



Research article

Solutions to a discrete resonance problem with eigenparameter-dependent boundary conditions

Chenghua Gao*, Enming Yang and Huijuan Li

College of Mathematics and Statistics, Northwest Normal University, Lanzhou 730070, China

* **Correspondence:** Email: gaokuguo@163.com.

Abstract: In this paper, we considered the existence of solutions to a discrete second-order resonance problem with eigenparameter-dependent boundary conditions. We first transformed the resonance problem into its corresponding equivalent system using the Lyapunov-Schmidt method. In addition, using Schauder’s fixed-point theorem and the connectivity theories of the solution set of compact vector fields, we obtained the existence and multiplicity of solutions to the second-order discrete resonance problem with eigenparameter-dependent boundary conditions.

Keywords: discrete resonance problem; eigenparameter-dependent boundary conditions; Lyapunov-Schmidt method; existence; multiplicity

1. Introduction

In 2005, Rodríguez [1] used the Lyapunov-Schmidt method and Brower fixed-point theorem to discuss the following discrete Sturm-Liouville boundary value problem

$$\begin{cases} \Delta[p(t-1)\Delta y(t-1)] + q(t)y(t) + \lambda y(t) = f(y(t)), & t \in [a+1, b+1]_{\mathbb{Z}}, \\ a_{11}y(a) + a_{12}\Delta y(a) = 0, & a_{21}y(b+1) + a_{22}\Delta y(b+1) = 0, \end{cases}$$

where λ is the eigenvalue of the corresponding linear problem and the nonlinearity f is bounded.

Furthermore, in 2007, Ma [2] studied the following discrete boundary value problem

$$\begin{cases} \Delta[p(t-1)\Delta y(t-1)] + q(t)y(t) + \lambda y(t) = f(t, y(t)) + h(t), & t \in [a+1, b+1]_{\mathbb{Z}}, \\ a_{11}y(a) + a_{12}\Delta y(a) = 0, & a_{21}y(b+1) + a_{22}\Delta y(b+1) = 0, \end{cases}$$

where f is subject to the sublinear growth condition

$$|f(t, s)| \leq A|s|^\alpha + B, \quad s \in \mathbb{R}$$

for some $0 \leq \alpha < 1$ and $A, B \in (0, \infty)$. Additional results to the existence of solutions to the related continuous and discrete problems on the nonresonance and the resonance can be found in [3–13] and the references therein. For example, Li and Shu [14] considered the existence of solutions to the continuous Sturm-Liouville problem with random impulses and boundary value problems using the Dhage's fixed-point theorem and considered the existence of upper and lower solutions to a second-order random impulsive differential equation in [15] using the monotonic iterative method.

Inspired by the above literature, we use the solution set connectivity theory of compact vector field [16] to consider the existence of solutions to discrete resonance problems

$$\begin{cases} -\Delta[p(t-1)\Delta y(t-1)] + q(t)y(t) = \lambda_k r(t)y(t) \\ \quad + f(t, y(t)) + \gamma\psi_k(t) + \bar{g}(t), \quad t \in [1, T]_{\mathbb{Z}}, \\ (a_0\lambda_k + b_0)y(0) = (c_0\lambda_k + d_0)\Delta y(0), \\ (a_1\lambda_k + b_1)y(T+1) = (c_1\lambda_k + d_1)\nabla y(T+1), \end{cases} \quad (1.1)$$

where $p : [0, T]_{\mathbb{Z}} \rightarrow (0, \infty)$, $q : [1, T]_{\mathbb{Z}} \rightarrow \mathbb{R}$, $\bar{g} : [1, T]_{\mathbb{Z}} \rightarrow \mathbb{R}$, $r(t) > 0$, $t \in [1, T]_{\mathbb{Z}}$, (λ_k, ψ_k) is the eigenpair of the corresponding linear problem

$$\begin{cases} -\Delta[p(t-1)\Delta y(t-1)] + q(t)y(t) = \lambda r(t)y(t), \quad t \in [1, T]_{\mathbb{Z}}, \\ (a_0\lambda + b_0)y(0) = (c_0\lambda + d_0)\Delta y(0), \\ (a_1\lambda + b_1)y(T+1) = (c_1\lambda + d_1)\nabla y(T+1). \end{cases} \quad (1.2)$$

It is worth noting that the difference between the problem (1.1) and the above questions is the eigenvalue that not only appears in the equation but also in the boundary conditions, which causes us considerable difficulties. Furthermore, it should be noted that these problems also apply to a number of physical problems, including those involving heat conduction, vibrating strings, and so on. For instance, Fulton and Pruess [17] discussed a kind of heat conduction problem, which has the eigenparameter-dependent boundary conditions. However, to discuss this kind of problem, we should know the spectrum of the problem (1.2). Fortunately, in 2016, Gao and Ma [18] obtained the eigenvalue theory of problem (1.2) under the conditions listed as follows:

$$(A_1) \delta_0 := a_0d_0 - b_0c_0 < 0, \quad c_0 \neq 0, \quad d_1 - b_1 \neq 0,$$

$$(A_2) \delta_1 := a_1d_1 - b_1c_1 > 0, \quad c_1 \neq 0, \quad b_0 + d_0 \neq 0,$$

which laid a theoretical foundation for this paper.

Under the conditions (A_1) and (A_2) , we assume the following conditions hold:

(H_1) (Sublinear growth condition) $f : [1, T]_{\mathbb{Z}} \times \mathbb{R} \rightarrow \mathbb{R}$ is continuous and there exist $\alpha \in [0, 1)$ and $A, B \in (0, \infty)$, such that

$$|f(t, y)| \leq A|y|^\alpha + B,$$

(H_2) (Symbol condition) There exists $\omega > 0$, such that

$$yf(t, y) > 0, \quad t \in [1, T]_{\mathbb{Z}} \text{ for } |y| > \omega, \quad (1.3)$$

or

$$yf(t, y) < 0, \quad t \in [1, T]_{\mathbb{Z}} \text{ for } |y| > \omega, \quad (1.4)$$

(H₃) $\bar{g} : [1, T]_{\mathbb{Z}} \rightarrow \mathbb{R}$ satisfies

$$\sum_{s=1}^T \bar{g}(s)\psi_k(s) = 0, \quad (1.5)$$

(H₄) $f : [1, T]_{\mathbb{Z}} \times \mathbb{R} \rightarrow \mathbb{R}$ is continuous and

$$\lim_{|y| \rightarrow \infty} f(t, y) = 0$$

uniformly for $t \in [1, T]_{\mathbb{Z}}$.

The organization of this paper is as follows. In the second section, we construct a completely new inner product space. In the new inner product space, we discuss the basic self-adjointness of the corresponding linear operator and the properties of the eigenpair of (1.2). Finally, under the above properties, the Lyapunov-Schmidt method is used to decompose the inner product space and transform our problem to an equivalent system, that is to say, finding the solutions of (1.1) is equivalent to finding the solutions of this system. Under the sublinear condition and sign conditions on nonlinear terms, an existence result of solutions to the problem (1.1) is obtained using Schauder's fixed-point theorem and the connectivity theories of the solution set of compact vector fields. Based on the first result, the existence of two solutions to the problem (1.1) is also obtained in this section.

2. Preliminaries

Definition 2.1. ([19]) A linear operator P from the linear space X to itself is called the projection operator, if $P^2 = P$.

Lemma 2.2. ([16]) Let C be a bounded closed convex set in Banach space E , $T : [\alpha, \beta] \times C \rightarrow C$ ($\alpha < \beta$) be a continuous compact mapping, then the set

$$S_{\alpha, \beta} = \{(\rho, x) \in [\alpha, \beta] \times C | T(\rho, x) = x\}$$

contains a connected branch connecting $\{\alpha\} \times C$ and $\{\beta\} \times C$.

Lemma 2.3. ([20])(Schauder) Let D be a bounded convex closed set in E , $A : D \rightarrow D$ is completely continuous, then A has a fixed point in D .

First, we construct the inner product space needed in this paper.

Let

$$Y := \{u | u : [1, T]_{\mathbb{Z}} \rightarrow \mathbb{R}\},$$

then Y is a Hilbert space under the following inner product

$$\langle y, z \rangle_Y = \sum_{t=1}^T y(t)z(t)$$

and its norm is $\|y\|_Y := \sqrt{\langle y, y \rangle_Y}$.

Furthermore, consider the space $H := Y \oplus \mathbb{R}^2$. Define the inner product as follows:

$$\langle [y, \alpha, \beta]^{\top}, [z, \zeta, \rho]^{\top} \rangle = \langle y, z \rangle_Y + \frac{p(0)}{|\delta_0|} \alpha \zeta + \frac{p(T)}{|\delta_1|} \beta \rho,$$

which norm is defined as

$$\|y^*\| = \langle [y, \alpha, \beta]^\top, [y, \alpha, \beta]^\top \rangle^{\frac{1}{2}},$$

where \top is transposition to a matrix.

Let

$$y_{0,0} = b_0 y(0) - d_0 \Delta y(0), \quad y_{0,1} = a_0 y(0) - c_0 \Delta y(0)$$

and

$$y_{T+1,0} = b_1 y(T+1) - d_1 \nabla y(T+1), \quad y_{T+1,1} = a_1 y(T+1) - c_1 \nabla y(T+1).$$

For $y^* = [y, \alpha, \beta]^\top$, define an operator $L : D \rightarrow H$ as follows:

$$Ly^* = \begin{bmatrix} -\Delta[p(t-1)\Delta y(t-1)] + q(t)y(t) \\ -y_{0,0} \\ -y_{T+1,0} \end{bmatrix} := \begin{bmatrix} Ly \\ -y_{0,0} \\ -y_{T+1,0} \end{bmatrix},$$

where $D = \{[y, \alpha, \beta]^\top : y \in Y, y_{0,1} = \alpha, y_{T+1,1} = \beta\}$. Define $S : D \rightarrow H$ as follows:

$$Sy^* = S \begin{bmatrix} y \\ \alpha \\ \beta \end{bmatrix} = \begin{bmatrix} ry \\ \alpha \\ \beta \end{bmatrix}.$$

Then, the problem (1.2) is equivalent to the eigenvalue problem as follows:

$$Ly^* = \lambda Sy^*, \quad (2.1)$$

that is, if (λ_k, y) is the eigenpair of the problem (1.2), then (λ_k, y^*) is the eigenpair of the operator L . Conversely, if (λ_k, y^*) is the eigenpair of the operator L , then (λ_k, y) is the eigenpair of the problem (1.2).

Eventually, we define $\mathcal{A} : D \rightarrow H$ as follows:

$$\mathcal{A}y^* = \mathcal{F}(t, y^*) + [\gamma\psi_k + \bar{g}, 0, 0]^\top,$$

where $\mathcal{F}(t, y^*) = \mathcal{F}(t, [y, \alpha, \beta]^\top) = [f(t, y), 0, 0]^\top$. Obviously, the solution of the problem (1.1) is equivalent to the fixed point of the following operator

$$Ly^* = \lambda_k Sy^* + \mathcal{A}y^*. \quad (2.2)$$

It can be seen that there is a homomorphism mapping $(\lambda_k, y) \leftrightarrow (\lambda_k, y^*)$ between the problem (1.1) and the operator Eq (2.2).

Next, we are committed to obtaining the orthogonality of the eigenfunction.

Lemma 2.4. Assume that (λ, y^*) and (μ, z^*) are eigenpairs of L , then

$$\langle y^*, Lz^* \rangle - \langle Ly^*, z^* \rangle = (\mu - \lambda) \langle y^*, Sz^* \rangle.$$

Proof Let $y^* = [y, \alpha, \beta]^\top \in D$, $z^* = [z, \zeta, \rho]^\top \in D$, then

$$\begin{aligned} \langle y^*, Lz^* \rangle &= \langle [y, \alpha, \beta]^\top, [Lz, -z_{0,0}, -z_{T+1,0}]^\top \rangle \\ &= \langle y, Lz \rangle_Y + \frac{p(0)}{|\delta_0|} \alpha(-z_{0,0}) + \frac{p(T)}{|\delta_1|} \beta(-z_{T+1,0}) \\ &= \mu \langle y, rz \rangle_Y + \frac{p(0)}{|\delta_0|} \alpha(\mu\zeta) + \frac{p(T)}{|\delta_1|} \beta(\mu\rho) \\ &= \mu \langle y^*, Sz^* \rangle. \end{aligned} \quad (2.3)$$

Similarly, we have

$$\begin{aligned}
 \langle Ly^*, z^* \rangle &= \langle [Ly, -y_{0,0}, -y_{T+1,0}]^\top, [z, \zeta, \rho]^\top \rangle \\
 &= \langle Ly, z \rangle_Y + \frac{p(0)}{|\delta_0|} (-y_{0,0})\zeta + \frac{p(T)}{|\delta_1|} (-y_{T+1,0})\rho \\
 &= \lambda \langle ry, z \rangle_Y + \frac{p(0)}{|\delta_0|} \lambda \alpha \zeta + \frac{p(T)}{|\delta_1|} \lambda \beta \rho \\
 &= \lambda \langle y^*, Sz^* \rangle.
 \end{aligned} \tag{2.4}$$

It can be seen from (2.3) and (2.4)

$$\langle y^*, Lz^* \rangle - \langle Ly^*, z^* \rangle = (\mu - \lambda) \langle y^*, Sz^* \rangle.$$

□

Lemma 2.5. The operator L is the self-adjoint operator in H .

Proof For $y^* = [y, \alpha, \beta]^\top \in D$, $z^* = [z, \zeta, \rho]^\top \in D$, we just need to prove that $\langle y^*, Lz^* \rangle = \langle Ly^*, z^* \rangle$. By the definition of inner product in H , we obtain

$$\langle y^*, Lz^* \rangle = \langle y, Lz \rangle_Y + \frac{p(0)}{|\delta_0|} \alpha (-z_{0,0}) + \frac{p(T)}{|\delta_1|} \beta (-z_{T+1,0}),$$

and

$$\langle Ly^*, z^* \rangle = \langle Ly, z \rangle_Y + \frac{p(0)}{|\delta_0|} (-y_{0,0})\zeta + \frac{p(T)}{|\delta_1|} (-y_{T+1,0})\rho.$$

Therefore,

$$\begin{aligned}
 \langle y^*, Lz^* \rangle - \langle Ly^*, z^* \rangle &= \langle y, Lz \rangle_Y - \langle Ly, z \rangle_Y + \frac{p(0)}{|\delta_0|} [\alpha (-z_{0,0}) - (-y_{0,0})\zeta] \\
 &\quad + \frac{p(T)}{|\delta_1|} [\beta (-z_{T+1,0}) - (-y_{T+1,0})\rho],
 \end{aligned}$$

where

$$\begin{aligned}
 \langle y, Lz \rangle_Y &= \sum_{t=1}^T y(t) (-\Delta[p(t-1)\Delta z(t-1)] + q(t)z(t)) \\
 &= \sum_{t=1}^T y(t)p(t-1)\Delta z(t-1) - \sum_{t=1}^T y(t)p(t)\Delta z(t) + \sum_{t=1}^T q(t)y(t)z(t) \\
 &= \sum_{t=0}^{T-1} y(t+1)p(t)\Delta z(t) - \sum_{t=1}^T y(t)p(t)\Delta z(t) + \sum_{t=1}^T q(t)y(t)z(t) \\
 &= \sum_{t=0}^{T-1} p(t)\Delta y(t)\Delta z(t) + p(0)y(0)\Delta z(0) - p(T)y(T)\Delta z(T) \\
 &\quad + \sum_{t=1}^T q(t)y(t)z(t)
 \end{aligned}$$

and

$$\begin{aligned} \langle Ly, z \rangle_Y &= \sum_{t=0}^{T-1} p(t) \Delta y(t) \Delta z(t) + p(0) \Delta y(0) z(0) - p(T) \Delta y(T) z(T) \\ &\quad + \sum_{t=1}^T q(t) y(t) z(t). \end{aligned}$$

Moreover, from

$$\begin{aligned} \alpha(-z_{0,0}) - (-y_{0,0})\zeta &= [a_0 y(0) - c_0 \Delta y(0)] [d_0 \Delta z(0) - b_0 z(0)] \\ &\quad - [d_0 \Delta y(0) - b_0 y(0)] [a_0 z(0) - c_0 \Delta z(0)] \\ &= (a_0 d_0 - b_0 c_0) [y(0) \Delta z(0) - \Delta y(0) z(0)] \end{aligned}$$

and

$$\begin{aligned} \beta(-z_{T+1,0}) - (-y_{T+1,0})\rho &= [a_1 y(T+1) - c_1 \nabla y(T+1)] [-b_1 z(T+1) + d_1 \nabla z(T+1)] \\ &\quad - [-b_1 y(T+1) + d_1 \nabla y(T+1)] [a_1 z(T+1) - c_1 \nabla z(T+1)] \\ &= (a_1 d_1 - b_1 c_1) [y(T+1) \nabla z(T+1) - \nabla y(T+1) z(T+1)], \end{aligned}$$

we have

$$\begin{aligned} \langle y^*, Lz^* \rangle - \langle Ly^*, z^* \rangle &= p(0) \begin{vmatrix} y(0) & \Delta y(0) \\ z(0) & \Delta z(0) \end{vmatrix} - p(T) \begin{vmatrix} y(T) & \Delta y(T) \\ z(T) & \Delta z(T) \end{vmatrix} \\ &\quad - p(0) \begin{vmatrix} y(0) & \Delta y(0) \\ z(0) & \Delta z(0) \end{vmatrix} \\ &\quad + p(T) \begin{vmatrix} y(T+1) & \nabla y(T+1) \\ z(T+1) & \nabla z(T+1) \end{vmatrix} \\ &= 0. \end{aligned}$$

□

In order to obtain the orthogonality of the eigenfunction, we define a weighted inner product related to the weighted function $r(t)$ in H . First, we define the inner product in Y as $\langle y, z \rangle_r = \sum_{t=1}^T r(t) y(t) z(t)$.

Similarly, the inner product associated with the weight function $r(t)$ in the space H is defined as follows:

$$\langle [y, \alpha, \beta]^\top, [z, \zeta, \rho]^\top \rangle_r = \langle y, z \rangle_r + \frac{p(0)}{|\delta_0|} \alpha \zeta + \frac{p(T)}{|\delta_1|} \beta \rho.$$

Lemma 2.6. (Orthogonality theorem) Assume that (A_1) and (A_2) hold. If (λ, y^*) and (μ, z^*) are two different eigenpairs corresponding to L , then y^* and z^* are orthogonal under the weight inner product related to the weight function $r(t)$.

Proof Assume that (λ, y^*) and (μ, z^*) is the eigenpair of L , then it can be obtained from Lemmas 2.4 and 2.5

$$0 = (\mu - \lambda) \langle y^*, S z^* \rangle = (\mu - \lambda) \langle y^*, z^* \rangle_r.$$

Therefore, if $\lambda \neq \mu$, then $\langle y^*, z^* \rangle_r = 0$, which implies that y^* and z^* are orthogonal to the inner product defined by the weighted function $r(t)$. □

Lemma 2.7. ([18]) Suppose that (A_1) and (A_2) hold. Then (1.2) has at least T or at most $T + 2$ simple eigenvalues.

In this paper, we consider that λ_k is a simple eigenvalue, that is, the eigenspace corresponding to each eigenvalue is one-dimensional. Let $\psi_k^* = [\psi_k, \alpha, \beta]^\top \in D$ be the eigenfunction corresponding to λ_k , and assume that it satisfies

$$\langle \psi_k^*, \psi_k^* \rangle = 1. \quad (2.5)$$

Denote by $\mathcal{L} := L - \lambda_k S$, then the operator (2.2) is transformed into

$$\mathcal{L}y^* = \mathcal{A}y^*. \quad (2.6)$$

Define $\mathcal{P} : D \rightarrow D$ by

$$(\mathcal{P}x^*)(t) = \psi_k^*(t) \langle \psi_k^*(t), x^*(t) \rangle.$$

Lemma 2.8. \mathcal{P} is a projection operator and $\text{Im}(\mathcal{P}) = \text{Ker}(\mathcal{L})$.

Proof Obviously, \mathcal{P} is a linear operator, next, we need to prove $\mathcal{P}^2 = \mathcal{P}$.

$$\begin{aligned} (\mathcal{P}^2 x^*)(t) &= \mathcal{P}(\mathcal{P}x^*)(t) = \psi_k^*(t) \langle \psi_k^*(t), \mathcal{P}x^*(t) \rangle \\ &= \psi_k^*(t) \langle \psi_k^*(t), \psi_k^*(t) \langle \psi_k^*(t), x^*(t) \rangle \rangle \\ &= \psi_k^*(t) \langle \psi_k^*(t), x^*(t) \rangle \langle \psi_k^*(t), \psi_k^*(t) \rangle \\ &= \psi_k^*(t) \langle \psi_k^*(t), x^*(t) \rangle \\ &= (\mathcal{P}x^*)(t). \end{aligned}$$

It can be obtained from the Definition 2.1, \mathcal{P} is a projection operator. In addition, $\text{Im}(\mathcal{P}) = \text{span}\{\psi_k^*\} = \text{Ker}(\mathcal{L})$. \square

Define $\mathcal{H} : H \rightarrow H$ by

$$\mathcal{H} \left(\begin{bmatrix} y \\ \alpha \\ \beta \end{bmatrix} \right) = \begin{bmatrix} y \\ \alpha \\ \beta \end{bmatrix} - \left\langle \begin{bmatrix} y \\ \alpha \\ \beta \end{bmatrix}, \psi_k^* \right\rangle \psi_k^*.$$

Lemma 2.9. \mathcal{H} is a projection operator and $\text{Im}(\mathcal{H}) = \text{Im}(\mathcal{L})$.

Proof Obviously, \mathcal{H} is a linear operator, next, we need to prove that $\mathcal{H}^2 = \mathcal{H}$.

$$\begin{aligned}
 \mathcal{H}^2 \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} &= \mathcal{H} \left(\mathcal{H} \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} \right) = \mathcal{H} \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} - \left\langle \mathcal{H} \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^* \\
 &= \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} - \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^* \\
 &\quad - \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} - \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^*, \psi_k^* \right\rangle \psi_k^* \\
 &= \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} - 2 \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^* + \left\langle \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^*, \psi_k^* \right\rangle \psi_k^* \\
 &= \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} - 2 \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^* + \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \langle \psi_k^*, \psi_k^* \rangle \psi_k^* \\
 &= \mathcal{H} \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}.
 \end{aligned}$$

It can be obtained from Definition 2.1 that \mathcal{H} is a projection operator. On the one hand, for any $[y, \alpha, \beta]^\top \in H$, we have

$$\begin{aligned}
 \left\langle \mathcal{H} \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle &= \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix} - \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^*, \psi_k^* \right\rangle \\
 &= \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle - \left\langle \left\langle \begin{pmatrix} y \\ \alpha \\ \beta \end{pmatrix}, \psi_k^* \right\rangle \psi_k^*, \psi_k^* \right\rangle \\
 &= 0,
 \end{aligned}$$

thus, $\text{Im}(\mathcal{H}) \subset \text{Im}(\mathcal{L})$. On the other hand, for any $y^* \in \text{Im}(\mathcal{L})$, we have

$$\langle y^*, \psi_k^* \rangle = 0.$$

In summary, $\text{Im}(\mathcal{H}) = \text{Im}(\mathcal{L})$. □

Denote that I is a identical operator, then

$$D = \text{Im}(\mathcal{P}) \oplus \text{Im}(I - \mathcal{P}), \quad H = \text{Im}(\mathcal{H}) \oplus \text{Im}(I - \mathcal{H}).$$

The restriction of the operator \mathcal{L} on $\mathcal{L}|_{\text{Im}(I - \mathcal{P})}$ is a bijection from $\text{Im}(I - \mathcal{P})$ to $\text{Im}(\mathcal{H})$. Define $\mathcal{M} : \text{Im}(\mathcal{H}) \rightarrow \text{Im}(I - \mathcal{P})$ by

$$\mathcal{M} := (\mathcal{L}|_{\text{Im}(I - \mathcal{P})})^{-1}.$$

It can be seen from $\text{Ker } \mathcal{L} = \text{span}\{\psi_k^*\}$ that there is a unique decomposition for any $y^* = [y, \alpha, \beta]^\top \in D$

$$y^* = \rho\psi_k^* + x^*,$$

where $\rho \in \mathbb{R}$, $x^* = [x, \alpha, \beta]^\top \in \text{Im}(I - \mathcal{P})$.

Lemma 2.10. The operator Eq (2.6) is equivalent to the following system

$$x^* = \mathcal{MH}\mathcal{A}(\rho\psi_k^* + x^*), \quad (2.7)$$

$$\sum_{t=1}^T \psi_k(t)f(t, \rho\psi_k(t) + x(t)) = \gamma\left(\frac{p(0)}{|\delta_0|}\alpha^2 + \frac{p(T)}{|\delta_1|}\beta^2 - 1\right) := \theta, \quad (2.8)$$

where $\alpha = a_0\psi_k(0) - c_0\Delta\psi_k(0)$, $\beta = a_1\psi_k(T+1) - c_1\nabla\psi_k(T+1)$.

Proof (i) For any $y^* = \rho\psi_k^* + x^*$, we have

$$\begin{aligned} \mathcal{L}y^* = \mathcal{A}y^* &\iff \mathcal{H}(\mathcal{L}(\rho\psi_k^* + x^*) - \mathcal{A}(\rho\psi_k^* + x^*)) = 0 \\ &\iff \mathcal{L}x^* - \mathcal{H}\mathcal{A}(\rho\psi_k^* + x^*) = 0 \\ &\iff x^* = \mathcal{MH}\mathcal{A}(\rho\psi_k^* + x^*). \end{aligned}$$

(ii) Since $\langle \mathcal{L}y^*, \psi_k^* \rangle = 0$, we have $\langle \mathcal{A}y^*, \psi_k^* \rangle = 0$. Therefore,

$$\begin{aligned} &\langle f(t, y) + \gamma\psi_k + \bar{g}, \psi_k \rangle_Y \\ &= \sum_{t=1}^T f(t, \rho\psi_k(t) + x(t))\psi_k(t) + \sum_{t=1}^T \gamma\psi_k(t)\psi_k(t) + \sum_{t=1}^T \bar{g}(t)\psi_k(t) \\ &= 0. \end{aligned}$$

Combining (H₃) with (2.5), we have

$$\sum_{t=1}^T \psi_k(t)f(t, \rho\psi_k(t) + x(t)) = \gamma\left(\frac{p(0)}{|\delta_0|}\alpha^2 + \frac{p(T)}{|\delta_1|}\beta^2 - 1\right) = \theta,$$

where $\alpha = a_0\psi_k(0) - c_0\Delta\psi_k(0)$, $\beta = a_1\psi_k(T+1) - c_1\nabla\psi_k(T+1)$. □

3. Main results

Let

$$A^+ = \{t \in \{1, 2, \dots, T\} \text{ s.t. } \psi_k(t) > 0\},$$

$$A^- = \{t \in \{1, 2, \dots, T\} \text{ s.t. } \psi_k(t) < 0\}.$$

Obviously,

$$A^+ \cup A^- \neq \emptyset, \min\{|\psi_k(t)| \mid t \in A^+ \cup A^-\} > 0.$$

Lemma 3.1. Supposed that (H₁) holds, then there exist constants M_0 and M_1 , such that

$$\|x^*\| \leq M_1(|\rho| \|\psi_k\|_Y)^\alpha,$$

where (ρ, x^*) is the solution of (2.7) and satisfies $|\rho| \geq M_0$.

Proof Since

$$\mathcal{A}(\rho\psi_k^* + x^*) = \mathcal{F}(t, \rho\psi_k^* + x^*) + [\gamma\psi_k + \bar{g}, 0, 0]^\top = [f(t, \rho\psi_k + x) + \gamma\psi_k + \bar{g}, 0, 0]^\top,$$

we have

$$\begin{aligned} & \|x^*\| \\ & \leq \|\mathcal{M}\|_{\text{Im}(\mathcal{H}) \rightarrow \text{Im}(I-\mathcal{P})} \|\mathcal{H}\|_{H \rightarrow \text{Im}(\mathcal{H})} [\|\bar{g}\|_Y + \gamma\|\psi_k\|_Y + A(|\rho|\|\psi_k\|_Y + \|x\|_Y)^\alpha + B] \\ & = \|\mathcal{M}\|_{\text{Im}(\mathcal{H}) \rightarrow \text{Im}(I-\mathcal{P})} \|\mathcal{H}\|_{H \rightarrow \text{Im}(\mathcal{H})} \left[\|\bar{g}\|_Y + A(|\rho|\|\psi_k\|_Y)^\alpha \left(1 + \frac{\|x\|_Y}{|\rho|\|\psi_k\|_Y} \right)^\alpha + B - \theta \right] \\ & \leq \|\mathcal{M}\|_{\text{Im}(\mathcal{H}) \rightarrow \text{Im}(I-\mathcal{P})} \|\mathcal{H}\|_{H \rightarrow \text{Im}(\mathcal{H})} \left[\|\bar{g}\|_Y + A(|\rho|\|\psi_k\|_Y)^\alpha \left(1 + \frac{\alpha\|x\|_Y}{|\rho|\|\psi_k\|_Y} \right) + B - \theta \right] \\ & = \|\mathcal{M}\|_{\text{Im}(\mathcal{H}) \rightarrow \text{Im}(I-\mathcal{P})} \|\mathcal{H}\|_{H \rightarrow \text{Im}(\mathcal{H})} \left[\|\bar{g}\|_Y \right. \\ & \quad \left. + A(|\rho|\|\psi_k\|_Y)^\alpha \left(1 + \frac{\alpha}{(|\rho|\|\psi_k\|_Y)^{1-\alpha}} \frac{\|x\|_Y}{(|\rho|\|\psi_k\|_Y)^\alpha} \right) + B - \theta \right]. \end{aligned}$$

Denote that

$$\begin{aligned} D_0 &= \|\mathcal{M}\|_{\text{Im}(\mathcal{H}) \rightarrow \text{Im}(I-\mathcal{P})} \|\mathcal{H}\|_{H \rightarrow \text{Im}(\mathcal{H})} (\|\bar{g}\|_Y + B - \theta), \\ D_1 &= A \|\mathcal{M}\|_{\text{Im}(\mathcal{H}) \rightarrow \text{Im}(I-\mathcal{P})} \|\mathcal{H}\|_{H \rightarrow \text{Im}(\mathcal{H})}. \end{aligned}$$

Furthermore, we have

$$\begin{aligned} \frac{\|x^*\|}{(|\rho|\|\psi_k\|_Y)^\alpha} & \leq \frac{D_0}{(|\rho|\|\psi_k\|_Y)^\alpha} + D_1 + \frac{\alpha D_1}{(|\rho|\|\psi_k\|_Y)^{1-\alpha}} \frac{\|x\|_Y}{(|\rho|\|\psi_k\|_Y)^\alpha} \\ & \leq \frac{D_0}{(|\rho|\|\psi_k\|_Y)^\alpha} + D_1 + \frac{\alpha D_1}{(|\rho|\|\psi_k\|_Y)^{1-\alpha}} \frac{\|x^*\|}{(|\rho|\|\psi_k\|_Y)^\alpha}. \end{aligned}$$

So, if we let

$$\frac{\alpha D_1}{(|\rho|\|\psi_k\|_Y)^{1-\alpha}} \leq \frac{1}{2},$$

we have

$$|\rho| \geq \frac{(2\alpha D_1)^{\frac{1}{1-\alpha}}}{\|\psi_k\|_Y} := M_0.$$

Thus,

$$\frac{\|x^*\|}{(|\rho|\|\psi_k\|_Y)^\alpha} \leq \frac{2D_0}{(M_0\|\psi_k\|_Y)^\alpha} + 2D_1 := M_1.$$

This implies that

$$\|x^*\| \leq M_1 (|\rho|\|\psi_k\|_Y)^\alpha.$$

□

Lemma 3.2. Suppose that (H_1) holds, then there exist constants M_0 and Γ , such that

$$\|x^*\| \leq \Gamma (|\rho| \min\{\|\psi_k(t)\| | t \in A^+ \cup A^-\})^\alpha,$$

where (ρ, x^*) is the solution of (2.7) and satisfies $|\rho| \geq M_0$.

According to Lemma 3.2, choose constant ρ_0 , such that

$$\rho_0 > \max\{M_0, \Gamma(|\rho_0| \min\{\|\psi_k(t)\| | t \in A^+ \cup A^-\})^\alpha\}. \quad (3.1)$$

Let

$$K := \{x^* \in \text{Im}(I - \mathcal{P}) \mid x^* = \mathcal{MH}\mathcal{A}(\rho\psi_k^* + x^*), |\rho| \leq \rho_0\}.$$

Then, for sufficiently large $\rho \geq \rho_0$, there is

$$\rho\psi_k(t) + x(t) \geq \omega, \quad \forall t \in A^+, x^* \in K, \quad (3.2)$$

$$\rho\psi_k(t) + x(t) \leq -\omega, \quad \forall t \in A^-, x^* \in K, \quad (3.3)$$

and for sufficiently small $\rho \leq -\rho_0$, there is

$$\rho\psi_k(t) + x(t) \leq -\omega, \quad \forall t \in A^+, x^* \in K, \quad (3.4)$$

$$\rho\psi_k(t) + x(t) \geq \omega, \quad \forall t \in A^-, x^* \in K. \quad (3.5)$$

Theorem 3.3. Suppose that (A_1) , (A_2) and (H_1) – (H_3) hold, then there exists a non-empty bounded set $\Omega_{\bar{g}} \subset \mathbb{R}$, such that the problem (1.1) has a solution if and only if $\theta \in \Omega_{\bar{g}}$. Furthermore, $\Omega_{\bar{g}}$ contains $\theta = 0$ and has a non-empty interior.

Proof We prove only the case of (1.3) in (H_2) , and the case of (1.4) can be similarly proved.

From (1.3) and (3.2)–(3.5), it is not difficult to see that

$$f(t, \rho\psi_k(t) + x(t)) > 0, \quad \forall t \in A^+, x^* \in K,$$

$$f(t, \rho\psi_k(t) + x(t)) < 0, \quad \forall t \in A^-, x^* \in K,$$

for sufficiently large $\rho \geq \rho_0$ and for sufficiently small $\rho \leq -\rho_0$,

$$f(t, \rho\psi_k(t) + x(t)) < 0, \quad \forall t \in A^+, x^* \in K,$$

$$f(t, \rho\psi_k(t) + x(t)) > 0, \quad \forall t \in A^-, x^* \in K.$$

Therefore, if $\rho \geq \rho_0$ is sufficiently large,

$$\psi_k(t)f(t, \rho\psi_k(t) + x(t)) > 0, \quad \forall t \in A^+ \cup A^-, x^* \in K, \quad (3.6)$$

if $\rho \leq -\rho_0$ is sufficiently small,

$$\psi_k(t)f(t, \rho\psi_k(t) + x(t)) < 0, \quad \forall t \in A^+ \cup A^-, x^* \in K. \quad (3.7)$$

Let

$$C := \{x^* \in \text{Im}(I - \mathcal{P}) \mid \|x^*\| \leq \rho_0\}.$$

Define $\mathcal{T}_\rho : \text{Im}(I - \mathcal{P}) \rightarrow \text{Im}(I - \mathcal{P})$ by

$$\mathcal{T}_\rho := \mathcal{MH}\mathcal{A}(\rho\psi_k^* + x^*).$$

Obviously, \mathcal{T}_ρ is completely continuous. By (3.1), for $x^* \in C$ and $\rho \in [-\rho_0, \rho_0]$,

$$\begin{aligned} \|\mathcal{T}_\rho x^*\| &\leq \Gamma(|\rho| \min\{\|\psi_k(t)\| \mid t \in A^+ \cup A^-\})^\alpha \\ &\leq \Gamma(|\rho_0| \min\{\|\psi_k(t)\| \mid t \in A^+ \cup A^-\})^\alpha \\ &\leq \rho_0, \end{aligned}$$

i.e.,

$$\mathcal{T}_\rho(C) \subseteq C.$$

According to Schauder's fixed point theorem, \mathcal{T}_ρ has a fixed point on C , such that $\mathcal{T}_\rho x^* = x^*$. It can be seen from Lemma 2.10 that the problem (1.1) is equivalent to the following system

$$\Psi(s, x^*) = \theta, \quad (s, x^*) \in S_{\bar{g}},$$

where

$$S_{\bar{g}} := \{(\rho, x^*) \in \mathbb{R} \times \text{Im}(I - \mathcal{P}) \mid x^* = \mathcal{MHA}(\rho\psi_k^* + x^*)\},$$

$$\Psi(\rho, x^*) := \sum_{s=1}^T \psi_k(s) f(s, \rho\psi_k(s) + x(s)).$$

At this time, the $\Omega_{\bar{g}}$ in Theorem 3.3 can be given by $\Omega_{\bar{g}} = \Psi(S_{\bar{g}})$. There exists a solution to the problem (1.1) for $\theta \in \Omega_{\bar{g}}$.

From (3.6), (3.7) and $A^+ \cup A^- \neq \emptyset$, we can deduce that for any $x^* \in K$

$$\sum_{s=1}^T \psi_k(s) f(s, -\rho_0\psi_k(s) + x(s)) < 0, \quad \sum_{s=1}^T \psi_k(s) f(s, \rho_0\psi_k(s) + x(s)) > 0.$$

Thus,

$$\Psi(-\rho_0, x^*) < 0 < \Psi(\rho_0, x^*), \quad \forall x^* \in K. \quad (3.8)$$

According to Lemma 2.2, $S_{\bar{g}} \subset \mathbb{R} \times \bar{B}_{\rho_0}$ contains a connected branch $\xi_{-\rho_0, \rho_0}$ connecting $\{-\rho_0\} \times C$ and $\{\rho_0\} \times C$. Combined with (3.8), $\Omega_{\bar{g}}$ contains $\theta = 0$ and has a non-empty interior. \square

Theorem 3.4. Suppose that (A_1) , (A_2) , (H_2) – (H_4) hold. $\Omega_{\bar{g}}$ as shown in Theorem 3.3, then there exists a nonempty set $\Omega_{\bar{g}}^* \subset \Omega_{\bar{g}} \setminus \{0\}$, such that problem (1.1) has at least two solutions for $\theta \in \Omega_{\bar{g}}^*$.

Proof We prove only the case of (1.3), and the case of (1.4) can be similarly proved. Since the condition (H_4) implies that (H_1) , using Theorem 3.3, we know that there exists $\rho_0 > 0$, such that

$$\Psi(\rho_0, x^*) > 0, \quad \forall x^* \in K.$$

Let

$$\delta := \min\{\Psi(\rho_0, x^*) \mid x^* \in K\},$$

then $\delta > 0$.

Next, we prove that problem (1.1) has at least two solutions for any $\theta \in (0, \delta)$.

Let

$$S_{\bar{g}} := \{(\rho, x^*) \in \mathbb{R} \times \text{Im}(I - \mathcal{P}) \mid x^* = \mathcal{MHA}(\rho\psi_k^* + x^*)\},$$

$$\bar{K} := \{x^* \in \text{Im}(I - \mathcal{P}) \mid (\rho, x^*) \in S_{\bar{g}}\}.$$

By (H_4) , there exists a constant A_0 such that

$$\|x^*\| \leq A_0, \quad \forall x^* \in K.$$

Similar to the derivation of Theorem 3.3, there exists $\rho^* > \rho_0$ such that the following results hold:

(i) For $\rho \geq \rho^*$, there is

$$\psi_k(t)f(t, \rho\psi_k(t) + x(t)) > 0, \quad \forall t \in A^+ \cup A^-, \quad x^* \in \bar{K}, \quad (3.9)$$

(ii) For $\rho \leq -\rho^*$, there is

$$\psi_k(t)f(t, \rho\psi_k(t) + x(t)) < 0, \quad \forall t \in A^+ \cup A^-, \quad x^* \in \bar{K}. \quad (3.10)$$

Let

$$C^* := \{x^* \in \text{Im}(I - \mathcal{P}) \mid \|x^*\| \leq A_0\}.$$

According to (H₄), (3.9) and (3.10), we have

$$\lim_{|\rho| \rightarrow \infty} \sum_{s=1}^T \psi_k(s)f(s, \rho\psi_k(s) + x(s)) = 0$$

uniformly for $x^* \in \bar{K}$, i.e.

$$\lim_{|\rho| \rightarrow \infty} \Psi(\rho, x^*) = 0, \quad x^* \in \bar{K}.$$

Therefore, there exists a constant $l : l > \rho^* > \rho_0 > 0$ such that $S_{\bar{g}}$ contains a connected branch between $\{-l\} \times C^*$ and $\{l\} \times C^*$, and

$$\begin{aligned} & \max\{|\Psi(\rho, x^*)| \mid \rho = \pm l, (\rho, x^*) \in \xi_{-l,l}\} \\ & \leq \max\{|\Psi(\rho, x^*)| \mid (\rho, x^*) \in \{-l, l\} \times \bar{K}\} \\ & \leq \frac{\theta}{3}. \end{aligned}$$

It can be seen from the connectivity of $\xi_{-l,l}$ that there exist (ρ_1, x_1^*) and (ρ_2, x_2^*) in $\xi_{-l,l} \subset S_{\bar{g}}$, such that

$$\Psi(\rho_1, x_1^*) = \theta, \quad \Psi(\rho_2, x_2^*) = \theta,$$

where $\rho_1 \in (-l, \rho_0)$, $\rho_2 \in (\rho_0, l)$. It can be proved that $\rho_1\psi_k^* + x_1^*$ and $\rho_2\psi_k^* + x_2^*$ are two different solutions of problem (1.1). \square

4. Example

In this section, we give a concrete example of the application of our major results of Theorems 3.3 and 3.4. We choose $T = 3$, $a_0, d_0, b_1, c_1 = 0$ and $a_1, d_1, b_0, c_0 = 1$, which implies that the interval becomes $[1, 3]_{\mathbb{Z}}$ and the conditions (A₁), (A₂) hold.

First, we consider the eigenpairs of the corresponding linear problem

$$\begin{cases} -\Delta^2 y(t-1) = \lambda y(t), & t \in [1, 3]_{\mathbb{Z}}, \\ y(0) = \lambda \Delta y(0), & \lambda y(4) = \nabla y(4). \end{cases} \quad (4.1)$$

Define the equivalent matrix of (4.1) as follows,

$$A_\lambda = \begin{pmatrix} \lambda - 2 + \frac{\lambda}{1+\lambda} & 1 & 0 \\ 1 & \lambda - 2 & 1 \\ 0 & 1 & \lambda - 2 + \frac{1}{1-\lambda} \end{pmatrix}$$

Consequently, $A_\lambda y = 0$ is equivalent to (4.1). Let $|A_\lambda| = 0$, we have

$$\lambda_1 = -1.4657, \lambda_2 = 0.1149, \lambda_3 = 0.8274, \lambda_4 = 2.0911, \lambda_5 = 3.4324,$$

which are the eigenvalues of (4.1). Next, we choose $\lambda = \lambda_1 = -1.4657$, then we obtain the corresponding eigenfunction

$$\psi_1(t) = \begin{cases} 1, & t = 1, \\ 3.4657, & t = 2, \\ 3.4657^2 - 1, & t = 3. \end{cases}$$

Example 4.1. Consider the following problem

$$\begin{cases} -\Delta^2 y(t-1) = -1.4657y(t) + f(t, y(t)) + \psi_1(t) + \bar{g}(t), & t \in [1, 3]_{\mathbb{Z}}, \\ y(0) = -1.4657\Delta y(0), \quad -1.4657y(4) = \nabla y(4), \end{cases} \quad (4.2)$$

where

$$f(t, s) = \begin{cases} ts^3, & s \in [-1, 1], \\ t\sqrt[3]{s}, & s \in (-\infty, -1) \cup (1, +\infty), \end{cases}$$

and

$$\bar{g}(t) = \begin{cases} 0, & t = 1, \\ 3.4657^2 - 1, & t = 2, \\ -3.4657, & t = 3. \end{cases}$$

Then, for $f(t, y(t))$, we have $|f(t, y(t))| \leq 3|y(t)|^{\frac{1}{3}}$. If we choose $\omega = 1$, $yf(t, y) > 0$ for $|y(t)| > 1$. For $\bar{g}(t)$, we have $\sum_{s=1}^3 \bar{g}(s)\psi_1(s) = 0$.

Therefore, the problem (4.2) satisfies the conditions (A_1) , (A_2) , (H_1) – (H_3) , which implies that the problem (4.2) has at least one solution by Theorem 3.3.

Example 4.2. Consider the following problem

$$\begin{cases} -\Delta^2 y(t-1) = -1.4657y(t) + f(t, y(t)) + \psi_1(t) + \bar{g}(t), & t \in [1, 3]_{\mathbb{Z}}, \\ y(0) = -1.4657\Delta y(0), \quad -1.4657y(4) = \nabla y(4), \end{cases} \quad (4.3)$$

where

$$f(t, s) = \frac{ts}{e^{|s|}}, \quad t \in [1, 3]_{\mathbb{Z}}$$

and

$$\bar{g}(t) = \begin{cases} 0, & t = 1, \\ 1 - 3.4657^2, & t = 2, \\ 3.4657, & t = 3. \end{cases}$$

Then, for $f(t, y(t))$, we always have $yf(t, y) > 0$ for all $y(t) > 0$ or $y(t) < 0$, f is continuous and satisfies

$$\lim_{|y| \rightarrow \infty} f(t, y) = 0.$$

For $\bar{g}(t)$, we have $\sum_{s=1}^3 \bar{g}(s)\psi_1(s) = 0$.

Therefore, the problem (4.3) satisfies the conditions (A_1) , (A_2) , (H_2) – (H_4) , which implies that the problem (4.3) has at least two solutions by Theorem 3.4.

Use of AI tools declaration

The authors declare that they have not used Artificial Intelligence (AI) tools in the creation of this article.

Acknowledgments

Supported by National Natural Science Foundation of China [Grant No. 11961060] and Natural Science Foundation of Qinghai Province(No.2024-ZJ-931).

Conflict of interest

The authors declare that there are no conflicts of interest.

References

1. J. Rodríguez, Nonlinear discrete Sturm-Liouville problems, *J. Math. Anal. Appl.*, **308** (2005), 380–391. <https://doi.org/10.1016/j.jmaa.2005.01.032>
2. R. Y. Ma, Nonlinear discrete Sturm-Liouville problems at resonance, *Nonlinear Anal.*, **67** (2007), 3050–3057. <https://doi.org/10.1016/j.na.2006.09.058>
3. T. S. He, Y. T. Xu, Positive solutions for nonlinear discrete second-order boundary value problems with parameter dependence, *J. Math. Anal. Appl.*, **379** (2011), 627–636. <https://doi.org/10.1016/j.jmaa.2011.01.047>
4. S. Smirnov, Green's function and existence of solutions for a third-order boundary value problem involving integral condition, *Lith. Math. J.*, **62** (2022), 509–518. <https://doi.org/10.1007/s10986-022-09576-7>
5. G. Zhang, S. Ge, Existence of positive solutions for a class of discrete Dirichlet boundary value problems, *Appl. Math. Lett.*, **48** (2015), 1–7. <https://doi.org/10.1016/j.aml.2015.03.005>
6. C. W. Ha, C. C. Kuo, On the solvability of a two-point boundary value problem at resonance II, *Topol. Methods Nonlinear Anal.*, **11** (1998), 159–168. <https://doi.org/10.12775/tmna.1998.010>
7. D. Maroncelli, J. Rodríguez, On the solvability of nonlinear discrete Sturm-Liouville problems at resonance, *Int. J. Differ. Equations*, **12** (2017), 119–129.

8. C. H. Gao, Y. L. Wang, L. Lv, Spectral properties of discrete Sturm-Liouville problems with two squared eigenparameter-dependent boundary conditions, *Acta Math. Sci.*, **40** (2020), 755–781. <https://doi.org/10.1007/s10473-020-0312-5>
9. B. Freedman, J. Rodríguez, On nonlinear discrete Sturm-Liouville boundary value problems with unbounded nonlinearities, *J. Differ. Equations Appl.*, **28** (2022), 1–9. <https://doi.org/10.1080/10236198.2021.2017426>
10. J. Henderson, R. Luca, Existence and multiplicity for positive solutions of a second-order multi-point discrete boundary value problem, *J. Differ. Equations Appl.*, **19** (2013), 418–438. <https://doi.org/10.1080/10236198.2011.648187>
11. Z. J. Du, F. C. Meng, Solutions to a second-order multi-point boundary value problem at resonance. *Acta Math. Sci.*, **30** (2010), 1567–1576. [https://doi.org/10.1016/S0252-9602\(10\)60150-6](https://doi.org/10.1016/S0252-9602(10)60150-6)
12. N. Kosmatov, A singular non-local problem at resonance, *J. Math. Anal. Appl.*, **394** (2012), 425–431. <https://doi.org/10.1016/j.jmaa.2012.04.069>
13. D. Maroncelli, J. Rodríguez, Existence theory for nonlinear Sturm-Liouville problems with non-local boundary conditions, *Differ. Equations Appl.*, **10** (2018), 147–161.
14. Z. H. Li, X. B. Shu, T. Y. Miao, The existence of solutions for Sturm-Liouville differential equation with random impulses and boundary value problems, *Bound. Value Probl.*, **97** (2021). <https://doi.org/10.1186/s13661-021-01574-x>
15. Z. H. Li, X. B. Shu, F. Xu, The existence of upper and lower solutions to second order random impulsive differential equation with boundary value problem, *AIMS Math.*, **5** (2020), 6189–6210. <https://doi.org/10.3934/math.2020398>
16. R. Y. Ma, Y. R. Yang, Existence result for a singular nonlinear boundary value problem at resonance, *Nonlinear Anal.*, **68** (2008), 671–680. <https://doi.org/10.1016/j.na.2006.11.030>
17. C. T. Fulton, S. Pruess, Numerical methods for a singular eigenvalue problem with eigenparameter in the boundary conditions, *J. Math. Anal. Appl.*, **71** (1979), 431–462. [https://doi.org/10.1016/0022-247X\(79\)90203-8](https://doi.org/10.1016/0022-247X(79)90203-8)
18. C. H. Gao, R. Y. Ma, Eigenvalues of discrete Sturm-Liouville problems with eigenparameter dependent boundary conditions, *Linear Algebra Appl.*, **503** (2016), 100–119. <https://doi.org/10.1016/j.laa.2016.03.043>
19. N. Dunford, J. T. Schwartz, *Linear Operators, Part I: General Theory*, John Wiley & Sons, New York, 1988.
20. J. T. Schwartz, *Nonlinear Functional Analysis*, CRC Press, New York, 1969.



AIMS Press

©2024 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (<https://creativecommons.org/licenses/by/4.0>)