$\frac{1}{2}(\lambda - 1)i + \frac{3}{2}(2\lambda - 1)j + \frac{1}{2}(2\lambda - 1)k - \frac{3}{Q}\left(\frac{Q+6}{4}\lambda - 1\right)e = 0.$$
 (2.15)

Now solving the linear system (2.12), (2.13), (2.14) and (2.15). Denoted the coefficient matrix by M

$$M = \begin{pmatrix} \frac{1}{2} & \frac{1}{2} & \frac{1}{4} & -\frac{1}{Q} \\ \frac{1}{2} & 1 & \frac{1}{2} & -\frac{Q+6}{4Q} \\ \lambda - 1 & 2\lambda - 1 & 2\lambda - 1 & 1 - \frac{Q+6}{4}\lambda \\ \frac{1}{2}(\lambda - 1) & \frac{3}{2}(2\lambda - 1) & \frac{1}{2}(2\lambda - 1) & \frac{3}{Q}(1 - \frac{Q+6}{4}\lambda) \end{pmatrix}.$$

By taking elementary transformation to the matrix, we can deduce the determinent

$$\det M = \frac{(Q+2)(2-Q)}{64Q}\lambda(2\lambda-1).$$

If det  $M \neq 0$ , then by Cramer rule, the linear system has a unique solution and

$$e = \frac{d_0}{\det M}\lambda(2\lambda - 1) = \frac{64Qd_0}{(Q+2)(2-Q)},$$

which is a contradiction since Q := p + q > 2,  $d_0 > 0$  and e > 0.

Therefore  $\det M = 0$ , which implies that

$$\lambda = 0 \text{ or } \lambda = \frac{1}{2}.$$

If  $\lambda = \frac{1}{2}$ , then (2.13) becomes

$$-\frac{1}{2}i - \frac{Q-2}{8}e = 0,$$

which is also a contradiction since Q > 2, i > 0 and e > 0. Therefore  $\lambda = 0$  and we prove the claim. Hence  $I'(u_0, v_0) = (0, 0)$ .

## 3. Proof of Theorem 1.2

In this section, we prove Theorem 1.2. Our strategy is to prove that the minimization problem

$$d = \inf\{I(u, v) : (u, v) \in \mathcal{M}\}\$$

defined in (2.1) can be solved. And then using Lemma 2.7 to show that the minimizer is a positive solution to (1.1), which is radially symmetric. Keep the definition of the functional  $\mathcal{G}(u, v)$  in mind:

$$\mathcal{G}(u,v) := \frac{1}{2}\mathcal{A}(u,v) + C(u,v) + \frac{1}{2}\mathcal{B}(u,v) - \frac{Q+6}{4Q}\int |u|^p|v|^q dx,$$

where  $\mathcal{A}(u,v) = \int (a_1|\nabla u|^2 + a_2|\nabla v|^2) dx$ ,  $C(u,v) = \int (|u|^2 + |v|^2) dx$  and  $\mathcal{B}(u,v) = b_1 \left(\int |\nabla u|^2 dx\right)^2 + b_2 \left(\int |\nabla v|^2 dx\right)^2 + 2c \int |\nabla u|^2 dx \int |\nabla v|^2 dx$ . We will prove the following proposition.

**Proposition 3.1.** *Under the conditions of Theorem 1.2, the minimum d defined by (2.1) is achieved.* 

*Proof.* Let  $(u_n, v_n) \in \mathcal{M}$  be such that  $\mathcal{I}(u_n, v_n) \to d$  as  $n \to \infty$ . Then for n large enough, we have that

$$d + o(1) = \frac{1}{2}\mathcal{A}(u_n, v_n) + \frac{1}{2}C(u_n, v_n) + \frac{1}{4}\mathcal{B}(u_n, v_n) - \frac{1}{Q}\int |u_n|^p |v_n|^q dx; \tag{3.1}$$

$$\frac{1}{2}\mathcal{A}(u_n, v_n) + C(u_n, v_n) + \frac{1}{2}\mathcal{B}(u_n, v_n) - \frac{Q+6}{4Q} \int |u_n|^p |v_n|^q dx = 0.$$
 (3.2)

Combining (3.1) and (3.2), we deduce that

$$d + o(1) = \left(\frac{1}{2} - \frac{2}{Q+6}\right) \mathcal{A}(u_n, v_n) + \left(\frac{1}{2} - \frac{4}{Q+6}\right) C(u_n, v_n) + \left(\frac{1}{4} - \frac{2}{Q}\right) \mathcal{B}(u_n, v_n). \tag{3.3}$$

Since Q := p + q > 2, we know that  $\frac{1}{2} > \frac{2}{Q+6}$ ,  $\frac{1}{2} > \frac{4}{Q+6}$  and  $\frac{1}{4} > \frac{2}{Q}$ . Therefore  $\mathcal{A}(u_n, v_n) + C(u_n, v_n)$  is bounded from above.

Going if necessary to a subsequence, still denoted by  $\{(u_n, v_n)\}$ , we may assume that  $(u_n, v_n) \rightarrow (\tilde{u}, \tilde{v})$  weakly in H. Since  $H_r^1(\mathbb{R}^3) \rightarrow L_r^s(\mathbb{R}^3)$  is compact for any  $s \in (2, 6)$ , we have that

$$\int |u_n|^p |v_n|^q dx \to \int |\tilde{u}|^p |\tilde{v}|^q dx \quad \text{as} \quad n \to \infty.$$

Using  $\mathcal{G}(u_n, v_n) = 0$  and  $\mathcal{M}$  is bounded away from zero, we obtain that  $\tilde{u} \neq 0$  and  $\tilde{v} \neq 0$ . In the following, we will prove that  $(u_n, v_n) \to (\tilde{u}, \tilde{v})$  strongly in H and then  $(\tilde{u}, \tilde{v}) \in \mathcal{M}$ .

Denote  $w_n := u_n - \tilde{u}$  and  $z_n := v_n - \tilde{v}$ . Supposing that as  $n \to \infty$ ,  $\mathcal{A}(w_n, z_n) + C(w_n, z_n) \not\to 0$ , then from Brezis-Lieb lemma [3]

$$0 = \mathcal{G}(u_{n}, v_{n}) = \frac{1}{2}\mathcal{A}(w_{n}, z_{n}) + \frac{1}{2}\mathcal{A}(\tilde{u}, \tilde{v}) + C(w_{n}, z_{n}) + C(\tilde{u}, \tilde{v}) - \frac{Q+6}{4Q}\int |u_{n}|^{p}|v_{n}|^{q}dx$$

$$+ \frac{1}{2}\left(b_{1}\left(\int |\nabla u_{n}|^{2}dx\right)^{2} + b_{2}\left(\int |\nabla v_{n}|^{2}dx\right)^{2} + 2c\int |\nabla u_{n}|^{2}dx\int |\nabla v_{n}|^{2}dx\right)$$

$$\geq \frac{1}{2}\mathcal{A}(\tilde{u}, \tilde{v}) + C(\tilde{u}, \tilde{v}) + \frac{1}{2}\mathcal{B}(\tilde{u}, \tilde{v}) - \frac{Q+6}{4Q}\int |\tilde{u}|^{p}|\tilde{v}|^{q}dx$$

$$+ \frac{1}{2}\mathcal{A}(w_{n}, z_{n}) + \frac{1}{2}\mathcal{B}(w_{n}, z_{n}) + C(w_{n}, z_{n})$$

$$= \mathcal{G}(\tilde{u}, \tilde{v}) + \frac{1}{2}\mathcal{A}(w_{n}, z_{n}) + \frac{1}{2}\mathcal{B}(w_{n}, z_{n}) + C(w_{n}, z_{n}),$$

which implies that  $\mathcal{G}(\tilde{u}, \tilde{v}) < 0$ , then according to Proposition 2.2 and Lemma 2.5, we have a unique  $\tilde{t} := t(\tilde{u}, \tilde{v}) \in (0, 1)$  such that  $(\tilde{u}^{\tilde{t}}, \tilde{v}^{\tilde{t}}) \in \mathcal{M}$ , where  $\tilde{u}^{\tilde{t}}(x) := \tilde{t}^{\frac{1}{4}} \tilde{u}(\tilde{t}^{-\frac{1}{2}}x)$  and  $\tilde{v}^{\tilde{t}}(x) := \tilde{t}^{\frac{1}{4}} \tilde{v}(\tilde{t}^{-\frac{1}{2}}x)$ .

As  $\{(u_n, v_n)\}\subset \mathcal{M}$  is a minimizing sequence, we deduce from  $\mathcal{G}(u_n, v_n)=0$  that

$$\begin{split} d+o(1) = &I(u_n,v_n) = \frac{1}{4}\mathcal{A}(u_n,v_n) + \frac{Q-2}{8Q}\int |u_n|^p |v_n|^q dx \\ \geq &\frac{1}{4}\mathcal{A}(\tilde{u},\tilde{v}) + \frac{Q-2}{8Q}\int |\tilde{u}|^p |\tilde{v}|^q dx \\ > &\frac{1}{4}\tilde{t}\mathcal{A}(\tilde{u},\tilde{v}) + \frac{Q-2}{8Q}\tilde{t}^{\frac{p+6}{4}}\int |\tilde{u}|^p |\tilde{v}|^q dx \\ = &\frac{1}{4}\mathcal{A}(\tilde{u}^{\tilde{t}},\tilde{v}^{\tilde{t}}) + \frac{Q-2}{8Q}\int |\tilde{u}^{\tilde{t}}|^p |\tilde{v}^{\tilde{t}}|^q dx \\ = &I(\tilde{u}^{\tilde{t}},\tilde{v}^{\tilde{t}}), \end{split}$$

which is a contradiction as  $(\tilde{u}^{\tilde{t}}, \tilde{v}^{\tilde{t}}) \in \mathcal{M}$ .

Hence  $\mathcal{A}(w_n, z_n) + C(w_n, z_n) \to 0$  as  $n \to \infty$ . Therefore  $(\tilde{u}, \tilde{v}) \in \mathcal{M}$  and  $(\tilde{u}, \tilde{v})$  is a minimizer of  $I|_{\mathcal{M}}$ .

**Proof of Theorem 1.2.** From Proposition 3.1, we have a  $(\tilde{u}, \tilde{v}) \in \mathcal{M}$  such that  $d = I(\tilde{u}, \tilde{v})$ . By Lemma 2.7, the  $(\tilde{u}, \tilde{v})$  is a critical point of I and hence a solution to (1.1). Using a standard argument, we know that  $(\tilde{u}, \tilde{v})$  is a positive radially symmetric ground state solution to (1.1). The proof is complete.

## **4.** Asymptotical behavior of solutions as $(b_1)^2 + (b_2)^2 + c^2 \rightarrow 0$

From previous section, we know that for any  $a_1 > 0$ ,  $a_2 > 0$ ,  $b_1 > 0$ ,  $b_2 > 0$ ,  $c \ge 0$  and p > 1, q > 1 with  $Q := p + q \in (2, 6)$ , (1.1) has a positive radially symmetric ground state solution  $(\tilde{u}, \tilde{v})$ . In this section, for any fixed  $a_1 > 0$  and  $a_2 > 0$ , we will study how this ground state solution depends on  $b_1$ ,  $b_2$  and c. To emphasize the role of  $b_1$ ,  $b_2$  and c, we write the system (1.1) as

$$\begin{cases}
-\left(a_{1}+b_{1}\int_{\mathbb{R}^{3}}|\nabla u|^{2}dx+c\int_{\mathbb{R}^{3}}|\nabla v|^{2}dx\right)\Delta u+u=\frac{p}{Q}|u|^{p-2}u|v|^{q}, \\
-\left(a_{2}+b_{2}\int_{\mathbb{R}^{3}}|\nabla v|^{2}dx+c\int_{\mathbb{R}^{3}}|\nabla u|^{2}dx\right)\Delta v+v=\frac{q}{Q}|u|^{p}|v|^{q-2}v, \\
u:=u(x), \ v:=v(x)\in H^{1}_{r}(\mathbb{R}^{3}), \quad x\in\mathbb{R}^{3}.
\end{cases}$$
(1.1)<sub>b1,b2,c</sub>

The solution obtained in Theorem 1.2 is denoted by  $(u_{b_1,b_2,c},v_{b_1,b_2,c})$ . We write  $\mathcal{M}$  as  $\mathcal{M}_{b_1,b_2,c}$ , the functional  $\mathcal{I}$  as  $\mathcal{I}_{b_1,b_2,c}$ , and the functional  $\mathcal{G}$  as  $\mathcal{G}_{b_1,b_2,c}$ .

**Lemma 4.1.** For any fixed  $a_1 > 0$  and  $a_2 > 0$ , let  $b_1$ ,  $b_2$ ,  $c \in (0,1]$  and p > 1, q > 1 with  $Q := p + q \in (2,6)$ . Denoted by  $(u_{b_1,b_2,c}, v_{b_1,b_2,c})$  the solution obtained in Theorem 1.2. Then  $\{(u_{b_1,b_2,c}, v_{b_1,b_2,c})\}$  is uniformly bounded in H with respect to  $b_1$ ,  $b_2$ ,  $c \in (0,1]$ .

*Proof.* For any  $b_1, b_2, c \in (0,1]$ , choosing nonzero radial functions  $\phi, \psi \in H^1_r(\mathbb{R}^3) \cap C_0^\infty(\mathbb{R}^3)$  and

defining  $\phi^t(x) := t^{\frac{1}{4}}\phi(t^{-\frac{1}{2}}x), \psi^t(x) := t^{\frac{1}{4}}\psi(t^{-\frac{1}{2}}x), t > 0$ , then by direct computation, we obtain that

$$\mathcal{G}_{b_{1},b_{2},c}(\phi^{t},\psi^{t}) = \frac{t}{2}\mathcal{A}(\phi,\psi) + t^{2}C(\phi,\psi) + \frac{t^{2}}{2}\mathcal{B}(\phi,\psi) - \frac{Q+6}{4Q}t^{\frac{Q+6}{4}}\int |\phi|^{p}|\psi|^{q}dx 
\leq \frac{t}{2}\mathcal{A}(\phi,\psi) + t^{2}C(\phi,\psi) - \frac{Q+6}{4Q}t^{\frac{Q+6}{4}}\int |\phi|^{p}|\psi|^{q}dx 
+ \frac{t^{2}}{2}\left(\left(\int |\nabla\phi|^{2}dx\right)^{2} + \left(\int |\nabla\psi|^{2}dx\right)^{2} + 2\int |\nabla\phi|^{2}dx\int |\nabla\psi|^{2}dx\right).$$

Hence by the last inequality and  $\frac{Q+6}{4} > 2$ , we have a  $t_3 > 0$  such that  $\mathcal{G}_{b_1,b_2,c}(\phi^{t_3},\psi^{t_3}) < 0$ , where  $t_3$  is independent of  $b_1$ ,  $b_2$  and c. And then  $(\phi^{t_3},\psi^{t_3})$  is independent of  $b_1$ ,  $b_2$  and c, either. Denote  $w(x) := \phi^{t_3}(x)$  and  $z(x) := \psi^{t_3}(x)$ . By Proposition 2.2 and Lemma 2.5, we get a  $t_0 := t(w,z) \in (0,1)$  such that  $\mathcal{G}_{b_1,b_2,c}(w^{t_0},z^{t_0}) = 0$ , where  $w^{t_0}(x) := t_0^{\frac{1}{4}}w(t_0^{-\frac{1}{2}}x)$  and  $z^{t_0}(x) := t_0^{\frac{1}{4}}z(t_0^{-\frac{1}{2}}x)$ . From this, we deduce that

$$\begin{split} I_{b_1,b_2,c}(w^{t_0},z^{t_0}) &= \frac{1}{4} \mathcal{A}(w^{t_0},z^{t_0}) + \frac{Q-2}{8Q} \int |w^{t_0}|^p |z^{t_0}|^q dx \\ &= \frac{t_0}{4} \mathcal{A}(w,z) + \frac{Q-2}{8Q} t_0^{\frac{Q+6}{4}} \int |w|^p |z|^q dx \\ &< \frac{1}{4} \mathcal{A}(w,z) + \frac{Q-2}{8Q} \int |w|^p |z|^q dx := M_4. \end{split}$$

In here  $M_4$  is a positive constant. Since neither w nor z depends on  $b_1$ ,  $b_2$  and c,  $M_4$  does not depend on any one of  $b_1$ ,  $b_2$  and c, either.

Next let  $\{(u_{b_1,b_2,c},v_{b_1,b_2,c})\}$  be a minimizer of  $\mathcal{I}_{b_1,b_2,c}$  under the constraint of  $\mathcal{M}_{b_1,b_2,c}$ . Then  $\mathcal{I}_{b_1,b_2,c}(u_{b_1,b_2,c},v_{b_1,b_2,c}) \leq \mathcal{I}_{b_1,b_2,c}(w^{t_0},z^{t_0}) < M_4$ . Using  $\mathcal{G}_{b_1,b_2,c}(w^{t_0},z^{t_0}) = 0$ , we get that

$$\begin{split} M_4 > I_{b_1,b_2,c}(u_{b_1,b_2,c},v_{b_1,b_2,c}) &= \frac{Q+2}{2(Q+6)} \mathcal{A}(u_{b_1,b_2,c},v_{b_1,b_2,c}) \\ &+ \frac{Q-2}{2(Q+6)} C(u_{b_1,b_2,c},v_{b_1,b_2,c}) + \frac{Q-2}{4(Q+6)} \mathcal{B}(u_{b_1,b_2,c},v_{b_1,b_2,c}). \end{split}$$

As Q > 2, we deduce that  $\mathcal{A}(u_{b_1,b_2,c}, v_{b_1,b_2,c}) + C(u_{b_1,b_2,c}, v_{b_1,b_2,c})$  is uniformly bounded with respect to  $b_1, b_2, c \in (0, 1]$ . This proves the lemma.

**Proof of Theorem 1.3.** For the sequences  $b_1^{(n)}$ ,  $b_2^{(n)}$  and  $c^{(n)}$ , we may assume that for all  $n=1, 2, \cdots$ ,  $b_1^{(n)} < 1$ ,  $b_2^{(n)} < 1$  and  $c^{(n)} < 1$ . To simplify notations, we denote

$$u^{(n)}(x) := u_{b_1^{(n)}, b_2^{(n)}, c^{(n)}}(x) \text{ and } v^{(n)}(x) := v_{b_1^{(n)}, b_2^{(n)}, c^{(n)}}(x).$$

From Theorem 1.2 and Lemma 4.1, we know that  $\{(u^{(n)}, v^{(n)})\}$  is bounded in H. Going if necessary to a subsequence, we may assume that

$$(u^{(n)}(x), v^{(n)}) \rightarrow (U_0, V_0)$$
 weakly in  $H$ .

Hence  $(U_0, V_0)$  is a weak solution to

$$\begin{cases}
-a_1 \Delta u + u = \frac{p}{Q} |u|^{p-2} u |v|^q, \\
-a_2 \Delta v + v = \frac{q}{Q} |u|^p |v|^{q-2} v, \\
u, v \in H_r^1(\mathbb{R}^3), \quad x \in \mathbb{R}^3.
\end{cases}$$
(1.1)<sub>0,0,0</sub>

Since  $H_r^1(\mathbb{R}^3) \to L_r^s(\mathbb{R}^3)$  is compact for any 2 < s < 6, we have that

$$\int |u^{(n)}|^p |v^{(n)}|^q dx \to \int |U_0|^p |V_0|^q dx \quad \text{as} \quad n \to \infty.$$

Using  $I'_{0,0,0}(U_0, V_0) = 0$  and  $I'_{b_1^{(n)}, b_2^{(n)}, c^{(n)}}(u^{(n)}, v^{(n)}) = 0$ , we get that

$$0 = \langle (I'_{b_{1}^{(n)}, b_{2}^{(n)}, c^{(n)}}(u^{(n)}, v^{(n)}) - I'_{0,0,0}(U_{0}, V_{0})), (u^{(n)} - U_{0}, v^{(n)} - V_{0}) \rangle$$

$$= \int \left( a_{1} |\nabla u^{(n)} - \nabla U_{0}|^{2} + |u^{(n)} - U_{0}|^{2} \right) dx + b_{1}^{(n)} \int |\nabla u^{(n)}|^{2} dx \int \nabla u^{(n)} \nabla (u^{(n)} - U_{0}) dx$$

$$- \frac{p}{Q} \int |u^{(n)}|^{p-2} u^{(n)} (u^{(n)} - U_{0}) |v^{(n)}|^{q} dx + \frac{p}{Q} \int |U_{0}|^{p-2} U_{0} (u^{(n)} - U_{0}) |V_{0}|^{q} dx.$$

$$+ \int \left( a_{2} |\nabla v^{(n)} - \nabla V_{0}|^{2} + |v^{(n)} - V_{0}|^{2} \right) dx + b_{2}^{(n)} \int |\nabla v^{(n)}|^{2} dx \int \nabla v^{(n)} \nabla (v^{(n)} - V_{0}) dx$$

$$- \frac{q}{Q} \int |v^{(n)}|^{q-2} v^{(n)} (v^{(n)} - V_{0}) |u^{(n)}|^{p} dx + \frac{q}{Q} \int |V_{0}|^{q-2} V_{0} (v^{(n)} - V_{0}) |U_{0}|^{p} dx$$

$$+ c^{(n)} \int |\nabla v^{(n)}|^{2} dx \int \nabla u^{(n)} \nabla (u^{(n)} - U_{0}) dx + c^{(n)} \int |\nabla u^{(n)}|^{2} dx \int \nabla v^{(n)} \nabla (v^{(n)} - V_{0}) dx.$$

$$(4.1)$$

Note that as  $n \to \infty$ ,  $(b_1^{(n)})^2 + (b_2^{(n)})^2 + (c^{(n)})^2 \to 0$ ; both  $\int |\nabla u^{(n)}|^2 dx$  and  $\int |\nabla v^{(n)}|^2 dx$  are bounded; we obtain that

$$\begin{split} &\int |u^{(n)}|^{p-2}u^{(n)}(u^{(n)}-U_0)|v^{(n)}|^q dx \\ &\leq \left(\int |u^{(n)}|^{\mathcal{Q}}dx\right)^{\frac{p-1}{\mathcal{Q}}} \left(\int |u^{(n)}-U_0||^{\mathcal{Q}}dx\right)^{\frac{1}{\mathcal{Q}}} \left(\int |v^{(n)}|^{\mathcal{Q}}dx\right)^{\frac{q}{\mathcal{Q}}} \to 0; \\ &\int |U_0|^{p-2}U_0(u^{(n)}-U_0)|V_0|^q dx \\ &\leq \left(\int |U_0|^{\mathcal{Q}}dx\right)^{\frac{p-1}{\mathcal{Q}}} \left(\int |u^{(n)}-U_0||^{\mathcal{Q}}dx\right)^{\frac{1}{\mathcal{Q}}} \left(\int |V_0|^{\mathcal{Q}}dx\right)^{\frac{q}{\mathcal{Q}}} \to 0; \\ &\int |v^{(n)}|^{q-2}v^{(n)}(v^{(n)}-V_0)|u^{(n)}|^p dx \\ &\leq \left(\int |v^{(n)}|^{\mathcal{Q}}dx\right)^{\frac{q-1}{\mathcal{Q}}} \left(\int |v^{(n)}-V_0||^{\mathcal{Q}}dx\right)^{\frac{1}{\mathcal{Q}}} \left(\int |u^{(n)}|^{\mathcal{Q}}dx\right)^{\frac{p}{\mathcal{Q}}} \to 0 \\ &\int |U_0|^p |V_0|^{q-2}V_0(v^{(n)}-V_0)dx \\ &\leq \left(\int |V_0|^{\mathcal{Q}}dx\right)^{\frac{q-1}{\mathcal{Q}}} \left(\int |v^{(n)}-V_0||^{\mathcal{Q}}dx\right)^{\frac{1}{\mathcal{Q}}} \left(\int |U_0|^{\mathcal{Q}}dx\right)^{\frac{p}{\mathcal{Q}}} \to 0. \end{split}$$

and

Combining these with (4.1), we deduce that for n large enough,

$$0 = \mathcal{A}(u^{(n)} - U_0, v^{(n)} - V_0) + C(u^{(n)} - U_0, v^{(n)} - V_0) + o(1).$$

Hence we have proven that  $(u^{(n)}, v^{(n)}) \to (U_0, V_0)$  strongly in H. And  $(U_0, V_0)$  is a positive radially symmetric solution to  $(1.1)_{0,0,0}$ . The proof is complete.

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#### **Conflict of interest**

The authors declare no conflict of interest in this paper.

#### References

- 1. C. O. Alves, F. J. S. A. Correa, T. F. Ma, *Positive solutions for a quasilinear elliptic equation of Kirchhoff type*, Comput. Math. Appl., **49** (2005), 85–93.
- 2. G. Autuori, A. Fiscella, P. Pucci, *Stationary Kirchhoff problems involving a fractional elliptic operator and a critical nonlinearity*, Nonlinear Analysis, **125** (2015), 699–714.
- 3. H. Brezis, E. Lieb, A relation between pointwise convergence of functions and convergence of functionals, P. AM. Math. Soc., **88** (1983), 486–490.
- 4. M. Caponi, P. Pucci, Existence theorems for entire solutions of stationary Kirchhoff fractional p-Laplacian equations, Ann. Mat. Pura Appl., **195** (2016), 2099–2129.
- 5. L. D'Onofrio, A. Fiscella, G. Molica Bisci, *Perturbation methods for nonlocal Kirchhoff type problems*, Fract. Calc. Appl. Anal., **20** (2017), 829–853.
- 6. P. D'Ancona, S. Spagnoto, Global solvability for the degenerate Kirchhoff equation with real analytic data, Invent. Math., **108** (1992), 247–262.
- 7. Y. Deng, S. Peng, W. Shuai, *Existence and asymptotic behavior of nodal solutions for the Kirchhoff-type problems in*  $\mathbb{R}^3$ , J. Funct. Anal., **269** (2015), 3500–3527.
- 8. A. Arosio and S. Panizzi, *On the well-posedness of the Kirchhoff string*, Trans. Amer. Math. Soc., **348** (1996), 305–330.
- 9. G. M. Figueiredo, N. Ikoma, J. R. S. Júnior, *Existence and concentration result for the Kirchhoff type equations with general nonlinearities*, Arch. Rational Mech. Anal., **213** (2014), 931–979.
- 10. Z. Guo, *Ground states for Kirchhoff equations without compact condition*, J. Differ. Equations, **259** (2015), 2884–2902.
- 11. X. He, W. Zou, *Existence and concentration behavior of positive solutions for a Kirchhoff equation in*  $\mathbb{R}^3$ , J. Differ. Equations, **252** (2012), 1813–1834.
- 12. J. H. Jin, X. Wu, *Infinitely many radial solutions for Kirchhoff-type problems in*  $\mathbb{R}^N$ , J. Math. Anal. Appl., **369** (2010), 564–574.
- 13. G. Kirchhoff, Vorlesungen Uber Mechanik, Teubner, Leipzig, 1883.
- 14. G. Li, H. Ye, *Existence of positive ground state solutions for the nonlinear Kirchhoff type equations in*  $\mathbb{R}^3$ , J. Differ. Equations, **257** (2014), 566–600.

- 15. Z. P. Liang, F. Y. Li, J. P. Shi, *Positive solutions to Kirchhoff type equations with nonlinearity having prescribed asymptotic behavior*, Ann. Inst. H. Poincare Anal. Non Lineaire, **31** (2014), 155–167.
- 16. Y. H. Li, F. Y. Li, J. P. Shi, Existence of a positive solution to Kirchhoff type problems without compactness conditions, J. Differ. Equations, 253 (2012), 2285–2294.
- 17. J. L. Lions, On some questions in boundary value problems of mathematical physics, North-Holland Mathematics Studies, **30** (1978), 284–346.
- 18. Z. Liu, S. Guo, *Existence of positive ground state solutions for Kirchhoff type problems*, Nonlinear Analysis, **120** (2015), 1–13.
- 19. D. Lü, Existence and multiplicity results for perturbed Kirchhoff-type Schrodinger systems in  $\mathbb{R}^3$ , Comput. Math. Appl., **68** (2014), 1180–1193.
- 20. T. Matsuyama, M. Ruzhansky, *Global well-posedness of Kirchhoff systems*, J. Math. Pures Appl., **100** (2013), 220–240.
- 21. A. Ourraoui, *On a p– Kirchhoff problem involving a critical nonlinearity*, C. R. Math. Acad. Sci. Paris Ser. I., **352** (2014), 295–298.
- 22. P. Piersanti, P. Pucci, Entire solutions for critical p-fractional Hardy Schrödinger Kirchhoff equations, Publ. Mat., **62** (2018), 3–36.
- 23. D. Ruiz, *The Schrödinger-Poisson equation under the effect of a nonlinear local term*, J. Funct. Anal., **237** (2006), 655–674.
- 24. D. Sun, Z. Zhang, Uniqueness, existence and concentration of positive ground state solutions for Kirchhoff type problems in  $\mathbb{R}^3$ , J. Math. Anal. Appl., **461** (2018), 128–149.
- 25. X. Wu, Existence of nontrivial solutions and high energy solutions for Schrödinger-Kirchhoff-type equations in  $\mathbb{R}^3$ , Nonlinear Anal-Real World Applications, **12** (2011), 1278–1287.
- 26. F. Zhou, K. Wu, X. Wu, *High energy solutions of systems of Kirchhoff-type equations on*  $\mathbb{R}^N$ , Comput. Math. Appl., **66** (2013), 1299–1305.



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