



Research article

Least energy sign-changing solutions for a class of fractional (p, q)-Laplacian problems with critical growth in R^N

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Abstract: This paper considers the following fractional (p, q)-Laplacian equation:

(-Δ)\_p^s u + (-Δ)\_q^s u + V(x)(|u|^{p-2}u + |u|^{q-2}u) = λf(u) + |u|^{q\_s^\*-2}u in R^N,

where s ∈ (0, 1), λ > 0, 2 < p < q < N/s, (-Δ)\_t^s with t ∈ {p, q} is the fractional t-Laplacian operator, and potential V is a continuous function. Using constrained variational methods, a quantitative Deformation Lemma and Brouwer degree theory, we prove that the above problem has a least energy sign-changing solution u\_λ under suitable conditions on f, V and λ. Moreover, we show that the energy of u\_λ is strictly larger than two times the ground state energy.

Keywords: fractional (p, q)-Laplacian; sign-changing solutions; critical problem

Mathematics Subject Classification: 35J20, 35J65

1. Introduction and main results

In this paper, we investigate the existence of the least energy sign-changing solution for the following fractional (p, q)-Laplacian problem:

(-Δ)\_p^s u + (-Δ)\_q^s u + V(x)(|u|^{p-2}u + |u|^{q-2}u) = λf(u) + |u|^{q\_s^\*-2}u in R^N, (1.1)

where s ∈ (0, 1), 2 < p < q < N/s, λ > 0. The potential V ∈ C(R^N, R) and the operator (-Δ)\_t^s with t ∈ {p, q} is the fractional Laplacian which, up to a normalizing constant, may be defined for any

$u : \mathbb{R}^N \rightarrow \mathbb{R}$  smooth enough by setting

$$(-\Delta)_t^s u(x) = 2 \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}^N \setminus B_\varepsilon(x)} \frac{|u(x) - u(y)|^{t-2} (u(x) - u(y))}{|x - y|^{N+ts}} dy, \quad x \in \mathbb{R}^N$$

along functions  $u \in C_0^\infty(\mathbb{R}^N)$ , where  $B_\varepsilon(x)$  denotes the ball of  $\mathbb{R}^N$  centered at  $x \in \mathbb{R}^N$  and radius  $\varepsilon > 0$ .

When  $s = 1$ , problem (1.1) boils down to a  $(p, q)$ -Laplacian problem of the following type:

$$-\Delta_p u - \Delta_q u + V(x) (|u|^{p-2} u + |u|^{q-2} u) = f(u) \quad \text{in } \mathbb{R}^N. \quad (1.2)$$

As can be seen in [5] and [30], the applications in plasma physics, chemical process design, and biology have generated the majority of interest in this broad class of problems. In the last decade, many authors investigated problem (1.2), for example, Barile and Figueiredo [5] showed that (1.2) has a least energy sign-changing solution by using the deformation lemma and the Brouwer degree theory. For more interesting results involving  $(p, q)$ -Laplacian problems, we also mention [9, 22, 24, 30, 32, 39] and references therein.

When  $s \in (0, 1)$  and  $p = q = 2$ , problem (1.1) appears in the study of standing wave solutions, i.e., solutions of the form  $\psi(x, t) = u(x)e^{-i\omega t}$ , to the following fractional Schrödinger equation:

$$i\hbar \frac{\partial \psi}{\partial t} = \hbar^{2s} (-\Delta)^s \psi + W(x)\psi - f(|\psi|) \quad \text{in } \mathbb{R}^N \times \mathbb{R}, \quad (1.3)$$

where  $\hbar$  is the Planck constant,  $W : \mathbb{R}^N \rightarrow \mathbb{R}$  is an external potential and  $f$  is a suitable nonlinearity. Laskin [28, 29] first introduced the fractional Schrödinger equation due to its fundamental importance in the study of particles on stochastic fields modeled by Lévy processes. After that, fractional Schrödinger equations received a lot of attention, and a lot of interesting results were obtained. We direct the curious reader to [33] for a basic overview of this topic for more information. For the existence, multiplicity, and behavior of standing wave solutions to Eq (1.3), we refer to [10, 11, 14, 16, 21, 23, 36, 37] and the references therein.

When  $p = q \neq 2$ , problem (1.1) boils down to the following fractional Laplacian problem:

$$(-\Delta)_p^s u + V(x)|u|^{p-2} u = f(u) \quad \text{in } \mathbb{R}^N. \quad (1.4)$$

Problem (1.4) piques the interest of researchers because of its nonlocal character and the operator's nonlinearity. In [15], the authors obtained infinitely many sign-changing solutions of (1.4) by using descent flow with invariant sets. By applying the deformation Lemma and the Brouwer degree, they also proved that (1.4) has a least energy sign-changing solution. It is noteworthy that Wang and Zhou [37] used a similar method to obtain the least energy sign-changing of (1.4) with  $p = 2$ . In addition, for Eq (1.4), we refer to [2, 3, 18, 19, 34, 35] for existence and multiplicity results, to [13, 25] for regularity results.

However, only a few papers considered fractional  $(p, q)$ -Laplacian problems. For instance, the authors of [17] investigated the existence, nonexistence and multiplicity of solutions for a fractional  $(p, q)$ -Laplacian problem with subcritical growth. Alves et al [1] studied the following problem:

$$(-\Delta)_p^s u + (-\Delta)_q^s u + V(\varepsilon x) (|u|^{p-2} u + |u|^{q-2} u) = f(u) \quad \text{in } \mathbb{R}^N, \quad (1.5)$$

where the potential  $V(x)$  satisfies the Rabinowitz conditions. By virtue of the Ljusternik-Schnirelmann theory and minimax theorems, they explored the existence, multiplicity, and concentration of nontrivial solutions provided that  $\varepsilon$  is sufficiently small. Ambrosio and Rădulescu [4] considered the existence and concentration of positive solutions for (1.5) with the del Pino-Felmer type potential conditions. For the other work on (1.1) or similar problems, we refer the reader to [4, 20, 26, 40–44] and the references therein.

Motivated by the above results, it is natural to ask, whether the problem (1.1) had sign-changing solutions when the nonlinear term  $f$  has critical growth. To our knowledge, this question is open. In [23], the authors considered the following problem:

$$\begin{cases} (-\Delta)^s u = \lambda f(x, u) + |u|^{2_s^*-2} u & \text{in } \Omega, \\ u = 0 & \text{in } \mathbb{R}^N \setminus \Omega, \end{cases} \quad (1.6)$$

where  $\Omega \subset \mathbb{R}^N$  is a bounded domain,  $2_s^* = \frac{2N}{N-2s}$  and  $f$  satisfies some suitable conditions. By using the constrained variational methods, they proved the least energy sign-changing solution of (1.6) when  $\lambda$  sufficiently large. However, since (1.1) contains the nonlocal and nonlinear term  $(-\Delta)_p^s + (-\Delta)_q^s$ , the decomposition of functional  $I_\lambda$  (see the definition in (1.10)) is more complicated than that in [23]. Therefore, some difficulties arise in studying the existence of a least energy sign-changing solution for problem (1.1), and this makes the study interesting.

To study problem (1.1), we consider the following assumptions on  $V$  and  $f$ :

(V<sub>1</sub>)  $V(x) \in C(\mathbb{R}^N)$  and there exists  $V_0 > 0$  such that  $V(x) \geq V_0$  in  $\mathbb{R}^N$ . Moreover,  $\lim_{|x| \rightarrow \infty} V(x) = +\infty$ .

$$(f_1) \quad \lim_{|t| \rightarrow 0^+} \frac{f(t)}{|t|^{p-1}} = 0.$$

(f<sub>2</sub>)  $f$  has a “quasical growth” at infinity, namely,

$$\lim_{|t| \rightarrow +\infty} \frac{f(t)}{|t|^{q_s^*-1}} = 0.$$

We suppose that the function  $f$  satisfies the Ambrosetti-Rabinowitz condition:

(f<sub>3</sub>) There exists  $\theta \in (q, q_s^*)$  such that

$$0 < \theta F(t) = \theta \int_0^t f(s) ds \leq f(t)t \quad \text{for all } |t| > 0, \text{ where } F(t) := \int_0^t f(\tau) d\tau,$$

furthermore, we assume that:

(f<sub>4</sub>) The map  $f$  and its derivative  $f'$  satisfy

$$f'(t) > (q-1) \frac{f(t)}{t} \quad \text{for all } t \neq 0.$$

Clearly, (f<sub>4</sub>) implies that the map  $t \mapsto \frac{f(t)}{|t|^{q-1}}$  is strictly increasing for all  $|t| > 0$ .

Before starting our results, we recall some useful notations. Let  $1 \leq \zeta \leq \infty$ , we denote by  $|u|_\zeta$  the  $L^\zeta$ -norm of  $u : \mathbb{R}^N \rightarrow \mathbb{R}$  belonging to  $L^\zeta(\mathbb{R}^N)$ . For  $0 < s < 1$ , let us define  $\mathcal{D}^{s,\zeta}(\mathbb{R}^N) = \overline{C_c^\infty(\mathbb{R}^N)}^{|\cdot|_{s,\zeta}}$ , where

$$[u]_{s,\zeta} := \left[ \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^\zeta}{|x - y|^{N+s\zeta}} dx dy \right]^{\frac{1}{\zeta}}.$$

Let us denote by  $W^{s,\zeta}(\mathbb{R}^N)$  the set of functions  $u \in L^\zeta(\mathbb{R}^N)$  such that  $[u]_{s,\zeta} < \infty$ , endowed with the natural norm

$$\|u\|_{s,\zeta}^\zeta = [u]_{s,\zeta}^\zeta + |u|_\zeta^\zeta.$$

According to [33], let  $s \in (0, 1)$  and  $N > sq$ , there exists a sharp constant  $S_q > 0$  such that for any  $u \in \mathcal{D}^{s,q}(\mathbb{R}^N)$

$$|u|_{q_s^*}^q \leq S_q^{-1} [u]_{s,q}^q, \quad (1.7)$$

where  $q_s^* = \frac{Nq}{N-qs}$  is the Sobolev critical exponent. Moreover,  $W^{s,q}(\mathbb{R}^N)$  is continuously embedded in  $L^\gamma(\mathbb{R}^N)$  for any  $\gamma \in [q, q_s^*]$  and compactly in  $L^\gamma(B_R(0))$ , for all  $R > 0$  and for any  $\gamma \in [1, q_s^*)$ .

To ensure that problem (1.1) has a variational structure, we consider the following Sobolev space:

$$X = W^{s,p}(\mathbb{R}^N) \cap W^{s,q}(\mathbb{R}^N) \quad (1.8)$$

endowed with the norm

$$\|u\|_X := \|u\|_{W^{s,p}(\mathbb{R}^N)} + \|u\|_{W^{s,q}(\mathbb{R}^N)}.$$

Notice that  $W^{s,r}(\mathbb{R}^N)$  is a separable reflexive Banach space for all  $r \in (1, +\infty)$ , then  $X$  is also a separable reflexive Banach space. We also introduce the following Banach space

$$X_V := \left\{ u \in X : \int_{\mathbb{R}^N} V(x) (|u|^p + |u|^q) dx < +\infty \right\}, \quad (1.9)$$

endowed with the norm

$$\|u\| := \|u\|_{X_V} := \|u\|_{V,p} + \|u\|_{V,q},$$

where  $\|u\|_{V,t}^t := [u]_{s,t}^t + \int_{\mathbb{R}^N} V(x) |u|^t dx$  for  $t \in \{p, q\}$ . For the weak solution to (1.1), we mean a function  $u \in X_V$  such that

$$\begin{aligned} & \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{p-2} (u(x) - u(y)) (\varphi(x) - \varphi(y))}{|x - y|^{N+sp}} dx dy + \int_{\mathbb{R}^N} V(x) |u(x)|^{p-2} u(x) \varphi(x) dx \\ & + \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{q-2} (u(x) - u(y)) (\varphi(x) - \varphi(y))}{|x - y|^{N+sq}} dx dy + \int_{\mathbb{R}^N} V(x) |u(x)|^{q-2} u(x) \varphi(x) dx \\ & = \int_{\mathbb{R}^N} \lambda f(u(x)) \varphi(x) + |u(x)|^{q_s^*-2} u(x) \varphi(x) dx \end{aligned}$$

for all  $\varphi \in X_V$ .

Define the energy functional  $I_\lambda : X_V \rightarrow \mathbb{R}$  by

$$I_\lambda(u) = \frac{1}{p} \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy + \frac{1}{q} \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^q}{|x - y|^{N+qs}} dx dy + \frac{1}{p} \int_{\mathbb{R}^N} V(x) |u(x)|^p dx$$

$$+ \frac{1}{q} \int_{\mathbb{R}^N} V(x)|u(x)|^q dx - \lambda \int_{\mathbb{R}^N} F(u(x)) - \frac{1}{q_s^*} \int_{\mathbb{R}^N} |u(x)|^{q_s^*} dx. \quad (1.10)$$

By the similar arguments as in [1], we can deduce that  $I_\lambda(u) \in C^1(X_V, \mathbb{R})$ .

For convenience, we consider the operator  $\mathcal{A}_p : X_V \rightarrow X_V^*$  and  $\mathcal{A}_q : X_V \rightarrow X_V^*$  given by

$$\begin{aligned} \langle \mathcal{A}_p(u), v \rangle_{X_V^*, X_V} &= \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{p-2} (u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+ps}} dx dy \\ &\quad + \int_{\mathbb{R}^N} V(x)|u|^{p-2} u v dx, \quad \forall u, v \in X_V \end{aligned}$$

and

$$\begin{aligned} \langle \mathcal{A}_q(u), v \rangle_{X_V^*, X_V} &= \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{q-2} (u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+qs}} dx dy \\ &\quad + \int_{\mathbb{R}^N} V(x)|u|^{q-2} u v dx, \quad \forall u, v \in X_V, \end{aligned}$$

where  $X_V^*$  is the dual space of  $X_V$ . In this sequel, for simplicity, we denote  $\langle \cdot, \cdot \rangle_{X_V^*, X_V}$  by  $\langle \cdot, \cdot \rangle$ . Moreover, we denote the Nehari set  $\mathcal{N}_\lambda$  by

$$\mathcal{N}_\lambda = \{u \in X \setminus \{0\} : \langle I'_\lambda(u), u \rangle_{X_V^*, X_V} = 0\}. \quad (1.11)$$

Clearly,  $\mathcal{N}_\lambda$  contains all the nontrivial solutions of (1.1). Denote  $u^+(x) := \max\{u(x), 0\}$  and  $u^-(x) := \min\{u(x), 0\}$ . Then, the sign-changing solutions of (1.1) stay on the following set:

$$\mathcal{M}_\lambda = \{u \in X_V \setminus \{0\} : u^\pm \neq 0, \langle I'_\lambda(u), u^+ \rangle = 0, \langle I'_\lambda(u), u^- \rangle = 0\}. \quad (1.12)$$

Set

$$c := \inf_{u \in \mathcal{N}_\lambda} I(u), \quad (1.13)$$

and

$$c_\lambda := \inf_{u \in \mathcal{M}_\lambda} I(u). \quad (1.14)$$

The main results of this paper are stated in the following theorem.

**Theorem 1.1.** *Suppose that  $(f_1) - (f_4)$  are satisfied. Then there exists  $\Lambda > 0$  such that for all  $\lambda \geq \Lambda$ , the problem (1.1) possesses a least energy sign-changing solution  $u_\lambda$ . Moreover,  $c_\lambda > 2c$ .*

The proof of Theorem 1.1 is based on the arguments presented in [8]. First, we make sure that the minimum of functional  $I_\lambda$  restricted on set  $\mathcal{M}_\lambda$  can be achieved. Then, we demonstrate that it is a critical point of  $I_\lambda$  by applying a suitable variant of the quantitative deformation Lemma. However, one cannot obtain a corresponding equivalent definition of  $(-\Delta)_t^s$  by the harmonic extension approach because of the two fractional  $t$ -Laplacian operators  $(-\Delta)_t^s$  with  $s \in (0, 1)$  and  $t \in \{p, q\}$  (see [11]). Thus, we don't get the decomposition

$$I_\lambda(u) = I_\lambda(u^+) + I_\lambda(u^-) \quad \text{and} \quad \langle I'_\lambda(u), u^\pm \rangle = \langle I'_\lambda(u^\pm), u^\pm \rangle, \quad (1.15)$$

which is very useful to get sign-changing solutions of (1.1), see for instance [5–8, 12]. Furthermore, we could not adapt similar methods like in [23, 37] to conclude the set  $\mathcal{M}_\lambda$  is non empty. This is because for the linear operator  $(-\Delta)^s$ , one can easily deduce that

$$\int_{\mathbb{R}^{2N}} \frac{(u(x) - u(y))(u^+(x) - u^+(y))}{|x - y|^{N+2s}} dx dy = \int_{\mathbb{R}^{2N}} \frac{(u^+(x) - u^+(y))^2}{|x - y|^{N+2s}} dx dy - \int_{\mathbb{R}^{2N}} \frac{(u^+(x)u^-(y) + u^-(x)u^+(y))}{|x - y|^{N+2s}} dx dy,$$

which is important to prove that  $\mathcal{M}_\lambda$  is non-empty. But, for the nonlinear operators  $(-\Delta)_p^s$  and  $(-\Delta)_q^s$ , the above decomposition seems invalid. Fortunately, we find a new way to overcome those difficulties. We use another decomposition estimation by dividing  $\mathbb{R}^{2N}$  into several regions (see Lemma 2.2) as following:

$$\begin{aligned} & \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{t-2}(u(x) - u(y))(u^+(x) - u^+(y))}{|x - y|^{N+ts}} dx dy \\ &= \int_{(\mathbb{R}^N)^+ \times (\mathbb{R}^N)^+} \frac{|u^+(x) - u^+(y)|^t}{|x - y|^{N+ts}} dx dy + \int_{(\mathbb{R}^N)^+ \times (\mathbb{R}^N)^-} \frac{|u^+(x) - u^-(y)|^{t-1} u^+(x)}{|x - y|^{N+ts}} dx dy \\ &+ \int_{(\mathbb{R}^N)^- \times (\mathbb{R}^N)^+} \frac{|u^-(x) - u^+(y)|^{t-1} u^+(y)}{|x - y|^{N+ts}} dx dy, \end{aligned}$$

where  $(\mathbb{R}^N)^+ = \{x \in \mathbb{R}^N : u(x) \geq 0\}$  and  $(\mathbb{R}^N)^- = \{x \in \mathbb{R}^N : u(x) < 0\}$ . As we can see that it will also play an important role in proving  $\deg(\Psi_1, D, 0) = 1$  (see Section 4), and then we can get the minimizer  $u_\lambda$  of  $c_\lambda$  (that is,  $I_\lambda(u_\lambda) = c_\lambda$ ) is exactly a sign-changing solution of Problem (1.1). Besides, due to the critical growth of the nonlinear term, another difficulty arises in verifying the compactness of the minimizing sequence in  $X_V$ . Fortunately, thanks to the sharp constant  $S_q$ , we overcome this difficulty by choosing  $\lambda$  appropriately large to ensure the compactness of the minimizing sequence. Therefore, to obtain the least energy sign-changing solutions of (1.1), a more accurate investigation and meticulous calculations are needed in our setting.

The paper is organized as follows: Section 2 contains some compactness results and the decomposition characteristics of  $I_\lambda$ , which will be crucial to proving the main results. In Section 3, we provide several technical lemmas. The main results are proved in Section 4 by combining the reduced arguments with a variation of the Deformation Lemma and Brouwer degree theory.

Throughout this paper, we will use the following notations:  $L^\lambda(\mathbb{R}^N)$  denotes the usual Lebesgue space with norm  $|\cdot|_\lambda$ ;  $C, C_1, C_2, \dots$  will denote different positive constants whose exact values are not essential to the exposition of arguments.

## 2. Preliminaries

We provide the variational framework for the problem (1.1) in this section and provide some preliminary Lemmas. To begin with, we obtain the following compactness results by recalling the notion of fractional Sobolev space  $X_V$  in (1.9).

**Lemma 2.1.** *Suppose that  $(V_1)$  holds, then for all  $\gamma \in [p, q_s^*]$ , the embedding  $X_V \hookrightarrow L^\gamma(\mathbb{R}^N)$  is continuous. For all  $\gamma \in [p, q_s^*)$ , the embedding  $X_V \hookrightarrow L^\gamma(\mathbb{R}^N)$  is compact.*

*Proof.* Denote  $Y = L^\gamma(\mathbb{R}^N)$  and  $B_R = \{x \in \mathbb{R}^N : |x| < R\}$ ,  $B_R^c = \mathbb{R}^N \setminus \overline{B_R}$ . Denote  $X_p := \{u \in W^{s,p}(\mathbb{R}^N) : \int_{\mathbb{R}^N} V(x)|u|^p dx < +\infty\}$ .

For any  $p \leq \gamma \leq q_s^*$ , the space  $X_p$  is continuously embedded in  $Y$ , the space  $X_V$  is continuously embedded in  $X_p$ , so  $X_V \hookrightarrow Y$  is continuous.

For any  $p \leq \gamma < q_s^*$ , Let  $X_p(\Omega)$  and  $Y(\Omega)$  be the spaces of functions  $u \in X_p$ ,  $u \in Y$  restricted onto  $\Omega \subset \mathbb{R}^N$  respectively. Then, it follows from theorems 6.9, 6.10 and 7.1 in [33] that  $X_p(B_R) \hookrightarrow Y(B_R)$  is compact for any  $R > 0$ . Denote  $V_R = \inf_{x \in B_R^c} V(x)$ . By  $(V_1)$ , we deduce that  $V_R \rightarrow \infty$  as  $R \rightarrow \infty$ . Therefore, we have

$$\int_{B_R^c} |u|^\gamma dx \leq \frac{1}{V_R} \int_{B_R^c} V(x)|u|^\gamma dx \leq \frac{1}{V_R} \|u\|_{X_p}^\gamma,$$

which implies

$$\lim_{R \rightarrow +\infty} \sup_{u \in X \setminus \{0\}} \frac{\|u\|_{L^\gamma(B_R^c)}}{\|u\|_{X_p}} = 0.$$

By virtue of Theorem 7.9 in [27], we can see that  $X_p \hookrightarrow Y$  is compact, moreover,  $X_V \hookrightarrow X_p$  is compact, therefore, by interpolation inequality, the embedding  $X_V \hookrightarrow Y$  is compact for any  $p \leq \gamma < q_s^*$ .  $\square$

**Remark 2.1.** It follows from Lemma 2.1 and  $(f_1)$ ,  $(f_2)$  that  $I_\lambda$  is well-defined on  $X_V$ . Moreover,  $I_\lambda \in C^1(X_V, \mathbb{R}^N)$  and

$$\begin{aligned} \langle I'_\lambda(u), v \rangle &= \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{p-2} (u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x)|u|^{p-2} u v dx \\ &+ \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{q-2} (u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+qs}} dx dy + \int_{\mathbb{R}^N} V(x)|u|^{q-2} u v dx \\ &- \lambda \int_{\mathbb{R}^N} f(u) v dx - \int_{\mathbb{R}^N} |u|^{q_s^*-2} u v dx \end{aligned} \quad (2.1)$$

for all  $v \in X_V$ . Consequently, the critical point of  $I_\lambda$  is the weak solution of the problem (1.1).

Our goal is to find the sign-changing solution to the problem (1.1). As we saw in section 1, one of the challenges is the fact that the functional  $I_\lambda$  does not possess a decomposition like (1.15). Inspired by [15, 37], we have the following:

**Lemma 2.2.** Let  $u \in X_V$  with  $u^\pm \neq 0$ . Then,

- (i)  $I_\lambda(u) > I_\lambda(u^+) + I_\lambda(u^-)$ ,
- (ii)  $\langle I'_\lambda(u), u^\pm \rangle > \langle I'_\lambda(u^\pm), u^\pm \rangle$ .

*Proof.* Observe that

$$\begin{aligned} I_\lambda(u) &= \frac{1}{p} \|u\|_{V,p}^p + \frac{1}{q} \|u\|_{V,q}^q - \lambda \int_{\mathbb{R}^N} F(u) dx - \frac{1}{q_s^*} \int_{\mathbb{R}^N} |u|^{q_s^*} dx \\ &= \frac{1}{p} \langle \mathcal{A}_p(u), u^+ \rangle + \frac{1}{p} \langle \mathcal{A}_p(u), u^- \rangle + \frac{1}{q} \langle \mathcal{A}_q(u), u^+ \rangle + \frac{1}{q} \langle \mathcal{A}_q(u), u^- \rangle \\ &- \lambda \int_{\mathbb{R}^N} F(u^+) dx - \lambda \int_{\mathbb{R}^N} F(u^-) dx - \frac{1}{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx - \frac{1}{q_s^*} \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx. \end{aligned} \quad (2.2)$$

By density (see Theorem 2.4 in [33]), we can assume that  $u$  is continuous. Defining

$$(\mathbb{R}^N)_+ = \{x \in \mathbb{R}^N; u^+(x) \geq 0\} \text{ and } (\mathbb{R}^N)_- = \{x \in \mathbb{R}^N; u^-(x) \leq 0\}.$$

Then for  $u \in X_V$  with  $u^\pm \neq 0$ , by a straightforward computation, one can see that

$$\begin{aligned} \langle \mathcal{A}_p(u), u^+ \rangle &= \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{p-2} (u(x) - u(y)) (u^+(x) - u^+(y))}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u^+|^p dx \\ &= \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_+} \frac{|u^+(x) - u^+(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} u^+(x)}{|x - y|^{N+ps}} dx dy \\ &\quad + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} u^+(y)}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u^+|^p dx \\ &> \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_+} \frac{|u^+(x) - u^+(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u^+|^p dx \\ &\quad + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^p}{|x - y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^p}{|x - y|^{N+ps}} dx dy \\ &= \langle \mathcal{A}_p(u^+), u^+ \rangle \end{aligned} \tag{2.3}$$

and

$$\begin{aligned} \langle \mathcal{A}_p(u), u^- \rangle &= \int_{\mathbb{R}^{2N}} \frac{|u(x) - u(y)|^{p-2} (u(x) - u(y)) (u^-(x) - u^-(y))}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u^-|^p dx \\ &= \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_-} \frac{|u^-(x) - u^-(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} (-u^-(y))}{|x - y|^{N+ps}} dx dy \\ &\quad + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} (-u^-(x))}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u^-|^p dx \\ &> \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_-} \frac{|u^-(x) - u^-(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u^-|^p dx \\ &\quad + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^p}{|x - y|^{N+ps}} dx dy \\ &= \langle \mathcal{A}_p(u^-), u^- \rangle. \end{aligned} \tag{2.4}$$

Similarly, we also have

$$\langle \mathcal{A}_q(u), u^+ \rangle > \langle \mathcal{A}_q(u^+), u^+ \rangle \quad \text{and} \quad \langle \mathcal{A}_q(u), u^- \rangle > \langle \mathcal{A}_q(u^-), u^- \rangle. \tag{2.5}$$

Taking into account (2.3)–(2.5), we deduce that  $I_\lambda(u) > I_\lambda(u^+) + I_\lambda(u^-)$ . Analogously, one can prove (ii).  $\square$

The following Brézis-Lieb type Lemma will be very useful in this work, its proof is similar to Lemma 2.8 in [1] and we omit it here.



**Lemma 2.3.** Let  $\{u_n\} \subset X_V$  be a sequence such that  $u_n \rightharpoonup u$  in  $X_V$ . Set  $v_n = u_n - u$ , then we have:

- (i)  $[v_n]_{s,p}^p + [v_n]_{s,q}^q = ([u_n]_{s,p}^p + [u_n]_{s,q}^q) - ([u]_{s,p}^p + [u]_{s,q}^q) + o_n(1)$ ,
- (ii)  $\int_{\mathbb{R}^N} V(x) (|v_n|^p + |v_n|^q) dx = \int_{\mathbb{R}^N} V(x) (|u_n|^p + |u_n|^q) dx - \int_{\mathbb{R}^N} V(x) (|u|^p + |u|^q) dx + o_n(1)$ ,
- (iii)  $\int_{\mathbb{R}^N} (F(v_n) - F(u_n) + F(u)) dx = o_n(1)$ ,
- (iv)  $\sup_{\|w\| \leq 1} \int_{\mathbb{R}^N} |(f(v_n) - f(u_n) + f(u)) w| dx = o_n(1)$ .

### 3. Some technical lemmas

The purpose of this section is to prove some technical lemmas related to the existence of a least energy sign-changing solution. Firstly, we collect some preliminary lemmas which will be fundamental to prove our main results.

Now, fixed  $u \in X_V$  with  $u^\pm \neq 0$ , we define function  $\psi_u : [0, \infty) \times [0, \infty) \rightarrow \mathbb{R}$  and mapping  $T_u : [0, \infty) \times [0, \infty) \rightarrow \mathbb{R}^2$  by

$$\psi_u(\sigma, \tau) := I_\lambda(\sigma u^+ + \tau u^-)$$

and

$$T_u(\sigma, \tau) := (\langle I'_\lambda(\sigma u^+ + \tau u^-), \sigma u^+ \rangle, \langle I'_\lambda(\sigma u^+ + \tau u^-), \tau u^- \rangle).$$

**Lemma 3.1.** For any  $u \in X_V$  with  $u^\pm \neq 0$ , there exists a unique maximum point pair  $(\tau_u, \sigma_u)$  of the function  $\psi_u$  such that  $\tau_u u^+ + \sigma_u u^- \in \mathcal{M}_\lambda$ .

*Proof.* Our proof will be divided into three steps.

**Step 1:** For any  $u \in X_V$  with  $u^\pm \neq 0$ , in the following, we will prove the existence of  $\sigma_u$  and  $\tau_u$ . From  $(f_1)$ ,  $(f_2)$  and Lemma 2.2 we deduce that

$$\begin{aligned} \langle I'_\lambda(\sigma u^+ + \tau u^-), \sigma u^+ \rangle &\geq \langle I'_\lambda(\sigma u^+), \sigma u^+ \rangle \\ &= \sigma^p \|u^+\|_{V,p}^p + \sigma^q \|u^+\|_{V,q}^q - \lambda \int_{\mathbb{R}^N} f(\sigma u^+) \sigma u^+ dx - \sigma^{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx \\ &\geq \sigma^p \|u^+\|_{V,p}^p + \sigma^q \|u^+\|_{V,q}^q - \lambda \varepsilon \sigma^p \int_{\mathbb{R}^N} |u^+|^p dx \\ &\quad - \lambda C_\varepsilon \sigma^{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx - \sigma^{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx \\ &\geq (1 - \lambda C_\varepsilon) \sigma^p \|u^+\|_{V,p}^p + \sigma^q \|u^+\|_{V,q}^q - (\lambda C C_\varepsilon + C) \sigma^{q_s^*} \|u^+\|^{q_s^*}. \end{aligned} \quad (3.1)$$

Similarly, we have that

$$\begin{aligned} \langle I'_\lambda(\sigma u^+ + \tau u^-), \tau u^- \rangle &\geq \langle I'_\lambda(\tau u^-), \tau u^- \rangle \\ &\geq (1 - \lambda C_\varepsilon) \sigma^p \|u^-\|_{V,p}^p + \sigma^q \|u^-\|_{V,q}^q - (\lambda C C_\varepsilon + C) \sigma^{q_s^*} \|u^-\|^{q_s^*}. \end{aligned} \quad (3.2)$$

Choose  $\varepsilon > 0$  such that  $(1 - \lambda C_\varepsilon) > 0$ . Since  $p < q < q_s^*$ , there exists  $r > 0$  small enough such that

$$\langle I'_\lambda(r u^+ + \tau u^-), r u^+ \rangle > 0 \text{ for all } \tau > 0 \quad (3.3)$$

and

$$\langle I'_\lambda(\sigma u^+ + ru^-), ru^- \rangle > 0 \text{ for all } \sigma > 0. \quad (3.4)$$

On the other hand, by  $(f_3)$ , there exist  $D_1, D_2 > 0$  such that

$$F(t) \geq D_1 t^\theta - D_2 \text{ for } t > 0. \quad (3.5)$$

Then we have

$$\begin{aligned} & \langle I'(\sigma u^+ + \tau u^-), \sigma u^+ \rangle \\ & \leq \sigma^p \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_+} \frac{|u^+(x) - u^+(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma u^+(x) - \tau u^-(y)|^{p-1} \sigma u^+(x)}{|x - y|^{N+ps}} dx dy \\ & + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau u^-(x) - \sigma u^+(y)|^{p-1} \sigma u^+(y)}{|x - y|^{N+ps}} dx dy + \sigma^q \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_+} \frac{|u^+(x) - u^+(y)|^q}{|x - y|^{N+qs}} dx dy \\ & + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma u^+(x) - \tau u^-(y)|^{q-1} \sigma u^+(x)}{|x - y|^{N+qs}} dx dy \\ & + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau u^-(x) - \sigma u^+(y)|^{q-1} \sigma u^+(y)}{|x - y|^{N+qs}} dx dy \\ & + \sigma^p \int_{\mathbb{R}^N} V(x) |u^+|^p dx + \sigma^q \int_{\mathbb{R}^N} V(x) |u^+|^q dx - \lambda D_1 \sigma^\theta \int_{A^+} |u^+|^\theta dx + \lambda D_2 |A^+|, \end{aligned}$$

where  $A^+ \subset \text{supp}(u^+)$  is measurable set with finite and positive measure  $|A^+|$ . Due to the fact  $\theta > p$ , for  $R$  sufficiently large, we get

$$\langle I'_\lambda(Ru^+ + \tau u^-), Ru^+ \rangle < 0 \text{ for all } \tau \in [r, R]. \quad (3.6)$$

Similarly, we get

$$\langle I'_\lambda(\sigma u^+ + Ru^-), Ru^- \rangle < 0 \text{ for all } \sigma \in [r, R]. \quad (3.7)$$

Hence, by virtue of Miranda's Theorem [31], and taking (3.3), (3.4), (3.6) and (3.7) into account, we can see that there exists  $(\sigma_u, \tau_u) \in [r, R] \times [r, R]$  such that  $T_u(\sigma, \tau) = (0, 0)$ , i.e.,  $\sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda$ .

**Step 2:** Now we prove the uniqueness of the pair  $(\sigma_u, \tau_u)$ .

**Case 1:**  $u \in \mathcal{M}_\lambda$ .

If  $u \in \mathcal{M}_\lambda$ , we have that

$$\begin{aligned} & \|u^+\|_{V,p}^p + \|u^+\|_{V,q}^q - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^p}{|x - y|^{N+ps}} dx dy - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^p}{|x - y|^{N+ps}} dx dy \\ & - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^q}{|x - y|^{N+qs}} dx dy - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^q}{|x - y|^{N+qs}} dx dy \\ & + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} u^+(x)}{|x - y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} u^+(y)}{|x - y|^{N+ps}} dx dy \quad (3.8) \\ & + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{q-1} u^+(x)}{|x - y|^{N+qs}} dx dy + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{q-1} u^+(y)}{|x - y|^{N+qs}} dx dy \\ & = \lambda \int_{\mathbb{R}^N} f(u^+) u^+ dx + \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx \end{aligned}$$

and

$$\begin{aligned}
& \|u^-\|_{V,p}^p + \|u^-\|_{V,q}^q - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^p}{|x-y|^{N+ps}} dx dy - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^p}{|x-y|^{N+ps}} dx dy \\
& - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^q}{|x-y|^{N+qs}} dx dy - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^q}{|x-y|^{N+qs}} dx dy \\
& + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_-} \frac{|u^-(x) - u^+(y)|^{p-1} (-u^-(x))}{|x-y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} (-u^-(y))}{|x-y|^{N+ps}} dx dy \\
& + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_-} \frac{|u^-(x) - u^+(y)|^{q-1} (-u^-(x))}{|x-y|^{N+qs}} dx dy + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{q-1} (-u^-(y))}{|x-y|^{N+qs}} dx dy \\
& = \lambda \int_{\mathbb{R}^N} f(u^-) u^- dx + \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx.
\end{aligned} \tag{3.9}$$

We will show that  $(\sigma_u, \tau_u) = (1, 1)$  is the unique pair of numbers such that  $\sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda$ . Let  $(\sigma_u, \tau_u)$  be a pair of numbers such that  $\sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda$  with  $0 < \sigma_u \leq \tau_u$ , then one can see

$$\begin{aligned}
& \sigma_u^p \|u^+\|_{V,p}^p + \sigma_u^q \|u^+\|_{V,q}^q - \sigma_u^p \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^p}{|x-y|^{N+ps}} dx dy - \sigma_u^p \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^p}{|x-y|^{N+ps}} dx dy \\
& - \sigma_u^q \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^q}{|x-y|^{N+qs}} dx dy - \sigma_u^q \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^q}{|x-y|^{N+qs}} dx dy \\
& + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma_u u^+(x) - \tau_u u^-(y)|^{p-1} \sigma_u u^+(x)}{|x-y|^{N+ps}} dx dy \\
& + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau_u u^-(x) - \sigma_u u^+(y)|^{p-1} \sigma_u u^+(y)}{|x-y|^{N+ps}} dx dy \\
& + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma_u u^+(x) - \tau_u u^-(y)|^{q-1} \sigma_u u^+(x)}{|x-y|^{N+qs}} dx dy \\
& + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau_u u^-(x) - \sigma_u u^+(y)|^{q-1} \sigma_u u^+(y)}{|x-y|^{N+qs}} dx dy \\
& = \lambda \int_{\mathbb{R}^N} f(\sigma_u u^+) \sigma_u u^+ dx + \sigma_u^{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx
\end{aligned} \tag{3.10}$$

and

$$\begin{aligned}
& \tau_u^p \|u^-\|_{V,p}^p + \tau_u^q \|u^-\|_{V,q}^q - \tau_u^p \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^p}{|x-y|^{N+ps}} dx dy - \tau_u^p \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^p}{|x-y|^{N+ps}} dx dy \\
& - \tau_u^q \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^q}{|x-y|^{N+qs}} dx dy - \tau_u^q \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^q}{|x-y|^{N+qs}} dx dy \\
& + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau_u u^-(x) - \sigma_u u^+(y)|^{p-1} (-\tau_u u^-(x))}{|x-y|^{N+ps}} dx dy
\end{aligned}$$

$$\begin{aligned}
& + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma_u u^+(x) - \tau_u u^-(y)|^{p-1} (-\tau_u u^-(y))}{|x-y|^{N+ps}} dx dy \\
& + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau_u u^-(x) - \sigma_u u^+(y)|^{q-1} (-\tau_u u^-(x))}{|x-y|^{N+qs}} dx dy \\
& + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma_u u^+(x) - \tau_u u^-(y)|^{q-1} (-\tau_u u^-(y))}{|x-y|^{N+qs}} dx dy \\
& = \lambda \int_{\mathbb{R}^N} f(\tau_u u^-) \tau_u u^- dx + \tau_u^{q_s^*} \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx.
\end{aligned} \tag{3.11}$$

Since  $0 < \sigma_u \leq \tau_u$ , it follows from (3.11) that

$$\begin{aligned}
& \tau_u^{p-q} \|u^-\|_{V,p}^p + \|u^-\|_{V,q}^q \\
& - \tau_u^{p-q} \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^p}{|x-y|^{N+ps}} dx dy - \tau_u^{p-q} \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^p}{|x-y|^{N+ps}} dx dy \\
& - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^q}{|x-y|^{N+qs}} dx dy - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^q}{|x-y|^{N+qs}} dx dy \\
& + \tau_u^{p-q} \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} (-u^-(x))}{|x-y|^{N+ps}} dx dy \\
& + \tau_u^{p-q} \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} (-u^-(y))}{|x-y|^{N+ps}} dx dy \\
& + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{q-1} (-u^-(x))}{|x-y|^{N+qs}} dx dy \\
& + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{q-1} (-u^-(y))}{|x-y|^{N+qs}} dx dy \\
& \geq \lambda \int_{\mathbb{R}^N} \frac{f(\tau_u u^-) \tau_u u^-}{\tau_u^q} dx + \tau_u^{q_s^*-q} \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx.
\end{aligned} \tag{3.12}$$

If  $\tau_u > 1$ , by (3.9) and (3.12), we get

$$\begin{aligned}
& (\tau_u^{p-q} - 1) \left( \|u^-\|_{V,p}^p - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x)|^p}{|x-y|^{N+ps}} dx dy - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^-(y)|^p}{|x-y|^{N+ps}} dx dy \right) \\
& + (\tau_u^{p-q} - 1) \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} (-u^-(x))}{|x-y|^{N+ps}} dx dy \\
& + (\tau_u^{p-q} - 1) \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} (-u^-(y))}{|x-y|^{N+ps}} dx dy \\
& \geq \lambda \int_{\mathbb{R}^N} \left( \frac{f(\tau_u u^-)}{|\tau_u u^-|^{q-1}} - \frac{f(u^-)}{|u^-|^{q-1}} \right) |u^-|^q dx + (\tau_u^{q_s^*-q} - 1) \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx.
\end{aligned}$$

The left side of the above inequality is negative, which is absurd because the right side is positive. Therefore, we conclude that  $0 < \sigma_u \leq \tau_u \leq 1$ .

Similarly, by (3.10) and  $0 < \sigma_u \leq \tau_u$ , we have that

$$\begin{aligned} & (\sigma_u^{p-q} - 1) \left( \|u^+\|_{V,p}^p - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^p}{|x-y|^{N+ps}} dx dy - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^p}{|x-y|^{N+ps}} dx dy \right) \\ & + (\sigma_u^{p-q} - 1) \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} u^+(x)}{|x-y|^{N+ps}} dx dy \\ & + (\sigma_u^{p-q} - 1) \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} u^+(y)}{|x-y|^{N+ps}} dx dy \\ & \leq \lambda \int_{\mathbb{R}^N} \left( \frac{f(\sigma_u u^+)}{|\sigma_u u^+|^{q-1}} - \frac{f(u^+)}{|u^+|^{q-1}} \right) |u^+|^q dx + (\sigma_u^{q_s^* - q} - 1) \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx. \end{aligned}$$

This fact implies that  $\sigma_u \geq 1$ . Consequently,  $\sigma_u = \tau_u = 1$ .

**Case 2:**  $u \notin \mathcal{M}_\lambda$ .

Suppose that there exist  $(\tilde{\sigma}_1, \tilde{\tau}_1)$ ,  $(\tilde{\sigma}_2, \tilde{\tau}_2)$  such that

$$u_1 := \tilde{\sigma}_1 u^+ + \tilde{\tau}_1 u^- \in \mathcal{M}_\lambda \quad \text{and} \quad u_2 := \tilde{\sigma}_2 u^+ + \tilde{\tau}_2 u^- \in \mathcal{M}_\lambda.$$

Hence,

$$u_2 = \left( \frac{\tilde{\sigma}_2}{\tilde{\sigma}_1} \right) \tilde{\sigma}_1 u^+ + \left( \frac{\tilde{\tau}_2}{\tilde{\tau}_1} \right) \tilde{\tau}_1 u^- = \left( \frac{\tilde{\sigma}_2}{\tilde{\sigma}_1} \right) u_1^+ + \left( \frac{\tilde{\tau}_2}{\tilde{\tau}_1} \right) u_1^- \in \mathcal{M}_\lambda.$$

Since  $u_1 \in \mathcal{M}_\lambda$ , we deduce from case 1 that

$$\frac{\tilde{\sigma}_2}{\tilde{\sigma}_1} = \frac{\tilde{\tau}_2}{\tilde{\tau}_1} = 1,$$

which implies  $\tilde{\sigma}_1 = \tilde{\sigma}_2$ ,  $\tilde{\tau}_1 = \tilde{\tau}_2$ .

**Step 3:** We assert that  $(\sigma_u, \tau_u)$  is the unique maximum point of  $\psi_u$  on  $[0, +\infty) \times [0, +\infty)$ . In fact, by  $(f_3)$  we can see that

$$\begin{aligned} I_\lambda(\sigma u^+ + \tau u^-) &= \frac{1}{p} \|\sigma u^+ + \tau u^-\|_{V,p}^p + \frac{1}{q} \|\sigma u^+ + \tau u^-\|_{V,q}^q - \lambda \int_{\mathbb{R}^N} F(\sigma u^+ + \tau u^-) dx \\ &\quad - \frac{1}{q_s^*} \int_{\mathbb{R}^N} |\sigma u^+ + \tau u^-|^{q_s^*} dx \\ &\leq \frac{1}{p} \|\sigma u^+ + \tau u^-\|_{V,p}^p + \frac{1}{p} \|\sigma u^+ + \tau u^-\|_{V,q}^q - \frac{\sigma^{q_s^*}}{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx - \frac{\tau^{q_s^*}}{q_s^*} \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx, \end{aligned}$$

which implies that  $\lim_{|\sigma, \tau| \rightarrow \infty} \phi_u(\sigma, \tau) = -\infty$  due to  $q_s^* > q$ . Noticing that  $\sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda$ , we conclude that  $(\sigma_u, \tau_u)$  is the unique critical point of  $\psi_u$  in  $(0, +\infty) \times (0, +\infty)$ . Hence, it is sufficient to check that a maximum point cannot be achieved on the boundary of  $[0, +\infty) \times [0, +\infty)$ . By contradiction, we assume that  $(0, \tau_1)$  is a maximum point of  $\psi_u$  with  $\tau_1 \geq 0$ . Then, arguing as

Lemma 2.2, we have

$$\begin{aligned}
 \psi_u(\sigma, \tau_1) &= \frac{1}{p} \|\sigma u^+ + \tau_1 u^-\|_{V,p}^p + \frac{1}{q} \|\sigma u^+ + \tau_1 u^-\|_{V,q}^q - \lambda \int_{\mathbb{R}^N} F(\sigma u^+) dx \\
 &\quad - \lambda \int_{\mathbb{R}^N} F(\tau_1 u^-) dx - \frac{\sigma^{q_s^*}}{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx - \frac{\tau_1^{q_s^*}}{q_s^*} \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx \\
 &> \frac{\sigma^p}{p} \|u^+\|_{V,p}^p + \frac{\sigma^q}{q} \|u^+\|_{V,q}^q - \lambda \int_{\mathbb{R}^N} F(\sigma u^+) dx - \frac{\sigma^{q_s^*}}{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx \\
 &\quad + \frac{\tau_1^p}{p} \|u^-\|_{V,p}^p + \frac{\tau_1^q}{q} \|u^-\|_{V,q}^q - \lambda \int_{\mathbb{R}^N} F(\tau_1 u^-) dx - \frac{\tau_1^{q_s^*}}{q_s^*} \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx \\
 &= \psi_u(0, \tau_1) + \psi_u(\sigma, 0).
 \end{aligned} \tag{3.13}$$

On the other hand, by the growth condition  $(f_1)$  and  $(f_2)$ , one can easily check that  $\psi_u(\sigma, 0) > 0$  for  $\sigma$  sufficiently small. Combining this with (3.13), we see that

$$\psi_u(0, \tau_1) < \psi_u(0, \tau_1) + \psi_u(\sigma, 0) < \psi_u(\sigma, \tau_1)$$

if  $\sigma$  is small enough, which yields a contradiction. Similarly,  $\psi_u$  can not achieve its global maximum point at  $(\sigma_1, 0)$ , where  $\sigma_1 \geq 0$ . As a result, we complete the proof of Lemma 3.1.  $\square$

**Lemma 3.2.** For any  $u \in X_V$  with  $u^\pm \neq 0$ , such that  $\langle I'_\lambda(u), u^\pm \rangle \leq 0$ , the unique maximum point of  $\psi_u$  in  $[0, +\infty) \times [0, +\infty)$  satisfies  $0 < \sigma_u, \tau_u \leq 1$ .

*Proof.* If  $\sigma_u = 0$  or  $\tau_u = 0$ , according Lemma 3.1,  $\psi_u$  can not achieve maximum. Without loss of generality, we assume  $\sigma_u \geq \tau_u > 0$ . Since  $\sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda$ , there holds

$$\begin{aligned}
 &\sigma_u^p \|u^+\|_{V,p}^p + \sigma_u^q \|u^+\|_{V,q}^q - \sigma_u^p \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^p}{|x-y|^{N+ps}} dx dy - \sigma_u^p \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^p}{|x-y|^{N+ps}} dx dy \\
 &\quad - \sigma_u^q \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^q}{|x-y|^{N+qs}} dx dy - \sigma_u^q \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^q}{|x-y|^{N+qs}} dx dy \\
 &\quad + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma_u u^+(x) - \tau_u u^-(y)|^{p-1} \sigma_u u^+(x)}{|x-y|^{N+ps}} dx dy \\
 &\quad + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau_u u^-(x) - \sigma_u u^+(y)|^{p-1} \sigma_u u^+(y)}{|x-y|^{N+ps}} dx dy \\
 &\quad + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|\sigma_u u^+(x) - \tau_u u^-(y)|^{q-1} \sigma_u u^+(x)}{|x-y|^{N+qs}} dx dy \\
 &\quad + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|\tau_u u^-(x) - \sigma_u u^+(y)|^{q-1} \sigma_u u^+(y)}{|x-y|^{N+qs}} dx dy \\
 &= \lambda \int_{\mathbb{R}^N} f(\sigma_u u^+) \sigma_u u^+ dx + \sigma_u^{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx.
 \end{aligned} \tag{3.14}$$

On the other hand, by  $\langle I'_\lambda(u), u^+ \rangle \leq 0$ , we have

$$\begin{aligned}
& \|u^+\|_{V,p}^p + \|u^+\|_{V,q}^q - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^p}{|x-y|^{N+ps}} dx dy - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^p}{|x-y|^{N+ps}} dx dy \\
& - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^q}{|x-y|^{N+qs}} dx dy - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^q}{|x-y|^{N+qs}} dx dy \\
& + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} u^+(x)}{|x-y|^{N+ps}} dx dy + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} u^+(y)}{|x-y|^{N+ps}} dx dy \quad (3.15) \\
& + \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{q-1} u^+(x)}{|x-y|^{N+qs}} dx dy + \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{q-1} u^+(y)}{|x-y|^{N+qs}} dx dy \\
& \leq \lambda \int_{\mathbb{R}^N} f(u^+) u^+ dx + \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx.
\end{aligned}$$

Then it follows (3.14) and (3.15) that

$$\begin{aligned}
& (\sigma_u^{p-q} - 1) \left( \|u^+\|_{V,p}^p - \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x)|^p}{|x-y|^{N+ps}} dx dy - \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^+(y)|^p}{|x-y|^{N+ps}} dx dy \right) \\
& + (\sigma_u^{p-q} - 1) \int_{(\mathbb{R}^N)_+ \times (\mathbb{R}^N)_-} \frac{|u^+(x) - u^-(y)|^{p-1} u^+(x)}{|x-y|^{N+ps}} dx dy \\
& + (\sigma_u^{p-q} - 1) \int_{(\mathbb{R}^N)_- \times (\mathbb{R}^N)_+} \frac{|u^-(x) - u^+(y)|^{p-1} u^+(y)}{|x-y|^{N+ps}} dx dy \quad (3.16) \\
& \geq \lambda \int_{\mathbb{R}^N} \left( \frac{f(\sigma_u u^+)}{|\sigma_u u^+|^{q-1}} - \frac{f(u^+)}{|u^+|^{q-1}} \right) |u^+|^q dx + (\sigma_u^{q_s^*-q} - 1) \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx.
\end{aligned}$$

In view of  $(f_4)$ , we conclude that  $\sigma_u \leq 1$ . Thus, we have that  $0 < \sigma_u, \tau_u \leq 1$ .  $\square$

**Lemma 3.3.** *There exists  $\rho > 0$  such that  $\|u^\pm\| \geq \rho$  for all  $u \in \mathcal{M}_\lambda$ .*

*Proof.* For any  $u \in \mathcal{M}_\lambda$ , by  $(f_1)$ ,  $(f_2)$  and the Sobolev inequalities, we have that

$$\begin{aligned}
\|u^\pm\|_{V,p}^p + \|u^\pm\|_{V,q}^q & \leq \lambda \int_{\mathbb{R}^N} f(u^\pm) u^\pm dx + \int_{\mathbb{R}^N} |u^\pm|^{q_s^*} dx \\
& \leq \lambda \varepsilon C_1 \|u^\pm\|_{V,p}^p + \lambda C_2 C_\varepsilon \|u^\pm\|^{q_s^*} + C_3 \|u^\pm\|^{q_s^*}.
\end{aligned}$$

Thus we get

$$C'_0 \|u\|_{V,p}^p + \|u\|_{V,q}^q \leq \widetilde{C}_2 \|u\|^{q_s^*}, \quad (3.17)$$

where  $C'_0 = (1 - \lambda \varepsilon C_1)$ ,  $\widetilde{C}_2 = (C_3 + \lambda C_2 C_\varepsilon)$  with  $C$  is a Sobolev embedding constant. If  $0 < \|u\| < 1$ , then  $\|u\|_{V,p}, \|u\|_{V,q} < 1$  and by order relations between  $p$  and  $q$  and by (3.17) we have

$$\begin{aligned}
C'' \|u\|^q & \leq C'' \left( \|u\|_{V,p} + \|u\|_{V,q} \right)^q \leq C' \left( \|u\|_{V,p}^q + \|u\|_{V,q}^q \right) \\
& \leq C'_0 \|u\|_{V,p}^p + \|u\|_{V,q}^q \leq \widetilde{C}_2 \|u\|^{q_s^*},
\end{aligned}$$

where  $C' = \min \{C'_0, 1\}$  and  $C'' = \frac{C'}{2^{q-1}}$ . Hence, there exists a positive radius  $\rho_1 > 0$  such that  $\|u\| \geq \rho_1$  with  $\rho_1 = \left( \frac{C''}{C_\varepsilon} \right)^{\frac{1}{q_s^*-q}}$ . Clearly we can reason analogously if  $\|u\| \geq 1$  so that for some  $\rho > 0$  and for every  $u \in \mathcal{M}_\lambda$ , we get  $\rho \leq \|u\|$ .  $\square$

**Lemma 3.4.** Let  $c_\lambda = \inf_{u \in \mathcal{M}_\lambda} I_\lambda(u)$ , then we have that  $\lim_{\lambda \rightarrow \infty} c_\lambda = 0$ .

*Proof.* Since  $u \in \mathcal{M}_\lambda$ , we have  $\langle I'_\lambda(u), u \rangle = 0$  and then

$$\begin{aligned} I_\lambda(u) &= I_\lambda(u) - \frac{1}{\theta} \langle I'_\lambda(u), u \rangle \\ &\geq \left( \frac{1}{p} - \frac{1}{\theta} \right) \|u\|_{V,p}^p + \left( \frac{1}{q} - \frac{1}{\theta} \right) \|u\|_{V,q}^q, \end{aligned} \quad (3.18)$$

thus  $I_\lambda$  is bounded below on  $\mathcal{M}_\lambda$ , which implies  $c_\lambda$  is well-defined.

For any  $u \in X_V$  with  $u^\pm \neq 0$ , by Lemma 3.1, for each  $\lambda > 0$ , there exists  $\sigma_\lambda, \tau_\lambda$  such that  $\sigma_\lambda u^+ + \tau_\lambda u^- \in \mathcal{M}_\lambda$ , we have

$$\begin{aligned} 0 \leq c_\lambda &= \inf I_\lambda(u) \leq I_\lambda(\sigma_\lambda u^+ + \tau_\lambda u^-) \\ &\leq \frac{1}{p} \|\sigma_\lambda u^+ + \tau_\lambda u^-\|_{V,p}^p + \frac{1}{q} \|\sigma_\lambda u^+ + \tau_\lambda u^-\|_{V,q}^q - \int_{\mathbb{R}^N} F(\sigma_\lambda u^+ + \tau_\lambda u^-) dx \\ &\quad - \frac{1}{q_s^*} \int_{\mathbb{R}^N} |\sigma_\lambda u^+ + \tau_\lambda u^-|^{q_s^*} dx \\ &\leq \frac{2^{p-1}}{p} \sigma_\lambda^p \|u^+\|_{V,p}^p + \frac{2^{p-1}}{p} \tau_\lambda^p \|u^-\|_{V,p}^p + \frac{2^{q-1}}{q} \sigma_\lambda^q \|u^+\|_{V,q}^q + \frac{2^{q-1}}{q} \tau_\lambda^q \|u^-\|_{V,q}^q. \end{aligned}$$

Next, we will prove that  $\sigma_\lambda \rightarrow 0$  and  $\tau_\lambda \rightarrow 0$  as  $\lambda \rightarrow \infty$ .

Let  $Q_u = \{(\sigma_\lambda, \tau_\lambda) \in [0, +\infty) \times [0, +\infty) : T_u(\sigma_\lambda, \tau_\lambda) = (0, 0), \lambda > 0\}$ . Due to  $\sigma_\lambda u^+ + \tau_\lambda u^- \in \mathcal{M}_\lambda$ , there holds

$$\begin{aligned} &\sigma_\lambda^{q_s^*} \int_{\mathbb{R}^N} |u^+|^{q_s^*} dx + \tau_\lambda^{q_s^*} \int_{\mathbb{R}^N} |u^-|^{q_s^*} dx + \lambda \int_{\mathbb{R}^N} f(\sigma_\lambda u^+) (\sigma_\lambda u^+) dx + \lambda \int_{\mathbb{R}^N} f(\tau_\lambda u^-) (\tau_\lambda u^-) dx \\ &= \|\sigma_\lambda u^+ + \tau_\lambda u^-\|_{V,p}^p + \|\sigma_\lambda u^+ + \tau_\lambda u^-\|_{V,q}^q \\ &\leq 2^{p-1} \sigma_\lambda^p \|u^+\|_{V,p}^p + 2^{p-1} \tau_\lambda^p \|u^-\|_{V,p}^p + 2^{q-1} \sigma_\lambda^q \|u^+\|_{V,q}^q + 2^{q-1} \tau_\lambda^q \|u^-\|_{V,q}^q. \end{aligned}$$

Therefore,  $Q_u$  is bounded in  $\mathbb{R}^2$ . Let  $\{\lambda_n\} \subset (0, \infty)$  be such that  $\lambda_n \rightarrow \infty$  as  $n \rightarrow \infty$ . Then there exist  $\sigma_0$  and  $\tau_0$  such that  $(\sigma_{\lambda_n}, \tau_{\lambda_n}) \rightarrow (\sigma_0, \tau_0)$  as  $n \rightarrow \infty$ .

Now, we claim  $\sigma_0 = \tau_0 = 0$ . By contradiction, suppose that  $\sigma_0 > 0$  or  $\tau_0 > 0$  by  $\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^- \in \mathcal{M}_{\lambda_n}$ , then for any  $n \in \mathbb{N}$ , there holds

$$\begin{aligned} &\|\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-\|_{V,p}^p + \|\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-\|_{V,q}^q \\ &= \lambda_n \int_{\mathbb{R}^N} f(\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-) (\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-) dx + \int_{\mathbb{R}^N} |\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-|^{q_s^*} dx. \end{aligned} \quad (3.19)$$

Thanks to  $\sigma_{\lambda_n} u^+ \rightarrow \sigma_0 u^+$  and  $\tau_{\lambda_n} u^- \rightarrow \tau_0 u^-$  in  $X_V$ ,  $(f_1)$ ,  $(f_2)$  and the Lebesgue dominated convergence theorem, we deduce that

$$\int_{\mathbb{R}^N} f(\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-) (\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-) dx \rightarrow \int_{\mathbb{R}^N} f(\sigma_0 u^+ + \tau_0 u^-) (\sigma_0 u^+ + \tau_0 u^-) dx > 0 \quad (3.20)$$

as  $n \rightarrow \infty$ . It follows from  $\lambda_n \rightarrow \infty$  and (3.20) that the right hand side of (3.19) tends to infinity, which contradicts with the boundedness of  $\{\sigma_{\lambda_n} u^+ + \tau_{\lambda_n} u^-\}$  in  $X_V$ . Hence,  $\sigma_0 = \tau_0 = 0$ . As a result, we conclude that  $\lim_{\lambda \rightarrow \infty} c_\lambda = 0$ .  $\square$



**Lemma 3.5.** *There exists  $\lambda^* > 0$  such that for all  $\lambda \geq \lambda^*$ , the infimum  $c_\lambda$  is achieved.*

*Proof.* By the definition of  $c_\lambda = \inf_{u \in \mathcal{M}_\lambda} I_\lambda(u)$ , there exists a sequence  $\{u_n\} \subset \mathcal{M}_\lambda$  such that

$$\lim_{n \rightarrow \infty} I_\lambda(u_n) = c_\lambda.$$

Obviously,  $\{u_n\}$  is bounded in  $X_V$ . Up to a subsequence, still denoted by  $\{u_n\}$ , there exists  $u \in X_V$  such that  $u_n \rightharpoonup u$  weakly in  $X_V$ . Since the embedding  $X_V \hookrightarrow L^r(\mathbb{R}^N)$  is compact for all  $r \in [p, q_s^*]$ , we have  $u_n^\pm \rightarrow u^\pm$  in  $L^r(\mathbb{R}^N)$  for all  $r \in [p, q_s^*]$ ,  $u_n^\pm(x) \rightarrow u^\pm(x)$  a.e.  $x \in \mathbb{R}^N$ .

Denote  $\delta := \frac{s}{N} S \frac{N}{q}$ , according to Lemma 3.4, there is  $\lambda^* > 0$  such that  $c_\lambda < \delta$  for all  $\lambda \geq \lambda^*$ . Fix  $\lambda \geq \lambda^*$ , it follows from Lemma 3.1 that  $I_\lambda(\sigma u_n^+ + \tau u_n^-) \leq I_\lambda(u_n)$  for all  $\sigma, \tau \geq 0$ . Then by using Brézis-Lieb type Lemma 2.3 and the Fatou's Lemma, it follows that

$$\begin{aligned} & \liminf_{n \rightarrow \infty} I_\lambda(\sigma u_n^+ + \tau u_n^-) \\ &= \liminf_{n \rightarrow \infty} \left( \frac{1}{p} \|\sigma u_n^+ + \tau u_n^-\|_{V,p}^p + \frac{1}{q} \|\sigma u_n^+ + \tau u_n^-\|_{V,q}^q - \frac{1}{q_s^*} |\sigma u_n^+ + \tau u_n^-|_{q_s^*}^{q_s^*} \right) - \lambda \int_{\mathbb{R}^N} F(\sigma u_n^+ + \tau u_n^-) dx \\ &= \liminf_{n \rightarrow \infty} \left( \frac{1}{p} \|\sigma u_n^+ + \tau u_n^- - (\sigma u^+ + \tau u^-)\|_{V,p}^p + \frac{1}{q} \|\sigma u_n^+ + \tau u_n^- - (\sigma u^+ + \tau u^-)\|_{V,q}^q \right) \\ &\quad - \frac{\sigma^{q_s^*}}{q_s^*} \lim_{n \rightarrow \infty} |u_n^+ - u^+|_{q_s^*}^{q_s^*} - \frac{\tau^{q_s^*}}{q_s^*} \lim_{n \rightarrow \infty} |u_n^- - u^-|_{q_s^*}^{q_s^*} - \frac{1}{q_s^*} |\sigma u^+ + \tau u^-|_{q_s^*}^{q_s^*} \\ &\quad + \frac{1}{p} \|\sigma u^+ + \tau u^-\|_{V,p}^p + \frac{1}{q} \|\sigma u^+ + \tau u^-\|_{V,q}^q - \lambda \int_{\mathbb{R}^N} F(\sigma u^+ + \tau u^-) dx \\ &= I_\lambda(\sigma u^+ + \tau u^-) + \lim_{n \rightarrow \infty} \left( \frac{1}{p} \|\sigma u_n^+ - \sigma u^+\|_{V,p}^p + \frac{1}{p} \|\tau u_n^- - \tau u^-\|_{V,p}^p \right) \\ &\quad + \liminf_{n \rightarrow \infty} \left( \frac{1}{p} \|\sigma u_n^+ + \tau u_n^- - (\sigma u^+ + \tau u^-)\|_{V,p}^p - \frac{1}{p} \|\sigma u_n^+ - \sigma u^+\|_{V,p}^p - \frac{1}{p} \|\tau u_n^- - \tau u^-\|_{V,p}^p \right) \\ &\quad + \lim_{n \rightarrow \infty} \left( \frac{1}{q} \|\sigma u_n^+ - \sigma u^+\|_{V,q}^q + \frac{1}{p} \|\tau u_n^- - \tau u^-\|_{V,q}^q \right) \\ &\quad + \liminf_{n \rightarrow \infty} \left( \frac{1}{q} \|\sigma u_n^+ + \tau u_n^- - (\sigma u^+ + \tau u^-)\|_{V,q}^q - \frac{1}{q} \|\sigma u_n^+ - \sigma u^+\|_{V,q}^q - \frac{1}{q} \|\tau u_n^- - \tau u^-\|_{V,q}^q \right) \\ &\quad - \frac{\sigma^{q_s^*}}{q_s^*} \lim_{n \rightarrow \infty} |u_n^+ - u^+|_{q_s^*}^{q_s^*} - \frac{\tau^{q_s^*}}{q_s^*} \lim_{n \rightarrow \infty} |u_n^- - u^-|_{q_s^*}^{q_s^*} \\ &\geq I_\lambda(\sigma u^+ + \tau u^-) + \frac{1}{p} \sigma^p A_1 + \frac{1}{q} \sigma^q A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 + \frac{1}{p} \tau^p A_2 + \frac{1}{q} \tau^q A_4 - \frac{\tau^{q_s^*}}{q_s^*} B_2, \end{aligned}$$

where

$$\begin{aligned} A_1 &= \lim_{n \rightarrow \infty} \|\sigma u_n^+ - \sigma u^+\|_{V,p}^p, & A_2 &= \lim_{n \rightarrow \infty} \|\tau u_n^- - \tau u^-\|_{V,p}^p, & A_3 &= \lim_{n \rightarrow \infty} \|\sigma u_n^+ - \sigma u^+\|_{V,q}^q, \\ A_4 &= \lim_{n \rightarrow \infty} \|\tau u_n^- - \tau u^-\|_{V,q}^q, & B_1 &= \lim_{n \rightarrow \infty} |u_n^+ - u^+|_{q_s^*}^{q_s^*}, & B_2 &= \lim_{n \rightarrow \infty} |u_n^- - u^-|_{q_s^*}^{q_s^*}. \end{aligned}$$

Hence, we can see that for all  $\sigma \geq 0$  and  $\tau \geq 0$ , there holds

$$c_\lambda \geq I_\lambda(\sigma u^+ + \tau u^-) + \frac{1}{p} \sigma^p A_1 + \frac{1}{q} \sigma^q A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 + \frac{1}{p} \tau^p A_2 + \frac{1}{q} \tau^q A_4 - \frac{\tau^{q_s^*}}{q_s^*} B_2. \quad (3.21)$$

Now we divide the proof into three steps.

**Step 1:** We prove that  $u^\pm \neq 0$ . Here we only prove  $u^+ \neq 0$  since  $u^- = 0$  is similar, by contradiction, we suppose  $u^+ = 0$ . Then we have the following two cases.

**Case 1:**  $B_1 = 0$ . If  $A_1 = A_3 = 0$ , that is,  $u_n^+ \rightarrow u^+$  in  $X_V$ . According to Lemma 3.3, we obtain  $\|u^+\| > 0$ , which contradicts  $u^+ = 0$ . If  $A_1$  or  $A_3 > 0$ , By (3.21) we get  $\frac{1}{p}\sigma^p A_1 + \frac{\sigma^q}{q} A_3 < c_\lambda$  for all  $\sigma \geq 0$ , which is a contradiction.

**Case 2:**  $B_1 > 0$ . According to definition of  $S_q$ , we have that  $\delta := \frac{s}{N} S_q^{\frac{N}{sq}} \leq \frac{s}{N} \left( \frac{A_3}{(B_1)^{\frac{q}{q_s^*}}} \right)^{\frac{N}{sq}}$ , by direct calculation, we have that

$$\frac{s}{N} \left( \frac{A_3}{(B_1)^{\frac{q}{q_s^*}}} \right)^{\frac{N}{sq}} = \max_{\sigma \geq 0} \left\{ \frac{\sigma^q}{q} A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 \right\} \leq \max_{\sigma \geq 0} \left\{ \frac{\sigma^p}{p} A_1 + \frac{\sigma^q}{q} A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 \right\}.$$

Since  $c_\lambda \rightarrow 0$  as  $\lambda \rightarrow \infty$ , there exists  $\lambda^* > 0$  such that for all  $\lambda > \lambda^*$ ,  $c_\lambda \leq \delta$ . Then, without loss of generality, we can assume  $c_\lambda < \delta$ . Choosing  $\tau = 0$ , by (3.21) it follows that

$$\delta \leq \max_{\sigma \geq 0} \left\{ \frac{\sigma^q}{q} A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 \right\} \leq \max_{\sigma \geq 0} \left\{ \frac{\sigma^p}{p} A_1 + \frac{\sigma^q}{q} A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 \right\} < \delta,$$

which is impossible. From the above discussion, we have that  $u^+ \neq 0$ . Similarly, we obtain  $u^- \neq 0$ .

**Step 2:** we prove that  $B_1 = 0$ ,  $B_2 = 0$ . We just prove  $B_1 = 0$  (the proof of  $B_2 = 0$  is analogous). By contradiction, we suppose that  $B_1 > 0$ .

**Case 1:**  $B_2 > 0$ , Let  $\widehat{\sigma}_1$  and  $\widehat{\tau}_1$  satisfy

$$\left\{ \frac{\widehat{\sigma}_1^p}{p} A_1 + \frac{\widehat{\sigma}_1^q}{q} A_3 - \frac{\widehat{\sigma}_1^{q_s^*}}{q_s^*} B_1 \right\} = \max_{\sigma \geq 0} \left\{ \frac{\sigma^p}{p} A_1 + \frac{\sigma^q}{q} A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 \right\}$$

and

$$\left\{ \frac{\widehat{\tau}_1^p}{p} A_2 + \frac{\widehat{\tau}_1^q}{q} A_4 - \frac{\widehat{\tau}_1^{q_s^*}}{q_s^*} B_2 \right\} = \max_{\tau \geq 0} \left\{ \frac{\tau^p}{p} A_2 + \frac{\tau^q}{q} A_4 - \frac{\tau^{q_s^*}}{q_s^*} B_2 \right\}.$$

According to  $[0, \widehat{\sigma}_1] \times [0, \widehat{\tau}_1]$  is compact, there exists  $(\sigma_u, \tau_u) \in [0, \widehat{\sigma}_1] \times [0, \widehat{\tau}_1]$  such that  $\psi_u(\sigma_u, \tau_u) = \max_{(\sigma, \tau) \in [0, \widehat{\sigma}_1] \times [0, \widehat{\tau}_1]} \psi_u(\sigma, \tau)$ .

In the following, we prove that  $(\sigma_u, \tau_u) \in (0, \widehat{\sigma}_1) \times (0, \widehat{\tau}_1)$ . Obviously, if  $\tau$  is small enough, we have

$$\psi_u(\sigma, 0) < I_\lambda(\sigma u^+) + I_\lambda(\tau u^-) \leq I_\lambda(\sigma u^+ + \tau u^-) = \psi_u(\sigma, \tau), \quad \forall \sigma \in [0, \widehat{\sigma}_1].$$

Hence, there exists  $\tau_0$  such that  $\psi_u(\sigma, 0) \leq \psi_u(\sigma, \tau_0)$ , for all  $\sigma \in [0, \widehat{\sigma}_1]$ . That is,  $(\sigma_u, \tau_u) \notin [0, \widehat{\sigma}_1] \times \{0\}$ . Similarly, one can prove that  $(\sigma_u, \tau_u) \notin \{0\} \times [0, \widehat{\tau}_1]$ .

On the other hand, we can easily deduce that

$$\frac{\sigma^p}{p} A_1 + \frac{\sigma^q}{q} A_3 - \frac{\sigma^{q_s^*}}{q_s^*} B_1 > 0, \quad \sigma \in (0, \widehat{\sigma}_1] \quad (3.22)$$

and

$$\frac{\tau^p}{p} A_2 + \frac{\tau^q}{q} A_4 - \frac{\tau^{q_s^*}}{q_s^*} B_2, \quad \tau \in (0, \widehat{\tau}_1]. \quad (3.23)$$

Then, for all  $\sigma \in (0, \widehat{\sigma}_1]$  and  $\tau \in (0, \widehat{\tau}_1]$ , we get

$$\begin{aligned}\delta &\leq \frac{\widehat{\sigma}_1^p}{p}A_1 + \frac{\widehat{\sigma}_1^q}{q}A_3 - \frac{\widehat{\sigma}_1^{q_s^*}}{q_s^*}B_1 + \frac{\tau^p}{p}A_2 + \frac{\tau^q}{q}A_4 - \frac{\tau^{q_s^*}}{q_s^*}B_2, \\ \delta &\leq \frac{\widehat{\tau}_1^p}{p}A_2 + \frac{\widehat{\tau}_1^q}{q}A_4 - \frac{\widehat{\tau}_1^{q_s^*}}{q_s^*}B_2 + \frac{\sigma^p}{p}A_1 + \frac{\sigma^q}{q}A_3 - \frac{\sigma^{q_s^*}}{q_s^*}B_1.\end{aligned}$$

Together with (3.21), we obtain  $\psi_u(\sigma, \widehat{\tau}_1) \leq 0$ ,  $\psi_u(\widehat{\sigma}_1, \tau) \leq 0$ , for all  $\sigma \in [0, \widehat{\sigma}_1]$  and  $\tau \in [0, \widehat{\tau}_1]$ , which is absurd. Therefore,  $(\sigma_u, \tau_u) \notin [0, \widehat{\sigma}_1] \times \{\widehat{\tau}_1\}$  and  $(\sigma_u, \tau_u) \notin \{0, \widehat{\sigma}_1\} \times [0, \widehat{\tau}_1]$ .

In conclusion, we get  $(\sigma_u, \tau_u) \in (0, \widehat{\sigma}_1) \times (0, \widehat{\tau}_1)$ . Hence,  $\sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda$ . So, combining (3.21), (3.22) with (3.23), we have that

$$\begin{aligned}c_\lambda &\geq I_\lambda(\sigma_u u^+ + \tau_u u^-) + \frac{1}{p}\sigma_u^p A_1 + \frac{1}{q}\sigma_u^q A_3 - \frac{\sigma_u^{q_s^*}}{q_s^*}B_1 + \frac{1}{p}\tau_u^p A_2 + \frac{1}{q}\tau_u^q A_4 - \frac{\tau_u^{q_s^*}}{q_s^*}B_2 \\ &> I_\lambda(\sigma_u u^+ + \tau_u u^-) \geq c_\lambda.\end{aligned}$$

Therefore, we have a contradiction.

**Case 2:**  $B_2 = 0$ . In this case, we can maximize in  $[0, \widehat{\sigma}_1] \times [0, \infty)$ . Indeed, it is possible to show that there exists  $\widehat{\tau}_0 \in [0, \infty]$  such that  $I_\lambda(\sigma u^+ + \tau u^-) < 0$  for all  $(\sigma, \tau) \in [0, \widehat{\sigma}_1] \times [\widehat{\tau}_0, \infty)$ . Hence, there exists  $(\sigma_u, \tau_u) \in [0, \widehat{\sigma}_1] \times [0, \infty)$  that satisfies  $\psi_u(\sigma_u, \tau_u) = \max_{\sigma \in [0, \widehat{\sigma}_1] \times [0, \infty)} \psi_u(\sigma, \tau)$ .

Following, we prove that  $(\sigma_u, \tau_u) \in (0, \widehat{\sigma}_1) \times (0, \infty)$ .

Indeed, since  $\psi_u(\sigma, 0) \leq \psi_u(\sigma, \tau)$  for  $\sigma \in [0, \widehat{\sigma}_1]$  and  $\tau$  is small enough, we have  $(\sigma_u, \tau_u) \notin [0, \widehat{\sigma}_1] \times \{0\}$ . Analogously, we have  $(\sigma_u, \tau_u) \notin \{0\} \times [0, \infty)$ . On the other hand, for all  $\tau \in [0, \infty)$ , it is obvious that

$$\delta \leq \frac{\widehat{\sigma}_1^p}{p}A_1 + \frac{\widehat{\sigma}_1^q}{q}A_3 - \frac{\widehat{\sigma}_1^{q_s^*}}{q_s^*}B_1 + \frac{\tau^p}{p}A_2 + \frac{\tau^q}{q}A_4.$$

Hence, we have that  $\psi_u(\widehat{\sigma}_1, \tau) \leq 0$  for all  $\tau \in [0, \infty)$ . Thus,  $(\sigma_u, \tau_u) \notin \{\widehat{\sigma}_1\} \times [0, \infty)$ . In summary, we have  $(\sigma_u, \tau_u) \in (0, \widehat{\sigma}_1) \times (0, \infty)$ , namely,  $\sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda$ . Therefore, according to (3.22), we have that

$$\begin{aligned}c_\lambda &\geq I_\lambda(\sigma_u u^+ + \tau_u u^-) + \frac{1}{p}\sigma_u^p A_1 + \frac{1}{q}\sigma_u^q A_3 - \frac{\sigma_u^{q_s^*}}{q_s^*}B_1 + \frac{1}{p}\tau_u^p A_2 + \frac{1}{q}\tau_u^q A_4 \\ &> I_\lambda(\sigma_u u^+ + \tau_u u^-) \geq c_\lambda,\end{aligned}$$

which is a contradiction.

Therefore, from the above discussion, we deduce that  $B_1 = B_2 = 0$ .

**Step 3:** we prove that  $c_\lambda$  is achieved. Since  $u^\pm \neq 0$ , by Lemma 3.1, there exist  $\sigma_u, \tau_u > 0$  such that

$$\widetilde{u} = \sigma_u u^+ + \tau_u u^- \in \mathcal{M}_\lambda.$$

Furthermore,  $B_1 = B_2 = 0$  and Fatou's Lemma implies  $\langle I'_\lambda(u), u^\pm \rangle \leq 0$ . By Lemma 3.2, we obtain  $\sigma_u, \tau_u \leq 1$ . Since  $u_n \in \mathcal{M}_\lambda$ , then according to Lemma 3.1 there holds

$$I_\lambda(\sigma_u u_n^+ + \tau_u u_n^-) \leq I_\lambda(u_n^+ + u_n^-) = I_\lambda(u_n).$$

Due to  $\sigma_u, \tau_u \leq 1$ , arguing as Lemma 2.2, one has  $\|\sigma_u u^+ + \tau_u u^-\|_{V,p}^p \leq \|u\|_{V,p}^p$ . Then by (f<sub>4</sub>), Fatou's Lemma and a straightforward calculation, we deduce that

$$\begin{aligned} c_\lambda &\leq I_\lambda(\bar{u}) - \frac{1}{q} \langle I'_\lambda(\bar{u}), \bar{u} \rangle \\ &= \left(\frac{1}{p} - \frac{1}{q}\right) \|\bar{u}\|_{V,p}^p + \lambda \int_{\mathbb{R}^N} \left[ \frac{1}{q} f(\bar{u}) \bar{u} - F(\bar{u}) \right] dx + \left(\frac{1}{q} - \frac{1}{q_s^*}\right) \int_{\mathbb{R}^N} |\bar{u}|^{q_s^*} dx \\ &= \left(\frac{1}{p} - \frac{1}{q}\right) \|\sigma_u u^+ + \tau_u u^-\|_{V,p}^p + \lambda \int_{\mathbb{R}^N} \left[ \frac{1}{q} f(\sigma_u u^+) \sigma_u u^+ - F(\sigma_u u^+) \right] dx \\ &\quad + \lambda \int_{\mathbb{R}^N} \left[ \frac{1}{q} f(\tau_u u^-) \tau_u u^- - F(\tau_u u^-) \right] dx + \left(\frac{1}{q} - \frac{1}{q_s^*}\right) \int_{\mathbb{R}^N} |\sigma_u u^+|^{q_s^*} dx \\ &\quad + \left(\frac{1}{q} - \frac{1}{q_s^*}\right) \int_{\mathbb{R}^N} |\tau_u u^-|^{q_s^*} dx \\ &\leq \left(\frac{1}{p} - \frac{1}{q}\right) \|u\|_{V,p}^p + \lambda \int_{\mathbb{R}^N} \left[ \frac{1}{q} f(u) u - F(u) \right] dx + \left(\frac{1}{q} - \frac{1}{q_s^*}\right) \int_{\mathbb{R}^N} |u|^{q_s^*} dx \\ &\leq \liminf_{n \rightarrow \infty} \left[ I_\lambda(u_n) - \frac{1}{q} \langle I'_\lambda(u_n), u_n \rangle \right] \leq c_\lambda. \end{aligned}$$

Therefore,  $\sigma_u = \tau_u = 1$ , and  $c_\lambda$  is achieved by  $u_\lambda := u^+ + u^- \in \mathcal{M}_\lambda$ . This ends the proof of Lemma 3.5.  $\square$

#### 4. Proof of Theorem 1.1

*Proof of Theorem 1.1.* Since  $u_\lambda \in \mathcal{M}_\lambda$ , we have  $\langle I'_\lambda(u_\lambda), u_\lambda^+ \rangle = \langle I'_\lambda(u_\lambda), u_\lambda^- \rangle = 0$ . By Lemma 3.5, for  $(\sigma, \tau) \in (\mathbb{R}^+ \times \mathbb{R}^+) \setminus (1, 1)$ , we have

$$I_\lambda(\sigma u_\lambda^+ + \tau u_\lambda^-) < I_\lambda(u_\lambda^+ + u_\lambda^-) = c_\lambda. \quad (4.1)$$

Now we prove  $u_\lambda$  is a solution of (1.1). Arguing by contradiction, we assume that  $I'_\lambda(u_\lambda) \neq 0$ , then there exists  $\delta > 0$  and  $\kappa > 0$  such that

$$|I'_\lambda(v)| \geq \kappa, \text{ for all } \|v - u_\lambda\| \leq 3\delta.$$

Define  $D := [1 - \delta_1, 1 + \delta_1] \times [1 - \delta_1, 1 + \delta_1]$  and a map  $g : D \rightarrow X_V$  by

$$g(\sigma, \tau) := \sigma w^+ + \tau w^-,$$

where  $\delta_1 \in (0, \frac{1}{2})$  small enough such that  $\|g(\sigma, \tau) - w\| \leq 3\delta$  for all  $(\sigma, \tau) \in \bar{D}$ . Thus, by virtue of Lemma 3.5, we can see that

$$I(g(1, 1)) = c_\lambda, \quad I(g(\sigma, \tau)) < c_\lambda \text{ for all } (\sigma, \tau) \in D \setminus \{(1, 1)\}.$$

Therefore,

$$\beta := \max_{(\sigma, \tau) \in \partial D} I(g(\sigma, \tau)) < c_\lambda.$$

By using [38, Theorem 2.3] with

$$\mathcal{S}_\delta := \{v \in X : \|v - u_\lambda\| \leq \delta\}$$

and  $c := c_\lambda$ . Then, choosing  $\varepsilon := \min\left\{\frac{c_\lambda - \beta}{4}, \frac{\kappa\delta}{8}\right\}$ , we deduce that there exists a deformation  $\eta \in C([0, 1] \times X_V, X_V)$  such that:

- (i)  $\eta(t, v) = v$  if  $v \notin I^{-1}([c_\lambda - 2\varepsilon, c_\lambda + 2\varepsilon])$ ;
- (ii)  $I_\lambda(\eta(1, v)) \leq c_\lambda - \varepsilon$  for each  $v \in X_V$  with  $\|v - u\| \leq \delta$  and  $I_\lambda(v) \leq c_\lambda + \varepsilon$ ;
- (iii)  $I_\lambda(\eta(1, v)) \leq I_\lambda(v)$  for all  $u \in X_V$ .

By (ii) and (iii) we conclude that

$$\max_{(\sigma, \tau) \in \bar{D}} I_\lambda(\eta(1, g(\sigma, \tau))) < c_\lambda. \quad (4.2)$$

Therefore, to complete the proof of this Lemma, it suffices to prove that

$$\eta(1, g(\bar{D})) \cap \mathcal{M}_\lambda \neq \emptyset. \quad (4.3)$$

Indeed, if (4.3) holds true, then by the definition of  $c_\lambda$  and (4.2), we get a contradiction.

In the following, we will prove (4.3). To this end, for  $(\sigma, \tau) \in \bar{D}$ , let  $\gamma(\sigma, \tau) := \eta(1, g(\sigma, \tau))$  and

$$\begin{aligned} \Psi_0(\sigma, \tau) &:= (\langle I'_\lambda(g(\sigma, \tau)), u_\lambda^+ \rangle, \langle I'_\lambda(g(\sigma, \tau)), u_\lambda^- \rangle) \\ &:= (\langle I'_\lambda(\sigma u_b^+ + \tau u_\lambda^-), u_\lambda^+ \rangle, \langle I'_\lambda(\sigma u_b^+ + \tau u_\lambda^-), u_\lambda^- \rangle) := (\varphi_u^1(\sigma, \tau), \varphi_u^2(\sigma, \tau)) \end{aligned}$$

and

$$\Psi_1(\sigma, \tau) := \left( \frac{1}{\sigma} \langle I'_\lambda(\gamma(\sigma, \tau)), (\gamma(\sigma, \tau))^+ \rangle, \frac{1}{\tau} \langle I'_\lambda(\gamma(\sigma, \tau)), (\gamma(\sigma, \tau))^- \rangle \right).$$

Firstly, let us denote

$$\begin{aligned} A_p &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{p-2} |u_\lambda^+(x) - u_\lambda^+(y)|^2}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u_\lambda^+|^p dx, \\ A_q &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{q-2} |u_\lambda^+(x) - u_\lambda^+(y)|^2}{|x - y|^{N+qs}} dx dy + \int_{\mathbb{R}^N} V(x) |u_\lambda^+|^q dx, \\ B_p &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{p-2} |u_\lambda^-(x) - u_\lambda^-(y)|^2}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} V(x) |u_\lambda^-|^p dx, \\ B_q &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{q-2} |u_\lambda^-(x) - u_\lambda^-(y)|^2}{|x - y|^{N+qs}} dx dy + \int_{\mathbb{R}^N} V(x) |u_\lambda^-|^q dx, \\ C_p &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{p-2} (u_\lambda^-(x) - u_\lambda^-(y))(u_\lambda^+(x) - u_\lambda^+(y))}{|x - y|^{N+ps}} dx dy, \\ C_q &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{q-2} (u_\lambda^-(x) - u_\lambda^-(y))(u_\lambda^+(x) - u_\lambda^+(y))}{|x - y|^{N+qs}} dx dy, \end{aligned}$$

$$\begin{aligned}
D_p &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{p-2} (u_\lambda^+(x) - u_\lambda^+(y))(u_\lambda^-(x) - u_\lambda^-(y))}{|x - y|^{N+ps}} dx dy, \\
D_q &:= \int_{\mathbb{R}^{2N}} \frac{|u_\lambda(x) - u_\lambda(y)|^{q-2} (u_\lambda^+(x) - u_\lambda^+(y))(u_\lambda^-(x) - u_\lambda^-(y))}{|x - y|^{N+qs}} dx dy, \\
a_1 &:= \lambda \int_{\mathbb{R}^N} f'(u_\lambda^+) |u_\lambda^+|^2 dx, & a_2 &:= \lambda \int_{\mathbb{R}^N} f(u_\lambda^+) u_\lambda^+ dx, \\
b_1 &:= \lambda \int_{\mathbb{R}^N} f'(u_\lambda^-) |u_\lambda^-|^2 dx, & b_2 &:= \lambda \int_{\mathbb{R}^N} f(u_\lambda^-) u_\lambda^- dx, \\
c_1 &:= \int_{\mathbb{R}^N} |u_\lambda^+|^{q_s^*} dx, & c_2 &:= \int_{\mathbb{R}^N} |u_\lambda^-|^{q_s^*} dx.
\end{aligned}$$

Clearly,  $C_p = D_p > 0$ ,  $C_q = D_q > 0$ ,  $A_p, A_q, B_p, B_q > 0$  and notice that  $u_\lambda \in \mathcal{M}_\lambda$ , we can see that

$$A_p + C_p + A_q + C_q = a_2 + c_1, \quad B_p + D_p + B_q + D_q = b_2 + c_2. \quad (4.4)$$

Moreover,  $(f_4)$  guarantees

$$a_1 > (q - 1)a_2, \quad b_1 > (q - 1)b_2. \quad (4.5)$$

Then by direct computation, we have

$$\begin{aligned}
\frac{\partial \varphi_u^1}{\partial \sigma}(1, 1) &= (p - 1)A_p + (q - 1)A_q - a_1 - (q_s^* - 1)c_1 < 0, \\
\frac{\partial \varphi_u^2}{\partial \tau}(1, 1) &= (p - 1)B_p + (q - 1)B_q - b_1 - (q_s^* - 1)c_2 < 0
\end{aligned} \quad (4.6)$$

and

$$\frac{\partial \varphi_u^2}{\partial \tau}(1, 1) = \frac{\partial \varphi_u^2}{\partial \sigma}(1, 1) = (p - 1)C_p + (q - 1)C_q = (p - 1)D_p + (q - 1)D_q. \quad (4.7)$$

Let

$$M = \begin{bmatrix} \frac{\varphi_u^1(\sigma, \tau)}{\partial \sigma} \Big|_{1,1} & \frac{\varphi_u^2(\sigma, \tau)}{\partial \sigma} \Big|_{1,1} \\ \frac{\varphi_u^1(\sigma, \tau)}{\partial \tau} \Big|_{1,1} & \frac{\varphi_u^2(\sigma, \tau)}{\partial \tau} \Big|_{1,1} \end{bmatrix}.$$

So we have

$$\begin{aligned}
\det M &= \left[ (p - 1)A_p + (q - 1)A_q - a_1 - (q_s^* - 1)c_1 \right] \cdot \left[ (p - 1)B_p + (q - 1)B_q - b_1 - (q_s^* - 1)c_2 \right] \\
&\quad - \left[ (p - 1)C_p + (q - 1)C_q \right] \left[ (p - 1)D_p + (q - 1)D_q \right] \\
&> \left[ (q - 1)a_2 + (q_s^* - 1)c_1 - (p - 1)A_p - (q - 1)A_q \right] \cdot \\
&\quad \left[ (q - 1)b_2 + (q_s^* - 1)c_2 - (p - 1)B_p - (q - 1)B_q \right] \\
&\quad - \left[ (p - 1)C_p + (q - 1)C_q \right] \left[ (p - 1)D_p + (q - 1)D_q \right] \\
&= \left[ (q - p)A_p + (q - 1)C_p + (q - 1)C_q(q_s^* - q)c_1 \right] \cdot \\
&\quad \left[ (q - p)B_p + (q - 1)D_p + (q - 1)D_q + (q_s^* - q)c_2 \right] \\
&\quad - \left[ (p - 1)C_p + (q - 1)C_q \right] \left[ (p - 1)D_p + (q - 1)D_q \right] \\
&> 0.
\end{aligned} \quad (4.8)$$

Since  $\Psi_0(\alpha, \beta)$  is a  $C^1$  function and  $(1, 1)$  is the unique isolated zero point of  $\Psi_0$ , by using the degree theory, we deduce that  $\deg(\Psi_0, D, 0) = 1$ . Furthermore, combining (4.2) and (a), we get

$$g(\sigma, \tau) = \gamma(\sigma, \tau) \text{ on } \partial D.$$

Consequently, we deduce that  $\deg(\Psi_1, D, 0) = 1$ . Therefore,  $\Psi_1(\sigma_0, \tau_0) = 0$  for some  $(\sigma_0, \tau_0) \in D$  so that

$$\eta(1, g(\sigma_0, \tau_0)) = \gamma(\sigma_0, \tau_0) \in \mathcal{M}_\lambda,$$

which is contradicted to (4.2). From the above discussions, we deduce that  $u_\lambda$  is a sign-changing solution for the problem (1.1).

Next, we prove that the energy of  $u_b$  is strictly larger than two times the ground state energy.

Similar to proof of Lemma 3.1, there exists  $\lambda_1^* > 0$  such that for all  $\lambda \geq \lambda_1^* > 0$ , there exists  $v \in \mathcal{N}_\lambda$  such that  $I_\lambda(v) = c^* > 0$ . By standard arguments, the critical points of the functional  $I_\lambda$  on  $\mathcal{N}_\lambda$  are critical points of  $I_\lambda$  in  $X_V$ , we obtain  $\langle I'_\lambda(v), v \rangle = 0$ , that is,  $v$  is a ground state solution of (1.1).

According to Theorem 1.1, we know that the problem (1.1) has the least energy sign-changing solution  $u_b$  when  $\lambda \geq \lambda^*$ . Denote  $\Lambda := \max\{\lambda^*, \lambda_1^*\}$ . As Proof of Lemma 3.5, there exist  $\sigma_{u_\lambda^+} > 0$  and  $\tau_{u_\lambda^-} > 0$  such that

$$\sigma_{u_\lambda^+} u_\lambda^+ \in \mathcal{N}_\lambda, \quad \tau_{u_\lambda^-} u_\lambda^- \in \mathcal{N}_\lambda.$$

Furthermore, Lemma 3.2 implies that  $\sigma_{u_\lambda^+}, \tau_{u_\lambda^-} \in (0, 1)$ .

Therefore, in view of Lemma 3.1, we have that

$$2c \leq I_\lambda(\sigma_{u_\lambda^+} u_\lambda^+) + I_\lambda(\tau_{u_\lambda^-} u_\lambda^-) < I_\lambda(\sigma_{u_\lambda^+} u_\lambda^+ + \tau_{u_\lambda^-} u_\lambda^-) < I_\lambda(u_\lambda^+ + u_\lambda^-) = c_\lambda.$$

The proof is complete. □

## 5. Conclusions

This paper considers the least energy sign-changing solution for a class of fractional  $(p, q)$ -Laplacian problems with critical growth in  $\mathbb{R}^N$ . We use constrained variational methods, quantitative deformation lemma and Brouwer degree theory to prove that the above problem has a least energy sign-changing solution  $u_\lambda$  if  $\lambda$  is large enough. Moreover, we show that the energy of  $u_\lambda$  is strictly larger than two times the ground state energy.

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## Conflict of interest

The authors declare that there is no conflict of interest.

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**References**

1. C. O. Alves, V. Ambrosio, T. Isernia, Existence, multiplicity and concentration for a class of fractional  $p&q$  Laplacian problems in  $\mathbb{R}^N$ , *Commun. Pure Appl. Anal.*, **18** (2019), 2009–2045. <https://doi.org/10.3934/cpaa.2019091>
2. V. Ambrosio, Multiple solutions for a fractional  $p$ -Laplacian equation with sign-changing potential, preprint paper, arXiv: 1603.05282, 2016. <https://doi.org/10.48550/arXiv.1603.05282>
3. V. Ambrosio, T. Isernia, Multiplicity and concentration results for some nonlinear Schrödinger equations with the fractional  $p$ -Laplacian, *Discrete Contin. Dyn. Syst.*, **38** (2018), 5835–5881. <https://doi.org/10.3934/dcds.2018254>
4. V. Ambrosio, V. D. Rădulescu, Fractional double-phase patterns: concentration and multiplicity of solutions, *J. Math. Pures Appl.*, **142** (2020), 101–145. <https://doi.org/10.1016/j.matpur.2020.08.011>
5. S. Barile, G. M. Figueiredo, Existence of least energy positive, negative and nodal solutions for a class of  $p&q$ -problems with potentials vanishing at infinity, *J. Math. Anal. Appl.*, **427** (2015), 1205–1233. <https://doi.org/10.1016/j.jmaa.2015.02.086>
6. T. Bartsch, Z. Liu, T. Weth, Sign changing solutions of superlinear Schrödinger equations, *Commun. Part. Diff. Eq.*, **29** (2004), 25–42. <https://doi.org/10.1081/PDE-120028842>
7. T. Bartsch, T. Weth, Three nodal solutions of singularly perturbed elliptic equations on domains without topology, *Ann. Inst. Henri Poincaré, Anal. NonLinéaire*, **22** (2005), 259–281. <https://doi.org/10.1016/j.anihpc.2004.07.005>
8. T. Bartsch, T. Weth, M. Willem, Partial symmetry of least energy nodal solutions to some variational problems, *J. Anal. Math.*, **96** (2005), 1–18. <https://doi.org/10.1007/BF02787822>
9. G. Bonanno, G. Molica Bisci, V. Rădulescu, Quasilinear elliptic non-homogeneous Dirichlet problems through Orlicz-Sobolev spaces, *Nonlinear Anal.*, **75** (2012), 4441–4456. <https://doi.org/10.1016/j.na.2011.12.016>
10. X. Cabré, Y. Sire, Nonlinear equations for fractional Laplacians, I: Regularity, maximum principles, and Hamiltonian estimates, *Ann. Inst. Henri Poincaré, Anal. Non Linéaire*, **31** (2014), 23–53. <https://doi.org/10.1016/j.anihpc.2013.02.001>
11. L. Caffarelli, L. Silvestre, An extension problem related to the fractional Laplacian, *Commun. Part. Diff. Eq.*, **32** (2007), 1245–1260. <https://doi.org/10.1080/03605300600987306>
12. A. D. Castro, J. Cossio, J. M. Neuberger, A sign-changing solution for a superlinear Dirichlet problem, *Rocky Mt. J. Math.*, **27** (1997), 1041–1053. <https://doi.org/10.1216/rmj.1181071858>
13. A. D. Castro, T. Kuusi, G. Palatucci, Local behavior of fractional  $p$ -minimizers, *Ann. Inst. Henri Poincaré, Anal. Non Linéaire*, **33** (2016), 1279–1299. <https://doi.org/10.1016/j.anihpc.2015.04.003>
14. A. D. Castro, T. Kuusi, G. Palatucci, Nonlocal Harnack inequalities, *J. Funct. Anal.*, **267** (2014), 1807–1836. <https://doi.org/10.1016/j.jfa.2014.05.023>
15. X. Chang, Z. Nie, Z. Q. Wang, Sign-changing solutions of fractional  $p$ -laplacian problems, *Adv. Nonlinear Stud.*, **19** (2019), 29–53. <https://doi.org/10.1515/ans-2018-2032>



16. X. Chang, Z. Q. Wang, Nodal and multiple solutions of nonlinear problems involving the fractional Laplacian, *J. Differ. Equ.*, **256** (2014), 2965–2992. <https://doi.org/10.1016/j.jde.2014.01.027>
17. C. Chen, J. Bao, Existence, nonexistence, and multiplicity of solutions for the fractional  $p&q$ -Laplacian equation in  $\mathbb{R}^N$ , *Bound Value Probl.*, **2016** (2016), 153. <https://doi.org/10.1186/s13661-016-0661-0>
18. W. Chen, S. Deng, Existence, nonexistence, and multiplicity of solutions for the fractional  $p&q$ -Laplacian equation in  $\mathbb{R}^N$ , *Nonlinear Anal. Real World Appl.*, **27** (2016), 80–92. <https://doi.org/10.1016/j.nonrwa.2015.07.009>
19. W. Chen, C. Li, Maximum principles for the fractional  $p$ -Laplacian and symmetry of solutions, *Adv. Math.*, **335** (2018), 735–758. <https://doi.org/10.1016/j.aim.2018.07.016>
20. C. D. Filippis, G. Palatucci, Hölder regularity for nonlocal double phase equations, *J. Differ. Equ.*, **267** (2020), 547–586. <https://doi.org/10.1016/j.jde.2019.01.017>
21. P. Felmer, A. Quaas, J. Tan, Positive solutions of the nonlinear Schrödinger equation with the fractional Laplacian, *Proc. Roy. Soc. Edinburgh Sect. A*, **142** (2012), 1237–1262. <https://doi.org/10.1017/S0308210511000746>
22. G. M. Figueiredo, Existence of positive solutions for a class of  $p&q$  elliptic problems with critical growth on  $\mathbb{R}^N$ , *J. Math. Anal. Appl.*, **378** (2011), 507–518. <https://doi.org/10.1016/j.jmaa.2011.02.017>
23. R. F. Gabert, R. S. Rodrigues, Existence of sign-changing solution for a problem involving the fractional Laplacian with critical growth nonlinearities, *Complex Var. Elliptic Equ.*, **65** (2020), 272–292. <https://doi.org/10.1080/17476933.2019.1579208>
24. C. He, G. Li, The regularity of weak solutions to nonlinear scalar field elliptic equations containing  $p&q$ -Laplacians, *Ann. Acad. Sci. Fenn., Math.*, **33** (2008), 337–371.
25. A. Iannizzotto, S. Mosconi, M. Squassina, Global Hölder regularity for the fractional  $p$ -Laplacian, *Rev. Mat. Iberoam.*, **32** (2016), 1353–1392. <https://doi.org/10.4171/RMI/921>
26. T. Isernia, Fractional  $p&q$ -Laplacian problems with potentials vanishing at infinity, *Opusc. Math.*, **40** (2020), 93–110. <https://doi.org/10.7494/OpMath.2020.40.1.93>
27. I. Kuzin, S. Pohozaev, *Entire solutions of semilinear Elliptic equations*, Basel: Birkhäuser, 1995.
28. N. Laskin, Fractional quantum mechanics and Lévy path integrals, *Phys. Lett. A*, **268** (2000), 298–305. [https://doi.org/10.1016/S0375-9601\(00\)00201-2](https://doi.org/10.1016/S0375-9601(00)00201-2)
29. N. Laskin, Fractional Schrödinger equation, *Phys. Rev. E*, **66** (2002), 056108. <https://doi.org/10.1103/PhysRevE.66.056108>
30. G. Li, X. Liang, The existence of nontrivial solutions to nonlinear elliptic equation of  $p&q$ -Laplacian type on  $\mathbb{R}^N$ , *Nonlinear Anal.*, **71** (2009), 2316–2334. <https://doi.org/10.1016/j.na.2009.01.066>
31. C. Miranda, Un’osservazione su un teorema di Brouwer, *Boll Un Mat. Ital.*, **3** (1940), 5–7.
32. D. Mugnai, N. S. Papageorgiou, Wang’s multiplicity result for superlinear  $(p, q)$ -equations without the Ambrosetti-Rabinowitz condition, *Trans. Amer. Math. Soc.*, **366** (2014), 4919–4937. <https://doi.org/10.1090/S0002-9947-2013-06124-7>

33. E. D. Nezza, G. Palatucci, E. Valdinoci, Hitchhiker's guide to the fractional Sobolev spaces, *Bull. Sci. Math.*, **136** (2012), 521–573. <https://doi.org/10.1016/j.bulsci.2011.12.004>
34. G. Palatucci, The Dirichlet problem for the  $p$ -fractional Laplace equation, *Nonlinear Anal.*, **177** (2018), 699–732. <https://doi.org/10.1016/j.na.2018.05.004>
35. P. Pucci, M. Xiang, B. Zhang, Multiple solutions for nonhomogeneous Schrödinger-Kirchhoff type equations involving the fractional  $p$ -Laplacian in  $\mathbb{R}^N$ , *Calc. Var. Partial Differ. Equ.*, **54** (2015), 2785–2806. <https://doi.org/10.1007/s00526-015-0883-5>
36. S. Secchi, Ground state solutions for nonlinear fractional Schrödinger equations in  $\mathbb{R}^N$ , *J. Math. Phys.*, **54** (2013), 031501. <https://doi.org/10.1063/1.4793990>
37. Z. Wang, H. Zhou, Radial sign-changing solution for fractional Schrödinger equation, *Discrete Contin. Dyn. Syst.*, **36** (2016), 499–508. <https://doi.org/10.3934/dcds.2016.36.499>
38. M. Willem, Progress in nonlinear differential equations and their applications, In: *Minimax theorems*, Berlin: Springer, 1997.
39. M. Wu, Z. Yang, A class of  $p$ & $q$ -Laplacian type equation with potentials eigenvalue problem in  $\mathbb{R}^N$ , *Bound Value Probl.*, **2009** (2009), 185319. <https://doi.org/10.1155/2009/185319>
40. J. Zhang, W. Zhang, V. D. Rădulescu, Double phase problems with competing potentials: concentration and multiplication of ground states, *Math. Z.*, **301** (2022), 4037–4078. <https://doi.org/10.1007/s00209-022-03052-1>
41. W. Zhang, J. Zhang, Multiplicity and concentration of positive solutions for fractional unbalanced double-phase problems, *J. Geom. Anal.*, **32** (2022), 235. <https://doi.org/10.1007/s12220-022-00983-3>
42. W. Zhang, J. Zhang, V. D. Rădulescu, Concentrating solutions for singularly perturbed double phase problems with nonlocal reaction, *J. Differ. Equ.*, **347** (2023), 56–103. <https://doi.org/10.1016/j.jde.2022.11.033>
43. W. Zhang, S. Yuan, L. Wen, Existence and concentration of ground-states for fractional Choquard equation with indefinite potential, *Adv. Nonlinear Anal.*, **11** (2022), 1552–1578. <https://doi.org/10.1515/anona-2022-0255>
44. Y. Zhang, X. Tang, V. D. Rădulescu, Concentration of solutions for fractional double-phase problems: critical and supercritical cases, *J. Differ. Equ.*, **302** (2021), 139–184. <https://doi.org/10.1016/j.jde.2021.08.038>



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