



Research article

### Blow-up for degenerate nonlinear parabolic problem

W. Y. Chan\*

Department of Mathematics, Texas A&M University-Texarkana, Texarkana, TX 75503

\* **Correspondence:** Email: wychan@tamut.edu; Tel: 9033346679.

**Abstract:** In this paper, we deal with the existence, uniqueness, and finite time blow-up of the solution to the degenerate nonlinear parabolic problem:  $u_\tau = (\xi^r u^m u_\xi)_\xi / \xi^r + u^p$  for  $0 < \xi < a$ ,  $0 < \tau < \Gamma$ ,  $u(\xi, 0) = u_0(\xi)$  for  $0 \leq \xi \leq a$ , and  $u(0, \tau) = 0 = u(a, \tau)$  for  $0 < \tau < \Gamma$ , where  $u_0(\xi)$  is a positive function and  $u_0(0) = 0 = u_0(a)$ . In addition, we prove that  $u$  exists globally if  $a$  is small through constructing a global-exist upper solution, and  $u_\tau$  blows up in a finite time.

**Keywords:** blow-up; degenerate nonlinear parabolic problem; global existence

**Mathematics Subject Classification:** 35K55, 35K57, 35K60, 35K65

### 1. Introduction

Let  $\Gamma \in (0, \infty]$ ,  $r$  be a nonnegative constant less than 1,  $a$  and  $m$  be positive constants, and  $p$  be a positive constant greater than 1. We study the following degenerate nonlinear parabolic first initial-boundary value problem:

$$u_\tau = \frac{1}{\xi^r} (\xi^r u^m u_\xi)_\xi + u^p \text{ in } (0, a) \times (0, \Gamma), \tag{1.1}$$

$$u(\xi, 0) = u_0(\xi) \text{ on } [0, a], u(0, \tau) = 0 = u(a, \tau) \text{ for } \tau \in (0, \Gamma), \tag{1.2}$$

where  $u_0(\xi)$  is a positive function in  $(0, a)$  such that  $u_0^{m+1}(\xi) \in C^{2+\alpha}(\bar{D})$  for some  $\alpha \in (0, 1)$  and  $u_0(0) = 0 = u_0(a)$ .

Problems (1.1)–(1.2) describe the creeping gravity flow of a power-law liquid on a rigid horizontal surface. The solution  $u$  is the thickness of the current and  $r$  represents the Cartesian symmetry, see [5]. It also explains the radial spreading of an axisymmetric current with  $\xi$  and  $u^{m+1} / (m + 1)$  corresponding respectively to the radial coordinate and the integral of velocity profile of the current, see [7]. If  $u$  represents the temperature, then it can be interpreted as a nonlinear heat conduction problem with  $u^m$  being the thermal diffusivity, see [12, pp. 73–74]. When  $m = 0$  and  $r = 0.5$ , it exemplifies heat transfer into one face of a flat cylinder with a small ratio of depth to diameter, see [2, 15]. Problems (1.1)–(1.2)

can illustrate population dynamics when  $r = 0$ , see [6]. (1.1) is a degenerate equation because the thermal diffusivity  $u^m \rightarrow 0$  when  $\xi \rightarrow 0$  or  $\xi \rightarrow a$ .

Let  $\xi = ax$ ,  $\tau = a^2(m+1)t$ ,  $\Gamma = a^2(m+1)T$ ,  $D = (0, 1)$ ,  $\Omega = D \times (0, T)$ ,  $\bar{D} = [0, 1]$ ,  $\bar{\Omega} = \bar{D} \times [0, T)$ ,  $\partial D = \{0, 1\}$ , and  $\partial\Omega = (\bar{D} \times \{0\}) \cup (\partial D \times (0, T))$ . Then, the problems (1.1)–(1.2) are transformed into the degenerate nonlinear parabolic problem below,

$$u_t = (m+1) \frac{1}{x^r} (x^r u^m u_x)_x + a^2(m+1)u^p \text{ in } \Omega, \quad (1.3)$$

$$u(x, 0) = u_0(x) \text{ on } \bar{D}, \quad u(0, t) = 0 = u(1, t) \text{ for } t \in (0, T). \quad (1.4)$$

When  $r = 0$  and  $u_0(x) \geq 0$  on  $\bar{D}$ , the multi-dimensional version of the problems (1.3)–(1.4) have been studied by [4, 8, 11, 13, 14]. Let  $\mu_1$  be the first eigenvalue of the following Sturm-Liouville problem,

$$\varphi'' + \mu\varphi = 0 \text{ in } D, \quad \varphi(0) = 0 = \varphi(1).$$

When  $p = m + 1$ , Sacks [13] proved that if  $a^2(m+1) > \mu_1$ , the solution blows up in a finite time. If  $a^2(m+1) \leq \mu_1$  (that is, the domain size is sufficiently small), the problems (1.3)–(1.4) have a global solution (also see [14]). In the case of  $p > m + 1$ , the solution may or may not exist for all time which depends on the initial condition  $u_0$ , see [8, 11, 13]. Galaktionov [4] proved that the problems (1.3)–(1.4) have a global solution if  $p < m + 1$ .

This paper is organized as follows. In section 2, we prove the existence and uniqueness of the classical solution of the problems (1.1)–(1.2). In section 3, we show that  $u$  blows up in a finite time when  $p \geq m + 1$ . Then, we prove that there is a global solution when  $a$  is sufficiently small. Different from [13], our method does not require additional conditions on  $p$  and  $m$ . In section 4, we prove that  $u_t$  blows up in a finite time when  $u$  is unbounded.

## 2. Existence and uniqueness of the solution

We assume that the initial data  $u_0(x)$  satisfies the condition below,

$$\frac{d^2(u_0)^{m+1}}{dx^2} + \frac{r}{x} \frac{d(u_0)^{m+1}}{dx} + a^2(m+1)(u_0)^p \geq 0 \text{ in } D. \quad (2.1)$$

We note that  $u_0 = \left[ Kx \sin\left(\pi(1-x)^2/2\right) \right]^{1/(m+1)}$ , where  $K$  is a positive constant, satisfies (2.1) and  $u_0(x) = 0$  on  $\partial D$ . Let  $v = u^{m+1}$ , the problems (1.3)–(1.4) become

$$v_t = (m+1)v^{m/(m+1)} \left[ v_{xx} + \frac{r}{x}v_x + a^2(m+1)v^{p/(m+1)} \right] \text{ in } \Omega, \quad (2.2)$$

$$v(x, 0) = v_0(x) \text{ on } \bar{D}, \quad v(0, t) = 0 = v(1, t) \text{ for } t \in (0, T), \quad (2.3)$$

where  $v_0(x) = u_0^{m+1}(x)$ . To prove the existence of a solution, Chan and Chan [1] consider the following nonlinear parabolic problem with  $\varepsilon$  being a small positive number less than 1,

$$v_{\varepsilon t} = (m+1)v_{\varepsilon}^{m/(m+1)} \left[ v_{\varepsilon xx} + \frac{r}{x}v_{\varepsilon x} + a^2(m+1)v_{\varepsilon}^{p/(m+1)} \right] \text{ in } \Omega,$$

$$v_{\varepsilon}(x, 0) = v_0(x) + \varepsilon \text{ on } \bar{D}, \quad v_{\varepsilon}(0, t) = \varepsilon = v_{\varepsilon}(1, t) \text{ for } t \in (0, T).$$

They prove that  $v_\varepsilon \in C(\bar{\Omega}) \cap C^{2+\alpha, 1+\alpha/2}(D \times [0, T])$ , and the sequence of solutions:  $\{v_\varepsilon\}$  converges to  $v \in C(\bar{\Omega}) \cap C^{2+\alpha, 1+\alpha/2}(D \times [0, T])$  when  $\varepsilon \rightarrow 0$ . They also show that  $v > 0$  in  $D \times [0, T)$  and  $v(x, t) \geq v_0(x)$  on  $\bar{D} \times [0, T)$ . Using these results, they prove that the problems (1.3)–(1.4) have a solution  $u \in C(\bar{\Omega}) \cap C^{2+\alpha, 1+\alpha/2}(D \times [0, T])$ ,  $u > 0$  in  $D \times [0, T)$ , and  $u(x, t) \geq u_0(x)$  on  $\bar{D} \times [0, T)$ . By (2.1), they show that  $u_t \geq 0$  and  $v_t \geq 0$  in  $D \times [0, T)$ . Further, they prove that  $u$  is unbounded in  $D \times (0, T)$  if  $T < \infty$ . For ease of reference, let us state their Theorem 2.8 below.

**Theorem 2.1.** *Problems (1.3)–(1.4) have a solution  $u \in C(\bar{\Omega}) \cap C^{2+\alpha, 1+\alpha/2}(D \times [0, T])$ . If  $T < \infty$ , then  $u$  is unbounded in  $D \times (0, T)$ .*

Let  $Lv = v^{-m/(m+1)}v_t/(m+1) - v_{xx} - rv_x/x$  and  $\beta(x, t)$  be a bounded function on  $\bar{\Omega}$ . Here is a comparison theorem.

**Lemma 2.2.** *Suppose that  $y$  and  $s \in C(\bar{\Omega}) \cap C^{2,1}(\Omega)$ , and*

$$Ly - \beta y \geq Ls - \beta s \text{ in } \Omega, \quad y \geq s \text{ on } \partial\Omega. \tag{2.4}$$

Then,  $y \geq s$  on  $\bar{\Omega}$ .

*Proof.* If not, let us assume that  $s > y$  somewhere, say,  $(\bar{x}, \bar{t}) \in \Omega$ . By the continuity of  $s$  and  $y$  over  $\bar{\Omega}$ , there exists an interval  $(a_1, a_2) \subset D$  such that  $\bar{x} \in (a_1, a_2)$ ,  $s(a_1, \bar{t}) - y(a_1, \bar{t}) = 0$ ,  $s(a_2, \bar{t}) - y(a_2, \bar{t}) = 0$ ,  $s(x, \bar{t}) > y(x, \bar{t})$  for  $x \in (a_1, a_2)$ , and  $s \leq y$  in  $[a_1, a_2] \times [0, \bar{t})$ . Then,

$$\int_{a_1}^{a_2} (s^{1/(m+1)}(x, \bar{t}) - y^{1/(m+1)}(x, \bar{t})) dx > 0. \tag{2.5}$$

Let  $\tilde{\phi}(x)$  and  $\tilde{\lambda}$  be the first eigenfunction and eigenvalue of the following Sturm-Liouville problem,

$$(x^r w')' + \lambda x^r w = 0 \text{ in } D, \quad w(a_1) = 0 = w(a_2).$$

By Theorem 3.1.2 of Pao [9, p. 97],  $\tilde{\phi}(x)$  exists and  $\tilde{\lambda} > 0$ . Further,  $\tilde{\phi}(x) > 0$  in  $(a_1, a_2)$ . Let  $\gamma$  be a positive real number to be determined. By the above equation, we have

$$\int_0^{\bar{t}} \int_{a_1}^{a_2} (s - y) \tilde{\lambda} x^r \tilde{\phi} e^{\gamma t} dx dt = - \int_0^{\bar{t}} \int_{a_1}^{a_2} (s - y) (x^r \tilde{\phi}')' e^{\gamma t} dx dt. \tag{2.6}$$

Using integration by parts,  $\tilde{\phi}'(a_1) \geq 0$ , and  $\tilde{\phi}'(a_2) \leq 0$ , we have

$$\int_0^{\bar{t}} \int_{a_1}^{a_2} (s - y) (x^r \tilde{\phi}')' e^{\gamma t} dx dt \geq \int_0^{\bar{t}} \int_{a_1}^{a_2} [(s - y)_x x^r]_x \tilde{\phi} e^{\gamma t} dx dt.$$

By (2.4), we get

$$x^r [y^{1/(m+1)} - s^{1/(m+1)}]_t - \beta x^r (y - s) \geq -x^r (s_{xx} - y_{xx}) - r x^{r-1} (s_x - y_x) = -[(s - y)_x x^r]_x.$$

From this, we have

$$\int_0^{\bar{t}} \int_{a_1}^{a_2} (s - y) (x^r \tilde{\phi}')' e^{\gamma t} dx dt \geq - \int_0^{\bar{t}} \int_{a_1}^{a_2} [y^{1/(m+1)} - s^{1/(m+1)}]_t x^r \tilde{\phi} e^{\gamma t} dx dt + \int_0^{\bar{t}} \int_{a_1}^{a_2} \beta (y - s) x^r \tilde{\phi} e^{\gamma t} dx dt.$$

By (2.6), we obtain

$$\begin{aligned}
& - \int_0^{\bar{t}} \int_{a_1}^{a_2} (s-y) \tilde{\lambda} x^r \tilde{\phi} e^{\gamma t} dx dt \\
& \geq - \int_0^{\bar{t}} \int_{a_1}^{a_2} [y^{1/(m+1)} - s^{1/(m+1)}]_t x^r \tilde{\phi} e^{\gamma t} dx dt + \int_0^{\bar{t}} \int_{a_1}^{a_2} \beta (y-s) x^r \tilde{\phi} e^{\gamma t} dx dt \\
& = - \int_{a_1}^{a_2} [y^{1/(m+1)}(x, \bar{t}) - s^{1/(m+1)}(x, \bar{t})] x^r \tilde{\phi} e^{\gamma \bar{t}} dx \\
& + \int_{a_1}^{a_2} [y^{1/(m+1)}(x, 0) - s^{1/(m+1)}(x, 0)] x^r \tilde{\phi} dx \\
& + \int_0^{\bar{t}} \int_{a_1}^{a_2} [y^{1/(m+1)} - s^{1/(m+1)}] \gamma x^r \tilde{\phi} e^{\gamma t} dx dt + \int_0^{\bar{t}} \int_{a_1}^{a_2} \beta (y-s) x^r \tilde{\phi} e^{\gamma t} dx dt.
\end{aligned}$$

The above expression is equivalent to

$$\begin{aligned}
& \int_0^{\bar{t}} \int_{a_1}^{a_2} (\beta - \tilde{\lambda}) (s-y) x^r \tilde{\phi} e^{\gamma t} dx dt + \int_{a_1}^{a_2} [s^{1/(m+1)}(x, 0) - y^{1/(m+1)}(x, 0)] x^r \tilde{\phi} dx \\
& + \int_0^{\bar{t}} \int_{a_1}^{a_2} [s^{1/(m+1)} - y^{1/(m+1)}] \gamma x^r \tilde{\phi} e^{\gamma t} dx dt \\
& \geq \int_{a_1}^{a_2} [s^{1/(m+1)}(x, \bar{t}) - y^{1/(m+1)}(x, \bar{t})] x^r \tilde{\phi} e^{\gamma \bar{t}} dx.
\end{aligned}$$

By the mean value theorem, there exists an  $\zeta$  between  $s^{1/(m+1)}$  and  $y^{1/(m+1)}$  such that

$$\begin{aligned}
& \int_0^{\bar{t}} \int_{a_1}^{a_2} [s^{1/(m+1)} - y^{1/(m+1)}] [(m+1)(\beta - \tilde{\lambda}) \zeta^m + \gamma] x^r \tilde{\phi} e^{\gamma t} dx dt \\
& + \int_{a_1}^{a_2} [s^{1/(m+1)}(x, 0) - y^{1/(m+1)}(x, 0)] x^r \tilde{\phi} dx \\
& \geq \int_{a_1}^{a_2} [s^{1/(m+1)}(x, \bar{t}) - y^{1/(m+1)}(x, \bar{t})] x^r \tilde{\phi} e^{\gamma \bar{t}} dx.
\end{aligned}$$

By the Gronwall inequality (cf. Walter [16, pp. 14–15]),

$$\begin{aligned}
& \int_{a_1}^{a_2} [s^{1/(m+1)}(x, \bar{t}) - y^{1/(m+1)}(x, \bar{t})] x^r \tilde{\phi} e^{\gamma \bar{t}} dx \\
& \leq \int_{a_1}^{a_2} [s^{1/(m+1)}(x, 0) - y^{1/(m+1)}(x, 0)] x^r \tilde{\phi} dx \left[ 1 + \int_0^{\bar{t}} [(m+1)(\beta - \tilde{\lambda}) \zeta^m + \gamma] e^{\int_t^{\bar{t}} [(m+1)(\beta - \tilde{\lambda}) \zeta^m + \gamma] dt} dt \right].
\end{aligned}$$

As  $\beta$  is bounded, we choose  $\gamma$  such that  $\gamma \geq (m+1)(\tilde{\lambda} - \beta) \zeta^m$ . By  $y \geq s$  in  $[a_1, a_2] \times [0, \bar{t}]$ , we have

$$\int_{a_1}^{a_2} [s^{1/(m+1)}(x, \bar{t}) - y^{1/(m+1)}(x, \bar{t})] x^r \tilde{\phi} e^{\gamma \bar{t}} dx \leq 0.$$

Since  $x^r \tilde{\phi} e^{\gamma \bar{t}} > 0$  in  $(a_1, a_2)$ , the above inequality contradicts (2.5). Therefore,  $y \geq s$  in  $\Omega$ . As  $y \geq s$  on  $\partial\Omega$ ,  $y \geq s$  on  $\bar{\Omega}$ . The proof is complete.  $\square$

Let  $\mathcal{L}y = v^{-m/(m+1)}y_t / (m + 1) - y_{xx} - ry_x/x$ . Based on a similar computation of Lemma 2.2, we have the following result.

**Lemma 2.3.** *Suppose that  $y$  and  $s \in C(\bar{\Omega}) \cap C^{2,1}(\Omega)$ , and*

$$\mathcal{L}y - \beta y \geq \mathcal{L}s - \beta s \text{ in } \Omega, \quad y \geq s \text{ on } \partial\Omega.$$

Then,  $y \geq s$  on  $\bar{\Omega}$ .

By Theorem 2.1 and Lemma 2.2, we obtain the result of the existence and uniqueness of solution.

**Theorem 2.4.** *Problems (1.3)–(1.4) and (2.2)–(2.3) have the unique classical solution.*

### 3. Blow-up of the solution and global existence

Instead of using condition (2.1), let us assume that  $u_0$  satisfies the inequality below in the following two sections:

$$\frac{d^2(u_0)^{m+1}}{dx^2} + \frac{r}{x} \frac{d(u_0)^{m+1}}{dx} + a^2(m+1)(u_0)^p > 0 \text{ in } D. \tag{3.1}$$

Then, by (1.3) and  $u \in C(\bar{\Omega}) \cap C^{2+\alpha, 1+\alpha/2}(D \times [0, T])$ , we have  $u_t(x, 0) > 0$  ( $v_t(x, 0) > 0$ ) in  $D$ . We want to prove that  $v_t(x, t) > 0$  in  $D$  for  $t > 0$ . To achieve it, we have the following two results.

**Lemma 3.1.**  $v(x, t) > v_0(x)$  in  $\Omega$ .

*Proof.* From (3.1), we obtain

$$\frac{d^2v_0}{dx^2} + \frac{r}{x} \frac{dv_0}{dx} + a^2(m+1)v_0^{p/(m+1)} > 0 \text{ in } D.$$

As stated in section 2, we have  $v(x, t) \geq v_0(x)$  on  $\bar{D} \times [0, T]$ . Subtract the above inequality from (2.2), it gives

$$\begin{aligned} v^{-m/(m+1)}v_t &> (m+1) \left[ (v - v_0)_{xx} + \frac{r}{x} (v - v_0)_x + a^2(m+1)(v^{p/(m+1)} - v_0^{p/(m+1)}) \right] \\ &\geq (m+1) \left[ (v - v_0)_{xx} + \frac{r}{x} (v - v_0)_x \right]. \end{aligned}$$

Further, we know that  $v(x, t) = v_0(x) = 0$  on  $\partial D \times (0, T)$  and  $v(x, 0) = v_0(x)$  on  $\bar{D}$ . Suppose that  $v(\tilde{x}, t) = v_0(\tilde{x})$  for some  $\tilde{x} \in D$  and  $t > 0$ . Then, the set

$$\{t : v(x, t) = v_0(x) \text{ for some } x \in D \text{ and } t > 0\}$$

is non-empty. Let  $\tilde{t}$  denote its infimum. Suppose that  $\tilde{t} > 0$ . Then,  $v(\tilde{x}, \tilde{t}) = v_0(\tilde{x})$  and  $v(x, t) > v_0$  in  $D \times (0, \tilde{t})$ . Therefore,  $(v(\tilde{x}, \tilde{t}) - v_0(\tilde{x}))_t \leq 0$ . From section 2, we have  $v_t(\tilde{x}, \tilde{t}) \geq 0$ . Thus,  $(v(\tilde{x}, \tilde{t}) - v_0(\tilde{x}))_t = 0$ . Further,  $v(x, t) - v_0(x)$  attains its local minimum at  $(\tilde{x}, \tilde{t})$ . This implies that  $(v(\tilde{x}, \tilde{t}) - v_0(\tilde{x}))_x = 0$  and  $(v(\tilde{x}, \tilde{t}) - v_0(\tilde{x}))_{xx} > 0$ . Since  $v(\tilde{x}, \tilde{t}) > 0$ , we have

$$0 = v^{-m/(m+1)}(\tilde{x}, \tilde{t})v_t(\tilde{x}, \tilde{t}) > (m+1) \left[ (v(\tilde{x}, \tilde{t}) - v_0(\tilde{x}))_{xx} + \frac{r}{\tilde{x}}(v(\tilde{x}, \tilde{t}) - v_0(\tilde{x}))_x \right] > 0.$$

It leads to a contradiction. If  $\tilde{t} = 0$ , we have  $v(x, 0) = v_0(x)$  on  $\bar{D}$  and  $v(x, t) > v_0(x)$  for  $t > 0$  in  $D$ . Hence,  $v(x, t) > v_0(x)$  in  $\Omega$ . □

Let  $h$  be a small positive real number and  $q(x, t) = v(x, t + h)$ . Further,  $q$  is the solution of the following problem:

$$q^{-m/(m+1)} q_t = (m+1) \left[ q_{xx} + \frac{r}{x} q_x + a^2 (m+1) q^{p/(m+1)} \right] \text{ in } \Omega, \quad (3.2)$$

$$q(x, 0) = v(x, h) \text{ on } \bar{D}, q(0, t) = 0 = q(1, t) \text{ for } t \in (0, T). \quad (3.3)$$

We follow a similar calculation of Lemma 3.1 to obtain the corollary below.

**Corollary 3.2.**  $q(x, t) > v(x, t)$  in  $\Omega$ .

Having these two results, we prove  $v_t$  being positive in the domain.

**Lemma 3.3.**  $v_t > 0$  in  $\Omega$ .

*Proof.* From the result of section 2,  $v_t \geq 0$  in  $D \times [0, T)$ . Let us assume that  $v_t(\rho, \omega) = 0$  for some  $(\rho, \omega) \in \Omega$ . Then, there exists a neighborhood  $(a_3, a_4) \times (t_1, t_2) \subset \Omega$  such that  $(\rho, \omega) \in (a_3, a_4) \times (t_1, t_2)$ . We differentiate (2.2) with respect to  $t$  to obtain

$$(v_t)_t = \frac{m}{(m+1)} v^{-1} (v_t)^2 + (m+1) v^{m/(m+1)} \left[ (v_t)_{xx} + \frac{r}{x} (v_t)_x + a^2 p v^{(p-m-1)/(m+1)} v_t \right]. \quad (3.4)$$

Since  $v > 0$  in  $(a_3, a_4) \times (t_1, t_2)$ , it gives

$$(v_t)_t \geq (m+1) v^{m/(m+1)} \left[ (v_t)_{xx} + \frac{r}{x} (v_t)_x + a^2 p v^{(p-m-1)/(m+1)} v_t \right] \text{ in } (a_3, a_4) \times (t_1, t_2).$$

By the strong maximum principle (cf. Protter and Weinberger [10, pp. 168–169]),  $v_t \equiv 0$  in  $(a_3, a_4) \times (t_1, t_2)$ . This contradicts Corollary 3.2 that  $v$  is strictly increasing in  $t$  in  $\Omega$ . Therefore,  $v_t > 0$  in  $(a_3, a_4) \times (t_1, t_2)$ . Since  $(\rho, \omega)$  is arbitrary in  $\Omega$ ,  $v_t > 0$  in  $\Omega$ .  $\square$

To study the blow-up of the solution  $u$ , we let  $z^{1/(1-r)} = x$ . By a direct computation,

$$v_x = v_z \frac{1-r}{z^{r/(1-r)}},$$

$$v_{xx} = (1-r)^2 z^{-2r/(1-r)} v_{zz} - r(1-r) \frac{v_z}{z^{(1+r)/(1-r)}}.$$

Then, the problems (2.2)–(2.3) are transformed into

$$v_t = (m+1) v^{m/(m+1)} \left[ (1-r)^2 z^{-2r/(1-r)} v_{zz} + a^2 (m+1) v^{p/(m+1)} \right] \text{ in } \Omega, \quad (3.5)$$

$$v(z, 0) = v_0(z) \text{ on } \bar{D}, v(0, t) = 0 = v(1, t) \text{ for } t \in (0, T). \quad (3.6)$$

Let

$$F(t) = \frac{(m+1)^2}{p+1} \int_0^1 z^{2r/(1-r)} v^{(p+1)/(m+1)} dz. \quad (3.7)$$

Since  $v > 0$  in  $D \times [0, T)$ ,  $F(t) > 0$  over  $[0, T)$ . We modify Lemma 4.3 of Deng, Duan and Xie [3] to obtain the result below.

**Lemma 3.4.** If  $p \geq m+1$ , then

$$(F'(t))^2 \leq \frac{p+1}{2p} F(t) F''(t).$$

*Proof.* By a direct computation, the derivative of  $F(t)$  is given by

$$F'(t) = (m+1) \int_0^1 z^{2r/(1-r)} v^{(p-m)/(m+1)} v_t dz. \quad (3.8)$$

By  $v_t(x, 0) > 0$  in  $D$  and Lemma 3.3  $v_t > 0$  in  $\Omega$ , we have  $F'(t) > 0$  over  $[0, T)$ . By (3.5), (3.8) is rewritten as

$$F'(t) = (m+1)^2 \int_0^1 \left[ (1-r)^2 v_{zz} + a^2 (m+1) z^{2r/(1-r)} v^{p/(m+1)} \right] v^{p/(m+1)} dz.$$

Differentiating  $F'(t)$  with respect to  $t$  and by (3.5), we have

$$\begin{aligned} F''(t) &= p \int_0^1 v^{(p-2m-1)/(m+1)} z^{2r/(1-r)} (v_t)^2 dz + (m+1)^2 (1-r)^2 \int_0^1 v^{p/(m+1)} v_{zzt} dz \\ &\quad + a^2 (m+1)^2 p \int_0^1 z^{2r/(1-r)} v^{[2p-(m+1)]/(m+1)} v_t dz. \end{aligned}$$

Using integration by parts and  $p \geq m+1$ , we obtain

$$\begin{aligned} F''(t) &= p \int_0^1 v^{(p-2m-1)/(m+1)} z^{2r/(1-r)} (v_t)^2 dz + a^2 (m+1)^2 p \int_0^1 z^{2r/(1-r)} v^{[2p-(m+1)]/(m+1)} v_t dz \\ &\quad + (1-r)^2 p (m+1) \left( \frac{p}{m+1} - 1 \right) \int_0^1 v^{p/(m+1)-2} v_t (v_z)^2 dz + p (m+1) (1-r)^2 \int_0^1 v^{p/(m+1)-1} v_t v_{zz} dz. \end{aligned}$$

By (3.5), the above expression becomes

$$\begin{aligned} F''(t) &= p \int_0^1 v^{(p-2m-1)/(m+1)} z^{2r/(1-r)} (v_t)^2 dz + a^2 (m+1)^2 p \int_0^1 z^{2r/(1-r)} v^{[2p-(m+1)]/(m+1)} v_t dz \\ &\quad + (1-r)^2 p (m+1) \left( \frac{p}{m+1} - 1 \right) \int_0^1 v^{p/(m+1)-2} v_t (v_z)^2 dz \\ &\quad + p \int_0^1 v^{p/(m+1)-1} v_t (m+1) \left[ \frac{v^{-m/(m+1)} z^{2r/(1-r)} v_t}{(m+1)} - a^2 (m+1) z^{2r/(1-r)} v^{p/(m+1)} \right] dz \\ &= 2p \int_0^1 v^{(p-2m-1)/(m+1)} z^{2r/(1-r)} (v_t)^2 dz + (1-r)^2 p (m+1) \left( \frac{p}{m+1} - 1 \right) \int_0^1 v^{p/(m+1)-2} v_t (v_z)^2 dz. \end{aligned}$$

By assumption  $p \geq m+1$ , it yields

$$F''(t) \geq 2p \int_0^1 z^{2r/(1-r)} v^{(p-2m-1)/(m+1)} (v_t)^2 dz. \quad (3.9)$$

By (3.8) and the Cauchy-Schwartz inequality, we obtain

$$\begin{aligned} (F'(t))^2 &= (m+1)^2 \left[ \int_0^1 z^{2r/(1-r)} v^{(p-m)/(m+1)} v_t dz \right]^2 \\ &\leq (m+1)^2 \int_0^1 z^{2r/(1-r)} v^{(p+1)/(m+1)} dz \int_0^1 z^{2r/(1-r)} v^{(p-2m-1)/(m+1)} (v_t)^2 dz. \end{aligned}$$

Then, by (3.7) and (3.9), we have

$$(F'(t))^2 \leq \frac{p+1}{2p} F(t) F''(t). \quad (3.10)$$

This completes the proof.  $\square$

**Lemma 3.5.** *If  $p \geq m + 1$ , then the solution  $u$  blows up somewhere on  $\bar{D}$  in a finite time  $T$ .*

*Proof.* By a direct computation,

$$\begin{aligned} \frac{d^2}{dt^2} F^{-(p-1)/(p+1)}(t) &= -\frac{p-1}{p+1} \left[ \frac{-2p}{p+1} F^{-(3p+1)/(p+1)} (F')^2 + F^{-2p/(p+1)} F'' \right] \\ &= \frac{2p(p-1)}{(p+1)^2} F^{-(3p+1)/(p+1)} \left[ (F')^2 - \frac{p+1}{2p} F F'' \right]. \end{aligned}$$

By (3.10),  $p > 1$ , and  $F > 0$  over  $[0, T)$ , we have

$$\frac{d^2}{dt^2} F^{-(p-1)/(p+1)}(t) \leq 0.$$

We integrate the above inequality over  $(0, t)$  to get

$$\left( F^{-(p-1)/(p+1)}(t) \right)' - \left( F^{-(p-1)/(p+1)}(0) \right)' \leq 0.$$

Equivalently,

$$\left( F^{-(p-1)/(p+1)}(t) \right)' \leq -\frac{p-1}{p+1} F^{-2p/(p+1)}(0) F'(0).$$

Then, we integrate this inequality over  $(0, t)$  to obtain

$$F^{-(p-1)/(p+1)}(t) \leq -\frac{p-1}{p+1} F^{-2p/(p+1)}(0) F'(0) t + F^{-(p-1)/(p+1)}(0).$$

Since  $F(0) > 0$ ,  $F'(0) > 0$ , and  $p > 1$ , the right side of the above inequality is a decreasing function in  $t$  and is equal to zero in a finite time. Therefore, there exists some finite  $T$  such that  $F^{-(p-1)/(p+1)}(T) = 0$ . Hence,  $F(T) = \infty$ . It implies that  $v(z, t) \rightarrow \infty$  when  $t \rightarrow T$  for some  $z \in \bar{D}$ . Thus,  $u(x, t)$  blows up somewhere on  $\bar{D}$  in a finite time  $T$ .  $\square$

Now, we prove that  $u$  exists globally if  $a$  is sufficiently small. This can be achieved through constructing a global-exist upper solution of the problems (2.2)–(2.3). In this proof, we do not have additional conditions on  $p$  and  $m$ .

**Theorem 3.6.** *If  $a$  is small enough, then  $u$  exists globally.*

*Proof.* It suffices to prove that  $v(x, t)$  exists globally. Let  $V(x) = kx^{1-r}(1-x)$  where  $k$  is a positive constant. Then,  $V(x) \in C(\bar{D}) \cap C^2(D)$ . We choose  $k$  such that  $V(x) \geq v_0(x)$ . Clearly,  $V(x) = 0$  at  $x = 0$  and  $x = 1$ . The expression of  $V_x$  and  $V_{xx}$  is below

$$\begin{aligned} V_x &= k \left[ (1-r)x^{-r} - (2-r)x^{1-r} \right], \\ V_{xx} &= k \left[ -r(1-r)x^{-r-1} - (2-r)(1-r)x^{-r} \right]. \end{aligned}$$



By a direct computation,

$$\begin{aligned} & V_{xx} + \frac{r}{x}V_x + a^2(m+1)V^{p/(m+1)} \\ &= k \left[ -r(1-r)x^{-r-1} - (2-r)(1-r)x^{-r} + r(1-r)x^{-r-1} - r(2-r)x^{-r} \right] \\ &+ a^2(m+1)k^{p/(m+1)} \left[ x^{1-r}(1-x) \right]^{p/(m+1)} \\ &= -k(2-r)x^{-r} + a^2(m+1)k^{p/(m+1)} \left[ x^{1-r}(1-x) \right]^{p/(m+1)}. \end{aligned}$$

If  $a$  is sufficiently small, then  $V_{xx} + rV_x/x + a^2(m+1)V^{p/(m+1)} \leq 0 (= V_t)$ . By Lemma 2.2,  $V(x) \geq v(x, t)$  on  $\bar{D} \times [0, \infty)$ . Therefore,  $v$  exists globally which implies  $u$  exists globally.  $\square$

#### 4. Blow-up of $u_t$

In this section, we want to prove that  $u_t$  tends to infinity if  $u$  blows up. From Lemma 3.3,  $v_t > 0$  in  $\Omega$ . Let  $J(x, t) = v_t(x, t) - \varepsilon v(x, t)$  where  $\varepsilon$  is a small positive number. Then,  $J = 0$  on  $\partial D \times [0, T)$ . Let  $t_3 \in (0, T)$ . We choose  $\varepsilon$  such that  $J(x, t_3) \geq 0$  on  $\bar{D}$ .

**Lemma 4.1.** *If  $p \geq m + 1$ , then  $J \geq 0$  on  $\bar{D} \times [t_3, T)$ .*

*Proof.* By a direct computation,  $J_t = v_{tt} - \varepsilon v_t$ ,  $J_x = v_{tx} - \varepsilon v_x$ , and  $J_{xx} = v_{txx} - \varepsilon v_{xx}$ . From (3.4), we have

$$v_{tt} = \frac{m}{m+1}v^{-1}(v_t)^2 + (m+1)v^{m/(m+1)} \left[ J_{xx} + \varepsilon v_{xx} + \frac{r}{x}(J_x + \varepsilon v_x) + a^2 p v^{(p-m-1)/(m+1)} v_t \right].$$

By Lemma 3.3,  $J_t + \varepsilon v_t = v_{tt}$ , and (2.2), we have

$$\begin{aligned} J_t + \varepsilon v_t &> (m+1)v^{m/(m+1)} \left( J_{xx} + \frac{r}{x}J_x \right) + (m+1)v^{m/(m+1)} \varepsilon \left[ \frac{v^{-m/(m+1)}}{m+1}v_t - a^2(m+1)v^{p/(m+1)} \right] \\ &+ a^2(m+1)pv^{(p-1)/(m+1)}(J + \varepsilon v). \end{aligned}$$

Simplifying the above inequality and by  $p \geq m + 1$ , it gives

$$\begin{aligned} J_t &> (m+1)v^{m/(m+1)} \left( J_{xx} + \frac{r}{x}J_x \right) + a^2(m+1)pv^{(p-1)/(m+1)}J + \varepsilon a^2(m+1)[p - (m+1)]v^{(p+m)/(m+1)} \\ &\geq (m+1)v^{m/(m+1)} \left( J_{xx} + \frac{r}{x}J_x \right) + a^2(m+1)pv^{(p-1)/(m+1)}J. \end{aligned}$$

By Lemma 2.3, we have  $J \geq 0$  on  $\bar{D} \times [t_3, T)$ .  $\square$

Our main result below is immediately followed by Lemma 3.5 and Lemma 4.1.

**Theorem 4.2.** *If  $p \geq m + 1$  and  $u$  is unbounded somewhere on  $\bar{D}$  in a finite time  $T$ , then  $u_t$  blows up at  $T$ .*

#### 5. Conclusion

In this paper, we prove the existence and uniqueness of the solution of a degenerate nonlinear parabolic problem. This solution blows up in a finite time if  $p \geq m + 1$ . Then, we show that  $u_t$  blows up somewhere in the domain in a finite time.

## Acknowledgments

The author thanks the anonymous referee for careful reading. This research did not receive any specific grant funding agencies in the public, commercial, or not-for-profit sectors.

## Conflict of interest

The author declares that there are no conflicts of interest in this paper.

## References

1. C. Y. Chan, W. Y. Chan, *Existence of classical solutions of nonlinear degenerate parabolic problems*, Proc. Dynam. Systems Appl., **5** (2008), 85–91.
2. C. Y. Chan, C. S. Chen, *A numerical method for semilinear singular parabolic quenching problems*, Q. Appl. Math., **47** (1989), 45–57.
3. W. Deng, Z. Duan, C. Xie, *The blow-up rate for a degenerate parabolic equation with a non-local source*, J. Math. Anal. Appl., **264** (2001), 577–597.
4. V. A. Galaktionov, *Boundary-value problem for the nonlinear parabolic equation  $u_t = \Delta u^{\sigma+1} + u^\beta$* , Differ. Uravn., **17** (1981), 836–842.
5. J. Gratton, F. Minotti, S. M. Mahajan, *Theory of creeping gravity currents of a non-Newtonian liquid*, Phy. Rev. E., **60** (1999), 6960–6967.
6. M. E. Gurtin, R. C. MacCamy, *On the diffusion of biological populations*, Math. Biosci., **33** (1977), 35–49.
7. H. E. Huppert, *The propagation of two-dimensional and axisymmetric viscous gravity currents over a rigid horizontal surface*, J. Fluid Mech., **121** (1982), 43–58.
8. H. A. Levine, P. E. Sacks, *Some existence and nonexistence theorems for solutions of degenerate parabolic equations*, J. Differ. Equations, **52** (1984), 135–161.
9. C. V. Pao, *Nonlinear Parabolic and Elliptic Equations*, New York: Plenum Press, 1992.
10. M. H. Protter, H. F. Weinberger, *Maximum Principles in Differential Equations*, New York: Springer-Verlag, 1984.
11. M. I. Roux, *Numerical solution of nonlinear reaction diffusion processes*, SIAM J. Numer. Anal., **37** (2000), 1644–1656.
12. P. L. Sachdev, *Self-Similarity and Beyond: Exact Solutions of Nonlinear Problems*, Florida: Chapman and Hall/CRC, 2000.
13. P. E. Sacks, *Global behavior for a class of nonlinear evolution equations*, SIAM J. Math. Anal., **16** (1985), 233–250.
14. A. A. Samarskii, V. A. Galaktionov, S. P. Kurdyumov, et al. *Blow-up in Quasilinear Parabolic Equations*, New York: Walter de Gruyter, 1995.
15. A. D. Solomon, *Melt time and heat flux for a simple PCM body*, Sol. Energy, **22** (1979), 251–257.

---

16. W. Walter, *Differential and Integral Inequalities*, New York: Springer-Verlag, 1970.



AIMS Press

©2019 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (<http://creativecommons.org/licenses/by/4.0>)