
Research article

Convergence for global curve diffusion flows[†]

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Abstract: In this note we establish exponentially fast smooth convergence for global curve diffusion flows, and discuss open problems relating embeddedness to global existence (Giga's conjecture) and the shape of Type I singularities (Chou's conjecture).

Keywords: curve diffusion; surface diffusion; global existence; convergence; curvature; partial differential equation; geometric analysis

1. Introduction

The curve diffusion flow is the one-parameter family of immersed curves $\gamma : \mathbb{S}^1 \times [0, T) \rightarrow \mathbb{R}^2$ with normal velocity equal to $-\text{grad}_{H^{-1}}(L(\gamma))$, that is

$$\partial_t^\perp \gamma = -k_{ss} \nu. \quad (\text{CD})$$

Here, s denotes arclength, k is the scalar curvature, and ν is the normal vector.

In [10], the author proves that if the isoperimetric ratio and L^2 -normalised oscillation of curvature are close to their value on any circle (with an explicit estimate on the constant) then the flow exists globally and converges to a circle. However the rate of convergence was not established.

In this note we prove a general convergence principle that can be summarised as:

$$T = \infty \implies \gamma \text{ converges exponentially fast to an } \omega\text{-circle} \\ \text{in the smooth topology.}$$

An ω -circle is an immersion $\gamma : \mathbb{S}^1 \rightarrow \mathbb{R}^2$ with winding number of the unit normal around the origin equal to ω . Note that also $\omega = \frac{1}{2\pi} \int k ds$.

Theorem 1.1. Suppose $\gamma : \mathbb{S} \times [0, \infty) \rightarrow \mathbb{R}^2$ is a global curve diffusion flow with smooth initial data γ_0 that has winding number ω . Then γ converges exponentially fast to a round ω -circle in the smooth topology, with explicit estimates (here $m \in \mathbb{N}$)

$$\|k_{s^m}\|_2^2 \leq c_m e^{-\bar{k}_0^4 t},$$

with

$$\mathcal{D}(t) \leq \mathcal{D}_0 e^{-4\bar{k}_0^4 t} \quad \text{and} \quad K_{osc}(t) \leq c_0 e^{-2\bar{k}_0^4 t}.$$

The constants c_m depend only on the initial data γ_0 .

The quantities \mathcal{D} and K_{osc} are the isoperimetric deficit and the L^2 -normalised oscillation of curvature:

$$\mathcal{D} = L^2 - 4\omega\pi A \quad \text{and} \quad K_{osc} = L \int_{\gamma} (k - \bar{k})^2 ds$$

where L and A are the length and signed enclosed area of γ , and k, \bar{k} are the curvature scalar and the average of k respectively. Section 2 contains the proof of Theorem 1.1.

Remark 1.2. Note that no smallness condition is required.

Remark 1.3. The regularity hypothesis that we impose (initial data of class C^∞) is not optimal. There are in the literature well-posedness (and regularisation) results for curve and surface diffusion with initial data of class $C^{1,\alpha}$ [7] as well as initial data in a Besov space [4]. It remains an interesting problem to determine the weakest possible conditions on the initial data that allow generation of a unique, regularising solution to the flow.

Remark 1.4. Theorem 1.1 is stated for flows of immersed multiply-covered circles which makes it applicable to the case considered by Miura-Okabe [8]. However the curve diffusion flow starting from a multiply-covered circle is not expected to be stable in general – perturbations that unbalance the length of different leaves are should lead to finite-time singularities.* Miura-Okabe use a rotational symmetry assumption so that only perturbations that preserve the balance of the different leaves are studied, and the power of this hypothesis can be seen in the isoperimetric inequality for immersed rotationally symmetric curves that they establish. That is, preservation of the smallness of the isoperimetric ratio (see Lemma 2.4) is proved in [8] for rotationally symmetric curves *without assuming global existence*. This is the decisive new ingredient that allows their result on the curve diffusion flow to go through.

Yoshikazu Giga famously conjectured around ten years ago that:

Conjecture (Giga's Conjecture). Suppose $\gamma : \mathbb{S} \times [0, T) \rightarrow \mathbb{R}^2$ is a curve diffusion flow with smooth initial data γ_0 that has the property:

$$\gamma(\cdot, t) \text{ is an embedding for each } t \in [0, T).$$

Then $T = \infty$.

*Rigorously establishing this is an interesting open problem.

If this conjecture holds, our Theorem 1.1 implies: *an embedded curve diffusion flow converges smoothly and exponentially fast to a round circle.*

Note that in the case of the curve diffusion flow in homotopy classes outside ω -circles, shrinkers are expected and one explicit example is known (the Lemniscate of Bernoulli, see [5]). Indeed, Chou [2] suggests that parabolic rescalings of Type I singularities (a definition also due to Chou) are self-similar solutions to the curve diffusion flow:

Conjecture (Chou's Conjecture). *Suppose $\gamma : \mathbb{S} \times [0, T) \rightarrow \mathbb{R}^2$ is a curve diffusion flow with $T < \infty$ that satisfies the estimate*

$$\|k\|_2^2(t) \leq C(T - t)^{-1/4}, \quad (1.1)$$

for some $C \in \mathbb{R}$, and $t \in [0, T)$.

Then a parabolic rescaling (we assume the centre of mass of γ is the origin)

$$\eta(s, \tau) = (T - t)^{-\frac{1}{4}}\gamma(s, t), \quad \partial_\tau \eta = -k_{ss}^\eta v^\eta + \frac{1}{4}\eta$$

where $\tau = -\log(T - t)$ yields a self similar solution η to the curve diffusion flow, that is, η converges as $t \rightarrow \infty$ to a solution η^∞ of

$$\langle \eta^\infty, v^{\eta^\infty} \rangle = 4k_{ss}^{\eta^\infty}. \quad (1.2)$$

Curve diffusion flows satisfying (1.1) are said to be *Type I*. As suggested by Chou [2], this conjecture can be approached by studying all solutions to the shrinker Eq (1.2), and classifying all such solutions is an interesting problem. Since area is constant along the curve diffusion flow, and shrinkers must eventually enclose zero area, all shrinkers have zero signed enclosed area. This means that all solutions to (1.2) are non-embedded (and have zero signed enclosed area).

The maximal time of smooth existence for the curve diffusion flow with $\omega = 0$ is always finite. This is because along the flow the curvature scalar always has a zero, and so the Wirtinger inequality implies that the flow can not exist past time $T = \frac{L_0^4}{16\pi^4}$. This is not a sharp estimate, as equality in the estimates used by the proof implies $k = \text{const}$, and there is no such curve that also has $\omega = 0$. A reasonable conjecture is that the maximal existence time for a curve diffusion flow with winding number zero and initial length L_0 is bounded by that of the Lemniscate of Bernoulli starting with length L_0 .

As far as we know the connection between non-convexity and maximal existence time was observed first by Chou [2]. It was later used in [10] to estimate the total waiting time before a curve diffusion flow with small L^2 -oscillation of curvature and isoperimetric ratio becomes uniformly convex. Theorem 1.1 here combined with the waiting time estimate [10, Proposition 1.5] yields an estimate for the total waiting time before *any* global curve diffusion flow becomes uniformly convex.

2. Proof of Theorem 1.1

We complete the proof in a sequence of lemmata. The hypotheses of Theorem 1.1 are not restated, but assumed throughout.

Lemma 2.1. *There exists a subsequence of times $t_j \rightarrow \infty$ such that $\|k_s\|_2^2(t_j) \rightarrow 0$.*

Proof. Note that

$$L' = -\|k_s\|_2^2.$$

Integration and the monotonicity of length implies the result. \square

As far as we are aware, Lemma 2.1 was first observed by Chou [2].

Now we observe that the signed enclosed area of the curve under the flow can not be non-positive.

Lemma 2.2. *The signed enclosed area A of γ is strictly positive.*

Proof. Our convention is that the area of an ω -circle is positive (and equal to $\omega\pi r^2$, where r is its radius). If the area is not strictly positive, then the flow can exist for at most a finite amount of time. Since area is constant under the flow, if it is non-positive at any time, it must be so for all time. Then, since $A(t) < 0$ implies that $k(s_t, t) = 0$ for some point s_t , we use the Wirtinger and Poincaré inequalities to conclude

$$L'(t) \leq -\frac{\pi^2}{L^2} \|k\|_2^2 \leq -\frac{4\pi^4}{L^3}$$

which implies that the flow can exist at most for a finite time (in fact $T \leq \frac{L_0}{16\pi^4}$). This is a contradiction. \square

Now set

$$\mathcal{D} = L^2 - 4\omega\pi A.$$

Recall that \mathcal{D} is the *isoperimetric defect*. Note that $\mathcal{D} \geq 0$ with equality if and only if γ is an ω -circle. Since length is monotone decreasing and signed enclosed area is constant in time, we immediately have:

Lemma 2.3. *The isoperimetric defect satisfies*

$$\mathcal{D}'(t) \leq 0.$$

By combining Lemma 2.1, Lemma 2.2 and Lemma 2.3 we find the following.

Corollary 2.4. *Let $\varepsilon > 0$ be arbitrary. There exists a t_j such that for all $t > t_j$,*

$$\mathcal{I}(t) := \frac{L^2}{4\omega A} \leq 1 + \varepsilon.$$

Proof. The estimate

$$\mathcal{D} \leq \frac{L^2}{\sqrt{2\pi}} K_{osc}^{\frac{1}{2}}, \quad K_{osc} = L \int_{\gamma} (k - \bar{k})^2 ds \quad (2.1)$$

(note that $\bar{k} = 2\omega\pi/L$ is the average of k) together with the Poincaré inequality implies

$$\mathcal{D} \leq \frac{L^{\frac{7}{2}}}{2^{\frac{3}{2}}\omega\pi^{\frac{3}{2}}} \|k_s\|_2. \quad (2.2)$$

The estimate (2.1) follows using an elementary argument (see the Appendix).

Note that length is uniformly bounded below by the measure-theoretic isoperimetric inequality

$$L^2(t) \geq 4\pi A^+$$

where A^+ is the area of γ changed so that the orientation of every loop is positive (no segments of negative area). Clearly we always have $A^+ \geq A$ and so (using the constancy of area along the flow)

$$L^2(t) \geq 4\pi A(0)$$

follows. This is not a sharp estimate for $\omega \neq 1$ but it is enough to give a uniform lower bound on evolving area for any winding number. Note that the shrinking Lemniscare of Bernoulli has $A(0) = 0$ and $T < \infty$. Lemma 2.2 implies that $A(0) > 0$ and so in particular length is uniformly bounded from below by a universal constant.

Therefore the estimate (2.2) implies that $\mathcal{D}(t_j) \rightarrow 0$, which implies $\mathcal{I}(t_j) \rightarrow 1$. Now take j sufficiently large that $\mathcal{I}(t_j) \leq 1 + \varepsilon$, where $\varepsilon > 0$ is as given by hypothesis. Then Lemma 2.3 implies that $\mathcal{I}(t) \leq 1 + \varepsilon$ for all $t > 0$, as required. \square

Lemma 2.5. *There are uniform bounds on all derivatives of curvature.*

Proof. We recall the estimate (2.2) and

$$K_{osc} \leq \frac{L^3}{4\omega^2\pi^2} \|k_s\|_2^2, \quad (2.3)$$

which follows from the Poincaré inequality. Each of the estimates (2.1), (2.3) are uniform (from the uniform boundedness of length).

These estimates imply that the isoperimetric defect and K_{osc} converge to zero at times t_j as $j \rightarrow \infty$. Therefore for j sufficiently large, the hypotheses of [10, Proposition 3.7] are satisfied at time t_j . Given the eventual smallness of the isoperimetric ratio (Corollary 2.4), the argument in [10] (see also [8]) applies to yield uniform estimates on K_{osc} . This then implies bounds on all derivatives of curvature by the evolutionary interpolation inequalities of Dziuk-Kuwert-Schätzle [3].[†] \square

Lemma 2.6. *The isoperimetric defect decays exponentially fast:*

$$\mathcal{D}(t) \leq \mathcal{D}_0 e^{-4\bar{k}_0^4 t}.$$

Proof. Since

$$\mathcal{D}' = -2L \int_{\gamma} k_s^2 ds,$$

the estimate (2.1) implies

$$\mathcal{D}' \leq -\mathcal{D} \frac{64\omega^4\pi^4}{L^4} \leq -\mathcal{D} \frac{64\omega^4\pi^4}{L_0^4}.$$

Therefore

$$(\log \mathcal{D})' \leq -4\bar{k}_0^4,$$

where \bar{k}_0 is the initial average of the curvature scalar. Then

$$\mathcal{D}(t) \leq \mathcal{D}_0 e^{-4\bar{k}_0^4 t},$$

as required. \square

Remark 2.7. One may also attempt the same strategy with the isoperimetric ratio in place of the isoperimetric defect. However, the resultant sharp estimate looks like

$$(\mathcal{I}^{-1})' \geq \bar{k}^4.$$

This results in linear decay of the isoperimetric ratio – not exponential.

[†]The argument is by now standard in the literature, see for instance [1, Theorem 6.4] for a recent example.

Lemma 2.8. *All derivatives of curvature decay exponentially fast, with the explicit estimate*

$$\|k_{s^m}\|_2^2 \leq c_m e^{-\bar{k}_0^4 t}.$$

Here $k_{s^m} = \partial_s^m k$.

Proof. Since the previous lemma establishes uniform bounds for all derivatives of curvature, a Fourier series argument (taught to us by Ben Andrews many years ago and used in for example [1]; we give an explanation in the appendix). [‡]

$$K_{osc} \leq c \frac{\sqrt{\mathcal{D}}}{L}$$

where c depends on the estimates for k and k_s . Now the isoperimetric inequality (and constancy of area) implies that L is uniformly bounded from below, so the above implies that K_{osc} satisfies

$$K_{osc} \leq c_0 e^{-2\bar{k}_0^4 t}.$$

Integration by parts and the uniform estimates imply that for any m we have

$$\|k_{s^m}\|_2^2 \leq c_m e^{-\bar{k}_0^4 t},$$

as required. \square

The exponential decay implies convergence of the position vector γ to an ω -circle with a standard argument (for example this was used by Huisken [6] in his seminal work on mean curvature flow). Briefly, this is because we then have estimates on γ by simply integrating the evolution equation in time. We may convert from arclength derivatives to arbitrary ones in a manner analogous to [3, proof of Theorem 3.1]. Note that this integration in time can not be done with only linear decay estimates.

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[‡]The author recently learned that similar Fourier series arguments (and inequalities that imply the below) also appear in [9].

Conflict of interest

The author declares no conflict of interest.

References

1. B. Andrews, J. McCoy, G. Wheeler, V. M. Wheeler, Closed ideal planar curves, *Geom. Topol.*, **24** (2020), 1019–1049.
2. K. S. Chou, A blow-up criterion for the curve shortening flow by surface diffusion, *Hokkaido Math. J.*, **32** (2003), 1–19.
3. G. Dziuk, E. Kuwert, R. Schätzle, Evolution of elastic curves in \mathbb{R}^n : existence and computation, *SIAM J. Math. Anal.*, **33** (2002), 1228–1245.
4. J. Escher, P. B. Mucha, The surface diffusion flow on rough phase spaces, *Discrete Cont. Dyn. A*, **26** (2010), 431–453.
5. M. Edwards, A. Gerhardt-Bourke, J. McCoy, G. Wheeler, V. M. Wheeler, The shrinking figure eight and other solitons for the curve diffusion flow, *J. Elasticity*, **119** (2015), 191–211.
6. G. Huisken, Flow by mean-curvature of convex surfaces into spheres, *J. Differ. Geom.*, **20** (1984), 237–266.
7. J. LeCrone, Y. Z. Shao, G. Simonett, The surface diffusion and the Willmore flow for uniformly regular hypersurfaces, *Discrete Cont. Dyn. S*, **13** (2020), 3503–3524.
8. T. Miura, S. Okabe, On the isoperimetric inequality and surface diffusion flow for multiply winding curves, 2019, *arXiv:1909.08816*.
9. T. Nagasawa, K. Nakamura, Interpolation inequalities between the deviation of curvature and the isoperimetric ratio with applications to geometric flows, *Adv. Differential Equ.*, **24** (2019), 581–608.
10. G. Wheeler, On the curve diffusion flow of closed plane curves, *Ann. Mat. Pur. Appl.*, **192** (2013), 931–950.

A. Bounding oscillation of curvature by the isoperimetric defect

Both estimates used in this article are contained in the following statement.

Lemma A.1. *Let $\gamma : \mathbb{S} \rightarrow \mathbb{R}^2$ be a smooth immersed curve. Then*

$$\frac{2\pi}{L^4} \mathcal{D}^2 \leq K_{osc} \leq c \frac{\mathcal{D}^{\frac{1}{2}}}{L}. \quad (\text{A.1})$$

where c is such that

$$L^3 \|k_s\|_2^2 + L^{5/2} \|k\|_6^3 K_{osc}^{\frac{1}{2}} \leq c.$$

Note that when the upper estimate is used, curvature and its derivative(s) are uniformly bounded, so the hypothesis involving c is satisfied.

A.1. Lower estimate

Let us prove first the lower estimate. To begin, we translate the curve γ to $\tilde{\gamma}$ so that

$$\int \tilde{\gamma} ds = 0$$

and the estimate

$$|\tilde{\gamma}| \leq \frac{L}{\sqrt{2\pi}}$$

holds. The lower estimate in (A.1) is invariant under translation (this is true of both sides), so if we can prove the estimate for $\tilde{\gamma}$ it will be true for γ also.

Recall that

$$A = -\frac{1}{2} \int \langle \tilde{\gamma}, v \rangle ds.$$

Then integration by parts using $\langle \tilde{\gamma}, \tilde{\gamma}_s \rangle_s = 1 + k \langle \tilde{\gamma}, v \rangle$ implies

$$\begin{aligned} \mathcal{D} &= \int L + 2\omega\pi \langle \tilde{\gamma}, v \rangle ds \\ &= \int -kL \langle \tilde{\gamma}, v \rangle + 2\omega\pi \langle \tilde{\gamma}, v \rangle ds \\ &= \int (-kL + 2\omega\pi) \langle \tilde{\gamma}, v \rangle ds. \end{aligned}$$

Since $\bar{k} = \frac{1}{L} \int k ds = \frac{2\omega\pi}{L}$, we have $\bar{k} - k = \frac{1}{L}(-kL + 2\omega\pi)$. Therefore

$$\mathcal{D} \leq \int |-kL + 2\omega\pi| \frac{L}{\sqrt{2\pi}} ds \leq \frac{L^2}{\sqrt{2\pi}} \|k - \bar{k}\|_1.$$

The lower estimate in (A.1) follows now from Hölder's inequality.

Remark A.2. The estimate

$$\left| \langle \gamma, v \rangle - \frac{1}{L} \int \langle \gamma, v \rangle ds \right| \leq L$$

is false in general. Take for example a circle with unit radius centred at $(P, 0)$ in the plane. Then,

$$\sup \left| \langle \gamma, v \rangle - \frac{1}{L} \int \langle \gamma, v \rangle ds \right| \geq (P + 1) + \frac{2\pi}{2\pi} = P + 2.$$

Since P is arbitrary, this quantity is unbounded. This estimate is used in [9], although it is not a major issue (translation invariance as used here can also be used there).

A.2. Upper estimate

As in [1, Proof of Theorem 6.1] we work in the complex plane. Let us identify $\gamma(s) = (x(s), y(s))$ with $s \mapsto x(s) + iy(s)$. The idea is to use a Fourier series decomposition for γ to reduce the question of proving geometric inequalities to properties of infinite series of integers. Here we have γ smooth, but in order to express γ as a convergent Fourier series we only require $\gamma \in W^{1,2}$. Our manipulations of the series below require significantly more regularity; for instance, the highest order one needs $\gamma \in W^{4,2}$.

We write

$$\gamma(s) = \sum_{p \in \mathbb{Z}} \hat{\gamma}(p) c_p(s)$$

where the coefficients c_p are given by

$$c_p(s) = \frac{1}{\sqrt{L}} e^{2i\omega\pi ps/L},$$

with projection $\hat{\gamma}$ defined via

$$\hat{\gamma}(p) = \int_{\gamma} \gamma(s) \overline{c_p(s)} ds.$$

Curvature arises by differentiating γ , which is the same as differentiating the Fourier decomposition. The basic relation is:

Lemma A.3. *For $q \geq 2$ we have*

$$\sum_{p \in \mathbb{Z}} p^q |\hat{\gamma}(p)|^2 = -\frac{i}{(2i\omega\pi/L)^q} \int k Q_{q-1} ds = \frac{i^{-q-1} L^q}{(2\omega\pi)^q} \int k Q_{q-1} ds$$

where $Q_q = (\partial_s^{q-1} \gamma) \overline{\partial_s \gamma}$. For the $q = 1$ case we have

$$\sum_{p \in \mathbb{Z}} p |\hat{\gamma}(p)|^2 = \frac{iL}{2\omega\pi} \int Q_1 ds.$$

Proof. We provide a sketch with the essential steps. First, observe that for $q \geq 2$ we have

$$\int Q_q ds = i \int k Q_{q-1} ds. \quad (\text{A.2})$$

This is because $Q_q = ikQ_{q-1} + \partial_s Q_{q-1}$ (this equality can be proved by a straightforward induction argument). But now we use the series decomposition with orthonormality of c_p to conclude

$$\begin{aligned} \int Q_q ds &= \int (\partial_s^{q-1} \gamma) \overline{\partial_s \gamma} ds \\ &= \int \left(\sum_p (2i\omega\pi/L)^{q-1} p^{q-1} c_p(s) \hat{\gamma}(p) \right) \overline{\left(\sum_p (2i\omega\pi/L) p c_p(s) \hat{\gamma}(p) \right)} ds \\ &= -(2i\omega\pi/L)^q \sum_{p \in \mathbb{Z}} p^q |\hat{\gamma}(p)|^2. \end{aligned}$$

In the $q = 1$ case this finishes the proof; for $q \geq 2$ we combine this with (A.2) to finish the proof. \square

Lemma A.3 allows us to simply express the isoperimetric defect and the L^2 oscillation of curvature in terms of infinite series. To see this, we consider the cases of $q = 1, 2, 3, 4, 5, 6$ in Lemma A.3. For each of computation we make the notation

$$\mathcal{T} = \langle \gamma, \partial_s \gamma \rangle, \quad \text{and} \quad \mathcal{N} = \langle \gamma, \nu \rangle.$$

Note that $\partial_s \mathcal{T} = 1 + k\mathcal{N}$ and $\partial_s \mathcal{N} = -k\mathcal{T}$. We calculate

$$(q = 1) \quad \sum_{p \in \mathbb{Z}} p |\hat{\gamma}(p)|^2 = \frac{iL}{2\omega\pi} \int Q_1 ds = \frac{iL}{2\omega\pi} \int (x + iy)(x_s - iy_s) ds$$

$$= \frac{iL}{2\omega\pi} \int \mathcal{T} + i\mathcal{N} ds = \frac{iL}{2\omega\pi} \int i\mathcal{N} ds = \frac{LA}{\omega\pi}$$

$$(q = 2) \quad \sum_{p \in \mathbb{Z}} p^2 |\hat{\gamma}(p)|^2 = \frac{i^{-3}L^2}{(2\omega\pi)^2} \int kQ_1 ds$$

$$= i \frac{L^2}{(2\omega\pi)^2} \int k\mathcal{T} + ik\mathcal{N} ds = \frac{L^3}{(2\omega\pi)^2}$$

$$(q = 3) \quad \sum_{p \in \mathbb{Z}} p^3 |\hat{\gamma}(p)|^2 = \frac{i^{-4}L^3}{(2\omega\pi)^3} \int kQ_2 ds = \frac{L^3}{(2\omega\pi)^3} \int ik^2 Q_1 + k\partial_s Q_1 ds$$

$$= \frac{L^3}{(2\omega\pi)^3} \int (ik^2 + k\partial_s) \mathcal{T} + i(ik^2 + k\partial_s) \mathcal{N} ds$$

$$= \frac{L^3}{(2\omega\pi)^3} \int k ds = \frac{L^3}{(2\omega\pi)^2}$$

In fact, $Q_2 = (\partial_s \gamma)(\overline{\partial_s \gamma}) = 1$. We continue to calculate:

$$(q = 4) \quad \sum_{p \in \mathbb{Z}} p^4 |\hat{\gamma}(p)|^2 = \frac{i^{-5}L^4}{(2\omega\pi)^4} \int kQ_3 ds = -i \frac{L^4}{(2\omega\pi)^4} \int (ik^2 + k\partial_s) Q_2 ds$$

$$= -i \frac{L^4}{(2\omega\pi)^4} \int ik^2 ds = \frac{L^4}{(2\omega\pi)^4} \int k^2 ds$$

$$(q = 5) \quad \sum_{p \in \mathbb{Z}} p^5 |\hat{\gamma}(p)|^2 = \frac{i^{-6}L^5}{(2\omega\pi)^5} \int kQ_4 ds = -\frac{L^5}{(2\omega\pi)^5} \int (ik^2 + k\partial_s) Q_3 ds$$

$$= -\frac{L^5}{(2\omega\pi)^5} \int (ik^2 + k\partial_s)(ik) ds$$

$$= \frac{L^5}{(2\omega\pi)^5} \int k^3 ds.$$

$$(q = 6) \quad \sum_{p \in \mathbb{Z}} p^6 |\hat{\gamma}(p)|^2 = \frac{i^{-7}L^6}{(2\omega\pi)^6} \int kQ_5 ds = i \frac{L^6}{(2\omega\pi)^6} \int (ik^2 + k\partial_s)(-k^2 + ik_s) ds$$

$$= i \frac{L^6}{(2\omega\pi)^6} \int (-ik^4 - k^2 k_s - 2k^2 k_s + ik k_{ss}) ds$$

$$= \frac{L^6}{(2\omega\pi)^6} \int (k^4 + k_s^2) ds.$$

Although we won't need it here, one can continue to calculate more of these identities for greater values of q . They can be used to establish higher order inequalities – we include an example (inequality (A.3)). Here are the next two:

$$(q = 7) \quad \sum_{p \in \mathbb{Z}} p^7 |\hat{\gamma}(p)|^2 = \frac{i^{-8} L^7}{(2\omega\pi)^7} \int k Q_6 ds$$

$$= \frac{L^7}{(2\omega\pi)^7} \int (ik^2 + k\partial_s)(-ik^3 - 3kk_s + ik_{ss}) ds$$

$$= \frac{L^7}{(2\omega\pi)^7} \int k^5 - 6ik^3k_s - 4k^2k_{ss} - 3kk_s^2 + ikk_{sss} ds$$

$$= \frac{L^7}{(2\omega\pi)^7} \int k^5 + 5kk_s^2 ds.$$

$$(q = 8) \quad \sum_{p \in \mathbb{Z}} p^8 |\hat{\gamma}(p)|^2 = \frac{i^{-9} L^8}{(2\omega\pi)^8} \int k Q_7 ds$$

$$= -i \frac{L^8}{(2\omega\pi)^8} \int (ik^2 + k\partial_s)(k^4 - 6ik^2k_s - 4kk_{ss} - 3k_s^2 + ik_{sss}) ds$$

$$= -i \frac{L^8}{(2\omega\pi)^8} \int ik^6 + 10k^4k_s - 15ik^2k_s^2 - 10ik^3k_{ss}$$

$$- 10kk_s k_{ss} - 5k^2k_{sss} + ikk_{s^4} ds$$

$$= \frac{L^8}{(2\omega\pi)^8} \int k^6 + 15k^2k_s^2 + 10k^3k_{ss} + 10ikk_s k_{ss} + 5ik^2k_{sss} + k_{ss}^2 ds$$

$$= \frac{L^8}{(2\omega\pi)^8} \int k^6 - 15k^2k_s^2 + k_{ss}^2 ds.$$

This last identity can be used trivially to deduce the estimate

$$15 \int k^2 k_s^2 ds \leq \int k^6 + k_{ss}^2 ds. \quad (\text{A.3})$$

This estimate does not seem to follow from using more common methods.

Now we return to our aim of proving the upper estimate. First, we write key quantities in terms of infinite series.

Lemma A.4. *The equalities*

$$\mathcal{D} = \frac{(2\omega\pi)^2}{L} \sum_{p \in \mathbb{Z}} p(p-1) |\hat{\gamma}(p)|^2$$

and

$$K_{osc} = \frac{(2\omega\pi)^4}{L^3} \sum_{p \in \mathbb{Z}} p^2(p^2-1) |\hat{\gamma}(p)|^2 = \frac{(2\omega\pi)^4}{L^3} \sum_{p \in \mathbb{Z}} p^3(p-1) |\hat{\gamma}(p)|^2$$

hold.

Proof. We calculate

$$\mathcal{D} = L^2 - 4\omega\pi A = \frac{4\omega^2\pi^2}{L} \left(\frac{L^3}{(2\omega\pi)^2} - \frac{AL}{\omega\pi} \right).$$

Using the expressions for $q = 2$ and $q = 1$, this shows

$$\mathcal{D} = \frac{4\omega^2\pi^2}{L} \left(\sum_{p \in \mathbb{Z}} (p^2 - p) |\hat{\gamma}(p)|^2 \right),$$

as required. Now

$$K_{osc} = L \int (k - \bar{k})^2 ds = L \int k^2 ds - (2\omega\pi)^2$$

which, using the expressions for $q = 4$ and $q = 2$, gives

$$K_{osc} = \frac{(2\omega\pi)^4}{L^3} \sum_{p \in \mathbb{Z}} p^4 |\hat{\gamma}(p)|^2 - \frac{(2\omega\pi)^4}{L^3} \sum_{p \in \mathbb{Z}} p^2 |\hat{\gamma}(p)|^2 = \frac{(2\omega\pi)^4}{L^3} \sum_{p \in \mathbb{Z}} (p^4 - p^2) |\hat{\gamma}(p)|^2.$$

For the second term we could have used the equality for $q = 3$ instead, which gives the last equality and finishes the proof. \square

Finally, we use the Cauchy-Schwartz inequality in $\ell^2(\mathbb{Z})$ to conclude essentially the upper estimate.

Lemma A.5. *We have*

$$K_{osc}^2 \leq L\mathcal{D} \left(\|k_s\|_2^2 + L^{-1/2} \|k\|_6^3 K_{osc}^{\frac{1}{2}} \right).$$

Proof. We use the Cauchy-Schwartz inequality to estimate

$$\begin{aligned} K_{osc} &= \frac{(2\omega\pi)^4}{L^3} \sum_{p \in \mathbb{Z}} p^3 (p-1) |\hat{\gamma}(p)|^2 \\ &\leq \frac{(2\omega\pi)^4}{L^3} \left(\sum_{p \in \mathbb{Z}} p(p-1) |\hat{\gamma}(p)|^2 \right)^{\frac{1}{2}} \left(\sum_{p \in \mathbb{Z}} p^5 (p-1) |\hat{\gamma}(p)|^2 \right)^{\frac{1}{2}}. \end{aligned}$$

Now Lemma A.4 implies

$$K_{osc} \leq \frac{(2\omega\pi)^4}{L^3} \mathcal{D}^{\frac{1}{2}} \frac{L^{\frac{1}{2}}}{2\omega\pi} \left(\sum_{p \in \mathbb{Z}} p^5 (p-1) |\hat{\gamma}(p)|^2 \right)^{\frac{1}{2}}.$$

The expressions for $q = 5$ and $q = 6$ now imply

$$\begin{aligned} K_{osc}^2 &\leq \frac{(2\omega\pi)^8}{L^6} \mathcal{D} \frac{L}{(2\omega\pi)^2} \left(\frac{L^6}{(2\omega\pi)^6} \int k^4 + k_s^2 - \frac{2\omega\pi}{L} k^3 ds \right) \\ &= L\mathcal{D} \int k_s^2 + k^3 (k - \bar{k}) ds \\ &\leq L\mathcal{D} \left(\|k_s\|_2^2 + L^{-1/2} \|k\|_6^3 K_{osc}^{\frac{1}{2}} \right), \end{aligned}$$

as required. \square

Now we finish the proof of the upper estimate. Since by hypothesis we have

$$L^3\|k_s\|_2^2 + L^{5/2}\|k\|_6^3 K_{osc}^{\frac{1}{2}} \leq c,$$

Lemma A.5 implies

$$K_{osc} \leq c \frac{\mathcal{D}^{\frac{1}{2}}}{L}$$

as required.



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